# ON DUAL SERIES RELATIONS INVOLVING LAGUERRE POLYNOMIALS 

K. N. Srivastava

In this paper, we shall consider the problem determining the sequence $\{A n\}$, such that
$\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1)\} L_{n}^{\alpha}(x)=f_{1}(x), \quad 0 \leqq x<y$,
$\sum_{n=0}^{\infty}\left\{A n / \Gamma(n+\alpha+1 / 2\} L_{n}^{\alpha}(x)=f_{2}(x), y<x \leqq \infty, \alpha>-1 / 2\right.$, where $L_{n}^{\alpha}(x)$ is a Laguerre polynomial, the functions $f_{1}(x)$ and $f_{2}(x)$ being prescribed. By expressing the sequence $\{A n\}$ in terms of a sequence of integrals involving an unknown function $g(u)$ the problem is reduced to that of solving an Abel integral equation for $g(u)$.

In recent years, dual series relations involving Fourier-Bessel, Dini series, trigonometric series and series of Jacobi polynomials have been investigated by various workers [1, 2, 5 to 12]. Here we shall apply the method developed by Sneddon and Srivastav for obtaining a solution of the dual series relations involving Laguerre polynomials.

As pointed out by Sneddon and Srivastav [6], with a view to simplify the calculations, we split the problem posed by the pair of dual equations given above into two parts: Problem (a). Determine the constants $\{A n\}$ satisfying the dual series relations

$$
\begin{gather*}
\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1)\} L_{n}^{\alpha}(x)=f_{1}(x), \quad 0 \leqq x<y,  \tag{1.1}\\
\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1 / 2)\} L_{n}^{\alpha}(x)=0, \quad y<x \leqq \infty, \alpha>-1 / 2 .
\end{gather*}
$$

Problem (b). Determine the constants $\{A n\}$ satisfying the dual series relations

$$
\begin{align*}
\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1)\} L_{n}^{\alpha}(x)=0, & 0 \leqq x<y  \tag{1.3}\\
\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1 / 2)\} L_{n}^{\alpha}(x)=f_{2}(x), & y<x \leqq \infty, \alpha>-1 / 2 \tag{1.4}
\end{align*}
$$

The solution of the general problem is obviously obtained merely by adding the solutions of problem (a) and (b). We suppose that functions $f_{1}(x)$ and $f_{2}(x)$ satisfy the following conditions:
(i) $F_{1}(x)=x^{\alpha} f_{1}(x)$ is finite and continuously differentiable for $0 \leqq x<y$,
(ii) $F_{2}(x)=\int_{x}^{\infty} e^{-x} f_{2}(x) d x$ is finite and continuously differentiable for $y<x \leqq \infty$.
As we shall presently see the classes of functions $f_{1}(x)$ and $f_{2}(x)$ for which the problem under discussion is solvable, must satisfy the above conditions,
2. In this section we list some results for ready reference. By combining the results [3, p. 292 (2), (3)], we have

$$
\begin{equation*}
\int_{0}^{\infty} x^{\alpha} e^{-x} L_{n}^{\alpha}(x) L_{m}^{\alpha}(x) d x=(n+1)_{\alpha} \cdot \delta_{m n}, \tag{2.1}
\end{equation*}
$$

where $\delta_{m n}$ is a Kronecker delta. From [4, p. 193 (27), (28)] we have

$$
\begin{gather*}
\frac{d}{d n}\left\{x^{\alpha} L_{n}^{\alpha}(x)\right\}=(n+\alpha) x^{\alpha-1} L_{n}^{\alpha-1}(x)  \tag{2.2}\\
\int_{x}^{\infty} e^{-y} L_{n}^{\alpha}(y) d y=e^{-x} L_{n}^{\alpha-1}(x) \tag{2.3}
\end{gather*}
$$

We shall also require the following results which are easily derived from the more general results given in [3, p. 293 (5), p. 405 (20)]. For $\alpha>-1 / 2$

$$
\begin{gather*}
\int_{x}^{\infty}(y-x)^{-1 / 2} e^{-y} L_{n}^{\alpha}(y) d y=\Gamma(1 / 2) e^{-x} L_{n}^{\alpha-1 / 2}(x),  \tag{2.4}\\
\int_{0}^{x}(x-y)^{-1 / 2} y^{\alpha} L_{n}^{\alpha}(y) d y=\frac{\Gamma(n+\alpha+1) \Gamma\left({ }^{1 / 2}\right)}{\Gamma(n+\alpha+3 / 2)} x^{\alpha+1 / 2} L_{n}^{\alpha+1 / 2}(x) \tag{2.5}
\end{gather*}
$$

We also note that if $f(x)$ is continuously differentiable then Abel integral equation

$$
\begin{equation*}
f(x)=\int_{0}^{x} \frac{\phi(y)}{(x-y)^{1 / 2}} d y \tag{2.6}
\end{equation*}
$$

has a continuous solution given by the equation

$$
\begin{equation*}
\phi(y)=\frac{1}{\Pi} \frac{d}{d y} \int_{0}^{y} \frac{f(x)}{(y-x)^{1 / 2}} d x \tag{2.7}
\end{equation*}
$$

Furthermore, if $f(x)$ is continuously differentiable then the integral equation

$$
\begin{equation*}
f(x)=\int_{x}^{\infty} \frac{\phi(y)}{(y-x)^{1 / 2}} d y \tag{2.8}
\end{equation*}
$$

has a continuous solution

$$
\begin{equation*}
\phi(y)=-\frac{1}{\Pi} \frac{d}{d y} \int_{y}^{\infty} \frac{f(x)}{(x-y)^{1 / 2}} d x \tag{2.9}
\end{equation*}
$$

This can be easily established by simple methods given in [13, p. 229]. The analysis given here is purely formal and no attempt is made to justify the interchange of various limiting processes.
3. Solution of the problem (a). Let us suppose that for $0 \leqq$ $x<y$

$$
\begin{equation*}
\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1 / 2)\} L_{n}^{\alpha}(x)=-e^{x} \frac{d}{d x} \int_{x}^{y} \frac{g_{1}(u)}{(u-x)^{1 / 2}} d u \tag{3.1}
\end{equation*}
$$

Using the orthogonal property (2.1), it can be shown that

$$
\begin{equation*}
A n=-\frac{\Gamma(n+\alpha+1 / 2) \Gamma(1 / 2)}{\Gamma(n+\alpha+1)} \int_{0}^{y} x^{\alpha} L_{n}^{\alpha}(x)\left(\frac{d}{d x} \int_{x}^{y} \frac{g_{1}(u)}{(u-x)^{1 / 2}} d u\right) d x \tag{3.2}
\end{equation*}
$$

Since

$$
\begin{equation*}
-\frac{d}{d x} \int_{x}^{y} \frac{g_{1}(u)}{(u-x)^{1 / 2}} d u=\frac{g_{1}(y)}{(y-x)^{1 / 2}}-\int_{x}^{y} \frac{d}{d u}\left\{g_{1}(u)\right\}, ~ d u \tag{3.3}
\end{equation*}
$$

we obtain with the help of (2.5), the equation
(3.4) $\quad A n=\Gamma(n+1) \Gamma(1 / 2) \int_{0}^{y} g_{1}(u) u^{\alpha-1 / 2} L_{n}^{\alpha-1 / 2}(u) d u, \quad n=0,1,2, \cdots$

If in the equation (1.1), we substitute for the coefficients $A n$ from (3.4), on interchanging the order of summation and integration, we get

$$
\begin{equation*}
f_{1}(x)=\int_{0}^{y} g_{1}(u) u^{\alpha-1 / 2} K_{1}(u, x) d u, \quad 0 \leqq x<y, \tag{3.5}
\end{equation*}
$$

where

$$
\begin{equation*}
K_{1}(u, x)=\sum_{n=0}^{\infty} \frac{\Gamma(n+1) \Gamma(1 / 2)}{\Gamma(n+\alpha+1)} L_{n}^{\alpha-1 / 2}(u) L_{n}^{\alpha}(x) \tag{3.6}
\end{equation*}
$$

with the help of equations (2.1) and (2.4) it can be shown that

$$
\begin{equation*}
K_{1}(u, x)=e^{u} x^{-\alpha}(x-u)^{-1 / 2} H(x-u) \tag{3.7}
\end{equation*}
$$

where $H(t)$ is Heaviside's unit function. (2.7) is easily proved. Let

$$
K_{1}(u, x)=\sum_{n=0}^{\infty} a_{n} L_{n}^{\alpha}(x)
$$

where the coefficients $a_{n}$ are given by

$$
\begin{aligned}
a_{n} & =\frac{\Gamma(n+1)}{\Gamma(n+\alpha+1)} \int_{0}^{\infty} K_{1}(u, x) x^{\alpha} e^{-n} L_{n}^{\alpha}(x) \\
& =\frac{\Gamma(n+1)}{\Gamma(n+\alpha+1)} e^{u} \int_{u}^{\infty} e^{-x}(x-u)^{-1 / 2} L_{n}^{\alpha}(x) d x \\
& =\frac{\Gamma(n+1) \Gamma(1 / 2)}{\Gamma(n+\alpha+1)} L_{n}^{\alpha-1 / 2}(u)
\end{aligned}
$$

Thus the equation (3.5) is equivalent to

$$
\begin{equation*}
F_{1}(x)=x^{\alpha} f_{1}(x)=\int_{0}^{x} \frac{g_{1}(u) u^{\alpha-1 / 2} e^{u}}{(x-u)^{1 / 2}} d u, \quad 0 \leqq x<y \tag{3.8}
\end{equation*}
$$

This is Abel integral equation, since $F_{1}(x)$ is finite and continuously differentiable, its solution is given by

$$
\begin{equation*}
u^{\alpha-1 / 2} e^{u} g_{1}(u)=\frac{1}{I I} \frac{d}{d u} \int_{0}^{u} \frac{x^{\alpha} f_{1}(x)}{(u-x)^{1 / 2}} d x . \tag{3.9}
\end{equation*}
$$

The coefficients An may now be calculated with the help of the relations (3.4) and (3.9).
4. Solution of the problem (b). We start with the assumption that for $y<x \leqq \infty$

$$
\begin{equation*}
\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1)\} L_{n}^{\alpha}(x)=x^{-\alpha} \int_{y}^{x} \frac{g_{2}(u)}{(x-u)^{1 / 2}} d u . \tag{4.1}
\end{equation*}
$$

This is equivalent to assuming that

$$
\begin{equation*}
A n=\Gamma(n+1) \Gamma(1 / 2) \int_{y}^{\infty} g_{2}(u) e^{-u} L_{n}^{\alpha-1 / 2}(u) d u, n=0,1,2, \cdots \tag{4.2}
\end{equation*}
$$

If we multiply both sides of the equation (1.4) by $\exp (-x)$ and integrate with respect to $x$ from $x$ to $\infty, y<x \leqq \infty$, we obtain

$$
\begin{equation*}
F_{2}(x)=\int_{x}^{\infty} e^{-x} f_{2}(x) d x=\sum_{n=0}^{\infty}\{A n / \Gamma(n+\alpha+1)\} e^{-x} L_{n}^{\alpha-1}(x) . \tag{4.3}
\end{equation*}
$$

Substituting the values of the coefficients from (4.2) in the equation (4.3) we find on interchanging the order of summation and integration that

$$
\begin{equation*}
e^{x} F_{2}(x)=\int_{y}^{\infty} g_{2}(u) e^{-u} K_{2}(u, x) d u, y<x \leqq \infty, \tag{4.4}
\end{equation*}
$$

where

$$
\begin{equation*}
K_{2}(u, x)=\sum_{n=0}^{\infty} \frac{\Gamma(n+1) \Gamma(1 / 2)}{\Gamma(n+\alpha+1 / 2)} L_{n}^{\alpha-1 / 2}(u) L_{n}^{\alpha-1}(x) \tag{4.5}
\end{equation*}
$$

From the relations (2.1) and (2.5) it easily follows that

$$
\begin{equation*}
K_{2}(u, x)-e^{x} u^{-\alpha+1 / 2}(u-x)^{-1 / 2} H(u-x) . \tag{4.6}
\end{equation*}
$$

Consequently the equation (4.4) reduces to the integral equation

$$
\begin{equation*}
F_{2}(x)=\int_{y}^{\infty} \frac{g_{2}(u) u^{1 /--\alpha} e^{-u}}{(u-x)^{1 / 2}} d u, \quad y<x \leqq \infty \tag{4.7}
\end{equation*}
$$

Since $F_{2}(x)$ is finite and continuously differentiable, the solution of the above equation is given by

$$
\begin{equation*}
g_{2}(u)=-\frac{e^{u} u^{\alpha-1 / 2}}{\Pi} \frac{d}{d u} \int_{0}^{u} \frac{F_{2}(x)}{(x-u)^{1 / 2}} d x . \tag{4.8}
\end{equation*}
$$

The coefficients $A_{n}$ are given by the relations (4.2) and (4.8).

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M. A. College of Technology, Bhopal India

