

Regular Wavelets, Heat Semigroup and Application to the Magneto-hydrodynamic Equations with Data in Critical Triebel-Lizorkin Type Oscillation Spaces

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Abstract. In this paper, based on a Poisson type extension of Triebel-Lizorkin type oscillation spaces $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$, we establish a bilinear estimate on some new tent spaces $\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$ associated with $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$. As an application, we get the well-posedness and regularity of the fractional magneto-hydrodynamic equations and quasi-geostrophic equations with initial data in the critical $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$.

1. Introduction

Triebel-Lizorkin spaces $\dot{F}_{p,q}^s(\mathbb{R}^n)$ were introduced as generalizations of many classical function spaces such as Lebesgue spaces and Sobolev spaces. We refer the reader to Triebel [48, 49] for an overview of $\dot{F}_{p,q}^s(\mathbb{R}^n)$. In the research of harmonic analysis and partial differential equations, Triebel-Lizorkin spaces have been used extensively and attracted the attention of many mathematicians. See Meyer-Yang [40], Chae [4, 5], Kozono-Shimada [28], Chen-Miao-Zhang [8] and the reference therein.

In this paper, we apply wavelets and semigroup to study a class of Triebel-Lizorkin type oscillation spaces $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$ defined as

$$\sup_Q |Q|^{\gamma_2/n-1/p} \inf_{P_{Q,f} \in S_{p,q,f}^{\gamma_1,\gamma_2}} \|\varphi_Q(f - P_{Q,f})\|_{\dot{F}_p^{\gamma_1,q}} < +\infty,$$

where the supremum is taken over all cubes Q and $S_{p,q,f}^{\gamma_1,\gamma_2}$ denotes the set of all polynomials satisfying (2.2) below. See Definition 2.3. The spaces $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$ cover many classical function spaces such as Sobolev spaces $W^{p,\gamma_1}(\mathbb{R}^n)$, Triebel-Lizorkin spaces $\dot{F}_{p,q}^{\gamma_1}(\mathbb{R}^n)$, Morrey spaces $L^{p,\gamma_2}(\mathbb{R}^n)$, bounded mean oscillation spaces $BMO(\mathbb{R}^n)$, Q type spaces $Q_\alpha^\beta(\mathbb{R}^n)$ and so on. We give the following space structure table to clarify the relation between these spaces and $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$:

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$1 < p < \infty, q = 2, \gamma_1 \in \mathbb{R}, \gamma_2 = n/p$	$\dot{F}_{p,2}^{\gamma_1, n/p}(\mathbb{R}^n) = W^{p, \gamma_1}(\mathbb{R}^n)$
$1 < p, q < \infty, \gamma_1 \in \mathbb{R}, \gamma_2 = n/p$	$\dot{F}_{p,q}^{\gamma_1, n/p}(\mathbb{R}^n) = \dot{F}_{p,q}^{\gamma_1}(\mathbb{R}^n)$
$1 < p < \infty, q = 2, \gamma_1 = 0, \gamma_2 = n/p$	$\dot{F}_{p,2}^{0, \gamma_2}(\mathbb{R}^n) = L^{p, \gamma_2}(\mathbb{R}^n)$
$p = q = 2, \gamma_1 = 0, \gamma_2 = -n/2,$	$\dot{F}_{2,2}^{0, -n/2}(\mathbb{R}^n) = \text{BMO}(\mathbb{R}^n)$
$p = q = 2, \gamma_1 = \alpha - \beta + 1, \gamma_2 = \alpha + \beta - 1$	$\dot{F}_{2,2}^{\alpha - \beta + 1, \alpha + \beta - 1}(\mathbb{R}^n) = Q_\alpha^\beta(\mathbb{R}^n)$

In Theorem 2.5, we give a wavelet characterization of $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$. As a consequence, $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ coincide with $\dot{F}_{p,q}^{s, \tau}(\mathbb{R}^n)$ introduced by Yang-Yuan [62]. Theorem 2.5 and Lemma 2.7 imply that Calderón-Zygmund operators are bounded on $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$. Moreover, our wavelet characterization is independent of the choice of wavelet basis. See Theorem 2.8 and Corollary 2.9.

It is well-known that for any $f \in \text{BMO}(\mathbb{R}^n)$, the Poisson integral $P_t(f)$ gives a harmonic extension of f . Obviously, the Poisson kernel can be replaced with the fractional heat semigroup $\{e^{-t(-\Delta)^\beta}\}$. In a sense, the function spaces on \mathbb{R}^n can be characterized via the semigroups and the tent spaces on \mathbb{R}_+^{n+1} , see Coifman-Meyer-Stein [9] for example. In the last decade, such an idea has been applied to the study of the well-posedness of Navier-Stokes equations. In 2001, Koch-Tataru [26] obtained a semigroup characterizations of $\text{BMO}(\mathbb{R}^n)$. In 2007, by Hausdorff capacity, Xiao [59] gave a semigroup characterization of $Q_\alpha(\mathbb{R}^n)$. Li-Zhai [33] developed the idea of [26, 59] and obtained a semigroup characterization of $Q_\alpha^\beta(\mathbb{R}^n)$. For further information, we refer the reader to Cannone [2, 3], Dafni-Xiao [14], Lemarié-Rieusset [29], Li-Xiao-Yang [30], Lin-Yang [35] and Miao-Yuan-Zhang [42].

In Section 2.3, we introduce tent type spaces $\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$ as follows:

$$\begin{aligned} \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2} &= \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I} \cap \mathbb{F}_{p,q}^{\gamma_1, \gamma_2, II} \cap \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, III} \cap \mathbb{F}_{p,q,m'}^{\gamma_1, \gamma_2, IV} \\ &=: X_1 \cap X_2 \cap X_3 \cap X_4. \end{aligned}$$

Via fractional heat semigroup, we can establish a relation between $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ and $\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$. See Theorem 2.17 for the details.

Actually, Theorem 2.17 is not a simple generalization of the results in [26, 33, 59]. In the above structure table, $\text{BMO}(\mathbb{R}^n)$, $Q_\alpha(\mathbb{R}^n)$ and $Q_\alpha^\beta(\mathbb{R}^n)$ are all $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ spaces with $p = q = 2$. Technically, the Plancherel identity of Fourier transform plays an important role in the semigroup characterizations of these function spaces. For the cases $p \neq 2$, this method is not valid. To overcome this difficulty, we apply a new method. Let $\{\Phi_{j,k}^\varepsilon\}_{(\varepsilon, j, k) \in \Lambda_n}$ be a wavelet basis and

$$f = \sum_{(\varepsilon, j, k) \in \Lambda_n} a_{j,k}^\varepsilon \Phi_{j,k}^\varepsilon \in \dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n).$$

For any cube Q , based on the relation between j and the radius of Q , we decompose the function

$$F(x, t) := e^{t(-\Delta)^\beta} f(x)$$

into several parts such that every part belongs to some X_i . Such a decomposition provides a clear view of the local structures of $F(x, t)$ related to the spatial variable and the frequency. So the desired semigroup characterization of $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ can be obtained easily.

As an application, with initial data in critical $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$, we consider the well-posedness of the fractional magneto-hydrodynamic equations $(\text{FMHD})_\beta$:

$$(1.1) \quad \begin{cases} \partial_t u + (-\Delta)^\beta u + u \cdot \nabla u + \nabla p - b \cdot \nabla b = 0; \\ \partial_t b + (-\Delta)^\beta b + u \cdot \nabla b - b \cdot \nabla u = 0; \\ \nabla \cdot u = \nabla \cdot b = 0; \\ u|_{t=0} = u_0, \quad b|_{t=0} = b_0 \end{cases}$$

with $\beta > 1/2$ and $n \geq 3$.

If $\beta = 1$, equations (1.1) reduce to the classical incompressible magneto-hydrodynamic (MHD) equations:

$$(1.2) \quad \begin{cases} \partial_t u + (-\Delta)u + u \cdot \nabla u + \nabla p - b \cdot \nabla b = 0; \\ \partial_t b + (-\Delta)b + u \cdot \nabla b - b \cdot \nabla u = 0; \\ u|_{t=0} = u_0, \quad b|_{t=0} = b_0, \end{cases}$$

with $\nabla \cdot u = \nabla \cdot b = 0$. Further, if the magnetic field $b(t, x)$ identically equals to zero, equations (1.2) reduce to the incompressible fractional Navier-Stokes equations:

$$(1.3) \quad \begin{cases} \partial_t u + (-\Delta)^\beta u + u \cdot \nabla u + \nabla p - u \cdot \nabla u = 0; \\ \nabla \cdot u = 0; \\ u|_{t=0} = u_0. \end{cases}$$

As a basic model in the fluid mechanics, the well-posedness of equations (1.3) with $\beta = 1$ in classical function spaces has been studied by many mathematicians. For example, with initial data in Lebesgue space L^n and Sobolev spaces $\dot{H}^{n/2-1}(\mathbb{R}^n)$, the well-posed results are obtained by Kato [24] and Fujita-Kato [25], respectively. Cannone [2, 3] and Planchon [46] proved the global well-posedness for small initial data in the Besov spaces $\dot{B}_{p,\infty}^{n/p-1}(\mathbb{R}^n)$, $n < p < \infty$. In 2001, for $\beta = 1$, Koch-Tataru [26] proved the well-posedness of (1.3) with data in $\text{BMO}^{-1}(\mathbb{R}^n)$. For further progress on equations (1.3), see Deng-Yao [15], Germain-Pavlović-Staffilani [18], Giga-Miyakawa [19], Giga-Sawada [20], Lemarié-Rieusset [29], Li-Zhai [33], Li-Xiao-Yang [30], Lions [36], Miao-Yuan-Zhang [42], Miura-Sawada [43], Taylor [47], Wang-Xiao [50], Wu [54, 57], Xiao [59, 60] and Zhai [66] and the reference therein.

For equations (1.1), because of the coupling effect between the velocity $u(x, t)$ and the magneto field $b(x, t)$, the situation is more complicated. Kozono [27] proved the existence of a classical solution for the two-dimensional MHD system in a bounded domain. Alexseev [1] established the existence and uniqueness results for the ideal MHD system in Sobolev space $W^k(\mathbb{R}^n)$. For $\beta = 1$, Miao-Yuan-Zhang [41] obtained the well-posedness of equations (1.1) with data in $BMO^{-1}(\mathbb{R}^n)$. For more information on the MHD equations, we refer the reader to Alexseev [1], Wu [52], Wu [53, 58] and the reference therein.

Leaving more notations and terminologies later, let us outline roughly our approach of the main result. With data in Besov spaces or Triebel-Lizorkin spaces, one usual method to derive the well-posedness of (1.3) is the “energy inequality” based on some commutator estimates and Littlewood-Paley theory. See Wu [54]. We apply a new method. The key step is a decomposition of the bilinear operator

$$B_\ell(u, v)(t, x) = \int_0^t e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell}(uv) ds, \quad \ell = 1, 2, \dots, n$$

based on multi-resolution analysis. We decompose $B_\ell(u, v)$ as follows.

$$B_\ell(u, v)(t, x) := \int_0^t \sum_{j' \in \mathbb{Z}} \sum_{i=1}^4 I_{j'}^{i, \ell}(s, t, x) ds,$$

where $I_{j'}^{i, \ell}$ are composed of the wavelet coefficients of u and v . See (3.2) for details. By this decomposition, Lemma 2.11 enables us to convert the bilinear estimate of $B_\ell(u, v)$ into various efficient computations involved in the wavelets coefficients of u and v . Finally we obtain the desired bilinear estimate on $\mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2} \times \mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}$. See Sections 3.2–3.3. The well-posedness of equations (1.1) can be derived from such bilinear estimates together with a usual procedure. Our method can be also applied to the well-posedness of quasi-geostrophic equations (4.4). See Section 4.

We point out that when we consider the well-posedness of equation (1.1) with data in $\dot{F}_{p, q}^{\gamma_1, \gamma_2}$, the tent type space $\mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}$ has some distinct advantages. On one hand, for $F(x, t) \in \mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}$, the four parts of $\|F\|_{\mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}}$ have different meanings:

- the norms $\|F\|_{X_1}$ and $\|F\|_{X_2}$ denote the Bloch-parts of $F(x, t)$,
- the norms $\|F\|_{X_3}$ and $\|F\|_{X_4}$ denote the L^p -parts of $F(x, t)$.

Furthermore, the index m represents the regularities for the variable x . Compared with the results in [26, 33, 59], if m becomes bigger, the elements in $\mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}$ have higher regularities.

On the other hand, Lin-Yang [35] and Li-Yang [31] used another tent type spaces to study the well-posedness of equations (1.3) and (1.1) with data in the critical Besov- Q spaces $\dot{B}_{p, q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$, respectively. In [35] and [31], the tent type spaces consist of two parts:

the L^∞ -part and L^p -part. Hence the index β is restricted to the scope $(1/2, 1]$. In this paper, Bloch parts $\|F\|_{X_1}$ and $\|F\|_{X_2}$ are substituted for the L^∞ -part. This change will enable us to obtain the well-posedness of equations (1.3) and (1.1) with $\beta > 1/2$.

Remark 1.1. Interestingly, Federbush [16] employed the divergence-free wavelets to study the classical Navier-Stokes equations, while the wavelets used in this paper are classical Meyer wavelets. In addition, when constructing a contraction mapping, Federbush’s method was based on the estimates of “long wavelength residues”. Nevertheless, our wavelet approach is based on Lemmas 3.2–3.7 and the Cauchy-Schwarz inequality.

Remark 1.2. In recent years, some special cases of $\dot{F}_{2,2}^{\gamma_1,\gamma_2}$ have been applied to the study well-posedness of equations (1.3) by several authors. We refer the reader to Wang-Xiao [50,51], Xiao [60] for further information.

The rest of this paper is organized as follows. In Section 2, we state some preliminary knowledge, notations and terminologies. In Section 3, we introduce a class of tent spaces $\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$. For $f \in \dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$, we prove its Poisson type extension belongs to $\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$. In Sections 3.2–3.3, we establish a bilinear estimate on $\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2} \times \mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$. In Section 4, we prove the well-posedness of equations (1.1) and (4.4) with data in $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$ being small.

2. Notations and preliminaries

In this paper, the symbols \mathbb{Z} and \mathbb{N} denote the sets of all integers and natural numbers, respectively. For $n \in \mathbb{N}$, \mathbb{R}^n is the n -dimensional Euclidean space, with Euclidean norm denoted by $|x|$ and the Lebesgue measure denoted by dx . \mathbb{R}_+^{n+1} is the upper half-space $\{(t, x) \in \mathbb{R}_+^{n+1} : t > 0, x \in \mathbb{R}^n\}$ with Lebesgue measure $dt dx$. $B(x, r)$ denotes a ball in \mathbb{R}^n with center x , radius r and the volume $|B|$. Denote by Q a cube in \mathbb{R}^n with the sides parallel to the coordinate axes. The volume and side length of Q are denoted by $|Q|$ and $l(Q)$, respectively.

For convenience, the positive constants C may change from one line to another and usually depend on the dimension n , α , β and other fixed parameters. The Schwartz class of rapidly decreasing functions and its dual will be denoted by $\mathcal{S}(\mathbb{R}^n)$ and $\mathcal{S}'(\mathbb{R}^n)$, respectively. For $f \in \mathcal{S}(\mathbb{R}^n)$, \hat{f} means the Fourier transform of f .

2.1. Wavelets

In this paper, we use real valued tensor product orthogonal wavelets Φ^ϵ , which will be Daubechies wavelets or Meyer wavelets. Daubechies wavelets are used only to get a wavelet characterization of $\dot{F}_{p,q}^{\gamma_1,\gamma_2}$ and Meyer wavelets will be used throughout this paper. If Φ^ϵ is a Daubechies wavelet, we assume that there exists an integer m_0 which is greater than some constant depending on the index of $\dot{F}_{p,q}^{\gamma_1,\gamma_2}$ such that

(1) for any $\epsilon \in \{0, 1\}^n$, $\Phi^\epsilon \in C_0^{m_0}([-2^M, 2^M]^n)$;

(2) for any $\epsilon \in E_n$, Φ^ϵ has the vanishing moments up to the order $m_0 - 1$.

We present some preliminaries on Meyer wavelets. Let $\Psi^0 \in C_0^\infty([-4\pi/3, 4\pi/3])$ be an even function satisfying $\Psi^0(\xi) \in [0, 1]$ and $\Psi^0(\xi) = 1$, if $|\xi| \leq 2\pi/3$. Let $\Omega(\xi) = [(\Psi^0(\xi/2)^2 - (\Psi^0(\xi))^2)^{1/2}]$. Then $\Omega \in C_0^\infty([-8\pi/3, 8\pi/3])$ is an even function satisfying

$$\begin{cases} \Omega(\xi) = 0, & \text{if } |\xi| \leq 2\pi/3; \\ \Omega^2(\xi) + \Omega^2(2\xi) = \Omega^2(\xi) + \Omega^2(2\pi - \xi) = 1, & \text{if } \xi \in [2\pi/3, 4\pi/3]. \end{cases}$$

Let $\Psi^1(\xi) = \Omega(\xi)e^{-i\xi/2}$. We define

$$\begin{cases} \widehat{\Phi}^\epsilon = \prod_{i=1}^n \Psi^{\epsilon_i}(\xi_i), & \epsilon = (\epsilon_1, \dots, \epsilon_n) \in \{0, 1\}^n; \\ \Phi_{j,k}^\epsilon(x) = 2^{nj/2} \Phi^\epsilon(2^j x - k), & j \in \mathbb{Z}, k \in \mathbb{Z}^n. \end{cases}$$

In this paper, we denote $E_n := \{0, 1\}^n \setminus \{0\}$, $F_n := \{(\epsilon, k) : \epsilon \in E_n, k \in \mathbb{Z}^n\}$ and $\Lambda_n = \{(\epsilon, j, k), \epsilon \in E_n, j \in \mathbb{Z}, k \in \mathbb{Z}^n\}$. For further information about wavelets, we refer the reader to Meyer [38] and Meyer-Coifman [39]. The following result is well-known.

Lemma 2.1. $\{\Phi_{j,k}^\epsilon\}_{(\epsilon,j,k) \in \Lambda_n}$ is an orthogonal basis in $L^2(\mathbb{R}^n)$.

For a function f , denote by $f_{j,k}^\epsilon = \langle f, \Phi_{j,k}^\epsilon \rangle$ the wavelet coefficients of f , where $\epsilon \in \{0, 1\}^n$, $k \in \mathbb{Z}^n$. Let

$$P_j f(x) = \sum_{k \in \mathbb{Z}^n} f_{j,k}^0 \Phi_{j,k}^0(x) \quad \text{and} \quad Q_j f(x) = \sum_{(\epsilon,k) \in F_n} f_{j,k}^\epsilon \Phi_{j,k}^\epsilon(x).$$

By Lemma 2.1, we can see that P_j and Q_j are projection operators on $L^2(\mathbb{R}^n)$. In fact, for any two functions u and v , we have

$$\begin{aligned} uv &= \sum_{j \in \mathbb{Z}} P_{j-3} u Q_j v + \sum_{j \in \mathbb{Z}} Q_j u Q_j v + \sum_{0 < j-j' \leq 3} Q_j u Q_{j'} v \\ (2.1) \quad &+ \sum_{0 < j'-j \leq 3} Q_{j'} u Q_j v + \sum_{j \in \mathbb{Z}} Q_j u P_{j-3} v. \end{aligned}$$

2.2. Triebel-Lizorkin type oscillation spaces and wavelet characterization

As a generalization of Triebel-Lizorkin spaces and Morrey spaces, Triebel-Lizorkin type oscillation spaces have been introduced by Yang [61] via wavelets. We refer the reader to Yuan-Sickel-Yang [65] for further information on Triebel-Lizorkin type oscillation spaces and their generalizations. To make the mathematicians on PDE to understand such spaces better, we give here a direct definition by using classical Triebel-Lizorkin spaces.

Suppose that ϕ is a function on \mathbb{R}^n such that

$$\text{supp } \widehat{\phi} \subset \{\xi \in \mathbb{R}^n : |\xi| \leq 1\} \quad \text{and} \quad \widehat{\phi}(\xi) = 1 \text{ for } \{\xi \in \mathbb{R}^n : |\xi| \leq 1/2\}.$$

For $v \in \mathbb{Z}$, define

$$\phi_v(x) = 2^{n(v+1)}\phi(2^{v+1}x) - 2^{nv}\phi(2^v x)$$

as the Littlewood-Paley functions. See Frazier-Jawerth-Weiss [17] and Triebel [48, 49] for more information on the Littlewood-Paley theory. The Triebel-Lizorkin spaces are defined as follows.

Definition 2.2. Let $-\infty < r < \infty$, $1 \leq p < \infty$ and $1 \leq q \leq \infty$. A function $f \in \mathcal{S}'(\mathbb{R}^n)/\mathcal{P}(\mathbb{R}^n)$ belongs to $\dot{F}_{p,q}^r(\mathbb{R}^n)$ if

$$\|f\|_{\dot{F}_{p,q}^r} = \left\| \left(\sum_{v \in \mathbb{Z}} 2^{qvr} |\phi_v * f(x)|^q \right)^{1/q} \right\|_p < \infty.$$

Take $\varphi \in C_0^\infty(B(0, n))$ such that $\varphi(x) = 1$ for $x \in B(0, \sqrt{n})$. Denote by $Q(x_0, r)$ a cube centered at x_0 and with side length r . For simplicity, sometimes, we write $Q(x_0, r)$ as Q or $Q(r)$. Let $\varphi_Q(x) = \varphi\left(\frac{x-x_Q}{r}\right)$. For $0 < p, q \leq \infty$, $\gamma_1, \gamma_2 \in \mathbb{R}$, let $m_0 = m_{p,q}^{\gamma_1, \gamma_2}$ be sufficiently big positive real number. For any function f , let $S_{p,q,f}^{\gamma_1, \gamma_2}$ be the set of polynomials $P_{Q,f}$ such that for any $|\alpha| \leq m_0$,

$$(2.2) \quad \int x^\alpha \varphi_Q(x)(f(x) - P(x)) dx = 0.$$

The Triebel-Lizorkin type oscillation spaces are defined as follows.

Definition 2.3. Let $0 < p, q \leq \infty$, $\gamma_1, \gamma_2 \in \mathbb{R}$. We say that f belongs to Triebel-Lizorkin type oscillation spaces $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ provided

$$(2.3) \quad \sup_Q |Q|^{\gamma_2/n-1/p} \inf_{P_{Q,f} \in S_{p,q,f}^{\gamma_1, \gamma_2}} \|\varphi_Q(f - P_{Q,f})\|_{\dot{F}_p^{\gamma_1, q}} < +\infty,$$

where the supremum is taken over all the cubes with center x_Q and length r in \mathbb{R}^n .

It is easy to verify the following results.

Proposition 2.4. Let $0 < p, q \leq \infty$ and $\gamma_1, \gamma_2 \in \mathbb{R}$.

- (i) The definition of $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ is independent of the choice of Littlewood-Paley function ϕ .
- (ii) For $1 \leq p, q \leq \infty$, $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ are Banach spaces.

Applying Hausdorff capacity and Littlewood-Paley theory, Yang-Yuan [62] introduced a new class of function spaces $\dot{F}_{p,q}^{s,\tau}(\mathbb{R}^n)$, $p \in (0, \infty)$, which generalize many classical function spaces. For more information, we refer to Liang et al. [34], Yang-Yuan [63, 64] and Yuan-Sickel-Yang [65].

We call Φ^ϵ a regular Daubechies wavelet, if there exist two integers m_0 and M such that

$$(2.4) \quad \begin{cases} \text{for } \epsilon \in \{0, 1\}^n, \Phi^\epsilon \in C_0^m([-2^M, 2^M]^n); \\ \text{for } \epsilon \in E_n \text{ and } |\alpha| \leq m_0, \int x^\alpha \Phi^\epsilon(x) dx = 0. \end{cases}$$

Let $0 < p, q \leq \infty$ and $\gamma_1, \gamma_2 \in \mathbb{R}$. In Theorem 2.5 below, we take an integer $m_{p,q}^{\gamma_1, \gamma_2}$ large enough such that the constant m_0 in (2.4) satisfies $m_0 \geq m_{p,q}^{\gamma_1, \gamma_2}$. For any cube Q , let $\Lambda_n^Q = \{(\epsilon, j, k) \in \Lambda_n, Q_{j,k} \subset Q\}$. We have the following wavelet characterization of $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$. We omit the proof and refer the reader to Yuan-Sickel-Yang [65] and Li-Yang-Zheng [32].

Theorem 2.5. *Let $0 < p, q \leq \infty$, $\gamma_1, \gamma_2 \in \mathbb{R}$. $f = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^\epsilon \Phi_{j,k}^\epsilon \in \dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$ if and only if*

$$(2.5) \quad \sup_Q |Q|^{\gamma_2/n-1/p} \left\| \left(\sum_{(\epsilon, j, k) \in \Lambda_n^Q} 2^{qj(\gamma_1+n/2)} |a_{j,k}^\epsilon|^q \chi(2^j \cdot -k) \right)^{1/q} \right\|_{L^p} < +\infty,$$

where the supremum is taken over all the dyadic cubes Q in \mathbb{R}^n .

Remark 2.6. If $\gamma_2 = n/p$, $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n) = \dot{F}_p^{\gamma_1, q}(\mathbb{R}^n)$. Moreover, for $\gamma_2 > n/p$, if $f \in \dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$, $f_{j,k}^\epsilon = 0$. Hence f is only a polynomial.

Now we state some preliminaries about Calderón-Zygmund operators and refer to Meyer [38] and Meyer-Coifman [39] for further information. For $x \neq y$, let $K(x, y)$ be a smooth function such that there exists a sufficiently large $N_0 \leq m_0$ satisfying that

$$(2.6) \quad \left| \partial_x^\alpha \partial_y^\beta K(x, y) \right| \leq C |x - y|^{-(n+|\alpha|+|\beta|)}, \quad \forall |\alpha| + |\beta| \leq N_0.$$

A linear operator T is said to be a Calderon-Zygmund operator if

- (i) T is continuous from $C^1(\mathbb{R}^n)$ to $(C^1(\mathbb{R}^n))'$;
- (ii) $Tf(x) = \int K(x, y)f(y) dy$, where $K(x, y)$ satisfies (2.6);
- (iii) $Tx^\alpha = T^*x^\alpha = 0$ for $\alpha \in \mathbb{N}^n$ with $|\alpha| \leq N_0$.

We denote $T \in \text{CZO}(N_0)$. By Schwartz kernel theorem, $K(\cdot, \cdot)$ is a distribution in $S'(\mathbb{R}^{2n})$. For $(\epsilon, j, k), (\epsilon', j', k') \in \Lambda_n$, let

$$a_{j,k,j',k'}^{\epsilon, \epsilon'} = \left\langle K(\cdot, \cdot), \Phi_{j,k}^\epsilon \Phi_{j',k'}^{\epsilon'} \right\rangle.$$

If T is a Calderon-Zygmund operator, then its kernel $K(\cdot, \cdot)$ and the related coefficients satisfy the following relations.

Lemma 2.7. [39, Section 8.3, Proposition 1]

(i) If $T \in \text{CZO}(N_0)$, then for $(\epsilon, j, k), (\epsilon', j', k') \in \Lambda_n$, $\{a_{j,k,j',k'}^{\epsilon,\epsilon'}\}$ satisfy the following condition:

$$(2.7) \quad \left| a_{j,k,j',k'}^{\epsilon,\epsilon'} \right| \leq C 2^{-|j-j'|(n/2+N_0)} \left(\frac{2^{-j} + 2^{-j'}}{2^{-j} + 2^{-j'} + |k2^{-j} - k'2^{-j'}|} \right)^{n+N_0}.$$

(ii) If $a_{j,k,j',k'}^{\epsilon,\epsilon'}$ satisfies the above condition (2.7), then

$$K(x, y) = \sum_{(\epsilon,j,k) \in \Lambda_n} \sum_{(\epsilon',j',k') \in \Lambda_n} a_{j,k,j',k'}^{\epsilon,\epsilon'} \Phi_{j,k}^\epsilon(x) \Phi_{j',k'}^{\epsilon'}(y)$$

in the sense of distributions and for any small positive real number δ , $T \in \text{CZO}(N_0 - \delta)$.

Following the ideas of Meyer-Yang [40], we can prove the following result. We omit the proof and refer the reader to Yuan-Sickel-Yang [65].

Theorem 2.8. Let $1 < p, q < \infty$ and $\gamma_1, \gamma_2 \in \mathbb{R}$. The Calderón-Zygmund operators are bounded on $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$.

Let $\{\Phi_{j,k}^{i,\epsilon}\}_{(\epsilon,j,k) \in \Lambda_n}$, $i = 1, 2$, be two regular orthogonal wavelets basis. For $(\epsilon, j, k), (\epsilon', j', k') \in \Lambda_n$, denote $a_{j,k,j',k'}^{\epsilon,\epsilon'} = \langle \Phi_{j,k}^{1,\epsilon}, \Phi_{j',k'}^{2,\epsilon'} \rangle$. We know that $\{a_{j,k,j',k'}^{\epsilon,\epsilon'}\}_{(\epsilon,j,k),(\epsilon',j',k') \in \Lambda_n}$ satisfies the condition (2.7). The following result is a corollary of Theorem 2.8. We refer to Yuan-Sickel-Yang [65] for a proof.

Corollary 2.9. Theorem 2.5 is also true for Meyer wavelets.

For $A > 0$ and a sequence $f := \{f_j\}$, we define the vector-valued maximal function $M_A(f)$ as

$$M_A(f)(x) = \left(\sum_j M(|f_j|^A)(x) \right)^{1/A}.$$

For $\{a_{j,k}^\epsilon\}_{(\epsilon,j,k) \in \Lambda_n}$, we set $f_{j'} = \sum_{(\epsilon,j,k) \in \Lambda_n} 2^{j(s+n/2)} |a_{j,k}^\epsilon| \chi(2^j x - k)$. Let

$$g_{j,j'}^k = \begin{cases} \sum_{\epsilon',k'} \frac{2^{j'(s+n/2)} |a_{j',k'}^{\epsilon'}|}{(1 + |k' - 2^{j'-j}k|)^{n+\gamma}}, & j \geq j', k \in \mathbb{Z}^n; \\ \sum_{\epsilon',k'} \frac{2^{j'(s+n/2)} |a_{j',k'}^{\epsilon'}|}{(1 + |k - 2^{j-j'}k'|)^{n+\gamma}}, & j < j', k \in \mathbb{Z}^n. \end{cases}$$

Yang [61] obtained the following result.

Lemma 2.10. [61, Chapter 5, Lemma 3.2] *For any $\gamma > n/A + 1$ and $x \in Q_{j,k}$, we have*

$$g_{j,j'}^k = \begin{cases} CM_A(f_{j'})(x), & \text{if } j \geq j'; \\ C2^{n(j'-j)/A}M_A(f_{j'})(x), & \text{if } j < j'. \end{cases}$$

2.3. Semigroup extension of $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$

In this section, we will introduce some tent spaces associated with $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$, which establishes a clear relationship among time t , frequency ξ and position x by wavelet methods. We would like to remind the reader that, for wavelets $\{\Phi_{j,k}^\epsilon : (\epsilon, j, k) \in \Lambda_n\}$, 2^j represents the range of frequency ξ and $2^{-j}k$ represents the range of position x in some sense. In the sequel, we only use the classical tensorial Meyer wavelets.

In the sequel, we use $N > 0$ to denote a fixed real number which is large enough. For fixed $\beta > 0$, we may choose a radial function $\phi \in \mathcal{S}(\mathbb{R}^n)$ (cf. [17, Lemma 1.1], [38, Chap. 3, §2]) such that there exists $C_\beta > 0$ satisfying

- (i) $\int_{\mathbb{R}^n} x^\gamma \phi(x) dx = 0$ for all $\gamma \in \mathbb{N}^n$;
- (ii) $\int_0^\infty \left(\widehat{\phi}(t^{1/(2\beta)}\xi)\right)^2 \frac{dt}{t} = 1$ for all $\xi \neq 0$;
- (iii) $\int_0^\infty \widehat{\phi}(t^{1/(2\beta)})e^{-t} \frac{dt}{t} = \frac{1}{C_\beta}$.

Define $\phi_t^\beta(x) = t^{-n/(2\beta)}\phi(t^{-1/(2\beta)}x)$. Then $\widehat{\phi}_t^\beta(\xi) = \widehat{\phi}(t^{1/(2\beta)}\xi)$, and hence

$$f(t, x) := e^{-t(-\Delta)^\beta} f(x) = K_t^\beta * f(x).$$

Applying the inverse Fourier transform, we can get

$$\widehat{f}(\xi) = C_\beta \int_0^\infty \widehat{\phi}(t^{1/(2\beta)})e^{-t} \frac{dt}{t} \widehat{f}(\xi) = C_\beta \int_0^\infty \widehat{\phi}(t^{1/(2\beta)}|\xi|)e^{-t|\xi|^{2\beta}} \widehat{f}(\xi) \frac{dt}{t}$$

which implies that

$$(2.8) \quad f(x) = C_\beta \int_0^\infty \int_{\mathbb{R}^n} f(t, x - y)\phi_t^\beta(y) dy \frac{dt}{t} := \pi_\phi f(\cdot, x).$$

For $(\epsilon, j, k) \in \Lambda_n$, let $a_{j,k}^\epsilon(t) = \langle f(t, \cdot), \Phi_{j,k}^\epsilon \rangle$ and $a_{j,k}^\epsilon = \langle f, \Phi_{j,k}^\epsilon \rangle$. Then

$$f = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^\epsilon \Phi_{j,k}^\epsilon \quad \text{and} \quad f(t, \cdot) = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^\epsilon(t) \Phi_{j,k}^\epsilon.$$

We first represent $a_{j,k}^\epsilon(t)$ by $a_{j',k'}^{\epsilon'}$. If $f(t, x) = K_t^\beta * f(x)$, then

$$\begin{aligned} a_{j,k}^\epsilon(t) &= \sum_{\epsilon', |j-j'| \leq 3, k'} a_{j',k'}^{\epsilon'} \langle K_t^\beta \Phi_{j',k'}^{\epsilon'}, \Phi_{j,k}^\epsilon \rangle \\ &= \sum_{\epsilon', |j-j'| \leq 3, k'} a_{j',k'}^{\epsilon'} \int e^{-t2^{2j\beta}|\xi|^{2\beta}} \widehat{\Phi}^{\epsilon'}(2^{j-j'}\xi) \widehat{\Phi}^\epsilon(\xi) e^{-i(2^j - j'k' - k)\xi} d\xi. \end{aligned}$$

By a simple application of integration by parts, we could control $a_{j,k}^\epsilon(t)$ by $\{a_{j',k'}^{\epsilon'}\}$ as follows.

Lemma 2.11. *There exists a fixed small constant $\tilde{c} > 0$ depending only on β and on the size of support for the Fourier transformation of Meyer wavelets such that*

(i) For $t2^{2\beta j} \geq 1$,

$$|a_{j,k}^\epsilon(t)| \leq C e^{-\tilde{c}t2^{2j\beta}} \sum_{\epsilon', |j-j'| \leq 3, k'} |a_{j',k'}^{\epsilon'}| \left(1 + |2^{j-j'}k' - k|\right)^{-N};$$

(ii) For $0 \leq t2^{2\beta j} \leq 1$,

$$|a_{j,k}^\epsilon(t)| \leq C \sum_{|j-j'| \leq 3} \sum_{\epsilon', k'} |a_{j',k'}^{\epsilon'}| \left(1 + |2^{j-j'}k' - k|\right)^{-N}.$$

Furthermore, if f is a function obtained by (2.8), then we could express $a_{j,k}^\epsilon$ by $\{a_{j',k'}^{\epsilon'}(t)\}$ as follows:

$$a_{j,k}^\epsilon = \int_{\mathbb{R}_+^{n+1}} \sum_{(\epsilon', j', k') \in \Lambda_n} a_{j',k'}^{\epsilon'}(t) (\phi_t^\beta * \Phi_{j',k'}^{\epsilon'}(x)) \Phi_{j,k}^\epsilon(x) dx \frac{dt}{t}.$$

Similarly, we apply integration by parts to obtain the following estimate.

Lemma 2.12.

$$|a_{j,k}^\epsilon| \leq C \sum_{|j-j'| \leq 3} \int_0^\infty \left(\max\{t2^{2j'\beta}, t^{-1}2^{-2j'\beta}\}\right)^{-N} \sum_{(\epsilon', k') \in F_n} \frac{|a_{j',k'}^{\epsilon'}(t)|}{(1 + |2^{j-j'}k' - k|)^N} \frac{dt}{t}.$$

We introduce some tent type spaces associated with $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$. For any $a(t, x)$ defined on \mathbb{R}_+^{n+1} , by wavelet theory there exists a family $\{a_{j,k}^\epsilon(t)\}_{(\epsilon, j, k) \in \Lambda_n}$ such that $a(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^\epsilon(t) \Phi_{j,k}^\epsilon(x)$. Let $\gamma_1, \gamma_2 \in \mathbb{R}$, $1 < p < \infty$, $m, \tilde{m} \in \mathbb{R}$, $m' > 0$ and $m'' \in \mathbb{Z}_+$. For any dyadic cube $Q_{j,k}$, we denote by $\chi_{j,k} = \chi(2^j \cdot -k)$ the characteristic function of $Q_{j,k}$.

The tent spaces associated with $\dot{F}_{p,q}^{\gamma_1, \gamma_2}$ consist of four parts. For the first two parts, we first take $\dot{F}_{p,q}^r$ -norms for space variable x , then take L^∞ -norm for variable t . We call them the Bloch type tent spaces. Precisely,

Definition 2.13. Given $\gamma_1, \gamma_2 \in \mathbb{R}$, $1 < p < \infty$ and $1 < q < \infty$. Let

$$\Lambda_{Q_r, 1} = \{(\epsilon, j, k) : j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}, Q_{j,k} \subset Q_r\}$$

and

$$\Lambda_{Q_r, 2} = \{(\epsilon, j, k) : -\log_2 r < j < -(\log_2 t)/(2\beta), Q_{j,k} \subset Q_r\}.$$

(i) We call $f \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,I}(\mathbb{R}_+^{n+1})$, if

$$\sup_{t>0} I_{p,q,m}^{\gamma_1,\gamma_2}(t) := \sup_{t,Q_r} |Q_r|^{\gamma_2/n-1/p} t^m \left\| \left(\sum_{(\epsilon,j,k) \in \Lambda_{Q_r,1}} 2^{qj(\gamma_1+n/2+2m\beta)} |a_{j,k}^\epsilon(t)|^q \chi_{j,k}(x) \right)^{1/q} \right\|_p < \infty.$$

(ii) We call $f \in \mathbb{F}_{p,q}^{\gamma_1,\gamma_2,II}(\mathbb{R}_+^{n+1})$, if for any $t > 0$

$$\sup_{t>0} II_{p,q}^{\gamma_1,\gamma_2}(t) := \sup_{t,Q_r} |Q_r|^{\gamma_2/n-1/p} \left\| \left(\sum_{(\epsilon,j,k) \in \Lambda_{Q_r,2}} 2^{qj(\gamma_1+n/2)} |a_{j,k}^\epsilon(t)|^q \chi_{j,k}(x) \right)^{1/q} \right\|_p < \infty.$$

For the other two parts, we first take L^p -norm for variable t , then take $\dot{F}_{p,q}^r$ -norm for variable x . We call them the L^p type tent spaces. Precisely,

Definition 2.14. Given $\gamma_1, \gamma_2 \in \mathbb{R}$, $1 < p < \infty$ and $1 < q < \infty$.

(i) $f \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}(\mathbb{R}_+^{n+1})$, if

$$III_{p,q,m}^{\gamma_1,\gamma_2} := \sup_{Q_r} |Q_r|^{\gamma_2/n-1/p} \left\| \left(\sum_{(\epsilon,j,k) \in \Lambda_Q^n} 2^{qj(\gamma_1+n/2+2m\beta)} \int_{2^{-2j\beta}}^{r^{2\beta}} |a_{j,k}^\epsilon(t)|^q t^{qm} \frac{dt}{t} \chi_{j,k}(x) \right) \right\|_p < \infty.$$

(ii) $f \in \mathbb{F}_{p,q}^{\gamma_1,\gamma_2,IV}(\mathbb{R}_+^{n+1})$, if

$$IV_{p,q,m'}^{\gamma_1,\gamma_2} := \sup_{Q_r} |Q_r|^{\gamma_2/n-1/p} \left\| \left(\sum_{(\epsilon,j,k) \in \Lambda_Q^n} 2^{qj(\gamma_1+n/2+2m'\beta)} \int_0^{2^{-2j\beta}} |a_{j,k}^\epsilon(t)|^q t^{qm'} \frac{dt}{t} \chi_{j,k}(x) \right) \right\|_p < \infty.$$

The tent spaces associated to $\dot{F}_{p,q}^{\gamma_1,\gamma_2}(\mathbb{R}^n)$ are defined as follows.

Definition 2.15. Let $\gamma_1, \gamma_2, m \in \mathbb{R}$, $m' > 0$ and $1 < p, q < \infty$. The tent spaces $\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$ are defined as

$$\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2} = \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,I} \cap \mathbb{F}_{p,q}^{\gamma_1,\gamma_2,II} \cap \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III} \cap \mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}.$$

The following lemma can be obtained immediately.

Lemma 2.16. Let $1 < p, q < \infty$, $\gamma_1, \gamma_2 \in \mathbb{R}$, $m > p$, $m' > 0$ and $a(x, t) \in \mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$.

(i) If $m > 0$, then $a(t, x)$ satisfies

$$\sup_{(\epsilon,j,k) \in \Lambda_n} \left\{ \sup_{t2^{2j\beta} \geq 1} (t2^{2j\beta})^m 2^{nj/2} 2^{j(\gamma_1-\gamma_2)} |a_{j,k}^\epsilon(t)| + \sup_{0 < t2^{2j\beta} \leq 1} 2^{nj/2} 2^{j(\gamma_1-\gamma_2)} |a_{j,k}^\epsilon(t)| \right\} < \infty.$$

(ii) If $-2\beta m < \gamma_1 - \gamma_2 < 0 < \beta$, then

$$\sup_{t>0} \sup_{j \in \mathbb{Z}, k \in \mathbb{Z}^n} t^{-(\gamma_1 - \gamma_2)/(2\beta)} 2^{nj/2} |\langle a(t, \cdot), \Phi_{j,k}^0 \rangle| < \infty.$$

Proof. (i) We divide the proof into two cases.

Case 1: $t2^{2j\beta} \geq 1$. Because $f \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I}$, for fixed j and k , we have

$$\begin{aligned} & |Q_r|^{\gamma_1/n - 1/p} t^m \left(\int 2^{pj(\gamma_1 + n/2 + 2m\beta)} |a_{j,k}^\epsilon(t)|^p \chi_{j,k}(x) dx \right)^{1/p} \\ & \lesssim |Q_r|^{\gamma_2/n - 1/p} t^m 2^{j(\gamma_1 + n/2 + 2m\beta)} |a_{j,k}^\epsilon(t)| 2^{-jn/p}. \end{aligned}$$

Take $r = 2^{-j}$. We can get $|a_{j,k}^\epsilon(t)| \lesssim 2^{j(\gamma_2 - \gamma_1 - n/2)} (t2^{2j\beta})^{-m}$.

Case 2: $0 < t2^{2j\beta} < 1$. Because $f \in \mathbb{F}_{p,q}^{\gamma_1, \gamma_2, II}$, we can obtain that

$$|Q_r|^{\gamma_2/n - 1/p} \left\| \left(\sum_{-\log_2 r < j < -(\log_2 t)/(2\beta)} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1 + n/2)} |a_{j,k}^\epsilon(t)|^q \chi_{j,k}(x) \right)^{1/q} \right\|_p \lesssim 1.$$

This implies that for fixed ϵ and k , $|a_{j,k}^\epsilon(t)| \lesssim 2^{-j(\gamma_1 - \gamma_2)} 2^{-jn/2}$.

Now we prove (ii). Because $\{\Phi_{j,k}^\epsilon\}$ are regular wavelets,

$$\langle a(t, \cdot), \Phi_{j,k}^0 \rangle = \left\langle \sum_{j' < j} a_{j',k'}^{\epsilon'}(t) \Phi_{j',k'}^{\epsilon'}, \Phi_{j,k}^0 \right\rangle.$$

For $0 < t2^{2j\beta} \leq 1$, a direct computation implies the desired result. For $t2^{2j\beta} > 1$, it suffices to divide into two subcases of $j' \leq -(\log_2 t)/(2\beta)$ and $-(\log_2 t)/(2\beta) < j' < j$. \square

In the rest of this paper, for any dyadic cube Q_r with side length r , we denote \tilde{Q}_r the dyadic cube which contains Q_r with side length $2^{\tau}r$. For all $w \in \mathbb{Z}^n$, write $\tilde{Q}_r^w := 2^{\tau}rw + \tilde{Q}_r$. Define

$$S_{Q_r}^{w,j'} = \left\{ (\epsilon', k') : (\epsilon', j', k') \in \Lambda_n \text{ and } Q_{j',k'} \subset \tilde{Q}_r^w \right\}.$$

Li-Yang-Zheng [32] established the following semigroup extension for $\dot{F}_{p,q}^{\gamma_1, \gamma_2}(\mathbb{R}^n)$.

Theorem 2.17. [32, Theorem 4.1] *Suppose that $\gamma_1, \gamma_2 \in \mathbb{R}$, $1 < p < m < \infty$, $\gamma_1 - \gamma_2 < 0 < \beta$, $m' > 0$ and $\tau + (\gamma_1 - \gamma_2)/(2\beta) > 0$.*

(i) *If $f \in \dot{F}_{p,q}^{\gamma_1, \gamma_2}$, then $f * K_t^\beta \in \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$;*

(ii) *The operator π_ϕ is a bounded and surjective operator from $\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$ to $\dot{F}_{p,q}^{\gamma_1, \gamma_2}$.*

For the Riesz transforms R_ℓ , $\ell = 1, 2, \dots, n$, let

$$a_{j,k,j',k'}^{\epsilon,\epsilon',\ell} = \left\langle \Phi_{j,k}^\epsilon, R_\ell \Phi_{j',k'}^{\epsilon'} \right\rangle, \quad (\epsilon', j', k') \in \Lambda_n.$$

If $|j - j'| \geq 2$, then $a_{j,k,j',k'}^{\epsilon,\epsilon',\ell} = 0$. By a similar method, we can obtain the boundedness of Riesz transforms on $\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$.

Theorem 2.18. *Given $1 < p, q < \infty$, $\gamma_1, \gamma_2 \in \mathbb{R}$, $m > p$ and $m' > 0$. The Riesz transforms R_1, R_2, \dots, R_n are bounded on $\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$.*

Proof.

$$R_\ell g(t, x) = \sum_{(\epsilon,j,k) \in \Lambda_n} g_{j,k}^\epsilon(t) R_\ell \Phi_{j,k}^\epsilon(x) := \sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^\epsilon(t) \Phi_{j,k}^\epsilon(x),$$

where

$$b_{j,k}^\epsilon(t) = \left\langle R_\ell g(t, \cdot), \Phi_{j,k}^\epsilon \right\rangle = \sum_{|j-j'| \leq 1} \sum_{\epsilon', k'} a_{j,k,j',k'}^{\epsilon,\epsilon'} g_{j',k'}^{\epsilon'}(t).$$

By Lemma 2.7, we have

$$\left| a_{j,k,j',k'}^{\epsilon,\epsilon'} \right| \lesssim 2^{-|j-j'|(n/2+N_0)} \left(\frac{2^{-j} + 2^{-j'}}{2^{-j} + 2^{-j'} + |2^{-j}k - 2^{-j'}k'|} \right)^{n+N_0}.$$

The rest of the proof is similar to that of Theorem 2.17. We omit the details and refer the reader to [32, Theorem 4.2]. □

3. Bilinear estimate on tent spaces associated with $\dot{F}_{p,q}^{\gamma_1,\gamma_2}$

3.1. Decomposition of bilinear operator

In Lemmas 3.2–3.7 below, we assume that the index $(\beta, p, q, \gamma_1, \gamma_2, m, m')$ satisfies the following conditions:

$$(3.1) \quad \left\{ \begin{array}{l} 1 < p < \infty, 1 < q < 2; \\ \beta > \frac{1}{2}, \gamma_1 = \gamma_2 - 2\beta + 1; \\ m > \max \left\{ p, \frac{n}{2\beta} \right\}; \\ 0 < m' < \min \left\{ 1, \frac{p}{2\beta} \right\}; \\ \frac{2\beta-2}{q} + \frac{n}{p} \left(1 - \frac{1}{q} \right) < \gamma_2 \leq \frac{n}{p} \end{array} \right. \quad \text{or} \quad \left\{ \begin{array}{l} 1 < p < \infty, 2 \leq q < \infty; \\ \beta > \frac{1}{2}, \gamma_1 = \gamma_2 - 2\beta + 1; \\ m > \max \left\{ p, \frac{n}{2\beta} \right\}; \\ 0 < m' < \min \left\{ 1, \frac{p}{2\beta} \right\}; \\ \beta - 1 + \frac{n}{2p} < \gamma_2 \leq \frac{n}{p}. \end{array} \right.$$

Here and henceforth, let

$$u(t, x) = \sum_{(\epsilon,j,k) \in \Lambda_n} u_{j,k}^\epsilon(t) \Phi_{j,k}^\epsilon(x) \quad \text{and} \quad v(t, x) = \sum_{(\epsilon,j,k) \in \Lambda_n} v_{j,k}^\epsilon(t) \Phi_{j,k}^\epsilon(x).$$

We can see that

$$e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell}(uv)(s, t, x) = \sum_{j' \in \mathbb{Z}} \sum_{i=1}^5 I_{j'}^{i, \ell}(s, t, x),$$

where

(3.2)

$$\begin{aligned} I_{j'}^{1, \ell}(u, v)(s, t, x) &= \sum_{\epsilon'} \sum_{k'} u_{j', k'}^{\epsilon'}(s) v_{j'-3, k''}^0(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j', k'}^{\epsilon'}(x) \Phi_{j'-3, k''}^0(x) \right), \\ I_{j'}^{2, \ell}(u, v)(s, t, x) &= \sum_{\epsilon'} \sum_{\epsilon'', k''} u_{j', k'}^{\epsilon'}(s) v_{j', k''}^{\epsilon''}(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j', k'}^{\epsilon'}(x) \Phi_{j', k''}^{\epsilon''}(x) \right), \\ I_{j'}^{3, \ell}(u, v)(s, t, x) &= \sum_{0 < |j' - j''| \leq 3} \sum_{\epsilon'} \sum_{\epsilon'', k''} u_{j', k'}^{\epsilon'}(s) v_{j'', k''}^{\epsilon''}(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j', k'}^{\epsilon'}(x) \Phi_{j'', k''}^{\epsilon''}(x) \right), \\ I_{j'}^{4, \ell}(u, v)(s, t, x) &= \sum_{\epsilon'} \sum_{k'} v_{j', k'}^{\epsilon'}(s) u_{j'-3, k''}^0(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j', k'}^{\epsilon'}(x) \Phi_{j'-3, k''}^0(x) \right). \end{aligned}$$

Write

$$I_\ell^i(s, t, x) = \sum_{j' \in \mathbb{Z}} I_{j'}^{i, \ell}(s, t, x) \quad \text{and} \quad I_\ell^i(u, v)(t, x) = \int_0^t I_\ell^i(s, t, x) ds.$$

For $\ell = 1, 2, \dots, n$, the bilinear terms $B_\ell(u, v)$ can be decomposed as

$$\begin{aligned} (3.3) \quad B_\ell(u, v)(t, x) &= \int_0^t e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell}(uv) ds =: \int_0^t \sum_{j' \in \mathbb{Z}} \sum_{i=1}^4 I_{j'}^{i, \ell}(s, t, x) ds \\ &=: \sum_{i=1}^4 \int_0^t I_\ell^i(s, t, x) ds =: \sum_{i=1}^4 I_\ell^i(u, v)(t, x). \end{aligned}$$

In order to estimate $B(u, v)$, we further decompose the terms $I_\ell^i(u, v)(t, x)$, $i = 1, 2, 3, 4$, respectively.

Decomposition of $I_\ell^1(u, v)$. The term $I_\ell^1(u, v)$ is decomposed according to two cases.

Case $[I_\ell^1]_1$: $t \geq 2^{-2j'\beta}$. For this case, we write $I_\ell^1(u, v)$ as the sum of the following three terms:

$$\begin{aligned} I_\ell^1(u, v)(t, x) &= \sum_{\epsilon', j', k'} \sum_{k''} \left(\int_0^{2^{-1-2j'\beta}} + \int_{2^{-1-2j'\beta}}^{t/2} + \int_{t/2}^t \right) \\ &\quad \times \left\{ u_{j', k'}^{\epsilon'}(s) v_{j'-3, k''}^0(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j', k'}^{\epsilon'}(x) \Phi_{j'-3, k''}^0(x) \right) \right\} ds \\ &=: I_\ell^{1,1}(u, v)(t, x) + I_\ell^{1,2}(u, v)(t, x) + I_\ell^{1,3}(u, v)(t, x). \end{aligned}$$

For $i = 1, 2, 3$, denote

$$I_\ell^{1,i}(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j, k}^{\epsilon, i}(t) \Phi_{j, k}^\epsilon(x).$$

Case $[I_\ell^1]_2$: $t < 2^{-2j\beta}$. For this case, we denote $a_{j,k}^{\epsilon,4}(t) = a_{j,k}^\epsilon(t)$ and then have

$$I_\ell^1(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^{\epsilon,4}(t) \Phi_{j,k}^\epsilon(x).$$

Decomposition of $I_\ell^2(u, v)$. We decompose $I_\ell^2(u, v)$ based on the relation between t and $2^{-2j\beta}$.

Case $[I_\ell^2]_1$: $t \geq 2^{-2j\beta}$. Naturally, $I_\ell^2(u, v)(t, x)$ can be divided into the following three terms:

$$\begin{aligned} I_\ell^2(u, v)(t, x) &= \sum_{j'} \sum_{\epsilon', k'} \sum_{\epsilon'', k''} \left(\int_0^{2^{-1-2j'\beta}} + \int_{2^{-1-2j'\beta}}^{t/2} + \int_{t/2}^t \right) \\ &\quad \times \left\{ u_{j',k'}^{\epsilon'}(s) v_{j',k''}^{\epsilon''}(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j',k'}^{\epsilon'}(x) \Phi_{j',k''}^{\epsilon''}(x) \right) \right\} ds \\ &=: I_\ell^{2,1}(u, v)(t, x) + I_\ell^{2,2}(u, v)(t, x) + I_\ell^{2,3}(u, v)(t, x). \end{aligned}$$

Case $[I_\ell^2]_2$: $t \leq 2^{-2j\beta}$. $I_\ell^2(u, v)$ can be decomposed into the sum of $II^4(u, v)$ and $II^5(u, v)$, where

$$\begin{aligned} I_\ell^2(u, v)(t, x) &= \sum_{j'} \sum_{\epsilon', k'} \sum_{\epsilon'', k''} \left(\int_0^{2^{-2j'\beta}} + \int_{2^{-2j'\beta}}^t \right) \\ &\quad \times \left\{ u_{j',k'}^{\epsilon'}(s) v_{j',k''}^{\epsilon''}(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j',k'}^{\epsilon'}(x) \Phi_{j',k''}^{\epsilon''}(x) \right) \right\} ds \\ &=: I_\ell^{2,4}(u, v)(t, x) + I_\ell^{2,5}(u, v)(t, x). \end{aligned}$$

For $i = 1, 2, 3, 4, 5$, set

$$I_\ell^{2,i}(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} b_{j,k}^{\epsilon,i}(t) \Phi_{j,k}^\epsilon(x).$$

Decompositions of $I_\ell^3(u, v)$. Without loss of generality, we assume that $0 < j' - j'' \leq 3$. Similarly, we have the following two cases.

Case $[I_\ell^3]_1$: $t \geq 2^{-2j\beta}$. This $I_\ell^3(u, v)$ can be divided into the following three terms:

$$\begin{aligned} I_\ell^3(u, v)(t, x) &= \sum_{0 < j' - j'' \leq 3} \sum_{\epsilon', k'} \sum_{\epsilon'', k''} \left(\int_0^{2^{-1-2j'\beta}} + \int_{2^{-1-2j'\beta}}^{t/2} + \int_{t/2}^t \right) \\ &\quad \times \left\{ u_{j',k'}^{\epsilon'}(s) v_{j',k''}^{\epsilon''}(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j',k'}^{\epsilon'}(x) \Phi_{j',k''}^{\epsilon''}(x) \right) \right\} ds \\ &=: I_\ell^{3,1}(u, v)(t, x) + I_\ell^{3,2}(u, v)(t, x) + I_\ell^{3,3}(u, v)(t, x). \end{aligned}$$

Case $[I_\ell^3]_2$: $t \leq 2^{-2j\beta}$. We divide $I_\ell^3(u, v)$ as follows.

$$\begin{aligned} I_\ell^3(u, v)(t, x) &= \sum_{0 < j' - j'' \leq 3} \sum_{\epsilon', k'} \sum_{\epsilon'', k''} \left(\int_0^{2^{-2j'\beta}} + \int_{2^{-2j'\beta}}^t \right) \\ &\quad \times \left\{ u_{j', k'}^{\epsilon'}(s) v_{j'', k''}^{\epsilon''}(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j', k'}^{\epsilon'}(x) \Phi_{j'', k''}^{\epsilon''}(x) \right) \right\} ds \\ &=: I_\ell^{3,4}(u, v)(t, x) + I_\ell^{3,5}(u, v)(t, x). \end{aligned}$$

For $i = 1, 2, 3, 4, 5$, denote

$$I_\ell^{3,i}(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} b_{j, k}^{\epsilon, i}(t) \Phi_{j, k}^\epsilon(x).$$

Decomposition of $I_j^{4,\ell}(u, v)$. It is easy to see that the terms $I_j^{1,\ell}(u, v)$ and $I_j^{4,\ell}(u, v)$ are symmetric associated with u and v . Hence for $I_\ell^4(u, v)$, we have a similar decomposition.

Case $[I_\ell^4]_1$: $t \geq 2^{-2j\beta}$. For this case, we write $I_\ell^4(u, v)$ as the sum of the following three terms:

$$\begin{aligned} I_\ell^4(u, v)(t, x) &= \sum_{\epsilon', j', k'} \sum_{k''} \left(\int_0^{2^{-1-2j'\beta}} + \int_{2^{-1-2j'\beta}}^{t/2} + \int_{t/2}^t \right) \\ &\quad \times \left\{ v_{j', k'}^{\epsilon'}(s) u_{j'' - 3, k''}^0(s) e^{-(t-s)(-\Delta)^\beta} \frac{\partial}{\partial x_\ell} \left(\Phi_{j', k'}^{\epsilon'}(x) \Phi_{j'' - 3, k''}^0(x) \right) \right\} ds \\ &=: I_\ell^{4,1}(u, v)(t, x) + I_\ell^{4,2}(u, v)(t, x) + I_\ell^{4,3}(u, v)(t, x). \end{aligned}$$

For $i = 1, 2, 3$, denote

$$I_\ell^{4,i}(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j, k}^{\epsilon, i}(t) \Phi_{j, k}^\epsilon(x).$$

Case $[I_\ell^4]_2$: $t < 2^{-2j\beta}$. For this case, we denote $a_{j, k}^{\epsilon, 4}(t) = a_{j, k}^\epsilon(t)$ and then have

$$I_\ell^4(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j, k}^{\epsilon, 4}(t) \Phi_{j, k}^\epsilon(x).$$

The above decompositions of $B(u, v)$ provide convenience of the bilinear estimates on $\mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}$. It is easy to see that the argument for $I_\ell^4(u, v)$ is similar to that for $I_\ell^1(u, v)$. Also the treatments of $I_\ell^3(u, v)$ is similar to that of $I_\ell^2(u, v)$. So we only show that

$$(3.4) \quad \begin{cases} I_\ell^1(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j, k}^\epsilon(t) \Phi_{j, k}^\epsilon(x) \in \mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}, \\ I_\ell^2(u, v)(t, x) = \sum_{(\epsilon, j, k) \in \Lambda_n} b_{j, k}^\epsilon(t) \Phi_{j, k}^\epsilon(x) \in \mathbb{F}_{p, q, m, m'}^{\gamma_1, \gamma_2}. \end{cases}$$

By the decompositions obtained above, (3.4) is equivalent to the following estimates:

$$\left\{ \begin{array}{l} \text{For } i = 1, 2, 3, 4, \quad \sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^{\epsilon, i}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}; \\ \text{For } i = 1, 2, 3, 4, 5, \quad \sum_{(\epsilon, j, k) \in \Lambda_n} b_{j,k}^{\epsilon, i}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}. \end{array} \right.$$

The demonstration will be concluded by the following Lemmas 3.2–3.7.

Remark 3.1. In the proofs of Lemmas 3.2–3.7 below, we only give the details for the case $q \geq 2$. The case $1 < q < 2$ can be dealt with similarly.

3.2. Bilinear estimate on Bloch type tent spaces

Lemma 3.2. *Let $(\beta, p, q, \gamma_1, \gamma_2, m, m')$ satisfy (3.1) and $u, v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$. Then*

- (i) $\sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^{\epsilon, i}(t) \Phi_{j,k}^\epsilon \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I}, i = 1, 2, 3;$
- (ii) $\sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^{\epsilon, 4}(t) \Phi_{j,k}^\epsilon \in \mathbb{F}_{p,q}^{\gamma_1, \gamma_2, II}.$

Proof. For simplicity, we assume $\|u\|_{\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}} = \|v\|_{\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}} = 1$. At first we prove (i). The proof of this part is divided into three parts. For $i = 1, 2, 3$, write

$$I_{a,i}(t) := |Q_r|^{\gamma_2/n-1/p} t^m \times \left\| \left(\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m\beta)} |a_{j,k}^{\epsilon, i}(t)|^q \chi_{j,k}(x) \right)^{1/q} \right\|_p.$$

Case 3.2.1: $\sum_{(\epsilon, j, k) \in \Lambda_n} a_{j,k}^{\epsilon, 1}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I}$. Because $v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$, by Lemma 2.16, one has $|v_{j'-3, k''}^0(s)| \lesssim s^{-(\gamma_2-\gamma_1)/(2\beta)} 2^{-nj'/2}$ and

$$\begin{aligned} |a_{j,k}^{\epsilon, 1}(t)| &\lesssim 2^{nj/2+j} \sum_{|j-j'| \leq 1} \sum_{\epsilon', k'} \int_0^{2^{-1-2j'\beta}} e^{-ct2^{2j\beta}} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k' - k|)^N} \frac{|v_{j'-3, k''}^0(s)|}{(1+|2^{j-j'+3}k'' - k|)^N} ds \\ &\lesssim 2^j \sum_{|j-j'| \leq 1} \sum_{\epsilon', k'} \int_0^{2^{-1-2j'\beta}} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k' - k|)^N} e^{-ct2^{2j\beta}} s^{1/(2\beta)-1} ds. \end{aligned}$$

By $|j - j'| \leq 1$ and Hölder’s inequality, it is easy to see that

$$\begin{aligned} &\left(\int_0^{2^{-1-2j'\beta}} \sum_{\epsilon', k'} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k' - k|)^N} s^{1/(2\beta)-1} ds \right)^q \\ &\lesssim \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \sum_{Q_{j',k'} \subset Q_{j,k}^w} \frac{1}{(1+|2^{j-j'}k' - k|)^N} 2^{-qj'} \left(\int_0^{2^{-1-2j'\beta}} |u_{j',k'}^{\epsilon'}(s)|^q s^{qm'} \frac{ds}{s} \right). \end{aligned}$$

Let

$$g_{j'} = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_0^{2^{-1-2j'\beta}} \left| u_{j',k'}^{\epsilon'}(s) \right|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \chi_{j',k'}(x).$$

We apply Lemma 2.10 to obtain

$$\begin{aligned} I_{a,1}(t) &\lesssim |Q_r|^{\gamma_2/n-1/p} t^m \\ &\times \left\| \left[\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} \sum_{|j-j'| \leq 1} 2^{qj(\gamma_1+n/2+2m\beta)} 2^{qj} e^{-cqt2^{2j\beta}} \right. \right. \\ &\quad \left. \left. \times \sum_{|j-j'| \leq 1} \left(\int_0^{2^{-1-2j'\beta}} \sum_{\epsilon',k'} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} s^{1/(2\beta)-1} ds \right)^q \chi_{j,k}(x) \right]^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left[\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} (M_A(g_{j'})(x))^q \right]^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left[\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} |g_{j'}(x)|^q \right]^{1/q} \right\|_p \lesssim 1. \end{aligned}$$

Case 3.2.2: $\sum_{(\epsilon,j,k) \in \Lambda_n} a_{j,k}^{\epsilon,2}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,I}$. Because $|v_{j',-3,k'}^0(s)| \lesssim s^{1/(2\beta)-1} 2^{-nj'/2}$, we get

$$\left| a_{j,k}^{\epsilon,2}(t) \right| \lesssim 2^j \sum_{|j-j'| \leq 1} \sum_{\epsilon',k'} \int_{2^{-1-2j'\beta}}^{t/2} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} e^{-ct2^{2j\beta}} s^{1/(2\beta)-1} ds.$$

Then

$$\begin{aligned} I_{a,2}(t) &\lesssim |Q_r|^{\gamma_2/n-1/p} t^m \\ &\times \left\| \left(\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m\beta)} \right. \right. \\ &\quad \left. \left. \times \left(2^j \sum_{|j-j'| \leq 1} \sum_{\epsilon',k'} \int_{2^{-1-2j'\beta}}^{t/2} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} e^{-ct2^{2j\beta}} s^{1/(2\beta)-1} ds \right)^q \chi_{j,k}(x) \right)^{1/q} \right\|_p \\ &\lesssim |Q_r|^{\gamma_2/n-1/p} t^m \left\| \left(\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m\beta)} 2^{qj} e^{-cqt2^{2j\beta}} \right. \right. \\ &\quad \left. \left. \times \sum_{|j-j'| \leq 1} \left(\sum_{\epsilon',k'} \int_{2^{-1-2j'\beta}}^{t/2} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} s^{1/(2\beta)-1} ds \right)^q \chi_{j,k}(x) \right)^{1/q} \right\|_p. \end{aligned}$$

At first we assume $t > r^{2\beta}$. Because $u \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, l}$, let $r = 2^{-j}$ and $Q_r = Q_{j,k}^w$. For fixed j' ,

$$2^{-nj(\gamma_2/n-1/p)} s^m \left\| \left(\sum_{Q_{j',k'} \subset Q_{j,k}^w} 2^{qj'(\gamma_1+n/2+2m\beta)} \left| u_{j',k'}^{\epsilon'}(s) \right|^q \chi_{j',k'}(x) \right)^{1/q} \right\|_p \lesssim 1.$$

For $s \geq 2^{-1-2j'\beta}$, if $s2^{2j'\beta} \geq 1$, then $\left| u_{j',k'}^{\epsilon'}(s) \right| \lesssim 2^{j'(\gamma_2-\gamma_1-n/2)} (s2^{2j'\beta})^{-m}$. If $s2^{2j'\beta} < 1$,

$$\left| u_{j',k'}^{\epsilon'}(s) \right| \leq 2^{j'(\gamma_2-\gamma_1-n/2)} \leq 2^{j'(\gamma_2-\gamma_1-n/2)} (s2^{2j'\beta})^{-m}.$$

Because $|j - j'| \leq 1$, the above estimate for $\left| u_{j',k'}^{\epsilon'}(s) \right|$ and Hölder's inequality imply that

$$\begin{aligned} I_{a,2}(t) &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1 + |w|)^{-N} \sum_{j' \geq \max\{-\log_2 r_w, -(\log_2 t)/(2\beta)-1\}} 2^{qj'(\gamma_1+n/2)} \\ &\quad \times (t2^{2j'\beta})^{m+1/q} e^{-ct2^{2j'\beta}} 2^{j'(\gamma_2-\gamma_1-n/2)} \left\| \left(\sum_{Q_{j',k'} \subset Q_r^w} \chi_{j',k'}(x) \right)^{1/q} \right\|_p \\ &\lesssim |Q_r|^{\gamma_2/n} \sum_{w \in \mathbb{Z}^n} (1 + |w|)^{-N} t^{-\gamma_2/(2\beta)} \sum_{j' \geq -(\log_2 t)/(2\beta)-1} (t2^{2j'\beta})^{m+1/q+\gamma_2/(2\beta)} e^{-ct2^{2j'\beta}} \lesssim 1, \end{aligned}$$

where we have used the fact that $t > r^{2\beta}$ in the last step.

Now we deal with the case $t \leq r^{2\beta}$. Let

$$g_{j'} = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_{2^{-1-2j'\beta}}^{t/2} \left| u_{j',k'}^{\epsilon'}(s) \right|^q (s2^{2j'\beta})^{qm} \frac{ds}{s} \right)^{1/q} \chi_{j',k'}(x).$$

Note that $m \gg q/(2\beta)$. By $t < r^{2\beta}$, we use Lemma 2.10 and Hölder's inequality to obtain

$$\begin{aligned} I_{a,2}(t) &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1 + |w|)^{-N} \\ &\quad \times \left\| \left[\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} \sum_{|j-j'| \leq 1} 2^{q(j-j')} \right. \right. \\ &\quad \times \left. \sum_{Q_{\epsilon',k'} \subset Q_{j,k}^w} \left(1 + \left| 2^{j-j'} k' - k \right| \right)^{-N} \left(\int_{2^{-1-2j'\beta}}^{t/2} \left| u_{j',k'}^{\epsilon'}(s) \right|^q (s2^{2j'\beta})^{qm} \frac{ds}{s} \right) \chi_{j,k}(x) \right]^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^N} \left\| \left(\sum_{j' \geq -\log_2 r_w} (M_A(g_{j'})(x))^q \right)^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^N} \left\| \left(\sum_{j' \geq -\log_2 r_w} |g_{j'}(x)|^q \right)^{1/q} \right\|_p \lesssim 1. \end{aligned}$$

Case 3.2.3: $\sum_{(\epsilon, j, k) \in \Lambda_n} a_{j, k}^{\epsilon, 3}(t) \Phi_{j, k}^\epsilon(x) \in \mathbb{F}_{p, q, m}^{\gamma_1, \gamma_2, I}$. We have

$$|a_{j, k}^{\epsilon, 3}(t)| \lesssim 2^j \sum_{|j-j'| \leq 1} \sum_{\epsilon', k', k''} \int_{t/2}^t e^{-c(t-s)2^{2j\beta}} \frac{|u_{j', k'}^{\epsilon'}(s)|}{(1 + |2^{j-j'} k' - k|)^N} s^{1/(2\beta)-1} ds.$$

By Hölder's inequality and the fact $|j - j'| \leq 1$, we can use the above estimate to deduce that

$$\begin{aligned} I_{a,3}(t) &\lesssim |Q_r|^{\gamma_2/n-1/p} t^m \left\| \left[\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m\beta)} 2^{qj} \right. \right. \\ &\quad \times \left. \sum_{|j-j'| \leq 1} \left| \sum_{\epsilon', k'} \int_{t/2}^t \frac{|u_{j', k'}^{\epsilon'}(s)|}{(1 + |2^{j-j'} k' - k|)^N} e^{-c(t-s)2^{2j\beta}} s^{1/(2\beta)-1} ds \right|^q \chi_{j,k}(x) \right]^{1/q} \right\|_p \\ &\lesssim |Q_r|^{\gamma_2/n-1/p} \left\| \left[\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m\beta)} \right. \right. \\ &\quad \times \sum_{|j-j'| \leq 1} \sum_{\epsilon', k'} \left(1 + |2^{j-j'} k' - k|\right)^{-N} \\ &\quad \times \left. \left. \left(\int_{t/2}^t |u_{j', k'}^{\epsilon'}(s)|^q e^{-c(t-s)2^{2j\beta}} s^{qm} \frac{ds}{s} \right) (t2^{2j\beta})^{q/(2\beta)-(q-1)} \chi_{j,k}(x) \right]^{1/q} \right\|_p. \end{aligned}$$

Let

$$\begin{aligned} g_{j'} &= 2^{j'(\gamma_1+n/2+2m\beta)} \\ &\quad \times \sum_{Q_{j', k'} \subset Q_r^w} \left[(t2^{2j'\beta})^{q/(2\beta)-(q-1)} \left(\int_{t/2}^t |u_{j', k'}^{\epsilon'}(s)|^q e^{-c(t-s)2^{2j'\beta}} s^{qm} \frac{ds}{s} \right) \right]^{1/q} \chi_{j', k'}(x). \end{aligned}$$

Hence Lemma 2.10 implies that

$$\sum_{\epsilon', k'} \frac{t^{q/(2\beta)-(q-1)}}{(1 + |2^{j-j'} k' - k|)^N} \left(\int_{t/2}^t e^{-c(t-s)2^{2j\beta}} |u_{j', k'}^{\epsilon'}(s)|^q s^{qm} \frac{ds}{s} \right) \lesssim \sum_{w \in \mathbb{Z}^n} \frac{1}{(1 + |w|)^N} (M_A(g_{j'})(x))^q.$$

We can get

$$\begin{aligned} I_{a,3}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^{N'}} \left\| \left[\sum_{j' \geq \max\{-\log_2 r_w, -(\log_2 t)/(2\beta)-1\}} 2^{qj'(\gamma_1+n/2+2m\beta)} (t2^{2j'\beta})^{q/(2\beta)-q+1} \right. \right. \\ &\quad \times \left. \sum_{Q_{j', k'} \subset Q_r^w} \int_{t/2}^t e^{-c(t-s)2^{2j'\beta}} |u_{j', k'}^{\epsilon'}(s)|^q s^{qm} \frac{ds}{s} \chi_{j', k'}(x) \right]^{1/q} \right\|_p. \end{aligned}$$

If $t > r^{2\beta}$, by Lemma 2.16, we can see that $|u_{j', k'}^{\epsilon'}(s)| \lesssim 2^{j'(\gamma_2-\gamma_1-n/2)} (s2^{2j'\beta})^{-m}$. This

gives

$$\begin{aligned}
 I_{a,3}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left[\sum_{j' \geq \max\{-\log_2 r_w, -(\log_2 t)/(2\beta)-1\}} (t2^{2j'\beta})^{q/(2\beta)-(q-1)-1} \right. \right. \\
 &\quad \left. \left. \times (t2^{2j'\beta})^{(q\gamma_2)/(2\beta)} t^{-(q\gamma_2)/(2\beta)} \sum_{Q_{j',k'} \subset Q_r^w} \chi_{j',k'}(x) \right] \right\|_p^{1/q} \\
 &\lesssim \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \sum_{j' \geq \max\{-\log_2 r_w, -(\log_2 t)/(2\beta)-1\}} (t2^{2j'\beta})^{1/(2\beta)-1+\gamma_2/(2\beta)} \lesssim 1,
 \end{aligned}$$

where we have used the facts $t > r^{2\beta}$ and $\gamma_1 < 0$, respectively.

For $t < r^{2\beta}$, because $u \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I}$, we have

$$|Q_r|^{\gamma_2/n-1/p} s^m \left\| \left(2^{qj'(\gamma_1+n/2+2m\beta)} \left| u_{j',k'}^{\epsilon'}(s) \right|^q \chi_{j',k'}(x) \right)^{1/q} \right\|_p \lesssim 1.$$

Because $Q_{j',k'} \subset Q_r^w$, that is, $2^{-j'} < r_w$, the above estimate implies that

$$\left| u_{j',k'}^{\epsilon'}(s) \right| (s2^{2j'\beta})^m \lesssim |Q_r|^{-\gamma_2/n} 2^{-j'(\gamma_1+n/2)}.$$

We can obtain

$$\begin{aligned}
 I_{a,3}(t) &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \left\| \left[\sum_{j' \geq \max\{-\log_2 r_w, -(\log_2 t)/(2\beta)-1\}} 2^{qj'(\gamma_1+2m\beta+n/2)} \right. \right. \\
 &\quad \times \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_{t/2}^t e^{-c(t-s)2^{2j'\beta}} s^{qm} 2^{-qj'(\gamma_1+n/2)} (s2^{2j'\beta})^{-qm} |Q_r|^{-q\gamma_2/n} \frac{ds}{s} \right) \\
 &\quad \left. \left. \times (t2^{2j'\beta})^{q/(2\beta)-(q-1)} \chi_{j',k'}(x) \right] \right\|_p^{1/q} \\
 &\lesssim |Q_r|^{-1/p} \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \\
 &\quad \times \left\| \left[\sum_{j' \geq \max\{-\log_2 r_w, -(\log_2 t)/(2\beta)-1\}} (t2^{2j'\beta})^{q/(2\beta)-q} \sum_{Q_{j',k'} \subset Q_r^w} \chi_{j',k'}(x) \right] \right\|_p^{1/q} \lesssim 1.
 \end{aligned}$$

Now we prove (ii). Write

$$II_{a,4}(t) := |Q_r|^{\gamma_2/n-1/p} \left\| \left(\sum_{-\log_2 r < j < -(\log_2 t)/(2\beta)} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} \left| a_{j,k}^{\epsilon,4}(t) \right|^q \chi_{j,k}(x) \right)^{1/q} \right\|_p.$$

For this case, we can see that

$$\begin{aligned} |a_{j,k}^{\epsilon,4}(t)| &\lesssim 2^{nj/2+j} \sum_{|j-j'|\leq 1} \sum_{\epsilon',k',k''} \int_0^t \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} \frac{|v_{j'-3,k''}^0(s)|}{(1+|k'-k''|)^N} e^{-c(t-s)2^{2j\beta}} ds \\ &\lesssim 2^j \sum_{|j-j'|\leq 1} \sum_{\epsilon',k'} \int_0^t \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} e^{-c(t-s)2^{2j\beta}} s^{1/(2\beta)-1} ds. \end{aligned}$$

By Hölder's inequality and the facts that $t \leq 2^{-2j\beta}$ and $|j - j'| \leq 1$, we obtain

$$\begin{aligned} &II_{a,4}(t) \\ &\lesssim |Q_r|^{\gamma_2/n-1/p} \left\| \left[\sum_{-\log_2 r < j' \leq -(\log_2 t)/(2\beta)} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} 2^{qj} \right. \right. \\ &\quad \times \left. \sum_{|j-j'|\leq 1} \left| \sum_{\epsilon',k'} \int_0^t e^{-c(t-s)2^{2j\beta}} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} s^{1/(2\beta)-1} ds \right|^q \chi_{j,k}(x) \right]^{1/q} \right\|_p \\ &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \left\| \left\{ \sum_{-\log_2 r < j' \leq -(\log_2 t)/(2\beta)} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} \sum_{|j-j'|\leq 1} \right. \right. \\ &\quad \times \left. \left[\sum_{Q_{\epsilon',k'} \subset Q_{j,k}^w} \frac{1}{(1+|2^{j-j'}k'-k|)^N} \left(\int_0^t |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \right]^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p. \end{aligned}$$

Let

$$g_{j'}(x) = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_0^t |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \chi_{j',k'}(x).$$

Finally, by Lemma 2.10, we have

$$\begin{aligned} II_{a,4}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \\ &\quad \times \left\| \left\{ \sum_{-\log_2 r < j' \leq -(\log_2 t)/(2\beta)} \sum_{|j-j'|\leq 1} 2^{q(j-j')(\gamma_1+n/2)} [M_A(g_{j'})(x)]^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{-\log_2 r w \leq j' \leq -(\log_2 t)/(2\beta)-1} |g_{j'}(x)|^q \right\}^{1/q} \right\|_p \lesssim 1, \end{aligned}$$

where in the last inequality, we have used the fact $t \leq 2^{-2j'\beta}$. □

Lemma 3.3. *Let $(\beta, p, q, \gamma_1, \gamma_2, m, m')$ satisfy (3.1) and $u, v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$. Then*

$$\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,i}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,I}, \quad i = 1, 2, 3.$$

Proof. For $i = 1, 2, 3$, write

$$I_{b,i}(t) := |Q_r|^{\gamma_2/n-1/p} t^m \times \left\| \left(\sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m\beta)} \left| a_{j,k}^{\varepsilon,i}(t) \right|^q \chi_{j,k}(x) \right)^{1/q} \right\|_p.$$

We divide the proof into three cases.

Case 3.3.1: Under $2 < q < \infty$, $\sum_{(\varepsilon,j,k) \in \Lambda_n} b_{j,k}^{\varepsilon,1}(t) \Phi_{j,k}^\varepsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I}$. For $0 < s < 2^{-1-2j'\beta}$ and $r = 2^{-j}$, the fact $v \in \mathbb{F}_{p,q}^{\gamma_1, \gamma_2, II}$ implies that

$$\left| v_{j',k''}^{\varepsilon''}(s) \right| \lesssim 2^{j(\gamma_2-n/p)} 2^{-j'(\gamma_1+n/2-n/p)}.$$

By Hölder's inequality, we have

$$\begin{aligned} \left| b_{j,k}^{\varepsilon,1}(t) \right| &\lesssim 2^{nj/2+j} e^{-ct2^{2j\beta}} \sum_{j < j'+2} \sum_{\varepsilon', k'} \sum_{\varepsilon'', k''} \int_0^{2^{-1-2j'\beta}} \frac{|u_{j',k'}^{\varepsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} \frac{|v_{j',k''}^{\varepsilon''}(s)|}{(1+|k'-k''|)^N} ds \\ &\lesssim \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} 2^{nj/2+j} 2^{j(\gamma_2-n/p)} e^{-ct2^{2j\beta}} \sum_{j < j'+2} 2^{-j'(\gamma_1+n/2-n/p)} 2^{-2j'\beta(1-1/q)} \\ &\quad \times \sum_{Q_{j',k'} \subset Q_{j,k}^w} \left(1+|2^{j-j'}k'-k|\right)^{-N} \left(\int_0^{2^{-1-2j'\beta}} |u_{j',k'}^{\varepsilon'}(s)|^q ds \right)^{1/q}, \end{aligned}$$

Let

$$g_{j'} = \sum_{Q_{j',k'} \subset Q_{j,k}^w} \left(1+|2^{j-j'}k'-k|\right)^{-N} \left(\int_0^{2^{-1-2j'\beta}} |u_{j',k'}^{\varepsilon'}(s)|^q ds \right)^{1/q}.$$

Hence Cauchy-Schwartz's inequality gives

$$\begin{aligned} I_{b,1}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} 2^{qj(\gamma_1+n+1+2m\beta)} (t2^{2j\beta})^{qm} \right. \right. \\ &\quad \left. \left. \times e^{-ct2^{2j\beta}} 2^{qj(\gamma_2-n/p)} \left[\sum_{j < j'+2} 2^{-j'(\gamma_1+n/2-n/p)} 2^{-2j'\beta(1-1/q)} |g_{j'}| \right]^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p. \end{aligned}$$

Let

$$f_{j'} := \sum_{Q_{j',k'} \subset Q_{j,k}^w} \frac{2^{j'(\gamma_1+n/2-2\beta/q)}}{(1+|2^{j-j'}k'-k|)^N} \left(\int_0^{2^{-1-2j'\beta}} |u_{j',k'}^{\varepsilon'}(s)|^q (s2^{2j'\beta})^{m'} ds \right)^{1/q} \chi_{j',k'}(x).$$

Then

$$\begin{aligned}
 I_{b,1}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} 2^{qj(\gamma_1+n+1+2m\beta+\gamma_2-n/p)} \right. \right. \\
 &\quad \times \left. \sum_{j < j'+2} 2^{(\delta+qn)(j'-j)} 2^{-qj'(2\gamma_1+n-n/p+2\beta)} (M_A(f_{j'})(x))^q \chi_{j,k}(x) \right\}^{1/q} \Big\|_p \\
 &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \sum_{j < j'+2} 2^{(j'-j)[\delta/q-(\gamma_1+\gamma_2+1-n/p)]} \\
 &\quad \times \left\| \left\{ \sum_{j' \geq -\log_2 r} \sum_{Q_{j',k'} \subset Q_r^w} 2^{qj'(\gamma_1+n/2)} \int_0^{2^{-1-2j'\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{m'} ds \chi_{j',k'}(x) \right\}^{1/q} \right\|_p \\
 &\lesssim \|u\|_{\mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}}.
 \end{aligned}$$

Case 3.3.2: Under $2 < q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,2}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,I}$. At first, we assume that $t > r^{2\beta}$. Because $|v_{j',k''}^{\epsilon''}(s)| \lesssim 2^{j(\gamma_2-n/p)} 2^{-j'(\gamma_1+n/2-n/p)} (s2^{2j'\beta})^{-m}$, we obtain that

$$\begin{aligned}
 |b_{j,k}^{\epsilon,2}(t)| &\lesssim 2^{nj/2+j} e^{-ct2^{2j\beta}} \sum_{j < j'+2} \sum_{\epsilon',k'} \sum_{\epsilon'',k''} \int_{2^{-1-2j'\beta}}^{t/2} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} \frac{|v_{j',k''}^{\epsilon''}(s)|}{(1+|k'-k''|)^N} ds \\
 &\lesssim 2^{nj/2+j} e^{-ct2^{2j\beta}} \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \sum_{w' \in \mathbb{Z}^n} (1+|w-w'|)^{-N} 2^{2j(\gamma_2-n/p)} \\
 &\quad \times \sum_{j < j'+2} 2^{-2j'(\gamma_1+n/2-n/p)} 2^{n(j'-j)} \int_{2^{-1-2j'\beta}}^{t/2} (s2^{2j'\beta})^{-2m} ds \\
 &\lesssim 2^{-nj/2+j} 2^{2\gamma_2j} 2^{-2j\gamma_1} e^{-ct2^{2j\beta}} 2^{-2j\beta},
 \end{aligned}$$

where we have used the fact that $\gamma_1 + \beta - n/p > 0$. By $\gamma_2 < n/p$, the above estimate implies that

$$I_{b,2}(t) \lesssim |Q_r|^{\gamma_2/n-1/p} t^m \sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} 2^{j(2\gamma_2-\gamma_1+2m\beta+1-2\beta)} e^{-ct2^{2j\beta}} 2^{-nj/p} \lesssim 1.$$

Now, we suppose $t \leq r^{2\beta}$. Because $v \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III} \cap \mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}$, if $Q_{j',k'} \subset Q_{j,k}^w$,

$$\left(\int_{2^{-1-2j'\beta}}^{r^{2\beta}} |v_{j',k''}^{\epsilon''}(s)|^q (s2^{2j'\beta})^{qm} \frac{ds}{s} \right)^{1/q} \lesssim r^{2j(\gamma_2-n/p)} 2^{-j'(\gamma_1+n/2-n/p)}.$$

It can be deduced that

$$\begin{aligned} |b_{j,k}^{\epsilon,2}(t)| &\lesssim 2^{j(n/2+1)} e^{-ct2^{2j\beta}} \sum_{j < j'+2} \sum_{\epsilon', k'} \sum_{\epsilon'', k''} \int_{2^{-1-2j'\beta}}^{t/2} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} \frac{|v_{j',k''}^{\epsilon''}(s)|}{(1+|k'-k''|)^N} ds \\ &\lesssim 2^{j(n/2+1+\gamma_2-n/p)} e^{-ct2^{2j\beta}} \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \sum_{j < j'+2} 2^{-j'(2\beta+\gamma_1+n/2-n/p)} \\ &\quad \times \sum_{Q_{j',k'} \subset Q_{j,k}^w} \left(1+|2^{j-j'}k'-k|\right)^{-N} \left(\int_{2^{-2j'\beta}}^{r^{2\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm} \frac{ds}{s}\right)^{1/q}. \end{aligned}$$

Let

$$g_{j'} := \sum_{Q_{j',k'} \subset Q_{j,k}^w} \left(1+|2^{j-j'}k'-k|\right)^{-N} \left(\int_{2^{-2j'\beta}}^{r^{2\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm} \frac{ds}{s}\right)^{1/q}.$$

This implies that

$$\begin{aligned} I_{b,2}(t) &\lesssim t^m \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n+1+2m\beta)} \right. \right. \\ &\quad \left. \left. \times 2^{qj(\gamma_2-n/p)} e^{-ct2^{2j\beta}} \sum_{j < j'+2} 2^{\delta(j'-j)} 2^{-qj'(2\beta+\gamma_1+n/2-n/p)} |g_{j'}|^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p \\ &\lesssim t^m \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} 2^{qj(\gamma_1+n+1)} 2^{qj(\gamma_2-n/p)} \right. \right. \\ &\quad \left. \left. \times \sum_{j < j'+2} 2^{\delta(j'-j)} 2^{-qj'(\gamma_1+n/2-n/p+2\beta)} |g_{j'}|^q \right\}^{1/q} \right\|_p. \end{aligned}$$

Let

$$f_{j'} = \sum_{Q_{j',k'} \subset Q_r^w} 2^{j'(\gamma_1+n/2)} \left(\int_{2^{-1-2j'\beta}}^{r^{2\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm} \frac{ds}{s}\right)^{1/q} \chi_{j',k'}(x).$$

By Lemma 2.10, we can see that $|g_{j'}|^q \lesssim 2^{-qj'(\gamma_1+n/2)} 2^{qm(j'-j)} (M_A(f_{j'})(x))^q$ and hence

$$\begin{aligned} I_{b,2}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \\ &\quad \times \left\| \left\{ \sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} \sum_{j < j'+2} 2^{(j'-j)[\delta-(\gamma_1+\gamma_2+1-n/p)]} (M_A(f_{j'})(x))^q \right\}^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \sum_{j < j'+2} 2^{(j'-j)[\delta-(\gamma_1+\gamma_2+1-n/p)]} \left\| \left\{ \sum_{j' \geq -\log_2 r_w} |f_{j'}(x)|^q \right\}^{1/q} \right\|_p \\ &\lesssim \|u\|_{\mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, III}} + \|u\|_{\mathbb{F}_{p,q,m'}^{\gamma_1, \gamma_2, IV}} \end{aligned}$$

Case 3.3.3: Under $2 < q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,3}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I}$. Because $v \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, I} \cap \mathbb{F}_{p,q}^{\gamma_1, \gamma_2, II}$, then

$$\left(\int_{t/2}^t |v_{j',k''}^{\epsilon''}(s)|^q ds \right)^{1/q} \lesssim 2^{j(\gamma_2-n/p)} 2^{-j'(\gamma_1+n/2-n/p)} (t2^{2j'\beta})^{-m} t^{1/q}.$$

By the above estimate and Hölder’s inequality, we can get

$$\begin{aligned} |b_{j,k}^{\epsilon,3}(t)| &\lesssim 2^{nj/2+j} \sum_{j < j'+2} \sum_{\epsilon',k'} \sum_{\epsilon'',k''} \int_{t/2}^t \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} \frac{|v_{j',k''}^{\epsilon''}(s)|}{(1+|k'-k''|)^N} e^{-c(t-s)2^{2j\beta}} ds \\ &\lesssim 2^{j(n/2+1+\gamma_2-n/p)} t^{1-1/q} \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \sum_{j < j'+2} 2^{-j'(\gamma_1+n/2-n/p)} (t2^{2j'\beta})^{-m} \\ &\quad \times \sum_{\epsilon',k'} (1+|2^{j-j'}k'-k|)^{-N} \left(\int_{t/2}^t |u_{j',k'}^{\epsilon'}(s)|^q ds \right)^{1/q}. \end{aligned}$$

By $\gamma_2 \leq n/p$ and $j \geq -\log_2 r$, we can get

$$\begin{aligned} I_{b,3}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} t^m \left\| \left\{ \sum_{j \geq \max\{-\log_2 r, -(\log_2 t)/(2\beta)\}} 2^{qj(\gamma_1+n+1+2m\beta+\gamma_2-n/p)} t^{q-1} \right. \right. \\ &\quad \times \sum_{j < j'+2} 2^{\delta(j'-j)} 2^{-qj'(\gamma_1+n/2-n/p)} (t2^{2j'\beta})^{-qm} \\ &\quad \times \left. \left[\sum_{\epsilon',k'} \frac{1}{(1+|2^{j-j'}k'-k|)^N} \left(\int_{t/2}^t |u_{j',k'}^{\epsilon'}(s)|^q ds \right)^{1/q} \right]^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p \\ &\lesssim \sum_{j' \geq -(\log_2 t)/(2\beta)-2} (t2^{2j'\beta})^{1-m} \sum_{j < j'+2} 2^{(j'-j)[\delta/q-(\gamma_1+\gamma_2+1-2n/p+2m\beta)]} \lesssim 1. \end{aligned}$$

This completes the proof. □

Lemma 3.4. *Let $(\beta, p, q, \gamma_1, \gamma_2, m, m')$ satisfy (3.1) and $u, v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$. Then*

$$\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,i}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q}^{\gamma_1, \gamma_2, II}, \quad i = 4, 5.$$

Proof. For $i = 4, 5$, write

$$II_{b,i}(t) := |Q_r|^{\gamma_2/n-1/p} \left\| \left(\sum_{-\log_2 r < j < -(\log_2 t)/(2\beta)} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} |b_{j,k}^{\epsilon,i}(t)|^q \chi_{j,k}(x) \right)^{1/q} \right\|_p.$$

We still divide the proof into two cases.

Case 3.4.1: Under $2 \leq q < \infty$, $\sum_{(\epsilon, j, k) \in \Lambda_n} b_{j, k}^{\epsilon, 4}(t) \Phi_{j, k}^\epsilon(x) \in \mathbb{F}_{p, q}^{\gamma_1, \gamma_2, II}$. We can see that

$$\begin{aligned} |b_{j, k}^{\epsilon, 4}(t)| &\lesssim 2^{nj/2+j} \sum_{j < j'+2} \sum_{\epsilon', k'} \sum_{\epsilon'', k''} \int_0^{2^{-2j'\beta}} \frac{|u_{j', k'}^{\epsilon'}(s)|}{(1 + |2^{j-j'}k' - k|)^N} \frac{|v_{j', k''}^{\epsilon''}(s)|}{(1 + |k' - k''|)^N} ds \\ &\lesssim 2^{nj/2+j} \sum_{j < j'+2} 2^{-2j'\beta(1-2/q)} \sum_{\epsilon', k'} \sum_{\epsilon'', k''} \frac{1}{(1 + |2^{j-j'}k' - k|)^N} \frac{1}{(1 + |k' - k''|)^N} \\ &\quad \times \left(\int_0^{2^{-2j'\beta}} |u_{j', k'}^{\epsilon'}(s)|^q ds \right)^{1/q} \left(\int_0^{2^{-2j'\beta}} |v_{j', k''}^{\epsilon''}(s)|^q ds \right)^{1/q}. \end{aligned}$$

Let

$$f_{j'} = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j', k'} \subset Q_r^w} \left(\int_0^{2^{-2j'\beta}} |u_{j', k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \chi_{j', k'}(x).$$

Because

$$\left(\int_0^{2^{-2j'\beta}} |v_{j', k''}^{\epsilon''}(s)|^q ds \right)^{1/q} \lesssim 2^{j\gamma_2-n/p} 2^{-j'(\gamma_1+n/2-n/p)} 2^{-2j'\beta/q},$$

by Lemma 2.10, we can obtain

$$\begin{aligned} II_{b, A}(t) &\lesssim |Q_r|^{\gamma_2/n-1/p} \left\| \left\{ \sum_{-\log_2 r < j < -(\log_2 t)/(2\beta)} \sum_{Q_{j, k} \subset Q_r} 2^{qj(\gamma_1+n+1+\gamma_2-n/p)} \right. \right. \\ &\quad \times \left. \left[\sum_{j < j'+2} \sum_{\epsilon', k'} \frac{2^{-j'(2\beta-2\beta/q+\gamma_1+n/2-n/p)}}{(1 + |2^{j-j'}k' - k|)^N} \left(\int_0^{2^{-2j'\beta}} |u_{j', k'}^{\epsilon'}(s)|^q ds \right)^{1/q} \right]^q \chi_{j, k}(x) \right\|^{1/q} \Big\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^{-N}} \left\| \left\{ \sum_{-\log_2 r < j < -(\log_2 t)/(2\beta)} 2^{qj(\gamma_1+\gamma_2+1+n-n/p)} \right. \right. \\ &\quad \times \left. \left. \sum_{j < j'+2} 2^{(\delta+qn)(j'-j)} 2^{-qj'[2\gamma_1+n-n/p+2\beta]} (M_A(f_{j'})(x))^q \right\}^{1/q} \right\|_p \\ &\lesssim \|u\|_{\mathbb{F}_{p, q, m'}^{\gamma_1, \gamma_2, IV}}. \end{aligned}$$

Case 3.4.2: Under $2 \leq q < \infty$, $\sum_{(\epsilon, j, k) \in \Lambda_n} b_{j, k}^{\epsilon, 5}(t) \Phi_{j, k}^\epsilon(x) \in \mathbb{F}_{p, q}^{\gamma_1, \gamma_2, II}$. By a simple computation, we have

$$\left(\int_{2^{-2j'\beta}}^t |u_{j', k'}^{\epsilon'}(s)|^q s^{qm} \frac{ds}{s} \right)^{1/q} \lesssim 2^{-2j'\beta m} \left(\int_{2^{-2j'\beta}}^t |u_{j', k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm} \frac{ds}{s} \right)^{1/q}$$

and

$$\left(\int_{2^{-2j'\beta}}^t |v_{j', k''}^{\epsilon''}(s)|^q s^{qm} \frac{ds}{s} \right)^{1/q} \lesssim 2^{-2j'\beta m} 2^{-j'(\gamma_1+n/2-n/p)} 2^{j(\gamma_2-n/p)}.$$

Similarly, the above estimate and Holder’s inequality imply that

$$\begin{aligned} |b_{j,k}^{\epsilon,5}(t)| &\lesssim 2^{nj/2+j} \sum_{j < j'+2} \sum_{\epsilon',k'} \sum_{\epsilon'',k''} \int_{2^{-2j'\beta}}^t \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1+|2^{j-j'}k'-k|)^N} \frac{|v_{j',k''}^{\epsilon''}(s)|}{(1+|k'-k''|)^N} ds \\ &\lesssim 2^{nj/2+j} 2^{j(\gamma_2-n/p)} \sum_{j < j'+2} 2^{-j'(\gamma_1+n/2-n/p)} 2^{-2j'\beta} \\ &\quad \times \sum_{\epsilon',k'} (1+|2^{j-j'}k'-k|)^{-N} \left(\int_{2^{-2j'\beta}}^t |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm} \frac{ds}{s} \right)^{1/q}. \end{aligned}$$

Let

$$f_{j'} = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_{2^{-2j'\beta}}^t |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \chi_{j',k'}(x).$$

Similar to Case 3.4.1, we can obtain

$$\begin{aligned} II_{b,5}(t) &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{-\log_2 r < j < -(\log_2 t)/(2\beta)} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+\gamma_2+1+n-n/p)} \right. \right. \\ &\quad \left. \left. \times \sum_{j < j'+2} 2^{(\delta+qn)(j'-j)} 2^{-qj'(2\gamma_1+n-n/p+2\beta)} (M_A(f_{j'})(x))^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p \\ &\lesssim 1, \end{aligned}$$

which completes the proof. □

3.3. Bilinear estimates on L^p type tent spaces

Lemma 3.5. *Let $(\beta, p, q, \gamma_1, \gamma_2, m, m')$ satisfy (3.1) and $u, v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$. Then*

- (i) $\sum_{(\epsilon,j,k) \in \Lambda_n} a_{j,k}^{\epsilon,i}(t) \Phi_{j,k}^\epsilon \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}$, $i = 1, 2, 3$;
- (ii) $\sum_{(\epsilon,j,k) \in \Lambda_n} a_{j,k}^{\epsilon,4}(t) \Phi_{j,k}^\epsilon \in \mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}$.

Proof. Now we prove (i). For $i = 1, 2, 3$, write

$$III_{a,i} := |Q_r|^{\gamma_2/n-1/p} \left\| \left(\sum_{(\epsilon,j,k) \in \Lambda_n^Q} 2^{qj(\gamma_1+n/2+2m\beta)} \int_{2^{-2j\beta}}^{r^{2\beta}} |a_{j,k}^{\epsilon,i}(t)|^q t^{qm} \frac{dt}{t} \chi_{j,k}(x) \right)^{1/q} \right\|_p.$$

We still divide the proof into three cases:

Case 3.5.1: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} a_{j,k}^{\epsilon,1}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}$;

Case 3.5.2: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} a_{j,k}^{\epsilon,2}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}$;

Case 3.5.3: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} a_{j,k}^{\epsilon,3}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}$.

For simplicity, we only give the proof of Case 3.5.1. The proofs of Cases 3.5.2 and 3.5.3 are similar.

Now we prove Case 3.5.1. By $u, v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2}$, we can see that

$$\begin{aligned} |a_{j,k}^\epsilon(t)| &\lesssim 2^{nj/2+j} \sum_{|j-j'| \leq 1} \sum_{\epsilon', k', k''} \int_0^{2^{-1-2j'\beta}} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1 + |2^{j-j'}k' - k|)^N} \frac{|v_{j'-3,k''}^0(s)|}{(1 + |2^{j-j'+3}k'' - k|)^N} e^{-ct2^{2j\beta}} ds \\ &\lesssim 2^j \sum_{|j-j'| \leq 1} \sum_{\epsilon', k', k''} \int_0^{2^{-1-2j'\beta}} |u_{j',k'}^{\epsilon'}(s)| e^{-ct2^{2j\beta}} (1 + |2^{j-j'}k' - k|)^{-N} s^{1/(2\beta)-1} ds. \end{aligned}$$

We can get

$$\begin{aligned} III_{a,1} &\lesssim |Q_r|^{\gamma_2/n-1/p} \left\| \left[\sum_{(\epsilon, j, k) \in \Lambda_n^Q} 2^{qj(\gamma_1+n/2+2m\beta)} 2^{qj} e^{-ct2^{2j\beta}} \int_{2^{-2j\beta}}^{r^{2\beta}} \sum_{|j-j'| \leq 1} \right. \right. \\ &\quad \left. \left. \times \left(\sum_{\epsilon', k'} \int_0^{2^{-1-2j'\beta}} \frac{|u_{j',k'}^{\epsilon'}(s)|}{(1 + |2^{j-j'}k' - k|)^N} s^{1/(2\beta)-1} ds \right)^q t^{qm} \frac{dt}{t} \chi_{j,k}(x) \right]^{1/q} \right\|_p. \end{aligned}$$

Let

$$g_{j'} = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_0^{2^{-1-2j'\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \chi_{j',k'}(x).$$

For $|j - j'| \leq 1$, we have, by Lemma 2.10 and Hölder's inequality for k' to deduce that

$$\begin{aligned} III_{a,1} &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1 + |w|)^{-N} \left\| \left\{ \sum_{j' \geq -\log_2 r_w} (M_A(g_{j'})(x))^q \right\}^{1/q} \right\|_p \\ &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1 + |w|)^{-N} \left\| \left\{ \sum_{j' \geq -\log_2 r_w} |g_{j'}(x)|^q \right\}^{1/q} \right\|_p \\ &\lesssim \|u\|_{\mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, III}} + \|u\|_{\mathbb{F}_{p,q,m'}^{\gamma_1, \gamma_2, IV}}. \end{aligned}$$

Now we prove (ii). Write

$$IV_{a,4} := |Q_r|^{\gamma_2/n-1/p} \left\| \left[\sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m'\beta)} \int_0^{2^{-2j\beta}} |a_{j,k}^{\epsilon,4}(t)|^q t^{qm'} \frac{dt}{t} \chi_{j,k}(x) \right]^{1/q} \right\|_p.$$

Take μ such that $qm' + q - 1 - q/(2\beta) \leq q\mu < q - 1$. Similarly we can get

$$\begin{aligned} IV_{a,4} &\lesssim |Q_r|^{\gamma_2/n-1/p} \left\| \left[\sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m'\beta)} \int_0^{2^{-2j\beta}} 2^{qj} \sum_{|j-j'| \leq 2} \right. \right. \\ &\quad \left. \left. \times \left[\sum_{\epsilon', k'} \frac{1}{(1 + |2^{j-j'}k' - k|)^N} \int_0^t |u_{j',k'}^{\epsilon'}(s)| e^{-c(t-s)2^{2j\beta}} s^{1/(2\beta)-1} ds \right]^q t^{qm'} \frac{dt}{t} \chi_{j,k}(x) \right]^{1/q} \right\|_p \end{aligned}$$

$$\begin{aligned} &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1 + |w|)^{-N} \\ &\times \left\| \left[\sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} \sum_{|j-j'| \leq 1} \sum_{Q_{j',k'} \subset Q_{j,k}^w} (1 + |2^{j-j'}k' - k|)^{-N} t^{q-1-q\mu} \right. \right. \\ &\times \left. \left. \int_0^{2^{-2j\beta}} |u_{j',k'}^{\epsilon'}(s)|^q s^{q(1/(2\beta)-1+\mu)} \left(\int_s^{2^{-2j\beta}} 2^{qj(1+2m'\beta)} t^{qm'+q-q\mu-2} dt \right) ds \chi_{j,k}(x) \right]^{1/q} \right\|_p. \end{aligned}$$

Let

$$g_{j'}(x) = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_0^{2^{-2j\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \chi_{j',k'}(x).$$

By $s2^{2j\beta'} \leq 1$ and $qm' + q - 1 - q/(2\beta) \leq q\mu$, a simple computation yields

$$\begin{aligned} IV_{a,4} &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^N} \left\| \left[\sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} \sum_{|j-j'| \leq 1} \sum_{Q_{j',k'} \subset Q_{j,k}^w} \frac{t^{q-1-q\mu}}{(1 + |2^{j-j'}k' - k|)^N} \right. \right. \\ &\times \left. \left. \int_0^{2^{-2j\beta}} |u_{j',k'}^{\epsilon'}(s)|^q s^{q(1/(2\beta)-1+\mu)+1} \frac{ds}{s} 2^{qj'(2\beta\mu+2\beta/q-2\beta+1)} \chi_{j,k}(x) \right]^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^N} \left\| \left\{ \sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} \sum_{|j-j'| \leq 1} 2^{q(j-j')(\gamma_1+n/2)} [M_A(g_{j'})(x)]^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p \\ &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^N} \left\| \left\{ \sum_{j' \geq -\log_2 r_w} |g_{j'}(x)|^q \right\}^{1/q} \right\|_p \\ &\lesssim \|u\|_{\mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}}. \end{aligned}$$

This completes the proof. □

Lemma 3.6. *Let $(\beta, p, q, \gamma_1, \gamma_2, m, m')$ satisfy (3.1) and $u, v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$. Then*

$$\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,i} \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}, \quad i = 1, 2, 3.$$

Proof. For $i = 1, 2, 3$, write

$$III_{b,i} := |Q_r|^{\gamma_2/n-1/p} \left\| \left\{ \sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2+2m\beta)} \int_{2^{-2j'\beta}}^{r^{2\beta}} |b_{j,k}^{\epsilon,1}(t)|^q t^{qm} \frac{dt}{t} \chi_{j,k}(x) \right\}^{1/q} \right\|_p.$$

We divide the proof into three cases:

Case 3.6.1: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,1}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}$;

Case 3.6.2: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,2}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, III}$;

Case 3.6.3: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,3}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m}^{\gamma_1, \gamma_2, III}$.

For simplicity, we only give the proof of Case 3.6.1. Cases 3.6.2 and 3.6.3 can be dealt with similarly. Now we prove Case 3.6.1. Because $s2^{2j'\beta} \leq 1$, we have

$$\left| u_{j',k'}^{\epsilon'}(s) \right| + \left| v_{j',k''}^{\epsilon''}(s) \right| \lesssim 2^{-(n/2+\gamma_1-\gamma_2)j'}$$

and

$$\left(\int_0^{2^{-1-2j'\beta}} \left| v_{j',k''}^{\epsilon''}(s) \right|^q ds \right)^{1/q} \lesssim 2^{-2j'\beta/q} 2^{-j'(\gamma_1+n/2-n/p)} 2^{j(\gamma_2-n/p)}.$$

Then Holder's inequality implies that

$$\begin{aligned} \left| b_{j,k}^{\epsilon,1}(t) \right| &\lesssim 2^{nj/2+j} e^{-ct2^{2j\beta}} \sum_{j < j'+2} \sum_{\epsilon',k'} \sum_{\epsilon'',k''} \int_0^{2^{-1-2j'\beta}} \frac{\left| u_{j',k'}^{\epsilon'}(s) \right|}{(1 + |2^{j-j'}k' - k|)^N} \frac{\left| v_{j',k''}^{\epsilon''}(s) \right|}{(1 + |k' - k''|)^N} ds \\ &\lesssim 2^{nj/2+j} e^{-ct2^{2j\beta}} 2^{j(\gamma_2-n/p)} \sum_{w \in \mathbb{Z}^n} (1 + |w|)^{-N} \sum_{j < j'+2} \sum_{\epsilon',k'} (1 + |2^{j-j'}k' - k|)^{-N} \\ &\quad \times 2^{-2j'\beta/q} 2^{-j'(\gamma_1+n/2-n/p)} 2^{-2j'\beta(1-2/q)} \left(\int_0^{2^{-1-2j'\beta}} \left| u_{j',k'}^{\epsilon'}(s) \right|^q ds \right)^{1/q}. \end{aligned}$$

The above estimate implies that

$$\begin{aligned} III_{b,1} &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1 + |w|)^N} \left\| \left\{ \sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n+1)} \int_{2^{-2qj\beta}}^{r^{2\beta}} e^{-cqt2^{2j\beta}} 2^{j(\gamma_2-n/p)} \right. \right. \\ &\quad \times \sum_{j < j'+2} 2^{\delta(j'-j)} 2^{-qj'(\gamma_1+n/2-n/p)} 2^{-2j'\beta(q-1)} \left[\sum_{Q_{j',k'} \subset Q_{j,k}^w} (1 + |2^{j-j'}k' - k|)^{-N} \right. \\ &\quad \left. \left. \times \left(\int_0^{2^{-1-2j'\beta}} \left| u_{j',k'}^{\epsilon'}(s) \right|^q ds \right)^{1/q} \right]^q (t2^{2j\beta})^{qm} \frac{dt}{t} \chi_{j,k}(x) \right\} \right\|_p. \end{aligned}$$

Let

$$g_{j'} = \left[\sum_{Q_{j',k'} \subset Q_{j,k}^w} (1 + |2^{j-j'}k' - k|)^{-N} \left(\int_0^{2^{-1-2j'\beta}} \left| u_{j',k'}^{\epsilon'}(s) \right|^q ds \right)^{1/q} \right]^q$$

and

$$f_{j'} = 2^{j'(\gamma_1+n/2)} \sum_{Q_{j',k'} \subset Q_r^w} \left(\int_0^{2^{-1-2j'\beta}} \left| u_{j',k'}^{\epsilon'}(s) \right|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q}.$$

Then

$$\begin{aligned} |g_{j'}| &\lesssim 2^{-2j'\beta} \left[\sum_{Q_{j',k'} \subset Q_{j,k}^w} \frac{1}{(1 + |2^{j-j'}k' - k|)^N} \left(\int_0^{2^{-1-2j'\beta}} \left| u_{j',k'}^{\epsilon'}(s) \right|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \right]^q \\ &\lesssim 2^{-2j'\beta} 2^{-qj'(\gamma_1+n/2)} 2^{qn(j'-j)} (M_A(f_{j'})(x))^q. \end{aligned}$$

Hence

$$\begin{aligned}
 III_{b,1} &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{j \geq \log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n+1+\gamma_2-n/p)} \sum_{j < j'+2} 2^{(\delta+qn)(j'-j)} \right. \right. \\
 &\quad \left. \left. \times 2^{-qj'(\gamma_1+n/2-n/p+2\beta+\gamma_1+n/2)} (M_A(f_{j'})(x))^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p \\
 &\lesssim \sum_{w \in \mathbb{Z}^n} \frac{|Q_r|^{\gamma_2/n-1/p}}{(1+|w|)^N} \left\| \left\{ \sum_{j' \geq -\log_2 r} \sum_{j < j'+2} 2^{q(j'-j)[\delta/q-(\gamma_1+\gamma_2+1-n/p)]} (M_A(f_{j'})(x))^q \right\}^{1/q} \right\|_p \\
 &\lesssim \|u\|_{\mathbb{F}_{p,q,m}^{\gamma_1,\gamma_2,III}}.
 \end{aligned}$$

The proof is finished. □

Lemma 3.7. *Let $(\beta, p, q, \gamma_1, \gamma_2, m, m')$ satisfy (3.1) and $u, v \in \mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2}$. Then*

$$\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,i}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}, \quad i = 4, 5.$$

Proof. For $i = 4, 5$, write

$$II_{b,i} := |Q_r|^{\gamma_2/n-1/p} \left\| \left\{ \sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n/2)} \int_0^{2^{-2j'\beta}} |b_{j,k}^{\epsilon,i}(t)|^q (t2^{2j'\beta})^{qm'} \frac{dt}{t} \chi_{j,k}(x) \right\}^{1/q} \right\|_p.$$

We still divide the proof into two cases:

Case 3.7.1: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,4}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}$;

Case 3.7.2: Under $2 \leq q < \infty$, $\sum_{(\epsilon,j,k) \in \Lambda_n} b_{j,k}^{\epsilon,5}(t) \Phi_{j,k}^\epsilon(x) \in \mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}$.

We only prove Case 3.7.1. The treatment for Case 3.7.2 is similar. We have

$$\begin{aligned}
 |b_{j,k}^{\epsilon,4}(t)| &\lesssim 2^{nj/2+j2j(\gamma_2-n/p)} \sum_{j < j'+2} 2^{-2j'\beta(1-2/q)} 2^{-j'(\gamma_1+n/2-n/p)} 2^{-4j'\beta/q} \\
 &\quad \times \sum_{\epsilon',k'} \frac{1}{(1+|2^{j-j'}k' - k|)^N} \left(\int_0^{2^{-2j'\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q}.
 \end{aligned}$$

Then by $\int_0^{2^{-2j'\beta}} (t2^{2j'\beta})^{qm'} \frac{dt}{t} \lesssim 1$, we apply Lemma 2.10 to derive that

$$\begin{aligned}
 II_{b,4} &\lesssim |Q_r|^{\gamma_2/n-1/p} \sum_{w \in \mathbb{Z}^n} (1+|w|)^{-N} \left\| \left\{ \sum_{j \geq -\log_2 r} \sum_{Q_{j,k} \subset Q_r} 2^{qj(\gamma_1+n+1+\gamma_2-n/p)} \right. \right. \\
 &\quad \times \sum_{j < j'+2} 2^{\delta(j'-j)} 2^{-2qj'\beta(1-2/q)} 2^{-qj'(\gamma_1+n/2-n/p)} 2^{-4j'\beta} \\
 &\quad \left. \left. \times \left[\sum_{Q_{j',k'} \subset Q_r^w} \frac{1}{(1+|2^{j-j'}k' - k|)^N} \left(\int_0^{2^{-2j'\beta}} |u_{j',k'}^{\epsilon'}(s)|^q (s2^{2j'\beta})^{qm'} \frac{ds}{s} \right)^{1/q} \right]^q \chi_{j,k}(x) \right\}^{1/q} \right\|_p \\
 &\lesssim \|u\|_{\mathbb{F}_{p,q,m'}^{\gamma_1,\gamma_2,IV}}.
 \end{aligned}$$

This completes the proof of Lemma 3.7. □

4. Well-posedness of magneto-hydrodynamic equations with data in $\dot{F}_{p,q}^{\gamma_1, \gamma_2}$

4.1. Generalized magneto-hydrodynamic equations

In this section, we consider the incompressible fractional magneto-hydrodynamic (FMHD) $_{\beta}$ equations in \mathbb{R}_+^{n+1} , $n \geq 2$:

$$(4.1) \quad \begin{cases} \partial_t u + (-\Delta)^{\beta} u + u \cdot \nabla u + \nabla p - b \cdot \nabla b = 0, \\ \partial_t b + (-\Delta)^{\beta} b + u \cdot \nabla b - b \cdot \nabla u = 0, \\ u|_{t=0} = u_0, \quad b|_{t=0} = b_0 \end{cases}$$

with $\nabla \cdot u = \nabla \cdot b = 0$ and $\beta \in (1/2, 1)$. Here $(-\Delta)^{\beta}$ is the fractional Laplacian with respect to x defined by

$$(-\Delta)^{\beta} u(t, \xi) = |\xi|^{2\beta} \widehat{u}(t, \xi).$$

In equations (4.1), $u(t, x) = (u_1(t, x), \dots, u_n(t, x))$ and $b(t, x) = (b_1(t, x), \dots, b_n(t, x))$ are the flow velocity and the magnetic field at point (x, t) , respectively. $u_0(x)$ and $b_0(x)$ are the initial velocity and magnetic field and satisfy $\nabla \cdot u_0 = 0$ and $\nabla \cdot b_0 = 0$, respectively.

Theorem 4.1. *Let*

$$(4.2) \quad \begin{cases} 1 < p, q < \infty, \\ \beta > 1/2, \quad \gamma_1 = \gamma_2 - 2\beta + 1, \\ m > \max \{p, n/(2\beta)\}, \quad 0 < m' < \min \{1, p/(2\beta)\}. \end{cases}$$

If the index (β, p, γ_2) obeys

$$1 < q \leq 2 \quad \& \quad 2\beta - \frac{2}{q} + \frac{n}{p} \left(1 - \frac{1}{q}\right) < \gamma_2 \leq \frac{n}{p}$$

or

$$2 < q < \infty \quad \& \quad \beta - 1 + \frac{n}{2p} < \gamma_2 \leq \frac{n}{p},$$

then (4.1) has a unique global mild solution in $(\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2})^n$ for any initial data (u_0, b_0) with $\|(u_0, b_0)\|_{(\dot{F}_{p,q}^{\gamma_1, \gamma_2})^n}$ being small.

Proof. By Lemmas 3.2–3.7, we have known that the bilinear operator

$$(4.3) \quad B(u, v) = \int_0^t e^{-(t-s)(-\Delta)^{\beta}} P \nabla \cdot (u \otimes v) ds$$

is bounded from $(\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2})^n \times (\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2})^n$ to $(\mathbb{F}_{p,q,m,m'}^{\gamma_1, \gamma_2})^n$.

The solution (u, b) to equations (4.1) can be written as

$$u(t, x) = e^{-t(-\Delta)^{\beta}} u_0(x) - B(u, u) + B(b, b) := F_1(u, b)$$

and

$$b(t, x) = e^{-(\Delta)^\beta} b_0(x) - B(u, b) + B(b, u) := F_2(u, b).$$

We rewrite the solution (u, b) as

$$\begin{pmatrix} u \\ b \end{pmatrix} = \begin{pmatrix} F_1(u, b) \\ F_2(u, b) \end{pmatrix} := F(u, b).$$

If (4.3) holds, we have

$$\|F_1(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \leq C_1 \|u_0\|_{(\dot{F}_{p,q}^{\gamma_1,\gamma_2})^n} + C_2 \|(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n}^2$$

and

$$\begin{aligned} & \|F_1(u, b)(t) - F_1(u', b')(t)\|_{(X_{p,m,m'}^{\gamma_1,\gamma_2,\tau})^n} \\ & \lesssim \|B(u, u) - B(u', u')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} + \|B(b, b) - B(b', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \\ & \lesssim \left(\|u\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} + \|u'\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \right) \|u - u'\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \\ & \quad + \left(\|b\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} + \|b'\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \right) \|b - b'\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \\ & \lesssim \|(u, b) - (u', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \left(\|(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} + \|(u', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \right). \end{aligned}$$

Similarly, it can be obtained that

$$\|F_2(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \leq C_1 \|b_0\|_{(\dot{F}_{p,q}^{\gamma_1,\gamma_2})^n} + C_2 \|(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n}^2$$

and

$$\begin{aligned} & \|F_2(u, b)(t) - F_2(u', b')(t)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \\ & \lesssim \|(u, b) - (u', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \left(\|(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} + \|(u', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \right). \end{aligned}$$

Thus, we have

$$\|F(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \leq C \|(u_0, b_0)\|_{(\dot{B}_{p,p}^{\gamma_1,\gamma_2})^n} + C \|(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n}^2$$

and

$$\begin{aligned} & \|F(u, b) - F(u', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \\ & \lesssim \|(u, b) - (u', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \left(\|(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} + \|(u', b')\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \right). \end{aligned}$$

Let R be a constant satisfying $R \leq 2C \|(u_0, b_0)\|_{(\dot{B}_{p,p}^{\gamma_1,\gamma_2})^n}$. Then F is a contraction mapping from

$$E = \left\{ (u, b) \in (\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n : \|(u, b)\|_{(\mathbb{F}_{p,q,m,m'}^{\gamma_1,\gamma_2})^n} \leq R \right\}$$

into itself. Thus, there exists a fixed point (u, b) of the operator F . □

4.2. Quasi-geostrophic equations

In this section, we study the well-posedness and regularity of quasi-geostrophic equation

$$(4.4) \quad \begin{cases} \partial_t \theta = -(-\Delta)^\beta \theta + \partial_1(\theta R_2 \theta) - \partial_2(\theta R_1 \theta), \\ \theta(0, x) = \theta_0(x), \end{cases}$$

with initial data in the space $\dot{B}_{p,p}^{\gamma_1, \gamma_2}$, where $\beta \in (1/2, 1)$.

The equations $(DQG)_\beta$ are important models in the atmosphere and ocean fluid dynamics. It was proposed by P. Constantin and A. Majda, etc. that the equations $(DQG)_\beta$ can be regarded as low-dimensional models for mathematical study of singularity in smooth solutions of unforced incompressible three-dimensional fluid equations. See e.g. [12, 21, 23, 44, 45] and the references therein.

Owing to the importance in mathematical and geophysical fluid dynamics mentioned above, the equations $(DQG)_\beta$ have been intensively studied. Many important progress has been made. We refer the reader to [5, 6, 8, 10, 11, 13, 22, 37, 55, 56] for details.

The solution to equations (4.4) can be represented as

$$u(t, \cdot) = e^{-t(-\Delta)^\beta} u_0 + B(u, u),$$

where the bilinear form $B(u, v)$ is defined by

$$B(u, v) = \int_0^t e^{-(t-s)(-\Delta)^\beta} (\partial_1(v R_2 u) - \partial_2(v R_1 u)) ds.$$

By the space $\mathbb{F}_{p,m,m'}^{\gamma_1, \gamma_2, \tau}$ introduced in Section 3, we consider the well-posedness and regularity of the quasi-geostrophic equations (4.4).

Theorem 4.2. *Given $1/2 < \beta < 1$, $\gamma_1 = \gamma_2 - 2\beta + 1$. If the index $(p, \gamma_1, \gamma_2, m, m')$ satisfies the conditions (4.2), the quasi-geostrophic dissipative equation (4.4) has a unique global mild solution in $(\mathbb{F}_{p,m,m'}^{\gamma_1, \gamma_2})^2$ for all small initial data $a(x)$ with $\nabla \cdot a = 0$ and $\|a\|_{(\dot{F}_{p,p}^{\gamma_1, \gamma_2})^2}$ small enough.*

Proof. By Picard’s contraction principle, it suffices to verify that the bilinear operator

$$B(u, v) = \int_0^t e^{-(t-s)(-\Delta)^\beta} (\partial_1(v R_2 u) - \partial_2(v R_1 u)) ds$$

is bounded from $(\mathbb{F}_{p,m,m'}^{\gamma_1, \gamma_2})^2 \times (\mathbb{F}_{p,m,m'}^{\gamma_1, \gamma_2})^2$ to $(\mathbb{F}_{p,m,m'}^{\gamma_1, \gamma_2})^2$. The proof is similar to that of Theorem 4.1. We omit the details. □

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References

- [1] G. V. Alekseev, *Solvability of a homogeneous initial-boundary value problem for equations of magnetohydrodynamics of an ideal fluid*, *Dinamika Sploshn. Sredy*. No. **57** (1982), 3–20.
- [2] M. Cannone, *Ondelettes, Paraproducts et Navier-Stokes*, Diderot Editeur, Paris, 1995.
- [3] ———, *Harmonic analysis tools for solving the incompressible Navier-Stokes equations*, in *Handbook of Mathematical Fluid Dynamics, Vol. III*, 161–244, North-Holland, Amsterdam, 2004. [https://doi.org/10.1016/s1874-5792\(05\)80006-0](https://doi.org/10.1016/s1874-5792(05)80006-0)
- [4] D. Chae, *On the well-posedness of the Euler equations in the Triebel-Lizorkin spaces*, *Comm. Pure Appl. Math.* **55** (2002), no. 5, 654–678.
<https://doi.org/10.1002/cpa.10029>
- [5] ———, *The quasi-geostrophic equation in the Triebel-Lizorkin spaces*, *Nonlinearity* **16** (2003), no. 2, 479–495. <https://doi.org/10.1088/0951-7715/16/2/307>
- [6] D. Chae and J. Lee, *Global well-posedness in the super-critical dissipative quasi-geostrophic equations*, *Comm. Math. Phys.* **233** (2003), no. 2, 297–311.
<https://doi.org/10.1007/s00220-002-0750-z>
- [7] Q. Chen, C. Miao and Z. Zhang, *A new Bernstein's inequality and the 2D dissipative quasi-geostrophic equation*, *Comm. Math. Phys.* **271** (2007), no. 3, 821–838.
<https://doi.org/10.1007/s00220-007-0193-7>
- [8] ———, *On the well-posedness of the ideal MHD equations in the Triebel-Lizorkin spaces*, *Arch. Ration. Mech. Anal.* **195** (2010), no. 2, 561–578.
<https://doi.org/10.1007/s00205-008-0213-6>
- [9] R. R. Coifman, Y. Meyer and E. M. Stein, *Some new function spaces and their applications to harmonic analysis*, *J. Funct. Anal.* **62** (1985), no. 2, 304–335.
[https://doi.org/10.1016/0022-1236\(85\)90007-2](https://doi.org/10.1016/0022-1236(85)90007-2)

- [10] P. Constantin, *Geometric statistics in turbulence*, SIAM Rev. **36** (1994), no. 1, 73–98.
<https://doi.org/10.1137/1036004>
- [11] P. Constantin, D. Cordoba and J. Wu, *On the critical dissipative quasi-geostrophic equations*, Indiana Univ. Math. J. **50** (2001), Special Issue, 97–107.
<https://doi.org/10.1512/iumj.2001.50.2153>
- [12] P. Constantin, A. J. Majda and E. Tabak, *Formation of strong fronts in the 2-D quasigeostrophic thermal active scalar*, Nonlinearity **7** (1994), no. 6, 1495–1533.
<https://doi.org/10.1088/0951-7715/7/6/001>
- [13] P. Constantin and J. Wu, *Behavior of solutions of 2D quasi-geostrophic equations*, SIAM J. Math. Anal. **30** (1999), no. 5, 937–948.
<https://doi.org/10.1137/s0036141098337333>
- [14] G. Dafni and J. Xiao, *Some new tent spaces and duality theorems for fractional Carleson measures and $Q_\alpha(\mathbb{R}^n)$* , J. Funct. Anal. **208** (2004), no. 2, 377–422.
[https://doi.org/10.1016/s0022-1236\(03\)00181-2](https://doi.org/10.1016/s0022-1236(03)00181-2)
- [15] C. Deng and X. Yao, *Well-posedness and ill-posedness for the 3D generalized Navier-Stokes equations in $\dot{F}_{\frac{3}{\alpha-1}}^{-\alpha,r}$* , Discrete Contin. Dyn. Syst. **34** (2014), no. 2, 437–459.
<https://doi.org/10.3934/dcds.2014.34.437>
- [16] P. Federbush, *Navier and Stokes meet the wavelet*, Commun. Math. Phys. **155** (1993), no. 2, 219–248. <https://doi.org/10.1007/bf02097391>
- [17] M. Frazier, B. Jawerth and G. Weiss, *Littlewood-Paley Theory and the Study of Function Spaces*, CBMS Regional Conference Series in Mathematics **79**, Amer. Math. Soc., Providence, RI, 1991. <https://doi.org/10.1090/cbms/079>
- [18] P. Germain, N. Pavlović and G. Staffilani, *Regularity of solutions to the Navier-Stokes equations evolving from small data in BMO^{-1}* , Int. Math. Res. Not. IMRN **2007**, no. 21, Art. ID rnm087, 35 pp. <https://doi.org/10.1093/imrn/rnm087>
- [19] Y. Giga and T. Miyakawa, *Navier-Stokes flow in \mathbb{R}^3 with measures as initial vorticity and Morry spaces*, Comm. Partial Differential Equations **14** (1989), no. 5, 577–618.
<https://doi.org/10.1080/03605308908820621>
- [20] Y. Giga and O. Sawada, *On regularizing-decay rate estimates for solutions to the Navier-Stokes initial value problem*, in *Nonlinear Analysis and Applications*, Nonlinear Anal. Appl. **1** (2002), 549–562.

- [21] I. M. Held, R. T. Pierrehumbert, S. T. Garner and K. L. Swanson, *Surface quasi-geostrophic dynamics*, J. Fluid Mech. **282** (1995), 1–20.
<https://doi.org/10.1017/s0022112095000012>
- [22] N. Ju, *On the two dimensional quasi-geostrophic equations*, Indiana Univ. Math. J. **54** (2005), no. 3, 897–926. <https://doi.org/10.1512/iumj.2005.54.2518>
- [23] ———, *The 2D quasi-geostrophic equations in the Sobolev space*, in *Harmonic Analysis, Partial Differential Equations, and related topics*, 75–92, Contemp. Math. **428**, Amer. Math. Soc., Providence, RI, 2007. <https://doi.org/10.1090/conm/428/08213>
- [24] T. Kato, *Strong L^p -solutions of the Navier-Stokes in \mathbb{R}^m with applications to weak solutions*, Math. Z. **187** (1984), no. 4, 471–480. <https://doi.org/10.1007/bf01174182>
- [25] T. Kato and H. Fujita, *On the non-stationary Navier-Stokes system*, Rend. Semin. Mat. Univ. Padova. **30** (1962), 243–260.
- [26] H. Koch and D. Tataru, *Well-posedness for the Navier-Stokes equations*, Adv. Math. **157** (2001), no. 1, 22–35. <https://doi.org/10.1006/aima.2000.1937>
- [27] H. Kozono, *Weak and classical solutions of the two-dimensional magnetohydrodynamic equations*, Tohoku Math. J. (2) **41** (1989), no. 3, 471–488.
<https://doi.org/10.2748/tmj/1178227774>
- [28] H. Kozono and Y. Shimada, *Bilinear estimates in homogeneous Triebel-Lizorkin spaces and the Navier-Stokes equations*, Math. Nachr. **276** (2004), no. 1, 63–74.
<https://doi.org/10.1002/mana.200310213>
- [29] P. G. Lemarié-Rieusset, *Recent Developments in the Navier-Stokes Problem*, Chapman & Hall/CRC Research Notes in Mathematics **431**, Chapman & Hall/CRC Press, FL, 2002. <https://doi.org/10.1201/9781420035674>
- [30] P. Li, J. Xiao and Q. Yang, *Global mild solutions of fractional Navier-Stokes equations with small initial data in critical Besov- Q spaces*, Electron. J. Differ. Equ. **2014** (2014), no. 185, 1–37.
- [31] P. Li and Q. Yang, *Wavelets and the well-posedness of incompressible magnetohydrodynamic equations in Besov type Q -spaces*, J. Math. Anal. Appl. **405** (2013), no. 2, 661–686. <https://doi.org/10.1016/j.jmaa.2013.04.035>
- [32] P. Li, Q. Yang and B. Zheng, *Wavelets and Triebel type oscillation spaces*, arXiv: 1401.0274.

- [33] P. Li and Z. Zhai, *Well-posedness and regularity of generalized Navier-Stokes equations in some critical Q -spaces*, *J. Funct. Anal.* **259** (2010), no. 10, 2457–2519. <https://doi.org/10.1016/j.jfa.2010.07.013>
- [34] Y. Liang, Y. Sawano, T. Ullrich, D. Yang and W. Yuan, *New characterizations of Besov-Triebel-Lizorkin-Hausdorff spaces including coorbital and wavelets*, *J. Fourier Anal. Appl.* **18** (2012), no. 5, 1067–1111. <https://doi.org/10.1007/s00041-012-9234-5>
- [35] C. C. Lin and Q. Yang, *Semigroup characterization of Besov type Morrey spaces and well-posedness of generalized Navier-Stokes equations*, *J. Differential Equations* **254** (2013), no. 2, 804–846. <https://doi.org/10.1016/j.jde.2012.09.017>
- [36] J.-L. Lions, *Quelques méthodes de résolution des problèmes aux limites non linéaires*, Dunod/Gauthier, Villars, Paris, 1969.
- [37] F. Marchand and P. G. Lemarié-Rieusset, *Solutions auto-similaires non radiales pour l'équation quasi-géostrophique dissipative critique*, *C. R. Math. Acad. Sci. Paris* **341** (2005), no. 9, 535–538. <https://doi.org/10.1016/j.crma.2005.09.004>
- [38] Y. Meyer, *Wavelets and Operators*, Cambridge Studies in Advanced Mathematics **37**, Cambridge University Press, Cambridge, 1992.
- [39] Y. Meyer and R. Coifman, *Calderón-Zygmund and multilinear operators*, Cambridge Studies in Advanced Mathematics. **48**, Cambridge University Press, Cambridge, 1997.
- [40] Y. Meyer, Q. X. Yang, *Continuity of Calderón-Zygmund operators on Besov or Triebel-Lizorkin spaces*, *Anal. Appl.* **6** (2008), no. 1, 51–81. <https://doi.org/10.1142/s0219530508001055>
- [41] C. Miao, B. Yuan and B. Zhang, *Well-posedness for the incompressible magneto-hydrodynamic system*, *Math. Methods Appl. Sci.* **30** (2007), no. 8, 961–976. <https://doi.org/10.1002/mma.820>
- [42] ———, *Well-posedness of the Cauchy problem for the fractional power dissipative equations*, *Nonlinear Anal.* **68** (2008), no. 3, 461–484. <https://doi.org/10.1016/j.na.2006.11.011>
- [43] H. Miura and O. Sawada, *On the regularizing rate estimates of Koch-Tataru's solution to the Navier-Stokes equations*, *Asymptot. Anal.* **49** (2006), no. 1-2, 1–15.
- [44] K. Ohkitani and M. Yamada, *Inviscid and inviscid-limit behavior of a surface quasi-geostrophic flow*, *Phys. Fluids* **9** (1997), no. 4, 876–882. <https://doi.org/10.1063/1.869184>

- [45] J. Peetre, *New Thoughts on Besov Spaces*, Duke University, Durham, N.C., 1976.
- [46] F. Planchon, *Global strong solutions in Sobolev or Lebesgue spaces to the incompressible Navier-Stokes equations in \mathbb{R}^3* , Ann. Inst. H. Poincaré Anal. Non Linéaire **13** (1996), no. 3, 319–336.
- [47] M. E. Taylor, *Analysis on Morrey spaces and applications to Navier-Stokes and other evolution equations*, Comm. Partial Differential Equations **17** (1992), no. 9-10, 1407–1456. <https://doi.org/10.1080/03605309208820892>
- [48] H. Triebel, *Theory of Function Spaces*, Monographs in Mathematics **78**, Birkhäuser Verlag, Basel, 1983. <https://doi.org/10.1007/978-3-0346-0416-1>
- [49] ———, *Theory of Function Spaces II*, Monographs in Mathematics **84**, Birkhäuser Verlag, Basel, 1992. <https://doi.org/10.1007/978-3-0346-0419-2>
- [50] Y. Wang and J. Xiao, *Mild solutions of incompressible Navier-Stokes system with dissipation*, preprint, 2014.
- [51] ———, *Homogeneous Campanato-Sobolev classes*, Appl. Comput. Harmon. Anal. **39** (2015), no. 2, 214–247. <https://doi.org/10.1016/j.acha.2014.09.002>
- [52] G. Wu, *Regularity criteria for the 3D generalized MHD equations in terms of vorticity*, Nonlinear Anal. **71** (2009), no. 9, 4251–4258. <https://doi.org/10.1016/j.na.2009.02.115>
- [53] J. Wu, *Generalized MHD equations*, J. Differential Equations **195** (2003), no. 2, 284–312. <https://doi.org/10.1016/j.jde.2003.07.007>
- [54] ———, *The generalized incompressible Navier-Stokes equations in Besov spaces*, Dyn. Partial Differ. Equ. **1** (2004), no. 4, 381–400. <https://doi.org/10.4310/dpde.2004.v1.n4.a2>
- [55] ———, *Global solutions of the 2D dissipative quasi-geostrophic equation in Besov spaces*, SIAM J. Math. Anal. **36** (2004/05), no. 3, 1014–1030. <https://doi.org/10.1137/s0036141003435576>
- [56] ———, *The two-dimensional quasi-geostrophic equation with critical or supercritical dissipation*, Nonlinearity **18** (2005), no. 1, 139–154. <https://doi.org/10.1088/0951-7715/18/1/008>
- [57] ———, *Lower bounds for an integral involving fractional Laplacians and the generalized Navier-Stokes equations in Besov spaces*, Comm. Math. Phys. **263** (2005), no. 3, 803–831. <https://doi.org/10.1007/s00220-005-1483-6>

- [58] ———, *Regularity criteria for the generalized MHD equations*, Comm. Partial Differential Equations **33** (2008), no. 2, 285–306.
<https://doi.org/10.1080/03605300701382530>
- [59] J. Xiao, *Homothetic variant of fractional Sobolev space with application to Navier-Stokes system*, Dyn. Partial Differ. Equ. **4** (2007), no. 3, 227–245.
<https://doi.org/10.4310/dpde.2007.v4.n3.a2>
- [60] ———, *Homothetic variant of fractional Sobolev space with application to Navier-Stokes system revisited*, Dyn. Partial Differ. Equ. **11** (2014), no. 2, 167–181.
<https://doi.org/10.4310/dpde.2014.v11.n2.a3>
- [61] Q. Yang, *Wavelet and Distribution*, Beijing Science and Technology Press, Beijing, 2002.
- [62] D. Yang and W. Yuan, *A new class of function spaces connecting Triebel-Lizorkin spaces and Q spaces*, J. Funct. Anal. **255** (2008), no. 10, 2760–2809.
<https://doi.org/10.1016/j.jfa.2008.09.005>
- [63] ———, *New Besov-type spaces and Triebel-Lizorkin-type spaces including Q spaces*, Math. Z. **265** (2010), no. 2, 451–480. <https://doi.org/10.1007/s00209-009-0524-9>
- [64] ———, *Characterizations of Besov-type and Triebel-Lizorkin-type spaces via maximal functions and local means*, Nonlinear Anal. **73** (2010), no. 12, 3805–3820.
<https://doi.org/10.1016/j.na.2010.08.006>
- [65] W. Yuan, W. Sickel and D. Yang, *Morrey and Campanato Meet Besov, Lizorkin and Triebel*, Lecture Notes in Mathematics **2005**, Springer-Verlag, Berlin, 2010.
<https://doi.org/10.1007/978-3-642-14606-0>
- [66] Z. Zhai, *Well-posedness for fractional Navier-Stokes equations in critical spaces close to $\dot{B}_{\infty,\infty}^{-(2\beta-1)}(\mathbb{R}^n)$* , Dyn. Partial Differ. Equ. **7** (2010), no. 1, 25–44.
<https://doi.org/10.4310/dpde.2010.v7.n1.a2>

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