

## EIGENVALUE ESTIMATES OF THE BASIC DIRAC OPERATOR ON A RIEMANNIAN FOLIATION

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**Abstract.** On a foliated Riemannian manifold with a transverse spin structure, we give a lower bound of the square for the eigenvalues of the basic Dirac operator in terms of the smallest eigenvalue of the basic Yamabe operator and of the norm of an appropriate endomorphism of the normal bundle  $Q$  of  $\mathcal{F}$ . We study the limiting case.

### 1. INTRODUCTION

The Dirac operator on a Riemannian spin manifold was studied by many authors([2,3,6,8,9,15,16,17]). In particular, the first sharp estimate for the eigenvalues  $\lambda$  of the Dirac operator  $D$  was given by Th. Friedrich ([6]). Using a suitable Riemannian spin connection, he proved the inequality

$$(1.1) \quad \lambda^2 \geq \frac{n}{4(n-1)} \inf_M R$$

on manifolds  $(M^n, g)$  with positive scalar curvature  $R > 0$ . The inequality (1.1) has been improved in several directions by many authors ([8,9,10,11,12,15,16,17]).

Recently, S. D. Jung([10]) proved the lower bound for the eigenvalues  $\lambda$  of the basic Dirac operator  $D_b$  on a foliated Riemannian manifold with a transverse spin structure, which is introduced by J.Brüning and F.W.Kamber([4]). Namely, let  $(M, g_M, \mathcal{F})$  be a Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  and

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a bundle-like metric  $g_M$ . Let  $S(\mathcal{F})$  be a foliated spinor bundle on  $M$ . Then the eigenvalue  $\lambda$  of the basic Dirac operator  $D_b$  satisfies the inequality

$$(1.2) \quad \lambda^2 \geq \frac{q}{4(q-1)} \inf_M (\sigma^\nabla + |\kappa|^2),$$

where  $q = \text{codim} \mathcal{F}$ ,  $\sigma^\nabla$  is the transversal scalar curvature and  $\kappa$  is the mean curvature form of  $\mathcal{F}$ . In the limiting case, it is proved that  $\mathcal{F}$  is minimal, transversally Einsteinian with constant transversal scalar curvature  $\sigma^\nabla$ .

In [12], S. D. Jung et al. improved (1.2) by the smallest eigenvalue of the basic Yamabe operator  $Y_b$ , which is defined by

$$(1.3) \quad Y_b = 4 \frac{q-1}{q-2} \Delta_B + \sigma^\nabla,$$

where  $\Delta_B$  is the basic Laplacian acting on basic functions. In fact, the eigenvalue  $\lambda$  of  $D_b$  satisfies

$$(1.4) \quad \lambda^2 \geq \frac{q}{4(q-1)} (\mu_1 + \inf_M |\kappa|^2),$$

where  $\mu_1$  is the smallest eigenvalue of the basic Yamabe operator  $Y_b$ .

In this paper, we prove a new estimate for the eigenvalues of the basic Dirac operator by using the appropriate endomorphism on the normal bundle  $Q$  of  $\mathcal{F}$ . This paper is corresponding to Hijazi's one([9]) for the foliations.

The paper is organized as follows. In Section 2, we review the known facts on the foliated Riemannian manifold. In Section 3, we study the basic properties of the transversal Dirac operators of transversally conformally related metrics. In Section 4, we introduce the new connection with an appropriate endomorphism on  $Q$  and estimate the lower bound for the eigenvalues of the basic Dirac operator. Namely, let  $(M, g_M, \mathcal{F})$  be a compact Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  of codimension  $q \geq 3$  and a bundle-like metric  $g_M$  such that  $\kappa$  is basic-harmonic. Then any eigenvalue  $\lambda$  of the basic Dirac operator corresponding to the eigenspinor  $\Psi \in \Gamma S(\mathcal{F})$  satisfies

$$(1.5) \quad \lambda^2 \geq \frac{1}{4} (\mu_1 + \inf_M |\kappa|^2) + \inf_M |l_\Psi|^2,$$

where  $\mu_1$  is the first eigenvalue of the basic Yamabe operator  $Y_b$  of  $\mathcal{F}$  and  $l_\Psi$  is an isomorphism on the normal bundle  $Q$ , which is defined in Section 4. In Section 5, we prove, in the limiting case that  $\mathcal{F}$  is a minimal foliation.

Throughout this paper, we consider the bundle-like metric  $\tilde{g}_M$  for  $(M, \mathcal{F})$  such that the mean curvature form of  $\mathcal{F}$  is basic-harmonic. The existence of the bundle-like metric  $g_M$  for  $(M, \mathcal{F})$  such that  $\kappa$  is basic, i.e.  $\kappa \in \Omega_B^1(\mathcal{F})$ , is proved in

[5]. In [19, 20], it is proved that for any bundle-like metric  $g_M$  with  $\kappa \in \Omega_B^1(\mathcal{F})$  there exists another bundle-like metric  $\tilde{g}_M$  for which the mean curvature form is basic-harmonic.

## 2. PRELIMINARIES AND KNOWN FACTS

Let  $(M, g_M, \mathcal{F})$  be a  $(p+q)$ -dimensional Riemannian manifold with a foliation  $\mathcal{F}$  of codimension  $q$  and a bundle-like metric  $g_M$  with respect to  $\mathcal{F}$ .

We recall the exact sequence

$$(2.1) \quad 0 \rightarrow L \rightarrow TM \xrightarrow{\pi} Q \rightarrow 0$$

determined by the tangent bundle  $L$  and the normal bundle  $Q = TM/L$  of  $\mathcal{F}$  ([14]). The assumption of  $g_M$  to be a bundle-like metric means that the induced metric  $g_Q$  on the normal bundle  $Q \cong L^\perp$  satisfies the holonomy invariance condition  $\overset{\circ}{\nabla} g_Q = 0$ , where  $\overset{\circ}{\nabla}$  is the Bott connection in  $Q$ .

For a distinguished chart  $\mathcal{U} \subset M$  the leaves of  $\mathcal{F}$  in  $\mathcal{U}$  are given as the fibers of a Riemannian submersion  $f : \mathcal{U} \rightarrow \mathcal{V} \subset N$  onto an open subset  $\mathcal{V}$  of a model Riemannian manifold  $N$ .

For overlapping charts  $U_\alpha \cap U_\beta$ , the corresponding local transition functions  $\gamma_{\alpha\beta} = f_\alpha \circ f_\beta^{-1}$  on  $N$  are isometries. Further, we denote by  $\nabla$  the canonical connection of the normal bundle  $Q$  of  $\mathcal{F}$ . To show that such connections exist, consider a Riemannian metric  $g_M$  on  $M$ . Then  $TM$  splits orthogonally as

$$TM = L \oplus L^\perp.$$

This means that there is a bundle map  $\sigma : Q \rightarrow L^\perp$  splitting the exact sequence (2.1), i.e., satisfying  $\pi \circ \sigma = \text{identity}$ . This metric  $g_M$  on  $TM$  is then a direct sum

$$g_M = g_L \oplus g_{L^\perp}.$$

With  $g_Q = \sigma^* g_{L^\perp}$ , the splitting map  $\sigma : (Q, g_Q) \rightarrow (L^\perp, g_{L^\perp})$  is a metric isomorphism. Then the adapted connection  $\nabla$  in  $Q$  is defined by ([14])

$$(2.2) \quad \begin{cases} \nabla_X s = \overset{\circ}{\nabla}_X s = \pi([X, Y_s]) & \text{for } X \in \Gamma L, \\ \nabla_X s = \pi(\nabla_X^M Y_s) & \text{for } X \in \Gamma L^\perp, \end{cases}$$

where  $s \in \Gamma Q$  and  $Y_s \in \Gamma L^\perp$  corresponding to  $s$  under the canonical isomorphism  $Q \cong L^\perp$ . The connection  $\nabla$  is metric and torsion free. It corresponds to the Riemannian connection of the model space  $N$ . The curvature  $R^\nabla$  of  $\nabla$  is defined by

$$R^\nabla(X, Y) = \nabla_X \nabla_Y - \nabla_Y \nabla_X - \nabla_{[X, Y]} \quad \text{for } X, Y \in TM.$$

Its curvature  $R^\nabla$  coincides with  $\overset{\circ}{R}$  for  $X, Y \in \Gamma L$ , hence  $R^\nabla(X, Y) = 0$  for  $X, Y \in \Gamma L$ . Since  $i(X)R^\nabla = \theta(X)R^\nabla = 0$  ([13, 14, 22]), we can define the (transversal) Ricci curvature  $\rho^\nabla : \Gamma Q \rightarrow \Gamma Q$  and the (transversal) scalar curvature  $\sigma^\nabla$  of  $\mathcal{F}$  by

$$\rho^\nabla(s) = \sum_a R^\nabla(s, E_a)E_a, \quad \sigma^\nabla = \sum_a g_Q(\rho^\nabla(E_a), E_a),$$

where  $\{E_a\}_{a=1, \dots, q}$  is an orthonormal basis of  $Q$ .  $\mathcal{F}$  is said to be (transversally) *Einsteinian* if the model space  $N$  is Einsteinian, that is,

$$(2.3) \quad \rho^\nabla = \frac{1}{q} \sigma^\nabla \cdot id$$

with constant transversal scalar curvature  $\sigma^\nabla$ . The *mean curvature form*  $\kappa$  for  $L$  is given by

$$(2.4) \quad \kappa(X) = g_Q \left( \sum_i \pi(\nabla_{E_i}^M E_i), X \right) \quad \text{for } X \in \Gamma Q,$$

where  $\{E_i\}_{i=1, \dots, p}$  is an orthonormal basis of  $L$ . The foliation  $\mathcal{F}$  is said to be *minimal* (or *harmonic*) if  $\kappa = 0$ .

Let  $\Omega_B^r(\mathcal{F})$  be the space of all *basic r-forms*, i.e.,

$$\Omega_B^r(\mathcal{F}) = \{ \phi \in \Omega^r(M) \mid i(X)\phi = 0, \theta(X)\phi = 0, \text{ for } X \in \Gamma L \}.$$

The foliation  $\mathcal{F}$  is said to be *isoparametric* if  $\kappa \in \Omega_B^1(\mathcal{F})$ . We already know that  $\kappa$  is closed, i.e.,  $d\kappa = 0$  if  $\mathcal{F}$  is isoparametric ([22]). Since the exterior derivative preserves the basic forms (that is,  $\theta(X)d\phi = 0$  and  $i(X)d\phi = 0$  for  $\phi \in \Omega_B^r(\mathcal{F})$ ), the restriction  $d_B = d|_{\Omega_B^*(\mathcal{F})}$  is well defined. Let  $\delta_B$  be the formal adjoint operator of  $d_B$ . Then it is well-known([1,10]) that

$$(2.5) \quad d_B = \sum_a \theta_a \wedge \nabla_{E_a}, \quad \delta_B = - \sum_a i(E_a) \nabla_{E_a} + i(\kappa_B),$$

where  $\kappa_B$  is the basic component of  $\kappa$ ,  $\{E_a\}$  is a local orthonormal basic frame in  $Q$  and  $\{\theta_a\}$  is its  $g_Q$ -dual 1-form.

The *basic Laplacian* acting on  $\Omega_B^*(\mathcal{F})$  is defined by

$$(2.6) \quad \Delta_B = d_B \delta_B + \delta_B d_B.$$

If  $\mathcal{F}$  is the foliation by points of  $M$ , the basic Laplacian is the ordinary Laplacian. In the more general case, the basic Laplacian and its spectrum provide information about the transverse geometry of  $(M, \mathcal{F})$  ([21]).

### 3. TRANSVERSAL DIRAC OPERATORS OF TRANSVERSALLY CONFORMALLY RELATED METRICS

Let  $(M, g_M, \mathcal{F}, S(\mathcal{F}))$  be a compact Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  of codimension  $q$ , a bundle-like metric  $g_M$  with respect to  $\mathcal{F}$  and a foliated spinor bundle  $S(\mathcal{F})$  ([4,7,10]). Then the transversal Dirac operator  $D_{tr}$  is locally defined by

$$(3.1) \quad D_{tr}\Psi = \sum_a E_a \cdot \nabla_{E_a} \Psi - \frac{1}{2} \kappa \cdot \Psi \quad \text{for } \Psi \in \Gamma S(\mathcal{F}),$$

where  $\{E_a\}$  is a local orthonormal basic frame of  $Q$ . We define the subspace  $\Gamma_B(S(\mathcal{F}))$  of *basic* or *holonomy invariant* sections of  $S(\mathcal{F})$  by

$$\Gamma_B(S(\mathcal{F})) = \{\Psi \in \Gamma S(\mathcal{F}) \mid \nabla_X \Psi = 0 \quad \text{for } X \in \Gamma L\}.$$

Trivially, we see that  $D_{tr}$  leaves  $\Gamma_B(S(\mathcal{F}))$  invariant if and only if the foliation  $\mathcal{F}$  is isoparametric, i.e.,  $\kappa \in \Omega_B^1(\mathcal{F})$ . Let  $D_b = D_{tr}|_{\Gamma_B(S(\mathcal{F}))} : \Gamma_B(S(\mathcal{F})) \rightarrow \Gamma_B(S(\mathcal{F}))$ . This operator  $D_b$  is called the *basic Dirac operator* on (smooth) basic sections. On an isoparametric transverse spin foliation  $\mathcal{F}$  with  $\delta\kappa = 0$ , it is well-known ([7,10]) that

$$(3.2) \quad D_{tr}^2 \Psi = \nabla_{tr}^* \nabla_{tr} \Psi + \frac{1}{4} K^\sigma \Psi,$$

where  $K^\sigma = \sigma^\nabla + |\kappa|^2$  and

$$(3.3) \quad \nabla_{tr}^* \nabla_{tr} \Psi = - \sum_a \nabla_{E_a, E_a}^2 \Psi + \nabla_\kappa \Psi.$$

The operator  $\nabla_{tr}^* \nabla_{tr}$  is non-negative and formally self-adjoint ([10]) such that

$$(3.4) \quad \int_M \langle \nabla_{tr}^* \nabla_{tr} \Phi, \Psi \rangle = \int_M \langle \nabla_{tr} \Phi, \nabla_{tr} \Psi \rangle$$

for all  $\Phi, \Psi \in \Gamma S(\mathcal{F})$ . Moreover, the curvature transform  $R^S$  on  $S(\mathcal{F})$  is given ([10,18]) as

$$(3.5) \quad R^S(X, Y)\Psi = \frac{1}{4} \sum_{a,b} g_Q(R^\nabla(X, Y)E_a, E_b) E_a \cdot E_b \cdot \Psi \quad \text{for } X, Y \in \Gamma TM.$$

Then we have the following lemma.

**Lemma 3.1.** ([10]) *On the foliated spinor bundle  $S(\mathcal{F})$ , we have the following equations*

$$(3.6) \quad \sum_{a < b} E_a \cdot E_b \cdot R^S(E_a, E_b)\Psi = \frac{1}{4} \sigma^\nabla \Psi,$$

$$(3.7) \quad \sum_a \bar{E}_a \cdot R^S(X, E_a)\Psi = -\frac{1}{2}\rho^\nabla(X) \cdot \Psi \quad \text{for } X \in \Gamma Q.$$

Now, we consider, for any real basic function  $u$  on  $M$ , the transversally conformal metric  $\bar{g}_Q = e^{2u}g_Q$ . Let  $\bar{P}_{so}(\mathcal{F})$  be the principal bundle of  $\bar{g}_Q$ -orthogonal frames. Locally, the section  $\bar{s}$  of  $\bar{P}_{so}(\mathcal{F})$  corresponding a section  $s = (E_1, \dots, E_q)$  of  $P_{so}(\mathcal{F})$  is  $\bar{s} = (\bar{E}_1, \dots, \bar{E}_q)$ , where  $\bar{E}_a = e^{-u}E_a$  ( $a = 1, \dots, q$ ). This isometry will be denoted by  $I_u$ . Thanks to the isomorphism  $I_u$  one can define a transverse spin structure  $\bar{P}_{spin}(\mathcal{F})$  on  $\mathcal{F}$  in such a way that the diagram

$$\begin{array}{ccc} P_{spin}(\mathcal{F}) & \xrightarrow{\bar{I}_u} & \bar{P}_{spin}(\mathcal{F}) \\ \downarrow & & \downarrow \\ P_{so}(\mathcal{F}) & \xrightarrow{I_u} & \bar{P}_{so}(\mathcal{F}) \end{array}$$

commutes.

Let  $\bar{S}(\mathcal{F})$  be the foliated spinor bundle associated with  $\bar{P}_{spin}(\mathcal{F})$ . For any section  $\Psi$  of  $S(\mathcal{F})$ , we write  $\bar{\Psi} \equiv I_u\Psi$ . If  $\langle \cdot, \cdot \rangle_{g_Q}$  and  $\langle \cdot, \cdot \rangle_{\bar{g}_Q}$  denote respectively the natural Hermitian metrics on  $S(\mathcal{F})$  and  $\bar{S}(\mathcal{F})$ , then for any  $\Phi, \Psi \in \Gamma S(\mathcal{F})$

$$(3.8) \quad \langle \Phi, \Psi \rangle_{g_Q} = \langle \bar{\Phi}, \bar{\Psi} \rangle_{\bar{g}_Q},$$

and the Clifford multiplication in  $\bar{S}(\mathcal{F})$  is given by

$$(3.9) \quad \bar{X} \cdot \bar{\Psi} = \overline{X \cdot \Psi} \quad \text{for } X \in \Gamma Q.$$

Let  $\bar{\nabla}$  be the metric and torsion free connection corresponding to  $\bar{g}_Q$ . Then we have ([12]) that for  $X, Y \in \Gamma TM$ ,

$$(3.10) \quad \bar{\nabla}_X \pi(Y) = \nabla_X \pi(Y) + X(u)\pi(Y) + Y(u)\pi(X) - g_Q(\pi(X), \pi(Y))grad_\nabla(u),$$

where  $grad_\nabla(u) = \sum_a E_a(u)E_a$  is a transversal gradient of  $u$  and  $X(u)$  is the Lie derivative of the function  $u$  in the direction of  $X$ . The formula (3.10) follows from that  $\bar{\nabla}$  is the metric and torsion free connection with respect to  $\bar{g}_Q$ . From (3.10), we have the following proposition.

**Proposition 3.2.** ([12]) *The connections  $\nabla$  and  $\bar{\nabla}$  acting respectively on the sections of  $S(\mathcal{F})$  and  $\bar{S}(\mathcal{F})$ , are related, for any vector field  $X$  and any spinor field  $\Psi$  by*

$$(3.11) \quad \bar{\nabla}_X \bar{\Psi} = \overline{\nabla_X \Psi} - \frac{1}{2}\pi(X) \cdot grad_\nabla(u) \cdot \bar{\Psi} - \frac{1}{2}g_Q(grad_\nabla(u), \pi(X))\bar{\Psi}.$$

Let  $\bar{D}_{tr}$  be the transversal Dirac operator associated with the metric  $\bar{g}_Q = e^{2u}g_Q$  and acting on the sections of the foliated spinor bundle  $\bar{S}(\mathcal{F})$ . Let  $\{E_a\}$  be a local frame of  $P_{so}(\mathcal{F})$  and  $\{\bar{E}_a\}$  a local frame of  $\bar{P}_{so}(\mathcal{F})$ . Locally,  $\bar{D}_{tr}$  is expressed by

$$(3.12) \quad \bar{D}_{tr}\bar{\Psi} = \sum_a \bar{E}_a \cdot \bar{\nabla}_{\bar{E}_a}\bar{\Psi} - \frac{1}{2}\kappa_{\bar{g}} \cdot \bar{\Psi},$$

where  $\kappa_{\bar{g}}$  is the mean curvature form associated with  $\bar{g}_Q$ , which satisfies  $\kappa_{\bar{g}} = e^{-2u}\kappa$ .

**Proposition 3.3.** ([12]) *Let  $\mathcal{F}$  be the transverse spin foliation of codimension  $q$ . Then the transverse Dirac operators  $D_{tr}$  and  $\bar{D}_{tr}$  satisfy*

$$(3.13) \quad \bar{D}_{tr}(e^{-\frac{q-1}{2}u}\bar{\Psi}) = e^{-\frac{q+1}{2}u}\overline{D_{tr}\Psi}$$

for any spinor field  $\Psi \in S(\mathcal{F})$ .

From Proposition 3.3, if  $D_{tr}\Psi = 0$ , then  $\bar{D}_{tr}\bar{\Phi} = 0$ , where  $\bar{\Phi} = e^{-\frac{q-1}{2}u}\Psi$ , and conversely. Hence the dimension of the space of the foliated harmonic spinors is a transversally conformal invariant.

Assume the mean curvature form  $\kappa$  of  $\mathcal{F}$  is basic- harmonic, i.e.,  $\kappa \in \Omega_B^1(\mathcal{F})$  and  $\delta_B\kappa = 0$ . Then we have the Lichnerowicz type formula ([12])

$$(3.14) \quad \bar{D}_{tr}^2\bar{\Psi} = \bar{\nabla}_{tr}^*\bar{\nabla}_{tr}\bar{\Psi} + \frac{1}{4}K_{\sigma}^{\bar{\nabla}}\bar{\Psi},$$

where

$$(3.15) \quad \bar{\nabla}_{tr}^*\bar{\nabla}_{tr}\bar{\Psi} = -\sum_a \bar{\nabla}_{\bar{E}_a}\bar{\nabla}_{\bar{E}_a}\bar{\Psi} + \bar{\nabla}_{\sum \bar{\nabla}_{\bar{E}_a}\bar{E}_a}\bar{\Psi} + \bar{\nabla}_{\kappa_{\bar{g}}}\bar{\Psi},$$

$$(3.16) \quad K_{\sigma}^{\bar{\nabla}} = \sigma^{\bar{\nabla}} + |\bar{\kappa}|^2 + 2(q-2)\kappa_{\bar{g}}(u).$$

Moreover, it is well-known[12] that for all  $\bar{\Phi}, \bar{\Psi} \in S(\mathcal{F})$

$$(3.17) \quad \int_M \langle \bar{\nabla}_{tr}^*\bar{\nabla}_{tr}\bar{\Psi}, \bar{\Phi} \rangle_{\bar{g}_Q} = \int_M \langle \bar{\nabla}_{tr}\bar{\Psi}, \bar{\nabla}_{tr}\bar{\Phi} \rangle_{\bar{g}_Q},$$

where  $\langle \bar{\nabla}_{tr}\bar{\Psi}, \bar{\nabla}_{tr}\bar{\Phi} \rangle_{\bar{g}_Q} = \sum_a \langle \bar{\nabla}_{\bar{E}_a}\bar{\Psi}, \bar{\nabla}_{\bar{E}_a}\bar{\Phi} \rangle_{\bar{g}_Q}$ .

#### 4. EIGENVALUE ESTIMATE OF THE BASIC DIRAC OPERATOR

Let  $(M, \tilde{g}_M, \mathcal{F}, S(\mathcal{F}))$  be a Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  of codimension  $q \geq 2$  and a bundle-like metric  $\tilde{g}_M$  such that  $\kappa$  is basic-harmonic.

Let  $l$  be a linear symmetric endomorphism of  $Q$ . For any tangent vector field  $X$  and any spinor field  $\Psi$ , we define the modified connection  $\overset{l}{\nabla}$  on  $S(\mathcal{F})$  by

$$(4.1) \quad \overset{l}{\nabla}_X \Psi = \nabla_X \Psi + l(\pi(X)) \cdot \Psi.$$

It is easy to see that the connection  $\overset{l}{\nabla}$  is metrical. By using the transversal divergence theorem ([23]), we have the following equation

$$(4.2) \quad \int_M \langle \overset{l}{\nabla}_{tr}^* \overset{l}{\nabla}_{tr} \Phi, \Psi \rangle = \int_M \langle \overset{l}{\nabla}_{tr} \Phi, \overset{l}{\nabla}_{tr} \Psi \rangle$$

for all  $\Phi, \Psi \in \Gamma S(\mathcal{F})$ . By a direct calculation, we have from (4.1) that for any linear symmetric endomorphism  $l$  of  $Q$

$$(4.3) \quad |\overset{l}{\nabla}_{tr} \Psi|^2 = |\nabla_{tr} \Psi|^2 - 2Re \sum_a \langle l(E_a) \cdot \nabla_{E_a} \Psi, \Psi \rangle + |l|^2 |\Psi|^2.$$

We now show that for an appropriate choice of the symmetric endomorphism  $l$ , one gets a sharp estimate for the first eigenvalue of the basic Dirac operator on compact Riemannian manifolds. For this, we need the following: On the complement of the set of zeros of a spinor field  $\Psi \in \Gamma_B S(\mathcal{F})$ , we define for any tangent vector fields  $X$  and  $Y$ , the symmetric bilinear tensor  $F_\Psi$  (see [9]) by

$$(4.4) \quad F_\Psi(X, Y) = \frac{1}{2} Re \langle \pi(X) \cdot \nabla_Y \Psi + \pi(Y) \cdot \nabla_X \Psi, \Psi / |\Psi|^2 \rangle.$$

Since  $\langle \kappa \cdot \Psi, \Psi \rangle$  is pure imaginary, we have

$$(4.5) \quad \text{tr} F_\Psi := \sum_a F_\Psi(E_a, E_a) = Re \langle D_b \Psi, \Psi / |\Psi|^2 \rangle.$$

Let  $l_\Psi$  be the symmetric endomorphism associated with  $F_\Psi$ . From (4.5), if  $D_b \Psi = \lambda \Psi$ , then

$$(4.6) \quad \text{tr} l_\Psi = \lambda.$$

On the other hand, since  $l$  is a linear symmetric endomorphism, we have

$$(4.7) \quad Re \sum_a \langle l(E_a) \cdot \nabla_{E_a} \Psi, \Psi / |\Psi|^2 \rangle = \sum_a \langle l(E_a), l_\Psi(E_a) \rangle = \langle l, l_\Psi \rangle.$$

From (4.3) and (4.7), we have the following equation;

$$(4.8) \quad |\overset{l}{\nabla}_{tr} \Psi|^2 = |\nabla_{tr} \Psi|^2 - f(l) |\Psi|^2,$$



where  $f(l) = 2 < l, l_\Psi > -|l|^2$ . Note that  $f(l)$  has maximum value  $|l_\Psi|^2$  at  $l = l_\Psi$  because  $l$  is a linear endomorphism. From (3.2), (3.4) and (4.8), we have that for any eigenspinor  $\Psi$  corresponding to an eigenvalue  $\lambda$

$$(4.9) \quad \int_M |\nabla_{tr}^{l_\Psi} \Psi|^2 = \int_M \left\{ \lambda^2 - \left( \frac{1}{4} K^\sigma + |l_\Psi|^2 \right) \right\} |\Psi|^2.$$

Hence we have the following theorem.

**Theorem 4.1.** *Let  $(M, \tilde{g}_M, \mathcal{F})$  be a compact Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  and a bundle-like metric  $\tilde{g}_M$ . Then any eigenvalue  $\lambda$  of the basic Dirac operator  $D_b$  corresponding to the eigenspinor  $\Psi \in \Gamma S(\mathcal{F})$  satisfies*

$$(4.10) \quad \lambda^2 \geq \inf_M \left( \frac{1}{4} K^\sigma + |l_\Psi|^2 \right),$$

where  $K^\sigma = \sigma^\nabla + |\kappa|^2$ .

Moreover, by the Cauchy-Schwarz inequality we have

$$(4.11) \quad \frac{1}{q} (\text{tr} l_\Psi)^2 \leq |l_\Psi|^2.$$

Hence we have the following corollary.

**Corollary 4.2.** (cf. [10]) *Under the same conditions as in Theorem 4.1, one has*

$$(4.12) \quad \lambda^2 \geq \frac{q}{4(q-1)} \inf_M K^\sigma.$$

On the other hand, we obtain that on the spinor bundle  $\bar{S}(\mathcal{F})$  associated with the metric  $\bar{g}_Q = e^{2u} g_Q$

$$(4.13) \quad F_{\bar{\Phi}}(\bar{X}, \bar{Y}) = e^{-u} F_\Phi(X, Y) = e^{-u} F_\Psi(X, Y), \quad X, Y \in \Gamma Q,$$

where  $\bar{\Phi} = e^{-\frac{q-1}{2}u} \Psi$ . From (4.13), we have the following identity

$$(4.14) \quad l_{\bar{\Phi}} = e^{-u} l_\Phi = e^{-u} l_\Psi.$$

From (4.14), we have the following proposition (see [9] for the details).

**Proposition 4.3.** *The following relations hold:*

$$(4.15) \quad |l_{\bar{\Phi}}|_{\bar{g}_Q}^2 = e^{-2u} |l_\Phi|^2 = e^{-2u} |l_\Psi|^2,$$

where  $\bar{\Phi} = e^{-\frac{q-1}{2}u} \Psi$ .

Now, we introduce a connection  $\bar{\nabla}^l$  on  $\bar{S}(\mathcal{F})$ , as

$$(4.16) \quad \bar{\nabla}_X^l \bar{\Psi} = \bar{\nabla}_X \bar{\Psi} + l(\pi(X)) \cdot \bar{\Psi} \quad \text{for } X \in TM,$$

where  $l$  is a linear symmetric endomorphism on  $Q$ . Trivially, the connection  $\bar{\nabla}^l$  is a metric connection and  $\bar{\nabla}_{tr}^{*l} \bar{\nabla}_{tr}^l$  is positive definite. By using the connection  $\bar{\nabla}^l$ , we can obtain the following equation

$$(4.17) \quad |\bar{\nabla}_{tr}^{l_{\bar{\Phi}}} \bar{\Phi}|_{\bar{g}_Q}^2 = |\bar{\nabla}_{tr} \bar{\Phi}|^2 - |l_{\bar{\Phi}}|_{\bar{g}_Q}^2 |\bar{\Phi}|_{\bar{g}_Q}^2.$$

Hence we have from (3.14), (3.17) and (4.17)

$$(4.18) \quad \int_M |\bar{\nabla}_{tr}^{l_{\bar{\Phi}}} \bar{\Phi}|_{\bar{g}_Q}^2 = \int_M \left\{ \langle \bar{D}_{tr}^2 \bar{\Phi}, \bar{\Phi} \rangle_{\bar{g}_Q} - \frac{1}{4} \langle K_{\sigma}^{\bar{\nabla}} \bar{\Phi}, \bar{\Phi} \rangle_{\bar{g}_Q} - |l_{\bar{\Phi}}|_{\bar{g}_Q}^2 |\bar{\Phi}|_{\bar{g}_Q}^2 \right\}.$$

Let  $D_b \Psi = \lambda \Psi$  ( $\Psi \neq 0$ ). From Proposition 3.3, we have

$$(4.19) \quad \bar{D}_b \bar{\Phi} = \lambda e^{-u} \bar{\Phi},$$

where  $\bar{\Phi} = e^{-\frac{q-1}{2}u} \Psi$ . From (4.15) and (4.18), we have that for any eigenspinor  $\Psi$  corresponding to the eigenvalue  $\lambda$

$$(4.20) \quad \int_M |\bar{\nabla}_{tr}^{l_{\bar{\Phi}}} \bar{\Phi}|_{\bar{g}_Q}^2 = \int_M e^{-2u} \left\{ \lambda^2 - \frac{1}{4} (e^{2u} K_{\sigma}^{\bar{\nabla}} + 4|l_{\Psi}|^2) \right\} |\bar{\Phi}|_{\bar{g}_Q}^2,$$

where  $\bar{\Phi} = e^{-\frac{q-1}{2}u} \Psi$ . Hence we have the following theorem.

**Theorem 4.4.** *Let  $(M, \tilde{g}_M, \mathcal{F})$  be a compact Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  of codimension  $q \geq 2$  and a bundle-like metric  $\tilde{g}_M$ . Assume that  $K_{\sigma}^{\bar{\nabla}} \geq 0$  for some transversally conformal metric  $\bar{g}_Q = e^{2u} g_Q$ . Then any eigenvalue  $\lambda$  of the basic Dirac operator  $D_b$  corresponding to the eigenspinor  $\Psi \in \Gamma S(\mathcal{F})$  satisfies*

$$(4.21) \quad \lambda^2 \geq \frac{1}{4} \sup_u \inf_M (e^{2u} K_{\sigma}^{\bar{\nabla}} + 4|l_{\Psi}|^2),$$

where  $K_{\sigma}^{\bar{\nabla}} = \sigma^{\bar{\nabla}} + |\bar{\kappa}|^2 + 2(q-2)\kappa_{\bar{g}}(u)$ .

From (4.11), we have the following corollary.

**Corollary 4.5.** (cf. [12]) *Under the same assumption as in Theorem 4.4, we have*

$$(4.22) \quad \lambda^2 \geq \frac{q}{4(q-1)} \sup_u \inf_M (e^{2u} K_\sigma^{\bar{\nabla}}).$$

The transversal Ricci curvature  $\rho^{\bar{\nabla}}$  of  $\bar{g}_Q = e^{2u} g_Q$  and the transversal scalar curvature  $\sigma^{\bar{\nabla}}$  of  $\bar{g}_Q$  are related to the transversal Ricci curvature  $\rho^{\nabla}$  of  $g_Q$  and the transversal scalar curvature  $\sigma^{\nabla}$  of  $g_Q$  by the following lemma.

**Lemma 4.6.** ([12]) *On a Riemannian foliation  $\mathcal{F}$ , we have that for any  $X \in Q$ ,*

$$(4.23) \quad e^{2u} \rho^{\bar{\nabla}}(X) = \rho^{\nabla}(X) + (2-q) \nabla_X \text{grad}_{\nabla}(u) + (2-q) |\text{grad}_{\nabla}(u)|^2 X \\ + (q-2) X(u) \text{grad}_{\nabla}(u) + \{\Delta_B u - \kappa(u)\} X.$$

$$(4.24) \quad e^{2u} \sigma^{\bar{\nabla}} = \sigma^{\nabla} + (q-1)(2-q) |\text{grad}_{\nabla}(u)|^2 + 2(q-1) \{\Delta_B u - \kappa(u)\}.$$

From (3.16) and (4.24), we have

$$(4.25) \quad e^{2u} K_\sigma^{\bar{\nabla}} = \sigma^{\nabla} + |\kappa|^2 + 2(q-1) \Delta_B u + (q-1)(2-q) |\text{grad}_{\nabla}(u)|^2 - 2\kappa(u).$$

On the other hand, for  $q \geq 3$ , if we choose the positive function  $h$  by  $u = \frac{2}{q-2} \ln h$ , then we have

$$(4.26) \quad \Delta_B u = \frac{2}{q-2} \{h^{-2} |\text{grad}_{\nabla}(h)|^2 + h^{-1} \Delta_B h\},$$

$$(4.27) \quad |\text{grad}_{\nabla}(u)|^2 = \left(\frac{2}{q-2}\right)^2 h^{-2} |\text{grad}_{\nabla}(h)|^2.$$

Hence we have

$$(4.28) \quad e^{2u} K_\sigma^{\bar{\nabla}} = h^{\frac{4}{q-2}} K_\sigma^{\bar{\nabla}} = h^{-1} Y_b h + |\kappa|^2 - \frac{4}{q-2} h^{-1} \kappa(h),$$

where

$$(4.29) \quad Y_b = 4 \frac{q-1}{q-2} \Delta_B + \sigma^{\nabla}$$

is a *basic Yamabe operator* of  $\mathcal{F}$  ([12]). Now we put  $\mathcal{K}_u = \{u \in \Omega_B^0(\mathcal{F}) | \kappa(u) = 0\}$ . If we choose  $u \in \mathcal{K}_u$ , then  $\kappa(h) = 0 = \kappa(u)$ . From (4.25) and (4.28), we have

$$(4.30) \quad e^{2u} K_\sigma^{\bar{\nabla}} = K^\sigma + 2(q-1) \Delta_B u + (q-1)(2-q) |\text{grad}_{\nabla}(u)|^2 = h^{-1} Y_b h + |\kappa|^2,$$

where  $K^\sigma = \sigma^{\nabla} + |\kappa|^2$ . From (4.21) and (4.30), we have the following theorem.

**Theorem 4.7.** *Let  $(M, \tilde{g}_M, \mathcal{F})$  be a compact Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  of codimension  $q \geq 3$  and a bundle-like metric  $\tilde{g}_M$ . Then, any eigenvalue  $\lambda$  of the basic Dirac operator corresponding to the eigenspinor  $\Psi \in \Gamma S(\mathcal{F})$  satisfies*

$$(4.31) \quad \lambda^2 \geq \frac{1}{4}(\mu_1 + \inf_M |\kappa|^2) + \inf_M |l\Psi|^2,$$

where  $\mu_1$  is the first eigenvalue of the basic Yamabe operator  $Y_b$  of  $\mathcal{F}$ .

From (4.11), we have the following corollary.

**Corollary 4.8.** (cf.[12]) *Let  $(M, \tilde{g}_M, \mathcal{F})$  be a compact Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  of codimension  $q \geq 3$  and a bundle-like metric  $\tilde{g}_M$ . If the transversal scalar curvature satisfies  $\sigma^\nabla \geq 0$ , then any eigenvalue  $\lambda$  of  $D_b$  satisfies*

$$(4.32) \quad \lambda^2 \geq \frac{q}{4(q-1)}(\mu_1 + \inf |\kappa|^2).$$

## 5. THE LIMITING CASE

In this section, we study the limiting foliations of (4.10) and (4.31). For any linear symmetric endomorphism  $l$  of  $Q$ , we define  $Ric_{\nabla}^l : \Gamma Q \otimes S(\mathcal{F}) \rightarrow S(\mathcal{F})$  ([9, 10]) by

$$(5.1) \quad Ric_{\nabla}^l(X \otimes \Psi) = \sum_a E_a \cdot R^l(X, E_a)\Psi,$$

where  $R^l$  is the curvature tensor with respect to  $\nabla$  defined by (4.1). Fix  $x \in M$  and choose an orthonormal basic frame  $\{E_a\}$  such that  $(\nabla E_a)_x = 0$  for all  $a$ . By a long calculation, for  $X \in \Gamma Q$  and  $\Psi \in \Gamma S(\mathcal{F})$  we have that at  $x$

$$(5.2) \quad Ric_{\nabla}^l(X \otimes \Psi) = -\frac{1}{2}\rho^\nabla(X) \cdot \Psi + \sum_a \{E_a \cdot dl(X, E_a) + E_a \cdot [l(X), l(E_a)]\} \cdot \Psi,$$

where  $[X, Y] = X \cdot Y - Y \cdot X$  and  $dl(X, Y) = (\nabla_X l)(Y) - (\nabla_Y l)(X)$ . By the definition of Clifford multiplication, we have

$$(5.3) \quad \sum_a E_a \cdot dl(X, E_a) = \sum_a E_a \wedge dl(X, E_a) - [X(\text{tr } l) - (\text{div}_{\nabla} l)(X)],$$

where  $(\operatorname{div}_{\nabla} l)(X) = \sum_a g_Q((\nabla_{E_a} l)(X), E_a)$ . Since  $l$  is symmetric,

$$(5.4) \quad \sum_a E_a \cdot l(E_a) = \sum_a l(E_a) \cdot E_a,$$

and then

$$(5.5) \quad \sum_a E_a \cdot l(E_a) = -\operatorname{tr} l.$$

From (5.4) and (5.5), we have

$$(5.6) \quad \sum_a E_a \cdot [l(X), l(E_a)] \cdot \Psi = 2(\operatorname{tr} l)l(X) \cdot \Psi - 2l^2(X) \cdot \Psi.$$

From (5.2), (5.3) and (5.6), we have that

$$(5.7) \quad \begin{aligned} \operatorname{Ric}_{\nabla}^l(X \otimes \Psi) &= -\frac{1}{2}\rho^{\nabla}(X) \cdot \Psi + \sum_a (E_a \wedge dl(X, E_a)) \cdot \Psi - 2l^2(X) \cdot \Psi \\ &\quad + 2(\operatorname{tr} l)l(X) \cdot \Psi - \{X(\operatorname{tr} l) - (\operatorname{div}_{\nabla} l)(X)\} \cdot \Psi. \end{aligned}$$

From (5.7), we have the following facts.

**Proposition 5.1.** *If  $M$  admits a non-zero spinor field  $\Psi \in \Gamma S(\mathcal{F})$  with  $\overset{l}{\nabla} \Psi = 0$ , then  $|\Psi|^2$  is constant and for any  $X \in \Gamma Q$*

$$(5.8) \quad \operatorname{grad}_{\nabla}(\operatorname{tr} l) = \operatorname{div}_{\nabla} l, \quad (\operatorname{tr} l)^2 = \frac{1}{4}\sigma^{\nabla} + |l|^2,$$

$$(5.9) \quad \frac{1}{2}\rho^{\nabla}(X) \cdot \Psi = \left[ \sum_a (E_a \wedge dl(X, E_a)) + 2(\operatorname{tr} l)l(X) - 2l^2(X) \right] \cdot \Psi.$$

*Proof.* Since  $\overset{l}{\nabla}$  is a metric connection, the condition  $\overset{l}{\nabla} \Psi = 0$  imply that  $|\Psi|^2$  is constant. If  $\overset{l}{\nabla}_X \Psi = 0$  for any  $X \in \Gamma Q$ , then  $\operatorname{Ric}_{\nabla}^l = 0$ . Hence from (5.7), we have

$$\begin{aligned} &< \{X(\operatorname{tr} l) - (\operatorname{div}_{\nabla} l)(X)\} \cdot \Psi, \Psi > \\ &= < \{-\frac{1}{2}\rho^{\nabla}(X) + \sum_a (E_a \wedge dl(X, E_a)) + 2(\operatorname{tr} l)l(X) - 2l^2(X)\} \cdot \Psi, \Psi >. \end{aligned}$$

The left-hand side is real, but the right-hand side is pure imaginary because  $\sum_a \langle E_a \wedge dl(X, E_a) \cdot \Psi, \Psi \rangle$  is pure imaginary. Therefore, both sides are zeros. Hence we get

$$(5.10) \quad X(\operatorname{tr} l) = (\operatorname{div}_{\nabla} l)(X),$$

$$(5.11) \quad \frac{1}{2} \rho^\nabla(X) \cdot \Psi = \left[ \sum_a (E_a \wedge dl(X, E_a)) + 2(\text{tr } l)l(X) - 2l^2(X) \right] \cdot \Psi.$$

Hence the first equation in (5.8) and (5.9) are proved from (5.10) and (5.11), respectively. For the second equation in (5.8), Clifford multiplication of (5.11) with  $E_b$  and for  $X = E_b$ , gives

$$\begin{aligned} \frac{1}{2} \sum_b E_b \cdot \rho^\nabla(E_b) \cdot \Psi &= \sum_{a,b} E_b \cdot (E_a \wedge dl(E_b, E_a)) \cdot \Psi + 2(\text{tr } l) \sum_b E_b \cdot l(E_b) \cdot \Psi \\ &\quad - 2 \sum_b E_b \cdot l^2(E_b) \cdot \Psi. \end{aligned}$$

From Lemma 3.1 and (5.5), we have

$$(5.12) \quad -\frac{1}{2} \sigma^\nabla \Psi = \sum_{a,b} E_a \cdot (E_b \wedge dl(E_a, E_b)) \cdot \Psi - 2(\text{tr } l)^2 \Psi + 2|l|^2 \Psi.$$

On the other hand,

$$(5.13) \quad \begin{aligned} \sum_{a,b} E_a \cdot (E_b \wedge dl(E_a, E_b)) \cdot \Psi &= \left( \sum_{a,b} E_a \wedge E_b \wedge dl(E_a, E_b) \right) \cdot \Psi \\ &\quad - \sum_{a,b} i(E_a)(E_b \wedge dl(E_a, E_b)) \cdot \Psi. \end{aligned}$$

The first term of the right-hand side of (5.13) is zero since  $l$  is symmetric, and the last term is the Clifford multiplication of  $\Psi$  with a vector field, which gives an imaginary function when taking its scalar product with  $\Psi$ . Thus we have

$$-\frac{1}{2} \sigma^\nabla |\Psi|^2 = -2(\text{tr } l)^2 |\Psi|^2 + 2|l|^2 |\Psi|^2.$$

Hence the proof of the second equation in (5.8) is completed. ■

Let  $\Psi_1$  be the eigenspinor corresponding to the eigenvalue  $\lambda_1^2 = \frac{1}{4} \inf_M (K^\sigma + 4|l_{\Psi_1}|^2)$ . From (4.9), we have

$$(5.14) \quad \nabla^{l_{\Psi_1}} \Psi_1 = 0, \quad K^\sigma = \text{constant}, \quad |l_{\Psi_1}| = \text{constant}.$$

From Proposition 5.1, we know that  $|\Psi_1|$  is constant and

$$(5.15) \quad \lambda_1^2 = \frac{1}{4} \sigma^\nabla + |l_{\Psi_1}|^2.$$

From (5.15), the transversal scalar curvature  $\sigma^\nabla$  is constant and we have

$$(5.16) \quad \inf |\kappa|^2 = 0.$$

Since  $\sigma^\nabla$  and  $K^\sigma = \sigma^\nabla + |\kappa|^2$  are constant,  $|\kappa|$  is constant and then  $|\kappa| = \inf_M |\kappa| = 0$ . This implies that  $\mathcal{F}$  is minimal. Hence we have the following theorem.

**Theorem 5.2.** *Under the same assumption as in Theorem 4.1, if there exists an eigenspinor  $\Psi (\neq 0)$  of the basic Dirac operator  $D_b$  for the eigenvalue  $\lambda^2 = \frac{1}{4} \inf_M (K^\sigma + 4|l_\Psi|^2)$ , then  $|\Psi|$  is constant and  $\mathcal{F}$  is minimal with the constant transversal scalar curvature  $\sigma^\nabla = 4(\lambda^2 - |l_\Psi|^2)$ .*

Next, we study the limiting case of (4.31). Similarly, we have

$$(5.17) \quad Ric_{\bar{\nabla}}^{l_{\bar{\Phi}}}(X \otimes \bar{\Phi}) = -\frac{1}{2} \rho^{\bar{\nabla}}(X) \cdot \bar{\Phi} + \sum (\bar{E}_a \wedge dl_{\bar{\Phi}}(X, \bar{E}_a)) \cdot \bar{\Phi} - 2l_{\bar{\Phi}}^2(X) \cdot \bar{\Phi} + 2(\text{tr}_{\bar{\Phi}})l_{\bar{\Phi}}(X) \cdot \bar{\Phi} - \{X(\text{tr}_{\bar{\Phi}}) - (\text{div}_{\bar{\nabla}} l_{\bar{\Phi}})(X)\} \cdot \bar{\Phi}.$$

From (5.17), we have the following proposition.

**Proposition 5.3.** *If  $M$  admits a non-zero spinor  $\Psi$  with  $\frac{l_{\bar{\Phi}}}{\bar{\nabla}} \bar{\Phi} = 0$ , where  $\bar{\Phi} = e^{-\frac{q-1}{2}u} \Psi$ , then  $|\Phi|$  is constant and*

$$(5.18) \quad \nabla_X \Psi = \frac{1}{2} \pi(X) \cdot \text{grad}_{\nabla}(u) \cdot \Psi + \frac{q}{2} X(u) \Psi - l_{\Psi}(\pi(X)) \cdot \Psi, \quad X \in TM$$

$$(5.19) \quad X(\text{tr } l_{\bar{\Phi}}) = (\text{div}_{\bar{\nabla}} l_{\bar{\Phi}})(X), \quad (\text{tr } l_{\bar{\Phi}})^2 = \frac{1}{4} \sigma^{\bar{\nabla}} + |l_{\bar{\Phi}}|^2.$$

*Proof.* Since  $\frac{l_{\bar{\Phi}}}{\bar{\nabla}}$  is metrical,  $|\bar{\Phi}|$  is constant. Moreover,  $\frac{l_{\bar{\Phi}}}{\bar{\nabla}} \bar{\Phi} = 0$  is equivalent to

$$(5.20) \quad \bar{\nabla}_X \bar{\Phi} + l_{\bar{\Phi}}(\pi(X)) \cdot \bar{\Phi} = 0.$$

From Proposition 3.2 and (5.20), we have that for  $\bar{\Phi} = e^{-\frac{q-1}{2}u} \Psi$

$$(5.21) \quad \nabla_X \bar{\Phi} = \frac{1}{2} \pi(X) \cdot \text{grad}_{\nabla}(u) \cdot \bar{\Phi} + \frac{1}{2} X(u) \bar{\Phi} - l_{\Psi}(\pi(X)) \cdot \bar{\Phi},$$

which gives (5.18). The proof of (5.19) is similar to the one in Proposition 5.1. ■

By a direct calculation together with (3.10), we obtain the following lemma.

**Lemma 5.4.** *For any vector field  $X \in \Gamma Q$  and any isomorphism  $l$ , we have*

$$(5.22) \quad (\text{div}_{\bar{\nabla}} l)(X) = (\text{div}_{\nabla} l)(X) + q g_Q(l(X), \text{grad}_{\nabla}(u)) - X(u) \text{tr } l,$$

where  $\bar{\nabla}$  is a Levi-Civita connection with respect to  $\bar{g}_Q = e^{2u}g_Q$ .

On the other hand, we have that for any  $\Phi = e^{-\frac{q-1}{2}u}\Psi$  and any  $X \in \Gamma Q$

$$(5.23) \quad (\operatorname{div}_{\bar{\nabla}} l_{\bar{\Phi}})(X) = e^{-u} \{ (\operatorname{div}_{\nabla} l_{\Psi})(X) - g_Q(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)) \}.$$

From (5.22) and (5.23), we have

$$(5.24) \quad \begin{aligned} & (\operatorname{div}_{\bar{\nabla}} l_{\bar{\Phi}})(X) \\ &= e^{-u} \{ (\operatorname{div}_{\nabla} l_{\Psi})(X) + (q-1)g_Q(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)) - X(u) \operatorname{tr} l_{\Psi} \}. \end{aligned}$$

Comparing with (5.19) and (5.24), we have the following proposition.

**Corollary 5.5.** *If  $M$  admits a non-zero spinor  $\Psi$  with  $\bar{\nabla} \bar{\Phi} = 0$ , where  $\bar{\Phi} = e^{-\frac{q-1}{2}u}\Psi$ , then for any  $X \in \Gamma Q$*

$$X(\operatorname{tr} l_{\Psi}) = (\operatorname{div}_{\nabla} l_{\Psi})(X) + (q-1)g_Q(l_{\Psi}(X), \operatorname{grad}_{\nabla}(u)).$$

Let  $D_b \Psi = \lambda \Psi$  with  $\lambda^2 = \frac{1}{4}(\mu_1 + \inf_M |\kappa|^2) + \inf_M |l_{\Psi}|^2$ . From (4.20) and (4.30), we have that for  $\Phi = e^{-\frac{q-1}{2}u}\Psi$

$$(5.25) \quad \bar{\nabla} \bar{\Phi} = 0, \quad |\kappa| = \text{constant}, \quad |l_{\Psi}| = \text{constant}.$$

Since  $|l_{\Psi}|$  is constant, we have from (4.15) and (5.19) that

$$(5.26) \quad \mu_1 + \inf_M |\kappa|^2 = e^{2u} \sigma^{\bar{\nabla}}.$$

From Lemma 4.6, we have that for  $u \in \mathcal{K}_u$

$$(5.27) \quad \mu_1 + \inf_M |\kappa|^2 = \sigma^{\nabla} + 2(q-1)\Delta_B u + (q-1)(2-q)|\operatorname{grad}_{\nabla}(u)|^2.$$

Let  $u$  be an eigenfunction of the basic Yamabe operator  $Y_b$ . Then from (4.30) we get

$$(5.28) \quad \inf_M |\kappa|^2 = 0.$$

Since  $|\kappa|$  is constant, we have that  $|\kappa| = 0$ , i.e.  $\mathcal{F}$  is minimal. Hence we have the following theorem.

**Theorem 5.6.** *Let  $(M, \tilde{g}_M, \mathcal{F})$  be a compact Riemannian manifold with a transverse spin foliation  $\mathcal{F}$  of codimension  $q \geq 3$  and a bundle-like metric  $\tilde{g}_M$ . Assume that an eigenvalue  $\lambda$  of  $D_b$  corresponding to the eigenspinor  $\Psi$  satisfies*

$$\lambda^2 = \frac{1}{4}(\mu_1 + \inf_M |\kappa|^2) + \inf_M |l_{\Psi}|^2.$$



Then  $|l_\Psi|$  is constant and the foliation  $\mathcal{F}$  is minimal. Moreover

$$(5.29) \quad (\operatorname{div}_\nabla l_\Psi)(X) = (1 - q)g_Q(l_\Psi(X), \operatorname{grad}_\nabla(u))$$

for any  $X \in \Gamma Q$ .

*Proof.* The equation (5.29) is trivial from corollary 5.5. ■

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