## ON THE CORRELATION COEFFICIENT OF A BIVARIATE, EQUAL VARIANCE, COMPLEX GAUSSIAN SAMPLE<sup>1</sup>

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Let  $u_n$  denote the sample correlation coefficient for n observations from a bivariate, equal variance, complex Gaussian distribution. In this note we derive the exact distribution of  $u_n$  by extending a method of Mehta and Gurland to the complex case. The asymptotic behavior of  $E|u_n|^k$  as  $n \to \infty$  is determined via the method of steepest descent. Applicability of the results to the analysis of certain estimators of spectral parameters of stationary time series is discussed.

1. Density of  $u_n$ . Let  $\xi_1, \xi_2, \dots, \xi_n$  be an independent sample from a zero mean, bivariate, equal variance, complex Gaussian distribution with correlation matrix

(1) 
$$\Sigma_{\xi} = E \xi \bar{\xi}' = \begin{pmatrix} \sigma^2 & \sigma^2 \rho \\ \sigma^2 \bar{\rho} & \sigma^2 \end{pmatrix}.$$

The Hermitian sample correlation matrix  $A = (A_{ij}) = n^{-1} \sum_{k=1}^{n} \xi_k \dot{\xi}_k'$ ,  $1 \le i$ ,  $j \le 2$ , then has the bivariate complex Wishart density

(2) 
$$p(A) = |A|^{n-2} [\pi \Gamma(n) \Gamma(n-1) |\Sigma_{\varepsilon}|^{n}]^{-1} \exp [-\operatorname{tr}(\Sigma_{\varepsilon}^{-1} A)].$$

(See Goodman (1963) for a detailed discussion of the complex Gaussian and complex Wishart distributions.) The function p(A), defined over the domain where A is Hermitian positive semi-definite, is a compact way of writing the joint density of the four real random variables  $A_{11}$ ,  $A_{22}$ ,  $A_{12R}$ , and  $A_{12I}$ . The usual estimator  $u_n$  of the complex correlation coefficient  $\rho$  is a function of the elements of A, namely

$$u_n = 2A_{12}[A_{11} + A_{22}]^{-1}$$
.

Throughout the remainder of this section we suppress the subscript n on  $u_n$ . The joint probability density of the real random variables  $u_R$  and  $u_I$  defined by the relation  $u=u_R+iu_I$  may be found by extending the method of Mehta and Gurland (1969) to the complex case as follows. First, two auxiliary variables  $v=A_{11}+A_{22}$  and  $w=A_{22}$  are introduced. A simple calculation reveals that the magnitude of the Jacobian of the transformation  $(A_{11},A_{22},A_{12R},A_{12I}) \rightarrow (u_R,u_I,v,w)$  equals  $(v/2)^2$ . It follows from (1), (2), and some algebra that  $p(u_R,u_I,v,w)=K_1(v/2)^2[(v-w)w-(|u|v/2)^2]^{n-2}\exp[-v\sigma^{11}+v\operatorname{Re}(\sigma^{12}\bar{u})]$ , where  $K_1^{-1}=\Gamma(n)\Gamma(n-1)\pi\sigma^{4n}(1-|\rho|^2)^n$ ,  $\sigma^{jk}$  is the (j,k)th element of  $\Sigma_{\varepsilon}^{-1}$  and the

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density vanishes outside the region  $0 \le w \le v < \infty$ ,  $|u| \le 2[(w/v) - (w/v)^2]^{\frac{1}{2}}$ . Although  $p(u_R, u_I)$  may be obtained by integrating  $p(u_R, u_I, v, w)$  over this region with respect to v and w, it proves more convenient first to introduce the variable t defined by

$$w = (v/2)[1 + t(1 - |u|^2)^{\frac{1}{2}}].$$

Then  $p(u_R, u_I, v, t) = p(u_R, u_I, v, w) |\partial w/\partial t|$  is supported on the cylinder  $0 \le |u| \le 1$ ,  $0 \le v < \infty$ ,  $-1 \le t \le 1$ , wherein it assumes the form

$$p(u_R, u_I, v, t) = K_1(1 - t^2)^{n-2}(1 - |u|^2)^{(2n-3)/2}(v/2)^{2n-1} \exp[-v\sigma^{11} + v \operatorname{Re}(\sigma^{12}\bar{u})].$$

Integrating over v and t yields

(3) 
$$p(u_R, u_I) = K_2(1 - |u|^2)^{(2n-3)/2} [\sigma^{11} - \operatorname{Re}(\sigma^{12}\bar{u})]^{-2n}, \quad |u| \leq 1,$$

where 
$$K_2 = \Gamma(2n)[\pi^{\frac{1}{2}}\Gamma(n)\Gamma(n-\frac{1}{2})\sigma^{4n}(1-|\rho|^2)^n2^{2n-1}]^{-1}$$
.

Subsequent discussion is facilitated by transforming (3) to polar coordinates. Letting  $u = |u|e^{i\theta_u}$  and  $\rho = |\rho|e^{i\theta_\rho}$ , and using the explicit expressions for  $\sigma^{11}$  and  $\sigma^{12}$  from (1), we obtain

(4) 
$$p(|u|, \theta_u) = C|u|(1 - |u|^2)^{(2n-3)/2}[1 - |u||\rho|\cos(\theta_u - \theta_\rho)]^{-2n}$$
$$(0 \le |u| \le 1, -\pi \le \theta_u \le \pi),$$

where

(5) 
$$C = (1 - |\rho|^2)^n \Gamma(2n) \left[\pi^{\frac{1}{2}} \Gamma(n) \Gamma(n - \frac{1}{2}) 2^{2n-1}\right]^{-1}.$$

Note that |u| and  $\theta_u$  are *not* statistically independent. Extensions of certain of the above results to the unequal variances case appear in Goodman (1957).

2. Asymptotic behavior of  $E|u_n|^k$ . The asymptotic behavior of  $E|u_n|^k$  for large n may be obtained as follows. From (4) we have

(6) 
$$E|u_n|^k = C \int_{-\pi}^{\pi} d\theta \int_0^1 dx \frac{x^{k+1}}{(1-x^2)^{\frac{3}{2}}} \left[ \frac{1-x^2}{(1-x|\rho|\cos\theta)^2} \right]^n.$$

For  $|\theta| < \pi/2$  the integrand has a high peak centered at  $x = |\rho| \cos \theta$  when n is large, while for  $\pi/2 < |\theta| < \pi$  the integrand contains the nth power of a factor that is smaller than 1 for all  $0 \le x \le 1$ . Accordingly, the error associated with limiting the range of integration to  $|\theta| < \pi/2$  is negligible when n is large. For each such  $\theta$  the inner integral in (6) is of the form

$$I = \int_0^1 g(x) \exp[-nf(x)] dx,$$

where  $g(x) = x^{k+1}(1-x^2)^{-\frac{3}{2}}$ ,  $f(x) = \log \left[ (1-bx)^2/(1-x^2) \right]$ , and  $b = |\rho| \cos \theta$ . Since f(x) has a minimum at x = b and  $f''(b) = 2(1-b^2)^{-2}$ , the method of steepest descent yields  $I \sim (\pi/n)^{\frac{1}{2}}b^{k+1}/(1-b^2)^{n+\frac{1}{2}}$ . Therefore

(7) 
$$E|u_n|^k \sim \left(\frac{\pi}{n}\right)^{\frac{1}{2}} C \int_{-\pi/2}^{\pi/2} \frac{(|\rho|\cos\theta)^{k+1}}{(1-|\rho|^2\cos^2\theta)^{n+\frac{1}{2}}} d\theta .$$

Since |
ho|<1 in nondegenerate cases, the integrand in (7) peaks sharply at heta=0

for large n. Steepest descent therefore is applicable again. In this regard we note that, to second order in  $\theta$ ,

$$(1 - |\rho|^2 \cos^2 \theta)^{-n}$$

$$= \exp\left[-n \log(1 - |\rho|^2 \cos^2 \theta)\right] \sim (1 - |\rho|^2)^{-n} \exp\left(-\frac{n|\rho|^2 \theta^2}{1 - |\rho|^2}\right).$$

It follows that

$$E|u_n|^k \sim \left(\frac{\pi}{n}\right)^{\frac{1}{2}} C \frac{|\rho|^{k+1}}{(1-|\rho|^2)^{n+\frac{1}{2}}} \int_{-\infty}^{\infty} \exp\left(-\frac{n|\rho|^2 \theta^2}{1-|\rho|^2}\right) d\theta = \frac{\pi C|\rho|^k}{n(1-|\rho|^2)^n}.$$

Substituting for C from (5) yields the desired result,

(8) 
$$E|u_n|^k \sim \frac{\pi^{\frac{1}{2}}\Gamma(2n)|\rho|^k}{n\Gamma(n)\Gamma(n-\frac{1}{2})2^{2n-1}}.$$

Asymptotic expansion of the gamma functions in (8) verifies the intuitively obvious fact that  $E|u_n|^k \to |\rho|^k$  as  $n \to \infty$ , whereupon the  $L_r$ -convergence theorem (Loève (1963)) implies that the  $|u_n|$  converge in rth mean to the constant  $|\rho|$  for all r > 0.

3. Applicability to estimation of spectral parameters. The above results find direct application to the analysis of certain radar estimates of spectral parameters of distributed-velocity media such as storm clouds and clear air turbulence. Specifically, the in-phase and quadrature signals returned from the portion of such a medium that is located at a fixed range from the transmitter may be modeled as the real and imaginary parts, respectively, of a zero mean complex Gaussian random process  $\{Z_i\}$ . It follows that range gating of the returns from a pair of narrow radar pulses spaced T seconds apart produces a bivariate complex Gaussian random variable  $\xi$ . If there is no clutter from ambiguous range cells, then  $\xi =$  $(Z_t, Z_{t+T})$ . Moreover, if  $\{Z_t\}$  is wide-sense stationary, which usually is the case, then  $Z_t$  and  $Z_{t+T}$  have equal variances. Their correlation matrix  $\Sigma_{\xi}$  then is of the form (1) with  $\sigma^2 = \int dF(\gamma)$  and  $\sigma^2 \rho = \int e^{i2\pi\gamma T} dF(\gamma)$ , where F is the spectral distribution function of  $\{Z_t\}$ . An independent sample  $\xi_j = (Z_{t,j}, Z_{t,j+T}), 1 \le 1$  $j \le n$ , of such bivariate, equal variance, complex Gaussian random variables may be obtained either by frequency-stepping the radar carrier frequency or by inserting sufficiently long delays between successive pulse pairs. In many applications it is of interest to estimate the centroid  $\gamma_0$  and spread  $\sigma_{\gamma}^2$  of F, which are defined by  $\int (\gamma - \gamma_0) dF(\gamma) = 0$  and  $\sigma^2 \sigma_{\gamma}^2 = \int (\gamma - \gamma_0)^2 dF(\gamma)$ . Reasonable estimators of these quantities proposed by Rummler (1968a) are expressible as functions of the complex statistic  $u=|u|e^{i\theta_u}$ , specifically  $\hat{\gamma}_0=(2\pi T)^{-1}\theta_u$  and  $\hat{\sigma}_{\gamma}{}^2=$  $(2\pi^2T^2)^{-1}(1-|u|)$ . Some properties of these estimators and of certain extensions of them to cases in which  $\{Z_t\}$  is corrupted by additive, independent, "white" receiver noise have been explored by Rummler (1968b), Hofstetter (1970), and Berger (1971).

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