# FINITE SUBGROUPS OF $S U_{2}$, DYNKIN DIAGRAMS AND AFFINE COXETER ELEMENTS 

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Dedicated to the memory of my friend Ernst Straus


#### Abstract

Using, among other things, some properties of affine Coxeter elements, for which we also present normal forms, we offer an explanation of the McKay correspondence, which associates to each finite subgroup of $S U_{2}$ an affine Dynkin diagram.


J. McKay [M] has observed that for each finite (Kleinian) subgroup $G$ of $S U_{2}$ the columns of the character table of $G$, one column for each conjugacy class, form a complete set of eigenvectors for the corresponding affine Cartan matrix (of type $A_{n}, D_{n}$ or $E_{n}$ ), the one that arises in connection with the resolution of the singularity of $\mathbf{C}^{2} / G$ at the origin (see 1(9) below). As he has observed, this follows at once from: if $\rho$ is the two-dimensional representation by which $G$ is defined, $\left\{\rho_{i}\right\}$ is the set of (complex) irreducible representations of $G$, and $\sum n_{i j} \rho_{j}$ denotes the decomposition of $\rho \otimes \rho_{i}$, then $C \equiv\left[c_{i j}\right] \equiv\left[2 \delta_{i j}-n_{i j}\right]$ is the relevant Cartan matrix. Partial explanations have been given by several authors (see [G], $[\mathbf{H}],[\mathbf{K}],\left[\mathbf{S}_{1}\right.$, Appendix III]). Here we shall give our own explanation of this and some related facts, including two normal forms for affine Coxeter elements which enter into our considerations. Section 1 details mainly with McKay's correspondence, Section 2 mainly with affine Coxeter elements. As general references for Kleinian groups, Kleinian singularities and root systems, we cite [C, Chapters 7, 11], [ $\mathbf{S}_{1}$, Section 6], [B], and the survey article $\left[\mathbf{S}_{2}\right]$.

1. In this section $G$ is a finite group, $\rho$ is a faithful (complex) representation of $G$ of finite dimension $d,\left\{\rho_{i}\right\}$ is the set of all irreducible representations of $G$ with $\rho_{0}$ the trivial one, $\sum n_{i j} \rho_{j}$ denotes the decomposition of $\rho \otimes \rho_{i}$, and $C$ is the matrix $\left[d \delta_{i j}-n_{i j}\right]$.
(1) The column $\left[\chi_{j}(x)\right.$ ( $x$ in $G$ fixed, $\left.j=1,2, \ldots\right)$ ] of the character table of $G$ is an eigenvector of $C$ with $d-\chi(x)=\chi(1)-\chi(x)$ as the corresponding eigenvalue. In particular $\left[d_{1}, d_{2}, \ldots\right]\left(d_{i}=\operatorname{dim} \rho_{i}\right)$ is an eigenvector corresponding to the eigenvalue 0 .

We have $\chi(x) \chi_{i}(x)=\sum n_{i j} \chi_{j}(x)$, whence the first statement. Then $x=1$ yields the second.
(2) The following equations hold.
(a) $n_{i j}=n_{j i}$ (bar denotes dual)
(b) $d d_{i}=\sum n_{i j} d_{j}$
(c) $d d_{i}=\sum n_{j i} d_{j}$
(d) $n_{i j}=n_{j i}$ for all $i$ and $j$ if and only if $\rho$ is self-dual.
(a) This follows from $n_{i j}=\left(\chi \chi_{i}, \chi_{j}\right)=$ Average $\chi \chi_{i} \bar{\chi}_{j}$.
(b) This is the second statement of (1). (c) $d d_{i}=d d_{i}=\sum n_{i j} d_{j}=$ $\sum n_{j i} d_{j}=\sum n_{j i} d_{j}$. (d) If $n_{i j}=n_{j i}$ always then $n_{0 j}=n_{j 0}=n_{0 j}$ and $\rho$ is self-dual. If $\rho$ is self-dual then $n_{i j}=n_{i j}=n_{j i}$.
(3) Now form a real vector space $V$ with a basis vector $\alpha_{i}$ for each $\rho_{i}$ and a scalar product given by $\left(\alpha_{i}, \alpha_{j}\right)=c_{i j} \equiv d \delta_{i j}-n_{i j}$. Then the line through $\sum d_{i} \alpha_{i}$ is the radical of $($,$) from the left and from the right, and$ it is also the radical of the quadratic form ( $\alpha, \alpha$ ) and this form is positive semidefinite.

By (2b) and (2c) the given line belongs to these radicals. It will be enough to prove the converse for the quadratic form since its radical contains the others. With $\alpha=\sum x_{i} \alpha_{i}$ arbitrary in $V$ we have

$$
\begin{aligned}
2(\alpha, \alpha) & =2 \sum c_{i j} x_{i} x_{j} \\
& =2 \sum\left(d-n_{i i}\right) x_{i}^{2}-2 \sum n_{i j} x_{i} x_{j} \quad(i \neq j) \\
& =\sum\left(n_{i j}+n_{j i}\right) d_{i}^{-1} d_{j} x_{i}^{2}-2 \sum n_{i j} x_{i} x_{j} \quad(i \neq j)
\end{aligned}
$$

by (2b) and (2c). For $i<j$ the pairs $i, j$ and $j, i$ together contribute

$$
\begin{aligned}
\left(n_{i j}\right. & \left.+n_{j i}\right)\left(d_{i}^{-1} d_{j} x_{i}^{2}+d_{j}^{-1} d_{i} x_{j}^{2}-2 x_{i} x_{j}\right) \\
& =\left(n_{i j}+n_{j i}\right) d_{i}^{-1} d_{j}^{-1}\left(d_{j} x_{i}-d_{i} x_{j}\right)^{2} \geq 0
\end{aligned}
$$

Thus ( $\alpha, \alpha$ ) is positive semidefinite. Now for each $i$ there exists a sequence $i_{1}, i_{2}, \ldots, i_{n}$ with $i_{1}=0$ corresponding to the trivial representation, $i_{n}=i$, and $n\left(i_{p}, i_{p+1}\right)=0$ for all $p$; this is because $\rho_{i}$ is necessarily contained in some tensor power of the faithful representation $\rho$. It follows from this and the above inequalities that if $(\alpha, \alpha)=0$ then $x_{i}=\left(x_{0} / d_{0}\right) d_{i}$ for all $i$, so that $\alpha$ is in the line of $\sum d_{i} \alpha_{i}$.

We now specialize to the case in which $\rho$ imbeds $G$ into $S U_{2}$. We assume that $G \neq\{1\}$.
(4) (a) $c_{i j}=c_{j i}$ (i.e. $n_{i j}=n_{j i}$ ) always.
(b) $c_{i i}=2$ (i.e. $n_{i i}=0$ ) always.
(c) If $G \neq\{ \pm 1\}$ and $i \neq j$ then $c_{i j}=0$ or -1 (i.e. $n_{i j}=0$ or 1 ).

In other words (a) $C$ is symmetric, (b) $\rho_{i}$ is disjoint from $\rho \otimes \rho_{i}$, and (c) $\rho \otimes \rho_{i}$ is multiplicity-free.
(a) This is by (2d): if $\alpha, \alpha^{-1}$ are the eigenvalues of $\rho(x)$ then $\alpha+\alpha^{-1}=\alpha+\bar{\alpha} \in \mathbf{R}$. (b) If $\rho$ is reducible it has the form $\sigma \dot{+} \bar{\sigma}$ with $\operatorname{dim} \sigma=1$. Thus $G$ is cyclic and all $\rho_{i}$ have dimension 1 . So, since $G$ is nontrivial, $\sigma \otimes \rho_{i}$ and $\bar{\sigma} \otimes \rho_{i}$ are different from $\rho_{i}$ and hence disjoint from it. If $\rho$ is irreducible then $\{ \pm 1\} \subseteq$ Center $(G)$. For then 2 divides $|G|$ and -1 is the unique element of order 2 of $S U_{2}$. Now if -1 in $G$ acts as (multiplication by) 1 (resp. -1) on $\rho_{i}$, it acts as -1 (resp. 1) on $\rho \otimes \rho_{i}$, hence also on its irreducible components, all of which must thus be different from $\rho_{i}$. (c) We have

$$
\sum_{j} n_{i j}^{2}=\left(\rho \otimes \rho_{i}, \rho \otimes \rho_{i}\right) \equiv \operatorname{Av}|\chi(x)|^{2}\left|\chi_{i}(x)\right|^{2} \leq \operatorname{Av} 4\left|\chi_{i}(x)\right|^{2}=4
$$

since $|\chi(x)| \leq 2$ for all $x$ in $G$ with equality only if $x= \pm 1$. If the strict inequality holds then $\sum_{j} n_{i j}^{2}<4$ and each $n_{i j}$ is 0 or 1 . If equality holds then $\chi_{i}(x)=0$ for all $x \neq \pm 1$. If also multiplicity occurs then $\rho \otimes \rho_{i}=2 \rho_{j}$ with $\rho_{j}$ irreducible, and $\rho \otimes \rho_{j}=2 \rho_{i}$ because of the values of $\chi_{i}$ and $\chi_{j}$. Now -1 , if it is in $G$, acts trivially on $\rho_{i}$ or on $\rho_{j}$, say on $\rho_{j}$. If $H$ denotes $G$ modulo its intersection with $\{ \pm 1\}$, then $\rho_{j}$ yields an irreducible representation of $H$ with character value $d_{j}$ at 1,0 elsewhere, whence $|H|=d_{j}^{2}$. However $|H|=\Sigma^{\prime} d_{j}^{2}$, summed over all irreducible representations of $H$. It follows that $\rho_{j}$ is the unique irreducible representation of $H$, hence that $H$ is trivial. Thus $G \subseteq\{ \pm 1\}$, contradicting our assumptions.

We now introduce a diagram $\Gamma$ with one vertex corresponding to each basis vector $\alpha_{i}$ of $V$ (or to each irreducible representation $\rho_{i}$ of $G$ ) and one edge for each pair $i, j$ such that $n_{i j}=1$ (i.e. $c_{i j}=-1$ ) in (4c), which is unambiguous by (4a). By (4b) no edge of $\Gamma$ is a loop.
(5) $\Gamma$ (resp. $C$ ) is the extended Dynkin diagram (resp. matrix) of a reduced, irreducible root system with all roots of one length (which we take to be $\sqrt{2}$ ) and $\Gamma^{\prime} \equiv\left\{\alpha_{i} \mid i \neq 0\right\}$ as a simple system and $\sum_{i \neq 0} d_{i} \alpha_{i}$ as the corresponding highest root.

By the argument at the end of (3) $\Gamma$ is connected. But then so is $\Gamma^{\prime}$ : We have $\rho \otimes \rho_{0}=\rho$. Thus if $\rho$ is irreducible it yields the unique vertex of $\Gamma$ joined to $\alpha_{0}$ and $\Gamma^{\prime}$ is connected. If $\rho$ is reducible then $\Gamma$ is a loop, as may be checked; thus again $\Gamma^{\prime}$ is connected. Now by (3), if $V^{\prime}$ is the subspace of $V$ generated by the $\alpha$ 's other than $\alpha_{0}$ and $L$ is the line $\mathbf{R} \sum d_{i} \alpha_{i}$, then $V^{\prime}$ projects isometrically onto $V / L$ and there (,) is positive definite. We identify the two spaces. Further $\left(\alpha_{i}, \alpha_{i}\right)=2$ and for $i \neq j\left(\alpha_{i}, \alpha_{j}\right)=-1$ or 0 according as $\alpha_{i}$ and $\alpha_{j}$ are or are not joined in $\Gamma$.

Thus $\left\{\alpha_{i} \mid i \neq 0\right\}$ is a simple system for an irreducible root system in which $(\alpha, \alpha)=2$ for every root and $\Gamma^{\prime}$ is its Dynkin diagram. Further $-\alpha_{0}$ is a root since $\left(-\alpha_{0},-\alpha_{0}\right)=2$, and it is dominant and hence the highest root since also $\left(-\alpha_{0}, \alpha_{i}\right)=n_{0 i} \geq 0$ for $i \neq 0$. Thus $\Gamma=\Gamma^{\prime} \cup\left\{\alpha_{0}\right\}$ is the corresponding extended Dynkin diagram. On $V^{\prime}$ we have $-\alpha_{0}=$ $-d_{0} \alpha_{0}=\sum_{i \neq 0} d_{i} \alpha_{i}$, whence the last point of (5).
(6) $G / G^{\prime}$ is isomorphic to $F$, the center of the simply connected complex Lie group $L$ whose extended Cartan matrix is $C$.

First the orders of the two groups are equal: $\left|G / G^{\prime}\right|$ is the number of 1 -dimensional representations of $G$, i.e., the number of $d_{i}$ 's equal to 1 , hence is $1+$ the number of coefficients that are 1 in the highest root, which is known to be $|F|$. An isomorphism is given by $x \in G / G^{\prime} \rightarrow$ $\Pi \alpha_{i}^{*}$ (det $\rho_{i}(x)$ ). Here $\alpha_{i}^{*}$ denotes the coroot of $\alpha_{i}$, viewed as a 1-parameter subgroup int $L$. All of this is relative to a choice of a maximal torus and an ordering of its character group. The proposed isomorphism is injective since if $x$ is in the kernel then det $\rho_{i}(x)=1$ for all $i$, whence $\rho_{i}(x)=1$ for all $i$ with $d_{i}=1$, and $x \in G^{\prime}$. The image is in $F$ since if $\alpha_{j}$ is any simple root then $\alpha_{j}($ image $)=\Pi\left(\operatorname{det} \rho_{i}(x)\right)^{c i(i, j)}=1$, as we see by taking determinants in $\rho \otimes \rho_{i}=\sum n_{i j} \rho_{j}$ and using det $\rho=1$, det $\rho_{0}=1$ and $c_{i j}=2 \delta_{i j}-$ $n_{i j}$.
(7) The (unextended) Dynkin diagram for $G^{\prime}$ can be gotten from that for $G$ by deleting all vertices $\alpha_{i}$ for which $d_{i}=1$.

At present this is only an empirical observation.
Because of (6) and (7) the derived series for $G$ can be written down easily in any given case. For example, $E_{7} \supset E_{6} \supset D_{4} \supset A_{1} \supset\{1\}$, with corresponding quotients $C_{2}, C_{3}, C_{2} \times C_{2}, C_{2}$.

If $G$ is reducible on $\mathbf{C}^{2}$ and hence cyclic, then $\Gamma$ is a cycle, hence of type $A_{n}$, as is mentioned above. Conversely if $\Gamma$ contains a cycle then by standard arguments $\Gamma$ must be a cycle, $\alpha_{0}$ (corresponding to the trivial representation) has two neighbors and $\rho=\rho \otimes \rho_{0}$ is reducible. Aside from these cases, the possibilities for $G$ are classified by numbers $p_{1} \geq q_{1} \geq r_{1}$ $\geq 2$ with $p_{1}^{-1}+q_{1}^{-1}+r_{1}^{-1}>1,\left(2 p_{1}, 2 q_{1}, 2 r_{1}\right.$ are the orders of the maximal cyclic subgroups of $G$, one subgroup for each conjugacy class), and so are the possibilities for $\Gamma^{\prime}$ (ordinary Dynkin diagram) ( $p, q, r$ are the branch lengths including the branch point).
(8) If $\Gamma^{\prime}(p, q, r)$ is the diagram coming from $G\left(p_{1}, q_{1}, r_{1}\right)$ then $(p, q, r)=\left(p_{1}, q_{1}, r_{1}\right)$. Thus McKay's correspondence is bijective.

Let $x, y, z$ in $G$ be such that their eigenvalues on $\mathbf{C}^{2}$ are $\exp ( \pm \pi i / p)$, $\exp ( \pm \pi i / q), \exp ( \pm \pi i / r)$. The elements $1,-1, x^{a}\left(1 \leq a<p_{1}\right), y^{b}$ $\left(1 \leq b<q_{1}\right), z^{c}\left(1 \leq c<r_{1}\right)$ form a system of representatives of the
conjugacy class of $G$, and on them $\chi$, the character of $\rho$, has the values 2 , $-2,2 \cos \pi a / p_{1}$, etc., resp. By (1) these are the eigenvalues of $2-C$ with $C$ the corresponding extended Cartan matrix. Thus $2,2,2 \cos 2 \pi a / p_{1}$, etc., are the eigenvalues of $(2-C)^{2}-2$. If we can show that this also holds with $p_{1}, q_{1}, r_{1}$ replaced by $p, q, r$ we will be done. Consider an affine Coxeter element $c$, the product of the reflections corresponding to the affine simple roots; these are the ordinary simple roots with $1-\mu$ adjoined ( $\mu$ is the highest root). Since $\Gamma$ has no circuits the conjugacy class of $c$ is independent of the order of the factors and the affine simple roots may be so ordered that the first few are mutually orthogonal as are the rest of them. Then in partitioned form we have

$$
2-C=\left[\begin{array}{cc}
0 & N \\
N^{\prime} & 0
\end{array}\right], \quad \text { whence }(2-C)^{2}-2=\left[\begin{array}{cc}
N N^{\prime}-2 & 0 \\
0 & N^{\prime} N-2
\end{array}\right] .
$$

On the other hand if $c=c_{1} c_{2}$ in accordance with this partition of the roots, then

$$
c_{1}=\left[\begin{array}{cc}
-1 & N \\
0 & 1
\end{array}\right] \quad \text { and } \quad c_{2}=\left[\begin{array}{cc}
1 & 0 \\
N^{\prime} & -1
\end{array}\right]
$$

in matrix form. It follows that $c+c^{-1}=(2-C)^{2}-2$. Thus by the above formulas the eigenvalues of $c$ are $1,1, \exp \left(2 \pi i a / p_{1}\right)\left(1 \leq a<p_{1}\right)$, etc., and those of $c^{\prime}$, the linear part of $c$, are the same with the first 1 deleted. We now invoke a result which will be proved in the next section (after (10) there).
(*) $c^{\prime}$, the linear part of $c$, is conjugate in the Weyl group to $c^{\prime \prime}$, the product of the ordinary simple reflections with the one at the branch point excluded. From (*) it follows that $c^{\prime \prime}$ has the same eigenvalues as $c^{\prime}$ as given above. However $c^{\prime \prime}$ is the product of three Coxeter elements of types $A_{n}(n=p-1, q-1, r-1)$ corresponding to the mutually orthogonal subsystems along the branches of $\Gamma^{\prime}$, and these, together with the branch root, contribute the eigenvalues $1, \exp (2 \pi i a / p)(1 \leq a<p)$, etc. Thus $(p, q, r)=\left(p_{1}, q_{1}, r_{1}\right)$, as required.
(9) Consider the minimal resolution of the singular surface $\mathbf{C}^{2} / G$. The singular fiber is a union of projective lines, and if we form a diagram by taking one node for each line and joining two nodes, by a simple bond, just when the corresponding lines intersect, we get an ordinary Dynkin diagram, of type $A_{n}, D_{n}$ or $E_{n}$ (see $\left[\mathbf{S}_{1}\right]$ ). It remains to show that this correspondence agrees with McKay's. Let $p, q, r$ be the branch lengths of the diagram just obtained. Type $A_{n}$ may be included by taking $q=r=1$ in what follows. Let $C^{\prime}=[c(i, j)]$ be the ordinary Cartan matrix. Then
the group $G$ is isomorphic to the abstract group defined by $n$ generators and the $n$ relations $\Pi x_{i}^{c(i, j)}=1 \quad(j=1,2, \ldots)$. (Thus $G / G^{\prime}$, the Abelianized group, is just $F$, as given in (6) above.) This result is due to Mumford [ $\mathbf{M u}$ ]. The relations yields, via an application of Van Kampen's Theorem, a presentation of the fundamental group of a "sphere" around the singular point of $\mathbf{C}^{2} / G$, and that group, quite clearly, is $G$ itself. Now let $x_{1}, x_{2}, \ldots, x_{p}$ be the generators along a branch of length $p$ towards the branch point $x_{p}$. The given relations yield $x_{1}^{2} x_{2}^{-1}=1, x_{1}^{-1} x_{2}^{2} x_{3}^{-1}=1, \ldots$, whence if $x_{1} \equiv x$ then $x_{a}=x^{a}$ for $(1 \leq a \leq p)$. Similarly on the other branches $y_{q}=y^{q}, z_{r}=z^{r}$ and $x_{p}=y_{q}=z_{r}$. The relation at the branch point yields $\left(x^{p}\right)^{2}=x^{p-1} y^{q-1} z^{r-1}$. Thus $x y z=x^{p}=y^{q}=z^{r}$. As is well known [C, 11.7, 7.4] this is a presentation of the Kleinian group of type ( $p, q, r$ ). Thus $G$ and the graph corresponding to it have the same type, as required.
(10) McKay's correspondence can be extended to yield Dynkin diagrams with multiple bonds in several different ways. One way, used in [H], is to start with representations over fields that are not algebraically closed. This yields most Dynkin diagrams, but not all of them. Another way, suggested in [S, App.III], which does yield all diagrams, is to start with certain pairs $G \triangleleft H$ of finite subgroups of $S U_{2}$, or, equivalently, with a single subgroup $G$ and an automorphism $\sigma$ of $G$ which stabilizes the defining representation of $G$, and then to associate a node to each $\sigma$-orbit of irreducible representations of $G$, or, dually, to each representation of $\langle G, \sigma\rangle$ induced by an irreducible representation of $G$. One can then carry out large parts of the above development in this new context (with weighted nodes, multiple bonds, etc.) or else notice that the coalescence of irreducible representations into $\sigma$-orbits corresponds exactly to the foldings of Dynkin diagrams according to their symmetries.
2. Affine Coexter elements. Our purpose is to develop the principal properties of these elements, including two normal forms and a proof of the property ( $*$ ) used in the proof of (8) above. II will be a simple system for an irreducible root system. We write $\lambda_{\alpha}$ for the fundamental weight corresponding to $\alpha$ and $\alpha^{*}$ for the coroot $2 \alpha /(\alpha, \alpha)$. We can decompose $\Pi$ into disjoint parts $\Pi_{1}$ and $\Pi_{2}$ so that each is an orthogonal set of roots. We exclude type $A_{n}$ mostly. Then $\mu$, the highest root, is orthogonal to all roots of $\Pi$ but one, so that the notation can be chosen so that $\mu$ is orthogonal to $\Pi_{2}$. We write $w_{1}$ (resp. $w_{2}$ ) for the product of the simple reflections in $\Pi_{1}$ (resp. $\Pi_{2}$ ), then $w_{i}=w_{1}$ (resp. $w_{2}$ ) when $i$ is odd (resp. even) (and similarly for $\Pi_{i}$ ), and finally $w^{1}=w_{1}, w^{2}=w_{2} w_{1}$,
$w^{3}=w_{3} w_{2} w_{1}, \ldots$. We observe that $w_{1}, w_{2}$ and each odd $w^{i}$ is an involution. Since type $A_{n}$ has been excluded, the Coxeter number, the order of $w_{1} w_{2}$, is even: $h=2 g$. For as can be easily proved or read off from the classification, $h$ is odd only for type $A_{2 n}$.
(1) We have $1<w^{1}<w^{2} \cdots<w^{2 g}$ and $w^{2 g}=w_{0}$, the element of the Weyl group that makes all positive roots negative.

Here $w<w^{\prime}$ means that $w$ is the terminal segment of a minimal expression for $w^{\prime}$ as a product of simple reflections, i.e., that the length of $w^{\prime}$ is the sum of those of $w$ and $w^{\prime} w^{-1}$. The more general Bruhat order could also be used in all that follows. The fact that $w^{2 g}=w_{0}$ is proved in [St]. In the expression $w_{0}=w_{2 g} \cdots w_{2} w_{1}$ with $w_{2}$ written as the product of the reflections for the roots in $\Pi_{i}$, the number of roots listed is $g|\Pi|=$ $(h / 2)|\Pi|$, which, as is known [St], is equal to the number of positive roots. It follows that the expression is minimal, as in each terminal segment, whence (1).
(2) Let $\lambda$ be a dominant (integral) weight, and $w<w^{\prime}$ in the Weyl group. Then $w \lambda \geq w^{\prime} \lambda$. Hence $\lambda \geq w^{\prime} \lambda \geq w^{2} \lambda \geq \cdots \geq w^{2 g} \lambda$.

It is enough, by induction, to prove this when $w^{\prime}=w_{\beta} w$, with $\beta>0$ and $w^{-1} \beta>0$. Here and elsewhere $w_{\beta}$ is the reflection relative to $\beta$. Now $w^{\prime} \lambda=w_{\beta} w \lambda=w \lambda-\left(w \lambda, \beta^{*}\right) \beta$, and $\left(w \lambda, \beta^{*}\right)=\left(\lambda, w^{-1} \beta^{*}\right) \geq 0$ since $\lambda$ is dominant and $w^{-1} \beta$ is positive, whence (2).
(3) Assume that $\lambda$ is dominant, $w_{0} \lambda=-\lambda$, and $\operatorname{Supp} \lambda \subseteq \Pi_{1}$. Then $w^{i} \lambda=-w^{2 g-1-i} \lambda$ for $0 \leq i<g$.

Here the third condition on $\lambda$ is that in its expression in terms of the fundamental weights $\lambda_{\alpha}$ only those with $\alpha \in \Pi_{1}$ are needed. We have

$$
\begin{aligned}
w^{i} \lambda & =w_{l} \cdots w_{2} w_{1} \lambda \\
& =w_{i+1} \cdots w_{2 g} \cdot w_{2 g} \cdots w_{i+1} w_{i} \cdots w_{1} \lambda=-w_{l+1} \cdots w_{2 g} \lambda \quad(\operatorname{by}(1)) \\
& =-w^{2 g-i-1} \lambda \quad \text { since } w_{2 g}=w_{2} \text { fixes } \lambda
\end{aligned}
$$

(4) In (3) $w^{g-1} \lambda$ is a nonnegative combination of roots in $\Pi_{g}$.

We have $w^{g-1} \lambda=\left(w^{g-1} \lambda-w^{g} \lambda\right) / 2$ by (3). This is $\geq 0$ by ( 2 ), and, since it equals $\left(1-w_{g}\right) w^{g-1} \lambda / 2$, it involves only the simple roots in $\Pi_{g}$.
(5) $w^{g-1} \mu$ is a simple root, an element of $\Pi_{g}$. It is the unique long simple root $b$ at which there is a branch point or a multiple bond. (Recall that $\mu$ is the highest root).

First, by (4), which is applicable since $\mu$ as a dominant weight has its support in $\Pi_{1}, w^{g-1} \mu=b$ is a nonnegative combination of roots in $\Pi_{g}$. Since $b$ is a root and the elements of $\Pi_{g}$ are mutually orthogonal, it easily follows that $b$ is an element of $\Pi_{g}$. And since $\mu$ is a long root, so is $b$.

Since type $A_{n}$ is being excluded, $\mu$ is connected to a unique simple root $\alpha$. Assume first that $\alpha$ is shorter than $\mu$. To prove (5) in this case we show that there is only one long simple root. In the extended Dynkin diagram $1-\mu$ is joined only to $\alpha$, by a multiple bond, and in the ordinary diagram $\alpha$ is connected to a nearest long root by a chain $C$, ending with a multiple bond. It is enough to show that $C$ is the full Dynkin diagram. If it were not, then $C \cup\{1-\mu\}$ would be a proper connected subdiagram of the extended Dynkin diagram, hence a Dynkin diagram in its own right, but one with two multiple bonds, namely those at its two ends, which is impossible. Now assume that $\alpha$ has the same lengths as $\mu$. We have

$$
\left(w^{g} \mu, w^{g+1} \mu\right)=\left(w^{2 g+1} \mu, \mu\right)=-\left(w_{1} \mu, \mu\right) \quad(\text { by }(3))=-|\mu|^{2} / 2
$$

since $w_{1} \mu=w_{\alpha} \mu$ and $\alpha$ and $\mu$ have equal length and form an angle of $60^{\circ}$ in the present case. Thus $\left|w^{g} \mu-w^{g+1} \mu\right|^{2}=3|\mu|^{2}$. However $w^{g} \mu-w^{g+1} \mu$ $=\left(1-w_{g+1}\right) w^{g} \mu=\left(1-w_{g+1}\right)(-b)=\Sigma\left(b, \gamma^{*}\right) \gamma$, summed over the elements of $\Pi_{g+1}$ that are not orthogonal to $b$, i.e., over the neighbors of $b$. Since $\mu$ and $b$ have the same length, the last two equations imply that $\Sigma\left(\gamma, b^{*}\right)\left(b, \gamma^{*}\right)=3$. Thus 3 bonds come together at $b$, which is therefore a branch point or a point with a multiple bond.

In the development in this section so far we have borrowed ideas from Kostant [Ko], who in turn has borrowed ideas from an earlier version of this paper. In that version, the transition from $\mu$ to $b$ was effected differently, namely by alternate applications of (1) the reflection corresponding to $b$, (2) the product of the other simple reflections ordered so as to move away from $b$. That method brings up other points of interest, but we shall not pursue them here.
(6) Assume as in (3) and (4) except that $\operatorname{Supp} \lambda \subseteq \Pi_{2}$. Then $w^{i} \lambda=$ $-w^{2 g+1-i} \lambda$ and $w^{g} \lambda$ is a nonnegative combination of the roots in $\Pi_{g+1}$.

This easily follows from (3) and (4) with the roles of $\Pi_{1}$ and $\Pi_{2}$ interchanged.
(7) Let $\lambda$ be a weight such that $w_{0} \lambda=-\lambda$, and $\lambda_{1}$ (resp. $\lambda_{2}$ ) the parts supported by the $\lambda_{\alpha}$ with $\alpha \in \Pi_{1}$ (resp. $\Pi_{2}$ ). Then $w^{g} \lambda_{1}$ (resp. $w^{g} \lambda_{2}$ ) is the part of $w^{g} \lambda$ supported by the roots of $\Pi_{g}\left(\right.$ resp. $\left.\Pi_{g+1}\right)$.

This follows from (4) and (6).
(8) Write $\mu=\Sigma_{1} n_{\alpha} \alpha+\Sigma_{2} n_{\beta} \beta$, the sums over $\Pi_{1}$ and $\Pi_{2}$. Then, with $b$ as in (5),
(a) $w^{g} \sum_{1} n_{\alpha} \alpha=-2 \lambda_{b}$.
(b) $w^{g} \sum_{2} n_{\beta} \beta=2 \lambda_{b}-b$.

By the orthogonality relations between the simple coroots and the fundamental weights we have $b=2 \lambda_{b}+\Sigma\left(b, \gamma^{*}\right) \lambda_{\gamma}$. Here $b$ is in $\Pi_{g}$ and
the sum, equal to $b-2 \lambda_{b}$, is over the neighbors of $b$, all in $\Pi_{g+1}$. Using (4) we get $-\mu=\left(w^{g}\right)^{-1}\left(2 \lambda_{b}\right)+\left(w^{g}\right)^{-1} \Sigma\left(b, \gamma^{*}\right) \lambda_{\gamma}$. By (7) with $w_{g}, w_{g+1}$, $\left(w^{g}\right)^{-1}$ in the roles of $w_{1}, w_{2}, w_{g}$, the first term on the right has support in $\Pi_{1}$, the second in $\Pi_{2}$, whence (a) and (b).
(9) With the notation as in (8)

$$
\sum_{1} n_{\alpha} \alpha=2 \sum_{1} n_{\alpha} \lambda_{\alpha}-2 \sum_{2} n_{\beta} \lambda_{\beta}
$$

If we dot with $\alpha^{*}$ we get $2 n_{\alpha}$ on both sides. If we dot with $\beta^{*}$ instead we get $-2 n_{\beta}$ since the left side equals $\mu-\sum_{2} n_{\beta} \beta$ and $\mu$ is orthogonal to all elements of $\Pi_{2}$.
(10) With the notation as in (8)
(a) $w^{g} \Sigma_{1} n_{\alpha} \lambda_{\alpha}=-\lambda_{b, g}$
(b) $w^{g} \sum_{2} n_{\beta} \lambda_{\beta}=\lambda_{b, g+1}$
(c) $w^{g}\left(\sum_{1} n_{\alpha} \lambda_{\alpha}-\sum_{2} n_{\beta} \lambda_{\beta}\right)=-\lambda_{b}$

Here $\lambda_{b, g}$, for example, denotes the part of $\lambda_{b}$ supported by $\Pi_{g}$, that one of $\Pi_{1}$ and $\Pi_{2}$ that contains $b$. By (8) and (9) we have (c). Then (a) and (b) follow from (7).

We turn now to our discussion of affine Coxeter elements. One of these is $c=w w_{2} w_{1}$ with $w_{1}$ and $w_{2}$ as above and $w=w_{1-\mu}$ the reflection corresponding to $1-\mu$. First we give the proof of (*) of $1(8)$, thus completing the proof of that result. We have to show that $c^{\prime}$, the linear part of $c$, is conjugate to the product of the reflections other than that at the branch point $b$. By (5), which is all that is needed, we have $w^{g} w_{\mu}\left(w^{g}\right)^{-1}$ $=w_{b}$. This can be written as $w^{g}\left(w_{\mu} w_{2} w_{1}\right)\left(w^{g}\right)^{-1}=w_{b} w_{g} w_{g+1}$. The left side is conjugate to $c^{\prime}$ and the right side equals the stated product since the $w_{b}$ in front cancels the $w_{b}$ that occurs as a factor of $w_{g}$.

Next we present a normal form for the affine Coxeter element $c=w w_{2} w_{1}$. Let $F$ denote the standard fundamental domain for the affine Weyl group, defined as the region (a simplex) where $\alpha \geq 0$ for all $\alpha \in \Pi$, and $1-\mu \geq 0$. Then $w_{1}$ is the reflection across the facet $F_{2}$ of $F$ where all $\alpha \in \Pi_{1}$ are 0 (since the elements of $\Pi_{1}$ are mutually orthogonal), and $w w_{2}$ is the reflection across the opposite facet $F_{1}$ where $1-\mu$ and all $\beta \in \Pi_{2}$ are 0 . We seek points $\gamma_{1}$ and $\gamma_{2}$ in $F_{1}$ and $F_{2}$ such that the line $L$ joining them is orthogonal to $F_{1}$ and to $F_{2}$. Then $c$ will be a screw displacement along $L$ (in 3 dimensions the motion of a screw whose axis is $L$ ): translation by the vector $2\left(\gamma_{1}-\gamma_{2}\right)$ in the direction of $L$ composed with a rotation around $L$ (i.e. an isometry fixing the points of $L$ ), the two factors necessarily commuting and being determined by the stated conditions. Since $c^{\prime}$, the linear part of $c$, has 1 as an eigenvalue of multiplicity 1 (with corresponding eigenvector in the direction of $L$ ), $L$, in the present case, may also be described as the set of points moved the least distance by $c$.
(11) Write $\mu=\sum_{1} n_{\alpha} \alpha+\sum_{2} n_{\beta} \beta$ as in (8) set $\delta=\sum_{1} n_{\alpha} \alpha$, so that $(\delta, \delta)=\sum_{1} n_{\alpha}^{2}(\alpha, \alpha)$. Then the solution to our problem is $\gamma_{1}=$ $2 \sum_{1} n_{\alpha} \lambda_{\alpha} /(\delta, \delta)$ and $\gamma_{2}=2 \sum_{2} n_{\beta} \lambda_{\beta} /(\delta, \delta)$. Thus $c$ is a rotation around the line joining these points composed with the translation by the vector $2\left(\gamma_{1}-\gamma_{2}\right)=2 \delta /(\delta, \delta)$, of length $2 /|\delta|$, along the line. Further $\gamma_{2}$ is the point on $L$ closest to the origin.

First $\delta$ is orthogonal to $F_{2}$ clearly and also to $F_{1}$ since $\delta=\mu-\sum_{2} n_{\beta} \beta$. Write the equation of (9) as $\delta=\delta_{1}-\delta_{2}$. Then by the definitions $\delta_{2} \in F_{2}$ and $\delta_{1} \in F_{1}$ except for the condition $\left(\mu, \delta_{1}\right)=1$. Now $\left(\mu, \delta_{1}\right)=$ $2 \sum_{1} n_{\alpha}^{2}\left(\alpha, \lambda_{\alpha}\right)=\sum_{1} n_{\alpha}^{2}(\alpha, \alpha)=(\delta, \delta)$. It follows that $\gamma_{1}=\delta_{1} /(\delta, \delta)$ is in $F_{1}$ and $\gamma_{2}$ is in $F_{2}$, and that $2\left(\gamma_{1}-\gamma_{2}\right)=2\left(\delta_{1}-\delta_{2}\right) /(\delta, \delta)=2 \delta /(\delta, \delta)$ is orthogonal to $F_{1}$ and to $F_{2}$. Finally $\gamma_{2}$ is the point of $L$ closest to the origin since it is orthogonal to the vector $\delta$ along $L$.

In the normal form for $c$ just given, the axis $L$ and the translational part can be calculated quite explicitly in any given case, but the same can not be said of the rotational part. Here is another normal form which remedies this deficiency.
(12) Let $b$ be the root in (5) above, and write $\lambda_{b}=\lambda_{b, g}+\lambda_{b, g+1}$ with $\lambda_{b, g}$ the part supported by $\Pi_{g}$ (which includes $b$ ) and $\lambda_{b, g+1}$ the part supported by $\Pi_{g+1}$. Let $L$ be the line through $\varepsilon=\lambda_{b, g+1} / 2\left(\lambda_{b}, \lambda_{b}\right)$ in the direction of $\lambda_{b}$. Then $\varepsilon$ is the point of $L$ closest to the origin. Form the rotation around $L$ whose linear part is the product of the simple reflections other than that for $b$ and compose it with the translation by $-\lambda_{b} /\left(\lambda_{b}, \lambda_{b}\right)$ in the direction of $L$. Then the result is an affine Coxeter element $c$ expressed in standard form as a screw displacement.

With $\gamma_{1}, \gamma_{2}, \delta$ as in (11) we have from (8) and (10) that $w^{g} \delta=-2 \lambda_{b}$, $w^{g} \gamma_{1}=-2 \lambda_{b, g} /(\delta, \delta)$ and $w^{g} \gamma_{2}=2 \lambda_{b, g+1} /(\delta, \delta)$. Thus (12) follows from (11).

We observe that since $\Pi-\{b\}$ is a union of systems of type $A_{n}$, the eigenvalues and eigenvectors of the linear part of $c$ above, as well as its order, can be easily determined. So can $\lambda_{b}$, hence the other items of (12) also, especially when all roots have the same length, so that $b$ is a branch point:
(13) If $\lambda_{b}=\sum m_{\alpha} \alpha(\alpha \in \Pi)$, then $m_{b}=\left(p^{-1}+q^{-1}+r^{-1}-1\right)^{-1}$ in terms of the branch lengths, and along the branch of length $p$, for example, starting at the end point the $m_{\alpha}$ 's are $p^{-1} m_{b}, 2 p^{-1} m_{b}, \ldots$

For, as is easily seen, the scalar product of the proposed vector with $b^{*}$ is 1 , with all other simple coroots is 0 .

We conclude our paper with some further remarks about the McKay correspondence that arise from the ideas of this section. $G$ will be a

Kleinian group, as in $\S 1, \Pi$ the corresponding simple root system, and the other notations as above.
(14) $2 \Sigma_{2} n_{\alpha}^{2}=g$, the order of $|G|$.

Regarding the $n_{\alpha}$ 's as the degrees of the irreducible representations of $G$ we have $\Sigma_{1} n_{\alpha}^{2}+\Sigma_{2} n_{\beta}^{2}+1$ (trivial representation) $=g$. The same for the group $G /\{ \pm 1\}$ yields $\Sigma_{2}+1=g / 2$ (see the proof of $1(4 \mathrm{~b})$ ), whence (14).
(15) $2 \Sigma_{1} n_{\alpha}^{2}=4\left(p^{-1}+q^{-1}+r^{-1}-1\right)^{-1}$ in terms of the branch lengths.

For, $(\delta, \delta)=\sum_{1} n_{\alpha}^{2}(\alpha, \alpha)$ in (11) and $\left(2 \lambda_{b}, 2 \lambda_{b}\right)=4\left(\lambda_{b}, \sum m_{\alpha} \alpha\right)=$ $2 m_{b}(b, b)$. These are equal by (8a), and then (15) follows from (13) and the equality of all root lengths.

These equations show that $m_{b}$ in (13) is just $g / 4$. They also lead to another, nontrigonometric, proof of $\S 1(8)$, with which we close our paper. We have $g=4\left(p_{1}^{-1}+q_{1}^{-1}+r_{1}^{-1}-1\right)$ since the decomposition of $G$ into conjugacy classes yields $g=1+1+\left(p_{1}-1\right) g / 2 p_{1}+\left(q_{1}-1\right) g / 2 q_{1}$ $+\left(r_{1}-1\right) g / 2 r_{1}$. Since $r=r_{1}=2$, this, (14) and (15) yield $p^{-1}+q^{-1}=$ $p_{1}^{-1}+q_{1}^{-1}$. But also $p+q=p_{1}+q_{1}$ since $p+q+r-1$ is the number of irreducible representations of $G$ and $p_{1}+q_{1}+r_{1}-1$ is the number of conjugacy classes. Dividing one equation by the other we get $p q=p_{1} q_{1}$. Thus $(p, q)=\left(p_{1}, q_{1}\right)$, as required.

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