

Research Article

On Bounds of Eigenvalues of Complex Sturm-Liouville Boundary Value Problems

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The paper is concerned with eigenvalues of complex Sturm-Liouville boundary value problems. Lower bounds on the real parts of all eigenvalues are given in terms of the coefficients of the corresponding equation and the bound on the imaginary part of each eigenvalue is obtained in terms of the coefficients of this equation and the real part of the eigenvalue.

1. Introduction

Consider the regular complex Sturm-Liouville problem

$$\tau y := -y'' + qy = \lambda \omega y \quad (1)$$

in $L^2_w[0, 1]$ associated with the Dirichlet boundary conditions

$$y(0) = y(1) = 0, \quad (2)$$

where q is a complex-valued function and ω is a real-valued function subjected to

$$\omega(x) > 0, \quad \text{a.e. } x \in [0, 1], \quad q, \omega \in L^1[0, 1]; \quad (3)$$

λ is a spectral parameter, $L^2_w[0, 1]$ is the weighted Hilbert space of all Lebesgue measurable complex-valued functions f on $[0, 1]$ satisfying $\int_0^1 \omega |f|^2 < \infty$ with the inner product $\langle f, g \rangle = \int_0^1 \omega f \bar{g}$ and the norm $\|f\| = \langle f, f \rangle^{1/2}$, and $L^1[0, 1]$ is the set of all Lebesgue measurable complex-valued functions f on $[0, 1]$ for which $\int_0^1 |f| < \infty$.

Equation (1) is formally self-adjoint, if and only if q is a real-valued function. Hence, (1) is formally non-self-adjoint, when the imaginary part of q is nonzero. The boundary value problems associated with (1) with real coefficients have been deeply studied (cf., e.g., [1–5] and their references). Similar to regular self-adjoint boundary value problems,

one can prove that (1) and (2) also have only countable eigenvalues without finite accumulation points using the spectral theory of compact operators in Hilbert spaces when (1) is formally non-self-adjoint. Unlike self-adjoint boundary value problems, (1) and (2) may have infinitely many nonreal eigenvalues, when the imaginary part of q is nonzero (see [6, Theorem 1.1]). The asymptotic behavior of eigenvalues of boundary value problems associated with (1) has been studied (cf. [7–9]). Sufficient conditions were given in [10] for all eigenvalues of (1) with the periodic, antiperiodic, Dirichlet, or Neumann boundary conditions to be simple. For more related results for non-self-adjoint differential expressions, the reader is referred to [11–14] and the references cited therein.

In this paper, we are interested in the bounds on the eigenvalues of (1) and (2). In the case where (1) is formally self-adjoint, the bounds on eigenvalues of (1) and (2) were constructed in terms of the coefficients of (1), the coefficients of a comparing equation, and the eigenvalues of the comparing eigenvalue problem using the comparison theorem in [4]. The Rayleigh-Ritz method was used in [3] to obtain the bounds of eigenvalues of (1) and (2), when q is positive and $\omega \equiv 1$. It is noted that the comparison theorem and the Rayleigh-Ritz method are not applicable to boundary value problems (1) and (2), for the case where (1) is formally non-self-adjoint. Here, we obtain lower bounds on the real parts of all eigenvalues of (1) and (2) in terms of the coefficients of (1)

and get the bound on the imaginary part of each eigenvalue in terms of the coefficients of (1) and the real part of the eigenvalue.

In the next section, we will present the main results of this paper.

2. The Bounds on the Eigenvalues

We denote by $\|\cdot\|_C$ the maximum norm of $C[0, 1]$, which is the set of the continuous functions on $[0, 1]$, by $\|\cdot\|_1$ the norm of the space $L^1[0, 1]$ and by $\|\cdot\|_2$ the norm of the space $L^2[0, 1]$. Let q_1 and q_2 be the real and imaginary parts of q , respectively; that is, $q_1 = \operatorname{Re} q$ and $q_2 = \operatorname{Im} q$. Then, $q = q_1 + iq_2$, where $i = \sqrt{-1}$. For convenience, we set

$$q_k^- := -\min\{q_k, 0\}, \quad q_k^+ := \max\{q_k, 0\}, \quad k = 1, 2. \quad (4)$$

In addition, we denote

$$S_1(\varepsilon) := \{x \in [0, 1] : w(x) < \varepsilon\}, \quad m_1(\varepsilon) := \operatorname{mes} S_1(\varepsilon), \quad (5)$$

$$S_2(\varepsilon) := \{x \in [0, 1] : w^2(x) < \varepsilon\}, \quad m_2(\varepsilon) := \operatorname{mes} S_2(\varepsilon). \quad (6)$$

Then the result below is one of the main results of the paper.

Theorem 1. *If λ is an eigenvalue of (1) and (2), then*

$$\operatorname{Re} \lambda \geq -\frac{8}{\varepsilon_1} (\|q_1^+\|_1 + 2\|q_1^-\|_1) \|q_1^-\|_1; \quad (7)$$

if λ is an eigenvalue of (1) and (2), with $\operatorname{Re} \lambda \leq 0$, then

$$|\operatorname{Im} \lambda| \leq \frac{8}{\varepsilon_1} \|q_2\|_1 \|q_1^-\|_1, \quad (8)$$

where $\varepsilon_1 > 0$ satisfies $8\|q_1^-\|_1^2 m_1(\varepsilon_1) < 1$; and if λ is an eigenvalue of (1) and (2), with $\operatorname{Re} \lambda > 0$, then

$$|\operatorname{Im} \lambda| \leq \frac{8}{\tilde{\varepsilon}_1} \|q_2\|_1 (\|q_1^-\|_1 + \operatorname{Re} \lambda \|w\|_1), \quad (9)$$

where $\tilde{\varepsilon}_1 > 0$ satisfies $8(\|q_1^-\|_1 + \operatorname{Re} \lambda \|w\|_1)^2 m_1(\tilde{\varepsilon}_1) < 1$.

Proof. We first consider the case where λ is an eigenvalue of (1) and (2), with $\operatorname{Re} \lambda \leq 0$. Let λ be such an eigenvalue and let $\phi(x)$ be the corresponding eigenfunction with $\|\phi\|_2 = 1$. Then, $\phi(0) = \phi(1) = 0$ and

$$-\phi'' + q\phi = \lambda w\phi. \quad (10)$$

By multiplying both sides of (10) by $\bar{\phi}$ and integrating over the interval $[0, 1]$, we have

$$\int_0^1 |\phi'|^2 + \int_0^1 q|\phi|^2 = \lambda \int_0^1 w|\phi|^2. \quad (11)$$

By separating the real and imaginary parts of both sides of (11), we get that

$$\operatorname{Re} \lambda \int_0^1 w|\phi|^2 = \int_0^1 |\phi'|^2 + \int_0^1 q_1|\phi|^2, \quad (12)$$

$$\operatorname{Im} \lambda \int_0^1 w|\phi|^2 = \int_0^1 q_2|\phi|^2. \quad (13)$$

From $\operatorname{Re} \lambda \leq 0$ and (12), it follows that

$$\int_0^1 |\phi'|^2 \leq \int_0^1 q_1^- |\phi|^2. \quad (14)$$

Now, let $Q_1(x) = \int_0^x q_1^-(t) dt$; then we have $\max_{x \in [0, 1]} |Q_1(x)| \leq \|q_1^-\|_1$ and

$$\begin{aligned} \int_0^1 q_1^- |\phi|^2 &= \int_0^1 Q_1' |\phi|^2 = - \int_0^1 Q_1 (\phi' \bar{\phi} + \phi \bar{\phi}') \\ &= -2 \operatorname{Re} \int_0^1 Q_1 \phi' \bar{\phi} \leq 2 \|q_1^-\|_1 \int_0^1 |\phi'| |\bar{\phi}| \\ &\leq 2 \|q_1^-\|_1 \|\phi'\|_2 \leq 2 \|q_1^-\|_1^2 + \frac{1}{2} \|\phi'\|_2^2 \end{aligned} \quad (15)$$

by the Cauchy-Schwarz inequality. It can be obtained from (14) and (15) that

$$\|\phi'\|_2^2 \leq 2 \|q_1^-\|_1^2 + \frac{1}{2} \|\phi'\|_2^2, \quad (16)$$

and, hence,

$$\|\phi'\|_2^2 \leq 4 \|q_1^-\|_1^2, \quad \int_0^1 q_1^- |\phi|^2 \leq 4 \|q_1^-\|_1^2. \quad (17)$$

Similar to (15), we can get that

$$\int_0^1 q_1^+ |\phi|^2 \leq 2 \|q_1^+\|_1 \|\phi'\|_2, \quad (18)$$

which, together with the first relation of (17), gives that

$$\int_0^1 q_1^+ |\phi|^2 \leq 4 \|q_1^+\|_1 \|q_1^-\|_1. \quad (19)$$

In addition, from $\phi(x) = \int_0^x \phi'(t) dt$ and the Cauchy-Schwarz inequality, we have

$$|\phi(x)|^2 = \left| \int_0^x \phi'(t) dt \right|^2 \leq x \int_0^x |\phi'(t)|^2 dt \leq \int_0^1 |\phi'|^2 = \|\phi'\|_2^2, \quad x \in [0, 1], \quad (20)$$

by which, together with (17), the definition of $m_1(\varepsilon_1)$ in (5), and $8\|q_1^-\|_1^2 m_1(\varepsilon_1) \leq 1$, we get that

$$\begin{aligned} \int_0^1 w|\phi|^2 &\geq \int_{[0, 1] \setminus S_1(\varepsilon_1)} w|\phi|^2 \geq \varepsilon_1 \int_{[0, 1] \setminus S_1(\varepsilon_1)} |\phi|^2 \\ &= \varepsilon_1 \left(1 - \int_{S_1(\varepsilon_1)} |\phi|^2 \right) \\ &\geq \varepsilon_1 \left(1 - \|\phi'\|_2^2 m_1(\varepsilon_1) \right) \geq \frac{\varepsilon_1}{2}. \end{aligned} \quad (21)$$

Using (12), (17), (19), and (21), we can easily conclude that, for every eigenvalue λ with $\operatorname{Re} \lambda \leq 0$,

$$|\operatorname{Re} \lambda| \frac{\varepsilon_1}{2} \leq 4 (\|q_1^+\|_1 + 2\|q_1^-\|_1) \|q_1^-\|_1, \quad (22)$$

which implies that (7) holds for all the eigenvalues of (1) and (2).

Let $Q_2(x) = \int_0^x q_2(t)dt$. Then, from (17),

$$\int_0^1 q_2|\phi|^2 \leq 2\|q_2\|_1\|\phi'\|_2 \leq 4\|q_2\|_1\|q_1^-\|_1, \tag{23}$$

which, together with (13) and (21), implies that (8) holds for every eigenvalue λ with $\text{Re } \lambda \leq 0$.

Now, let λ_0 be an eigenvalue of (1) and (2), with $\text{Re } \lambda_0 > 0$. Then, we consider the problem

$$-y'' + (q - \text{Re } \lambda_0 w) y = \lambda w y \tag{24}$$

with the Dirichlet boundary conditions (2). It can be easily verified that $\lambda_0 - \text{Re } \lambda_0$ is an eigenvalue of (24) and (2). Clearly, $\text{Re}(\lambda_0 - \text{Re } \lambda_0) = 0$. Hence, by (8), there exists $\bar{\varepsilon}_1$ satisfying $8(\|q_1^-\|_1 + \text{Re } \lambda_0 \|w\|_1)^2 m_1(\bar{\varepsilon}_1) < 1$ such that

$$|\text{Im } \lambda_0| = |\text{Im}(\lambda_0 - \text{Re } \lambda_0)| \leq \frac{8}{\bar{\varepsilon}_1} \|q_2\|_1 (\|q_1^-\|_1 + \text{Re } \lambda_0 \|w\|_1). \tag{25}$$

This completes the proof. \square

The following corollary is a direct consequence of Theorem 1.

Corollary 2. *If $q_1 \geq 0$, then all eigenvalues λ of (1) and (2) satisfy that $\text{Re } \lambda \geq 0$. Consequently, if there exists $q_0 \in \mathbf{R}$ such that $q_1 \geq q_0 w$ on $[0, 1]$, then all eigenvalues λ of (1) and (2) satisfy that $\text{Re } \lambda \geq q_0$.*

If $w \in AC[0, 1]$ and $w' \in L^2[0, 1]$, where $AC[0, 1]$ denotes the set of functions which are locally absolutely continuous on $[0, 1]$, then we have the following result.

Theorem 3. *Assume that $w \in AC[0, 1]$ and $w' \in L^2[0, 1]$. If λ is an eigenvalue of (1) and (2), then*

$$\text{Re } \lambda \geq -\frac{8}{\varepsilon_2} \|q_1^-\|_1^2 [\|w\|_C (1 + \|q_1\|_1) + \|w'\|_2]; \tag{26}$$

if λ is an eigenvalue of (1) and (2), with $\text{Re } \lambda \leq 0$, then

$$|\text{Im } \lambda| \leq \frac{8}{\varepsilon_2} \|q_1^-\|_1^2 (\|w\|_C \|q_2\|_1 + \|w'\|_2), \tag{27}$$

where $\varepsilon_2 > 0$ satisfies $8\|q_1^-\|_1^2 m_2(\varepsilon_2) < 1$; and if λ is an eigenvalue of (1) and (2), with $\text{Re } \lambda > 0$, then

$$|\text{Im } \lambda| \leq \frac{8}{\bar{\varepsilon}_2} (\|q_1^-\|_1 + \text{Re } \lambda \|w\|_1)^2 (\|w\|_C \|q_2\|_1 + \|w'\|_2), \tag{28}$$

where $\bar{\varepsilon}_2 > 0$ satisfies $8(\|q_1^-\|_1 + \text{Re } \lambda \|w\|_1)^2 m_2(\bar{\varepsilon}_2) < 1$.

Proof. Let λ be an eigenvalue of (1) and (2) with $\text{Re } \lambda \leq 0$ and let $\phi(x)$ be the corresponding eigenfunction with $\|\phi\|_2 = 1$. Then, $\phi(0) = \phi(1) = 0$ and (10) holds. By multiplying both

sides of (10) by $w\bar{\phi}$ and integrating over the interval $[0, 1]$, we have

$$\int_0^1 w|\phi'|^2 + \int_0^1 w'\phi'\bar{\phi} + \int_0^1 wq|\phi|^2 = \lambda \int_0^1 w^2|\phi|^2. \tag{29}$$

By separating the real and imaginary parts of both sides of (29), we get that

$$\text{Re } \lambda \int_0^1 w^2|\phi|^2 = \int_0^1 w|\phi'|^2 + \text{Re} \left(\int_0^1 w'\phi'\bar{\phi} \right) + \int_0^1 wq_1|\phi|^2, \tag{30}$$

$$\text{Im } \lambda \int_0^1 w^2|\phi|^2 = \text{Im} \left(\int_0^1 w'\phi'\bar{\phi} \right) + \int_0^1 wq_2|\phi|^2. \tag{31}$$

On the other hand, it follows from (17) and (20) that

$$\begin{aligned} \int_0^1 w|\phi'|^2 &\leq \|w\|_C \|\phi'\|_2^2 \leq 4\|w\|_C \|q_1^-\|_1^2, \\ \left| \int_0^1 wq_1|\phi|^2 \right| &\leq \|w\|_C \|\phi'\|_2^2 \|q_1\|_1 \leq 4\|w\|_C \|q_1^-\|_1^2 \|q_1\|_1, \\ \left| \int_0^1 w'\phi'\bar{\phi} \right| &\leq \|w'\|_2 \|\phi'\|_2^2 \leq 4\|w'\|_2 \|q_1^-\|_1^2, \\ \left| \int_0^1 wq_2|\phi|^2 \right| &\leq \|w\|_C \|q_2\|_1 \|\phi'\|_2^2 \leq 4\|w\|_C \|q_1^-\|_1^2 \|q_2\|_1. \end{aligned} \tag{32}$$

In addition, by (17), the definition of $m_2(\varepsilon_2)$ in (6), and $8\|q_1^-\|_1^2 m_2(\varepsilon_2) < 1$, it can be obtained that

$$\begin{aligned} \int_0^1 w^2|\phi|^2 &\geq \int_{[0,1] \setminus S_2(\varepsilon_2)} w^2|\phi|^2 \geq \varepsilon_2 \int_{[0,1] \setminus S_2(\varepsilon_2)} |\phi|^2 \\ &= \varepsilon_2 \left(1 - \int_{S_2(\varepsilon_2)} |\phi|^2 \right) \\ &\geq \varepsilon_2 \left(1 - 4\|q_1^-\|_1^2 m_2(\varepsilon_2) \right) \geq \frac{\varepsilon_2}{2}, \end{aligned} \tag{33}$$

which, together with (30) and (32), implies that (26) holds for every eigenvalue λ of (1) and (2), with $\text{Re } \lambda \leq 0$, and, hence, (26) holds for all eigenvalues of (1) and (2). Furthermore, if λ is an eigenvalue of (1) and (2), with $\text{Re } \lambda \leq 0$, then (27) follows from (31)–(33). With a similar argument to that in the proof of Theorem 1, (28) can be proved. This completes the proof. \square

Conflict of Interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

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