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Metastability of a random walk with catastrophes

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Abstract

We consider a random walk with catastrophes which was introduced to model population biology. It is known that this Markov chain gets eventually absorbed at 0 for all parameter values. Recently, it has been shown that this chain exhibits a metastable behavior in the sense that it can persist for a very long time before getting absorbed. In this paper we study this metastable phase by making the parameters converge to extreme values. We obtain four different limits that we believe shed light on the metastable phase.

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1 A random walk with catastrophes

Let $(X_n)_{n\geq 0}$ be the following discrete-time Markov chain. For every $n \geq 1$, X_n is a non-negative integer. The model has two parameters, $p \in (0,1)$ and $c \in (0,1)$. For $n \geq 0$, let $X_n = k \geq 0$. If $X_n = 0$ then $X_{n+1} = 0$. That is, 0 is an absorbing state. If $k \geq 1$, there are two possibilities:

- With probability p there is a birth. Then, $X_{n+1} = k + 1$.
- With probability 1 p there is a catastrophe. Then, $X_{n+1} = k B_n$, where B_n is a binomial random variable with parameters k and c. The random variables B_n are sampled independently of each other and of everything else.

In [2] it was shown that this Markov chain is eventually absorbed at 0 for all values of p and c in (0, 1). They have also shown through simulations and first moment computations that the time to absorption can be unusually long, in particular if p is close to 1 or c is close to 0. Before absorption the chain fluctuates in a narrow band around

$$n^* = \frac{p}{(1-p)c}.$$

That is, the chain seems to have reached some equilibrium but this equilibrium turns out to be unstable, see Figure 1. This is why we think of the chain as exhibiting metastable

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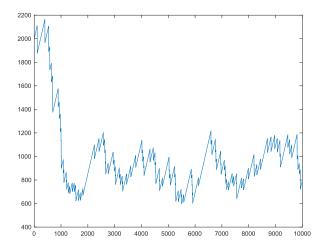


Figure 1: We ran this simulation for 10^4 steps, starting at $X_0 = 2000$, with p = 0.99 and c = 0.1. We see that the process drops very fast to $n^* = 990$ and then fluctuates around this value.

behavior. By looking at the limits of the process for $c \to 0$ and/or $p \to 1$ we will get more insight into this metastable phase. We find four different limit processes explicitly. In the expression of n^* , (1-p) and c play the same role. However, we will show that as capproaches 0 and 1-p approaches 0 the limiting processes are very different.

This model goes back to at least [5], see Section 2 there. The catastrophe distribution need not be binomial, see [3]. For a more recent survey of these models, see [1].

2 The limit as c approaches 0 and p is fixed

Let $(Y_n)_{n\geq 0}$ be the following discrete-time Markov chain. Assume $Y_n = k \geq 1$,

- With probability p, $Y_{n+1} = k + 1$.
- With probability 1 p, $Y_{n+1} = k P_n$, where P_n is a Poisson random variable with mean $\frac{p}{1-p}$. The random variables P_n are sampled independently of each other and of everything else.

Let $(X_n(L))_{n\geq 0}$ be the Markov chain defined in the previous section with c = 1/L and p in (0,1). With c = 1/L we get $n^*(L) = L\frac{p}{1-p}$.

Proposition 2.1. Let $X_0(L) = Y_0 = n^*(L)$ then for any fixed T and p,

$$\lim_{L \to +\infty} P(\exists n \le T : X_n(L) \ne Y_n) = 0.$$

In words, the process $(X_n(L))_{n \leq T}$ approaches $(Y_n)_{n \leq T}$ as L goes to infinity.

Proof. We introduce an auxiliary process (U_n) . If $U_n = k \ge 1$, then

- With probability p, $U_{n+1} = k + 1$.
- With probability 1 p, $U_{n+1} = k P'_n$, where P'_n is a Poisson random variable with mean kc. The random variables P'_n are sampled independently of each other and of everything else.

We proceed in two steps. First we show that

$$\lim_{L \to +\infty} P(\exists n \le T : X_n(L) \ne U_n) = 0,$$

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using that

$$P(\operatorname{bin}(x,c) \neq \operatorname{Poiss}(xc)) \leq \frac{1}{2}xc^2.$$

Then, we show that

$$\lim_{L \to +\infty} P(\exists n \le T : U_n \ne Y_n) = 0,$$

using that

$$P(\text{Poiss}(\lambda) \neq \text{Poiss}(\mu)) \le |\lambda - \mu|.$$

First step. Note that $U_n \leq n + n^* \leq T + n^*$ for $n \leq T$. Therefore, each time the process drops it drops from a state smaller than $T + n^*$. There are at most T drops up to time T. Hence, the sum of all drops up to time T is stochastically less than a sum of T i.i.d. Poisson variables with mean $(T + n^*)c$. The latter sum being itself a Poisson random variable with mean $T(T + n^*)c$.

Let

$$A = \{ \text{Poiss} (T(T + n^*)c) < T(T + n^*)c + M \},\$$

where M will be large. Let $\tau = \min\{n \ge 1 : X_n(L) \ne U_n\}$. We have,

$$P(\tau = n) \le P(X_{n-1}(L) = U_{n-1}; X_n(L) \ne U_n; A) + P(A^c).$$

On *A*, the process $(U_n)_{n < T}$ is a positive integer in the interval

$$I = [n^* - T(T + n^*)c - M, T + n^*].$$

Hence,

$$P(X_{n-1}(L) = U_{n-1}; X_n(L) \neq U_n; A) \le \sum_{x \in I} P(\mathsf{bin}(x, c) \neq \mathsf{Poiss}(xc)).$$

Therefore,

$$P(X_{n-1}(L) = U_{n-1}; X_n(L) \neq U_n; A)$$

$$\leq \sum_{x \in I} \frac{1}{2}xc^2$$

$$= \frac{1}{2}c^2\frac{1}{2}(2n^* + T - T(T+n^*)c - M)(T + T(T+n^*)c + M + 1).$$

Using that cn^* is a constant and that c goes to 0 as L goes to infinity,

$$\lim_{L \to \infty} P(X_{n-1}(L) = U_{n-1}; X_n(L) \neq U_n; A) = 0.$$

Hence,

$$\limsup_{L \to \infty} P(\tau = n) \le P(\operatorname{Poiss}(\lambda T) > \lambda T + M),$$

where $\lambda = cn^*$. Letting now M go to infinity we get $\lim_{L\to\infty} P(\tau = n) = 0$ for every $n \leq T$. This completes the first step.

Second step. Using the notation introduced in the first step,

$$P(U_{n-1} = Y_{n-1}; U_n \neq Y_n; A) \le \sum_{x \in I} P(\operatorname{Poiss}(xc) \neq \operatorname{Poiss}(\lambda)).$$

Therefore,

$$P(U_{n-1} = Y_{n-1}; U_n \neq Y_n; A) \le \sum_{x \in I} |cx - \lambda|.$$

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For x in I,

$$-T(Tc + \lambda)c - Mc \le cx - \lambda \le Tc.$$

Hence, for x in I and M large enough,

$$|cx - \lambda| \le (M + T\lambda)c.$$

Therefore,

$$P(U_{n-1} = Y_{n-1}; U_n \neq Y_n; A) \le (T(T+n^*)c + M + T + 1)(M + T\lambda)c.$$

As L goes to infinity the r.h.s. goes to 0 and we can conclude as in Step 1.

3 The limit as *p* approaches 1 and *c* is fixed

We now make $p = 1 - \frac{1}{L}$ and let $c \in (0, 1)$ be fixed. Let $y \ge 0$ and set $X_0 = \lfloor yL \rfloor$. In this context, it is convenient to describe (X_n) as follows. It alternates between the following two modes.

- Forward mode (F): The chain jumps to the right a geometric number of steps. Each jump takes a unit time. The mean of the geometric random variable is L 1.
- Backward mode (B): The chain jumps to the left in one unit time from the position it has gotten to, say z_L . The size of the jump is distributed as a Binomial random variable with mean cz_L and variance $c(1-c)z_L$.

The chain starts with Mode (F) with probability p or with Mode (B) with probability 1 - p. Then it alternates deterministically between the two modes.

We now introduce what will turn out to be the limiting process of the rescaled process $(\frac{1}{L}X_{tL})_{t>0}$.

Let $\mathcal{E}_0, \mathcal{E}_1, \ldots$ be independent mean 1 exponential random variables. Let $S_0 = 0$, and $S_n = \sum_{i=1}^n \mathcal{E}_i$, $n \ge 1$. Define recursively $Y_0 = y$ and, for $n \ge 1$,

$$Y_{S_n-} = Y_{S_{n-1}} + \mathcal{E}_n$$
$$Y_{S_n} = (1-c)Y_{S_n-}$$

For t in (S_{n-1}, S_n) , let Y_t be the linear interpolation of $Y_{S_{n-1}}$ and Y_{S_n} .

In words, starting from y, $(Y_t)_{t\geq 0}$ first moves to the right at speed 1 for a mean 1 exponential random time, after which it finds itself at $y + \mathcal{E}$, and then it jumps instantaneously to the left by $c(y + \mathcal{E})$ units. Forward and backward jumps keep alternating in a deterministic way.

Proposition 3.1. Let $p = 1 - \frac{1}{L}$ and let c fixed in (0,1) be fixed. Let $y \ge 0$ and set $X_0 = \lfloor yL \rfloor$. Then, $(\frac{1}{L}X_{tL})_{t\ge 0}$ (with the proper interpolated definition of X_t for t outside $\frac{1}{L}\mathbb{N}$) converges weakly as $L \to \infty$ to $(Y_t)_{t\ge 0}$ where $Y_0 = y$.

Proof. As L goes to infinity, p goes to 1. So the first mode taken by the chain after time 0 is (F). Therefore, let $n \in \frac{1}{L}\mathbb{N}$ then $X_{nL} = \lfloor yL \rfloor + nL$ for $nL \leq G_1$ where G_1 is a mean L-1 geometric random variable. Since G_1/L converges weakly to \mathcal{E}_1 (a mean 1 exponential random variable) we get the following weak convergence,

$$\lim_{L \to +\infty} \frac{1}{L} X_{nL} = y + n \text{ for } n < \mathcal{E}_1.$$

This shows the existence of $\lim_{L\to+\infty} \frac{1}{L} X_{tL}$ for t in $\frac{1}{L}\mathbb{N}$. For t outside this set we extend the definition of X_{tL} by linear interpolation. This gives

$$\lim_{L \to +\infty} \frac{1}{L} X_{tL} = y + t \text{ for all } t < \mathcal{E}_1^-.$$

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After the first mode (F) we switch to mode (B). Let $LT_L = G_1$, note that T_L converges weakly as L goes to infinity is \mathcal{E}_1 . Let $z_L = X_{LT_L}$ then

$$X_{LT_L+1} = z_L - \operatorname{bin}(z_L, c).$$

Since z_L goes to infinity with L, by the Law of Large Numbers and the weak convergence of G_1/L we get the following weak convergence,

$$\lim_{L \to +\infty} \frac{1}{L} X_{LT_L+1} = (1-c)(\mathcal{E}_1 + y).$$

At this point we have shown that $(\frac{1}{L}X_{tL})$ converges to (Y_t) for the first (F) and (B) modes. Using this method we can continue computing limits for the successive (F) and (B) modes. Since the latter process is non explosive, convergence in the usual Skorohod space of càdlàg trajectories readily follows.

Proposition 3.2. The process $(Y_t)_{t\geq 0}$ is ergodic. That is, it converges weakly to an invariant measure. Moreover, the distribution of the invariant measure is the same as the distribution of $\sum_{n\geq 0}(1-c)^n \mathcal{E}_n$, where $\mathcal{E}_0, \mathcal{E}_1, \ldots$ are mean 1 i.i.d. exponential random variables.

Proof. The infinitesimal generator of Y for $f \in C_c^1$, the continuously differentiable real functions on \mathbb{R}^+ with compact support, is given by

$$\Omega f(x) = f'(x) + f((1-c)x) - f(x).$$
(3.1)

To justify this, write

$$E_x[f(Y_t)] = f(x+t)e^{-t} + \int_0^t ds \, e^{-s} f((1-c)(x+s) + t - s) + o(t);$$

thus,

$$E_x[f(Y_t)] - f(x) = [f(x+t) - f(x)]e^{-t} + \int_0^t ds \, e^{-s}[f((1-c)(x+s) + t - s) - f(x)] + o(t),$$

and (3.1) follows by dividing by t and taking the limit as $t \to 0$.

In order to find an invariant distribution, let us suppose one such measure admits a continuous density ψ , which thus must satisfy

$$\int_0^\infty \Omega f(x)\psi(x)dx = \int_0^\infty f'(x)\psi(x)dx - \int_0^\infty dx\,\psi(x)\int_{(1-c)x}^x f'(y)dy$$
$$= \int_0^\infty f'(x)\left\{\psi(x) - \int_x^{ax}\psi(y)dy\right\} = 0$$

for all $f \in C_c^1$, where $a = (1 - c)^{-1}$. It follows that

$$\psi(x) = \int_x^{ax} \psi(y) dy$$

for all x > 0. By taking Laplace transforms, we readily find that $\varphi(\theta) = \int_0^\infty e^{-\theta x} \psi(x) dx$, $\theta > 0$, must satisfy

$$\varphi(\theta) = \frac{1}{1+\theta}\varphi((1-c)\theta), \ \theta > 0.$$

Iterating and taking the appropriate limit, we find that

$$\varphi(\theta) = \prod_{n=0}^{\infty} \frac{1}{1 + (1-c)^n \theta}, \, \theta > 0,$$

and the claimed form of the invariant measure is established in this case.

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To verify uniqueness, we resort to [4].

1. (Lebesgue)-irreducibility: It is enough to check the condition (midway at page 490) for *B* a finite nonempty open interval (a, b) with a > 0. The condition is clear for x < b; if $x \ge b$, then it is enough to establish that $P_x(\tau_{(0,a)} < \infty) > 0$, but this follows from the fact that after *n* jumps, our process is found at

$$W_n := (1-c)^n x + \sum_{i=1}^n (1-c)^{n-i+1} \mathcal{E}_i,$$
(3.2)

with $\mathcal{E}_1, \mathcal{E}_2, \ldots$ as above. It is enough now to have *n* large enough to make the first term less than a/2 and then $\mathcal{E}_1, \ldots, \mathcal{E}_n$ small so as to make the sum in the second term less than a/2, and event of positive probability.

- 2. Non-evanescence: If $Y_t \to \infty$ as $t \to \infty$, then $W_n \to \infty$ as $n \to \infty$, with W_n as in (3.2). But W_n is stochastically bounded by $Z := x + \sum_{n \ge 1} (1-c)^n \mathcal{E}_n$ uniformly in n. Since Z is a proper random variable, it follows that $P_x(\lim_{t\to\infty} Y_t = \infty) = 0$ for all x.
- 3. T-process property: We resort to Theorem 4.1 of [4]. We already have nonevanescence, so we want to argue that [0, L] is *petite* for every L > 0; we want to exhibit a probability measure $a = a_L$ and a nontrivial measure $\nu = \nu_L$ on \mathbb{R}^+ such that $\int_0^\infty da(t) P_x(Y_t \in \cdot) \ge \nu(\cdot)$ for all $x \in [0, L]$. We choose $a(t) = e^{-t+L} \mathbf{1}_{\{t>L\}}$ and thus

$$\begin{split} &\int_0^\infty da(t) P_x(Y_t \in \cdot) = \int_L^\infty dt \, e^{-t+L} P_x(Y_t \in \cdot) \\ &\geq \int_L^\infty dt \, e^{-t+L} P_x(\mathcal{E}_1 > L - x, \, Y_t \in \cdot) = \int_L^\infty dt \, e^{-t+L} e^{-L+x} P_L(Y_{t-L+x} \in \cdot) \\ &= \int_x^\infty dt \, e^{-s-L+2x} P_L(Y_s \in \cdot) \geq e^{-L} \int_L^\infty ds \, e^{-s} \, P_L(Y_s \in \cdot) =: \nu(\cdot) \end{split}$$

for all $x \in [0, L]$, and we have found our measure ν .

4. To conclude, we apply Theorem 3.2 in [4] to get that X is Harris recurrent. This is sufficient for uniqueness of the invariant distribution, as pointed out in [4] — see the penultimate sentence of the second paragraph page 491.

Remark 3.3. One amusing point related to the above proposition is as follows. The discrete time processes $M_n := Y_{S_n-}$ and $m_n := Y_{S_n}$, $n \ge 1$, represent local maxima and minima of the trajectory of Y. One would then perhaps be led to guess that the invariant distribution of Y should (strictly) dominate the invariant distribution of m_n , and be dominated by the invariant distribution of M_n . However, it equals the latter distribution (as one may easily check by computing the invariant distributions of m_n and M_n). The apparent contradiction is dispelled by the realization that (looking at the invariant distribution of Y as the limiting distribution of Y_t as $t \to \infty$) the interval (m_n, M_{n+1}) containing t is larger than typical (this is of course an instance of the *inspection paradox*), and in this case it is asymptotically 'twice' the size of a typical interval. Indeed, one can argue along this line to show that $\lim_{t\to\infty} Y_t = m_\infty + \mathcal{E} = M_\infty$ in distribution, where m_∞ and M_∞ are the invariant distributions of m_n and M_n respectively, with m_∞ and \mathcal{E} independent; actually, this provides an alternative proof of Proposition 3.1.

4 The limit as p approaches 1 and c approaches 0

We now make $p = 1 - \frac{1}{L^{\alpha}}$ and let $c = \frac{1}{L^{\beta}}$, $\alpha \in (0, 1)$, $\beta = 1 - \alpha$, and let $X_0 = R_L + \lfloor yL^{\alpha} \rfloor$, with $R_L = rL + o(L)$, $r \in \mathbb{R}^+$, $y \in \mathbb{R}$ fixed, and set $\bar{X}_t = X_t - R_L$.

The limiting process is defined by (using the notation of Section 3) $\bar{Y}_0 = y$ and, for $n \ge 1$, $\bar{Y}_{S_n-} = \bar{Y}_{S_{n-1}} + \mathcal{E}_n$, $\bar{Y}_{S_n} = \bar{Y}_{S_n-} - r$, and linear interpolation on (S_{n-1}, S_n) .

Proposition 4.1. The rescaled and centered process $\left(\frac{1}{L^{\alpha}}\bar{X}_{tL^{\alpha}}\right)_{t\geq 0}$ converges weakly as $L \to \infty$ to $(\bar{Y}_t)_{t\geq 0}$.

With our choice of parameters the metastable equilibrium n^* is of order L. The initial state is of order rL. Maybe surprisingly the limiting process drifts linearly with a speed 1 - r.

Remark 4.2. Figure 1 may be seen as an illustration of Proposition 4 in two cases. An appropriate choice of parameters would be L = 1000, $\alpha = 2/3$ and r = 2. In the early phase, up to time 1000, we see a negatively drifting process, roughly consistent with a drift of -1 predicted by the proposition (but the scales overlap in the simulation, which does not happen in the proposition). After time 1000, we are in the driftless, metastable regime where r = 1 of the proposition.

Proof. We follow the analysis done in Proposition 3.1. Again we start as L goes to infinity with a forward mode. Let T' be such that $T'L^{\alpha} = G'$ where G' is a geometric random variable with mean $L^{\alpha} - 1$. Then, $X_{T'L^{\alpha}} = X_0 + G'$. Since $X_0 = R_L + \lfloor yL^{\alpha} \rfloor$,

$$X_{T'L^{\alpha}} = R_L + \lfloor yL^{\alpha} \rfloor + G'.$$

Therefore, $\bar{X}_{T'L^{\alpha}} = \lfloor yL^{\alpha} \rfloor + G'$, and

$$\lim_{L \to \infty} \frac{1}{L^{\alpha}} \bar{X}_{T'L^{\alpha}} = y + \mathcal{E}_1,$$

where \mathcal{E}_1 is a mean 1 exponential random variable.

At time $T'L^{\alpha} + 1$ we switch to the backward mode with a single jump. Let $z'_L = X_{T'L^{\alpha}}$. Then,

$$X_{T'L^{\alpha}+1} = z'_L - \operatorname{bin}(z'_L, \frac{1}{L^{\beta}}).$$

Hence,

$$\frac{1}{L^{\alpha}}\bar{X}_{T'L^{\alpha}+1} = \frac{1}{L^{\alpha}}\left(\lfloor yL^{\alpha} \rfloor + G' - \operatorname{bin}(z'_{L}, \frac{1}{L^{\beta}})\right).$$

In order to prove that the limit of the l.h.s. is $y + \mathcal{E}_1 - r$ we need to show that

$$\lim_{L\to+\infty}\frac{1}{L^{\alpha}}\mathrm{bin}(z'_L,\frac{1}{L^{\beta}})=r.$$

We do this next. Note that $z'_L = rL + o(L)$ and let

$$J_L := \frac{1}{L^{\alpha}} \sum_{i=1}^{rL+o(L)} \xi_i,$$

where ξ_1, ξ_2, \ldots are iid Bernoullis with parameter $1/L^{\beta}$. Taking the Laplace transform, we find

$$E\left(e^{\theta J_L}\right) = \left(1 + \frac{1}{L^{\beta}}\left(e^{\theta/L^{\alpha}} - 1\right)\right)^{rL + o(L)} = \left(1 + \frac{\theta}{L}\frac{e^{\theta/L^{\alpha}} - 1}{\theta/L^{\alpha}}\right)^{rL + o(L)} \to e^{r\theta},$$

for all θ , and this shows that $J_L \to r$ in probability as $L \to \infty$.

Using this method we can continue computing limits for successive forward and backward modes. $\hfill \Box$

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5 A limit as p and c approach 0

We now take $p = \frac{1}{L^{\gamma}}$, $c = \frac{1}{L^{1+\gamma}}$, $\gamma > 0$. Let us make $X_0 = R_L + k$, with $R_L = rL + o(L)$, $r \in \mathbb{R}^+$, $k \in \mathbb{Z}$ fixed, and set $\bar{X}_t = X_t - R_L$. Notice that $\bar{X}_0 = k$.

Proposition 5.1. The rescaled and centered process $(\bar{X}(tL^{\gamma}))_{t\geq 0}$ converges weakly as $L \to \infty$ to a continuous time simple random walk on \mathbb{Z} with jump rate 1 + r, jumping to the right with probability $\frac{1}{1+r}$.

Proof. Let us describe the jump times and sizes of X_t starting at a location $Q_L = rL + o(L)$ as follows. Let Z_1^1, Z_2^1, \ldots be independent Bernoulli random variables with success parameter $1/L^{\gamma}$, and, independently, let Z_1^2, Z_2^2, \ldots be independent binomial random variables with Q_L trials and success parameter $1/L^{1+\gamma}$. Now set $T_L^i = \inf\{j \ge 1 : Z_j^i > 0\}$, i = 1, 2. Notice that T_L^1 and T_L^2 are independent geometric random variables. Moreover, for i = 1, 2, as $L \to \infty$ we have the following convergence in distribution, $\frac{1}{L^{\gamma}}T_L^i \to \mathcal{T}^i$ where \mathcal{T}^i is a rate λ_i exponential random variable and $\lambda_1 = 1$, $\lambda_2 = r$.

Hence, the time of the first jump is $T_L := T_L^1 \wedge T_L^2$. Note that $\frac{1}{L^{\gamma}}T_L$ converges in distribution to an exponential random variable with rate $\lambda_1 + \lambda_2 = 1 + r$.

We now turn to the the jump length. If $T_L^1 \leq T_L^2$ then the chain jumps one unit to the right. As $L \to \infty$ this has probability $\frac{1}{1+r}$. If $T_L^1 > T_L^2$ then the jump is equal $-Z_{T_L}^2$. Note that this is a strictly negative integer. We claim that

$$P(Z_{T_L}^2 \ge 2|T_L^1 > T_L^2) \to 0$$

as $L \to \infty$. It is enough to show that $P(Z^2_{T^2_L} \ge 2) \to 0$ as $L \to \infty$. The latter probability equals

$$\frac{P(Z_1^2 \ge 2)}{P(Z_1^2 \ge 1)}.$$

The denominator equals $1 - \left(1 - \frac{1}{L^{1+\gamma}}\right)^{Q_L} \ge 1 - e^{-\frac{r}{2L\gamma}} \ge \frac{r}{3L\gamma}$ for L large enough; and the numerator is bounded above by $(Q_L/L^{1+\gamma})^2 \le 2r^2/L^{2\gamma}$ for L large enough. The claim is established. This shows that in the limit when the process $(\bar{X}(tL^{\gamma}))_{t\ge 0}$ jumps to the left it jumps exactly one unit.

Remark 5.2. A note about the distinction between R_L and Q_L . The former quantity is part of the position of X at time 0, while the latter is meant for a generic position of the process after a fixed number of steps (independent of L) — the o(L) of R_L is fixed, and might have been written as $o_0(L)$, while that of Q_L varies from step to step.

References

- [1] J. R. Artalejo, A. Economou and M. J. Lopez-Herrero (2007). Evaluating growth measures in populations subject to binomial and geometric catastrophes. Math. Biosci. Eng. 4, 573–594. MR-2388812
- [2] I. Ben-Ari, A. Roitershtein, R. B. Schinazi (2019). A random walk with catastrophes. Electron. J. Probab. 24, 1–21. MR-3933207
- [3] P. J. Brockwell (1986). The extinction time of a general birth and death process with catastrophes. J. Appl. Probab. 23, 851–858. MR-0867182
- [4] S. P. Meyn and R. L. Tweedie (1993). Stability of Markovian processes II: continuous time processes and sampled chains. Adv. Appl. Prob. 25, 487–517. MR-1234294
- [5] M. F. Neuts (1994). An interesting random walk on the non-negative integers. J. Appl. Probab. 31, 48–58. MR-1260570