

A REMARK ON THE
LIMITING ABSORPTION PRINCIPLE
FOR THE REDUCED WAVE EQUATION
WITH TWO UNBOUNDED MEDIA

YOSHIMI SAITŌ

Eidus recently proved the limiting absorption principle for the reduced wave equation with two unbounded media, and he used it to show the limiting amplitude principle for the wave equation. In this paper we shall show that his limiting absorption principle can be improved so that it holds on the same weighted Sobolev spaces as were used in the case of the Schrödinger equation.

1. Introduction. Let us consider the reduced wave operator

$$(1.1) \quad Hu = -\mu(x)^{-1}\Delta$$

in \mathbf{R}^N , where Δ is the Laplacian in \mathbf{R}^N and $\mu(x)$ is a positive function in \mathbf{R}^N . The operator H can be regarded as a selfadjoint operator in the Hilbert space \mathcal{H} of all measurable functions $f(x)$ on \mathbf{R}^N such that $f(x)\sqrt{\mu(x)}$ is square integrable over \mathbf{R}^N (see §2). The reduced wave operator H is obtained from the wave equation

$$(1.2) \quad \mu(x)\frac{\partial^2 w}{\partial t^2} - \Delta w = 0$$

by separation of the time variable t .

Through this work it is assumed that the function $\mu(x)$ on \mathbf{R}^N with $N \geq 2$ has the form

$$(1.3) \quad \mu(x) = \mu_l \quad (x \in \Omega_l, l = 1, 2),$$

with positive constants μ_l ($\mu_1 \neq \mu_2$) and disjoint open sets Ω_l , $l = 1, 2$, given by

$$(1.4) \quad \begin{aligned} \Omega_1 &= \{x \in \mathbf{R}^N / x_N > \varphi(x')\} \quad \text{and} \\ \Omega_2 &= \{x \in \mathbf{R}^N / x_N < \varphi(x')\}, \end{aligned}$$

where $x = (x', x_N)$, $x' = (x_1, x_2, \dots, x_{N-1})$ and $\varphi \in C^1(\mathbf{R}^{N-1} \setminus \{0\})$. The separating surface S is defined by

$$(1.5) \quad S = \{x \in \mathbf{R}^N / x_N = \varphi(x')\}$$

and $n_l(x) = (n_{l1}(x), n_{l2}(x), \dots, n_{lN}(x))$, $l = 1, 2$, denote the outward unit normal of the boundary $\partial\Omega_l$ of Ω_l at $x \in S \setminus \{0\}$.

Eidus [4] proved that the limiting absorption principle and the limiting amplitude principle hold for the reduced wave operator H and the wave equation (1.2), respectively, under the following conditions on the surface S :

Assumption 1.1. The separating surface S has a “cone-like” shape in the following sense: Let $n_l(x)$ be as above. Then,

(i) the N th component $n_{lN}(x)$ of $n_l(x)$ satisfies

$$(1.6) \quad |n_{lN}(x)| > c_1 \quad (l = 1, 2, x \in S \setminus \{0\})$$

with a positive constant c_1

(ii) and we have

$$(1.7) \quad |x \cdot n_l(x)| \leq c_2 \quad (l = 1, 2, x \in S \setminus \{0\})$$

with a positive constant c_2 , where $x \cdot n_l(x)$ means the inner product of the vectors x and $n_l(x)$ in \mathbf{R}^N .

Let us define the resolvent $R(z)$ of the operator H by

$$(1.8) \quad R(z) = (H - z)^{-1}$$

for $z \in \mathbf{C} \setminus \mathbf{R}$. The resolvent $R(z)$ is a bounded linear operator on the Hilbert space \mathcal{H} which is now equivalent to the usual L_2 space $L_2(\mathbf{R}^N)$. Let us introduce the weighted L_2 space $L_{2,\beta}(\mathbf{R}^N)$ by

$$(1.9) \quad L_{2,\beta}(\mathbf{R}^N) = \left\{ f(x) / \int_{\mathbf{R}^N} (1+|x|)^{2\beta} |f(x)|^2 dx < \infty \right\}$$

with its norm

$$(1.10) \quad \|f\|_\beta = \left[\int_{\mathbf{R}^N} (1+|x|)^{2\beta} |f(x)|^2 dx \right]^{1/2}.$$

where β is a real number. (In Eidus [4], $L_{2,\beta}(\mathbf{R}^N)$ and $\|\cdot\|_\beta$ are denoted as $L_{2\beta}^2$ and $\|\cdot\|_{2\beta}$, respectively.) Then Eidus' result on the limiting absorption principle is stated as follows:

THEOREM 1.2 (Eidus [4], Theorem 3.2). *Let Assumption 1.1 be satisfied. Then there exist the limits*

$$(1.11) \quad \lim_{\eta \rightarrow \pm 0} R(\lambda + i\eta) = R^\pm(\lambda) \quad \text{in } B(L_{2,1}(\mathbf{R}^N), L_{2,-1}(\mathbf{R}^N))$$

for each $\lambda > 0$, where $B(X, Y)$ is the Banach space of all bounded, linear operators from X into Y . Furthermore, the limit $R^\pm(\lambda)$ are Hölder continuous in λ in the topology of $B(L_{2,1}(\mathbf{R}^N), L_{2,-1}(\mathbf{R}^N))$.

The limiting absorption principle for partial differential operators has been studied for about twenty years. Especially many works have been done for Schrödinger operators

$$(1.12) \quad T = -\Delta + V(x)$$

in \mathbf{R}^N to find various sufficient conditions on the potential $V(x)$ that the limiting absorption principle holds for T . Among them, we refer to Jäger [6], Saitō [10], Agmon [1] for the short-range potential, Ikebe & Saitō [5], Lavine [8] for the long-range potential. Mochizuki & Uchiyama [9], Devinatz & Rejto [2] for the oscillatory long-range potential. It is shown by these works that under certain conditions on $V(x)$ there exist the limits

$$(1.13) \quad \lim_{\eta \rightarrow \pm 0} R_T(\lambda + i\eta) = R_T^\pm(\lambda) \quad \text{in } B(L_{2,\delta}(\mathbf{R}^N), L_{2,-\delta}(\mathbf{R}^N)),$$

where $R_T(z) = (T - z)^{-1}$, λ belongs to the continuous spectrum of T , and δ is a constant such that

$$(1.14) \quad \delta > 1/2.$$

The condition (1.14) is in a sense best possible, because, in general, $u = R_T^\pm(\lambda)f$ does not belong to $L_{2,-\beta}(\mathbf{R}^N)$ for $\beta \leq 1/2$.

In this work we shall prove some new estimates (Propositions 3.3, 3.8, 3.10 etc.) for $u = R(z)f$ which, combined with the methods and results of Eidus [4], enable us to show, under Assumption 1.1, that for a set K in $\mathbf{C} \setminus \mathbf{R}$ of the form

$$(1.15) \quad K = \{z = \lambda + i\eta \in \mathbf{C} / \lambda_0 \leq \lambda \leq \lambda_1, 0 < |\eta| \leq \eta_0\}$$

with positive constants $\lambda_0 < \lambda_1$ and η_0 there exists a positive constant $C = C(K)$ depending only on K such that

$$(1.16) \quad \|R(z)f\|_{2,-\delta} \leq C\|f\|_\delta \quad (z \in K, f \in L_{2,\delta}(\mathbf{R}^N)),$$

where δ is a constant satisfying (1.14) and $\|\cdot\|_{2,-\delta}$ denotes the norm of the weighted Sobolev space $H_{-\delta}^2(\mathbf{R}^N)$ of all functions u such that all the derivatives up to the second order belong to the weighted L_2 space $L_{2,-\delta}(\mathbf{R}^N)$ (cf. Eidus [4], Corollary 2.2). It is easy to see from (1.16) and Theorem 1.1 that the limits

$$(1.17) \quad \text{s-lim}_{\eta \rightarrow \pm 0} R(\lambda + i\eta)f = R^\pm(\lambda)f \quad \text{in } H_{-\delta}^2(\mathbf{R}^N)$$

exist for all $f \in L_{2,\delta}(\mathbf{R}^N)$, where s-lim means the strong limit.

Let us explain our main idea. When we studied the limiting absorption principle for various Schrödinger operators, the classical Sommerfeld radiation condition

$$(1.18) \quad \frac{\partial u}{\partial r} - iku = \text{small at infinity}$$

or its modifications played a very important role (e.g., Jäger [6], Saitō [10], Ikebe & Saitō [5], Mochizuki & Uchiyama [9], Saitō [11], [12]). First, as is well-known, the radiation condition guarantees the uniqueness of the solution. At the same time it has been known that we can show the limiting absorption principle through a priori estimates of the radiation condition. It will be seen in our case that it is useful for getting the estimate (1.16) to introduce a “modified radiation condition-like” term which contains a surface integral over the separating surface S given by (1.5). The limiting absorption principle for the operator H will be obtained through estimating the above “modified radiation condition-like” term. However, it seems that we need further investigation to see whether our radiation condition fully guarantees the uniqueness of the solution.

In §2 some basic a priori estimates on $u = R(z)f$, which were obtained by Eidus [4], will be given. In §3 we shall prove some more a priori estimates on the “modified radiation condition-like” term and $u = R(z)f$. The estimate (1.16) and the limiting absorption principle will be shown in §4.

2. A priori estimates for $u = R(z)f$. In this section we shall give some a priori estimates for $u = R(z)f$ with $z \in \mathbf{C} \setminus \mathbf{R}$ and $f \in L_2(\mathbf{R}^N)$, $N \geq 2$. All these estimates except for Lemma 2.4 are proved in Eidus [4].

Let us now call some usual notation for some function spaces which will be used in the sequel. Let m and β be a nonnegative integer and a real number, respectively, and let G be an open set in \mathbf{R}^N . Then the weighted Sobolev space $H_\beta^m(G)$ is defined by

$$(2.1) \quad H_\beta^m(G) = \{v \in \mathcal{D}'(G) / (1 + |x|)^\beta \partial^\alpha v \in L_2(G), |\alpha| \leq m\},$$

where $\mathcal{D}'(G)$ is all distributions on G , $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_N)$ is a multi-index with $|\alpha| = \alpha_1 + \alpha_2 + \dots + \alpha_N$ and

$$(2.2) \quad \partial^\alpha v = \frac{\partial^{|\alpha|} v}{\partial_1^{\alpha_1} \partial_2^{\alpha_2} \dots \partial_N^{\alpha_N}}$$

with

$$(2.3) \quad \partial_j^{\alpha_j} = \left(\frac{\partial}{\partial x_j} \right)^{\alpha_j} \quad (j = 1, 2, \dots, N).$$

The space $H_\beta^m(G)$ is a Hilbert space with its inner product

$$(2.4) \quad (v, w)_{m,\beta,G} = \sum_{|\alpha| \leq m} \int_G (1 + |x|)^{2\beta} (\partial^\alpha v) \overline{(\partial^\alpha w)} dx$$

and norm

$$(2.5) \quad \|v\|_{m,\beta,G} = [(v, v)_{m,\beta,G}]^{1/2}.$$

We set

$$(2.6) \quad H_\beta^0(G) = L_{2,\beta}(G),$$

and the subscript 0 in $(\cdot, \cdot)_{0,\beta,G}$ or $\|\cdot\|_{0,\beta,G}$ will be omitted as in $(\cdot, \cdot)_{\beta,G}$ or $\|\cdot\|_{\beta,G}$. When $G = \mathbf{R}^N$, the subscript G will be also omitted as in $(\cdot, \cdot)_\beta$. Let \mathcal{H} be the Hilbert space defined by

$$(2.7) \quad \mathcal{H} = \left\{ f(x) \text{ on } \mathbf{R}^N / \int_{\mathbf{R}^N} |f(x)|^2 \mu(x) dx < \infty \right\}$$

with its inner product

$$(2.8) \quad (f, g)_\mathcal{H} = \int_{\mathbf{R}^N} f(x) \overline{g(x)} \mu(x) dx$$

and norm

$$(2.9) \quad \|f\|_\mathcal{H} = [(f, f)_\mathcal{H}]^{1/2}.$$

Since $\mu(x)$ is a positive step function on \mathbf{R}^N , the Hilbert space \mathcal{H} and $L_2(\mathbf{R}^N)$ are the same as sets.

Let us define the operator H in \mathcal{H} by

$$(2.10) \quad Hu = -\mu(x)^{-1} \Delta u,$$

$$(2.11) \quad D(H) = H^2(\mathbf{R}^N),$$

where $D(H)$ means the domain of H and the Sobolev space $H^2(\mathbf{R}^N)$ is regarded as a subset of \mathcal{H} . It is easy to see that H is a selfadjoint operator. We denote the resolvent $(H - z)^{-1}$ by $R(z)$.

In order to evaluate some integrals over the separating surface S let us prepare the following lemma due to Eidus [4].

LEMMA 2.1 (*Eidus [4], (2.9)*). *Suppose that $\mu(x)$ is a bounded measurable function on \mathbf{R}^N such that $\inf \mu(x) > 0$. Let H and $R(z)$ be as above. Let*

$$(2.12) \quad u = u(\cdot, z, f) = R(z)f$$

with $f \in L_2(\mathbf{R}^N)$ and $z = \lambda + i\eta \in \Pi$, Π being a bounded set in $\mathbf{C} \setminus \mathbf{R}$. Then there exists a constant $C = C(\Pi, \mu)$ depending only on Π and $\mu(x)$ such that

$$(2.13) \quad |\eta| \int_{\mathbf{R}^N} (|\nabla u|^2 + |u|^2) dx \leq C(|f|, |u|)_0.$$

Here $(\cdot, \cdot)_0$ is the inner product of $L_2(\mathbf{R}^N)$.

Proof. Let us first notice that $u = R(z)f$, $z \in \mathbf{C} \setminus \mathbf{R}$, satisfies the equation

$$(2.14) \quad -\Delta u - z\mu(x)u = \mu(x)f$$

and u belongs to $H^2(\mathbf{R}^N)$. Multiply both sides of (2.14) by \bar{u} , integrate over \mathbf{R}^N and take the imaginary part. Then the estimate for the term $|\eta| \int |u|^2 dx$ is obtained. If we multiply both sides of (2.14) by $|\eta|\bar{u}$ and take the real part, we get the estimate for the term $|\eta| \int |\nabla u|^2 dx$. \square

Let us now study the surface integrals of $u = R(z)f$ over the separating surface S . Suppose that $v \in H^2(\mathbf{R}^N)_{\text{loc}}$. Then the traces of v and $\partial_j v$ ($j = 1, 2, \dots, N$) on S are well-defined as elements of $L_2(S)_{\text{loc}}$. These traces on S will be denoted as v and $\partial_j v$ again. As usual, the inner product $(\cdot, \cdot)_S$ and norm $\|\cdot\|_S$ of $L_2(S)$ are defined by

$$(2.15) \quad (v, w)_S = \int_S v \bar{w} dS \quad \text{and} \quad \|v\|_S = [(v, v)_S]^{1/2}.$$

LEMMA 2.2 (*Eidus* [4], (2.14)). Assume (1.6). Let $u = R(z)f$ with $f \in L_2(\mathbf{R}^N)$ and $z = \lambda + i\eta \in \mathbf{C} \setminus \mathbf{R}$. Then $u \in L_2(S)$ and there exists a positive constant $C = C(\mu)$ depending only on $\mu(x)$ such that

$$(2.16) \quad |\lambda| \|u\|_S^2 \leq C(|f|, |u| + |\partial_N u|)_0.$$

Here $(\cdot, \cdot)_0$ means the inner product of $L_2(\mathbf{R}^N)$.

Proof. Multiply both sides of (2.14) by $\partial_N \bar{u}$, integrate over \mathbf{R}^N and take the real part. Then

$$(2.17) \quad -\operatorname{Re} \int_{\mathbf{R}^N} (\Delta u) \partial_N \bar{u} dx - \operatorname{Re} \int_{\mathbf{R}^N} z\mu(x)u(\partial_N \bar{u}) dx \\ = \operatorname{Re} \int_{\mathbf{R}^N} \mu(x)f(\partial_N \bar{u}) dx.$$

By the use of integration by parts, we have

$$(2.18) \quad \operatorname{Re} \int_{\mathbf{R}^N} (\partial_j^2 u)(\partial_N \bar{u}) dx = -\operatorname{Re} \int_{\mathbf{R}^N} (\partial_j u)(\partial_N \partial_j \bar{u}) dx \\ = -2^{-1} \int_{\mathbf{R}^N} \partial_N \{|\partial_j u|^2\} dx = 0$$

for $j = 1, 2, \dots, N$, which implies the first term in the left-hand side of (2.17) is identically zero. Let us next estimate the second term of the left-hand side of (2.17):

$$\begin{aligned}
 (2.19) \quad & -\operatorname{Re} \int_{\mathbf{R}^N} z\mu(x)u(\partial_N \bar{u}) \, dx \\
 & = \lambda \operatorname{Re} \int_{\mathbf{R}^N} \mu(x)u(\partial_N \bar{u}) \, dx \\
 & \quad - \eta \operatorname{Im} \int_{\mathbf{R}^N} \mu(x)u(\partial_N \bar{u}) \, dx \equiv I_1 - I_2.
 \end{aligned}$$

The terms I_1 and I_2 are estimated as follows:

$$\begin{aligned}
 (2.20) \quad |I_1| & = (|\lambda|/2) \left| \int_{\mathbf{R}^N} \mu(x)\partial_N(|u|^2) \, dx \right| \\
 & = 2^{-1}|\lambda(\mu_2 - \mu_1)| \left| \int_S |u|^2 n_{2N}(x) \, dS \right| \\
 & \geq 2^{-1}c_1|\lambda(\mu_2 - \mu_1)| \int_S |u|^2 \, dS,
 \end{aligned}$$

where we have used (1.6), and

$$(2.21) \quad |I_2| \leq |\eta|(|u|, |\partial_N u|)_0 \leq C_1(|f|, |u|)_0$$

with a constant C_1 , where we have used Lemma 2.1. The right-hand side of (2.17) is estimated as

$$(2.22) \quad \left| \operatorname{Re} \int_{\mathbf{R}^N} \mu(x)f(\partial_N \bar{u}) \, dx \right| \leq \operatorname{Max}(\mu_1, \mu_2)(|f|, |\partial_N u|)_0.$$

Combining these estimates, we obtain (2.16). □

Using Lemmas 2.1 and 2.2, Eidus proved the following proposition:

PROPOSITION 2.3 (Eidus [4], Lemma 2.1). *Let us assume (1.6). Let $f \in L_2(\mathbf{R}^N)$, $\gamma > 3/2$ for $N = 3$, $\gamma = 3/2$ for $N > 3$,*

$$(2.23) \quad 0 < |\eta| < 1 \quad \text{and} \quad -1 \leq \lambda \leq M,$$

where $z = \lambda + i\eta$ as above and M is some constant. Then

$$(2.24) \quad \|u\|_{-\gamma,1} \leq C(|f|, |u| + |\nabla u|)_0$$

for $u = R(z)f$, where the constant $C = C(M, \gamma)$ does not depend on f, λ, η (but may depend on M, γ).

For the proof see Eidus [4], p. 33–34. As for $\partial_j u$ on the separating surface S we have the following estimates:

LEMMA 2.4. Assume (1.6). Let $u = R(z)f$ with $f \in L_2(\mathbf{R}^N)$ and $z \in \Pi$, Π being a bounded set in $\mathbf{C} \setminus \mathbf{R}$. Then there exists a positive constant $C = C(\Pi, \mu)$ depending only on Π and $\mu(x)$ such that

$$(2.25) \quad \|\nabla u\|_S \leq C\{\|u\|_0 + \|f\|_0\},$$

where $\|\cdot\|_0$ is the norm of $L_2(\mathbf{R}^N)$ and

$$(2.26) \quad \|\nabla u\|_S^2 = \sum_{j=1}^N \|\partial_j u\|_S^2.$$

Proof. Let $B_R = \{x \in \mathbf{R}^N / |x| < R\}$ and $S_R = \{x \in \mathbf{R}^N / |x| = R\}$ for $R > 0$. By the Green formula we have

$$(2.27) \quad \int_{\Omega_1 \cap B_R} \partial_N(|\nabla u|^2) dx \\ = \int_{S \cap B_R} |\nabla u|^2 n_{1N}(x) dS + \int_{\Omega_1 \cap S_R} |\nabla u|^2 \tilde{x}_N dS,$$

where $n_{1N}(x)$ is the N th element of the outward normal of $\partial\Omega_1$ at x and $\tilde{x}_N = x_N/|x|$. By using (1.6) we obtain from (2.27)

$$(2.28) \quad c \int_{S \cap B_R} |\nabla u|^2 dS \\ \leq 2 \sum_{j=1}^N \int_{\Omega_1 \cap B_R} |\partial_N \partial_j u| |\partial_j u| dx + \int_{\Omega_1 \cap S_R} |\nabla u|^2 dS.$$

Let $R \rightarrow \infty$ along an appropriate sequence $\{R_m\}$ so that the second term of the right-hand side of (2.28) converges to zero. Then we have

$$(2.29) \quad \int_S |\nabla u|^2 dS \leq C_2 \|u\|_{2,0}^2$$

with a constant C_2 , where $\|\cdot\|_{2,0}$ is the norm of $H^2(\mathbf{R}^N)$. The estimate (2.25) follows from (2.29), the well-known estimate

$$(2.30) \quad \|u\|_{2,0} \leq C_3\{\|\Delta u\|_0 + \|u\|_0\}$$

with a constant C_3 and the equation (2.14). \square

3. More a priori estimates. In this section we shall introduce a “modified radiation condition-like” term for the solution u of the equation (2.14), i.e., $-\Delta u - z\mu(x)u = \mu(x)f$, where $z \in \mathbf{C} \setminus \mathbf{R}$. Some estimates for it will be proved. These estimates can be regarded as modifications of the estimates of the usual radiation conditions given in, e.g., [6], [5], [11], although we treat here only the case that $z \in \mathbf{C} \setminus \mathbf{R}$.

By using these estimates, we shall also show some estimates for the norm $\|u\|_{-\tau}$ with $\tau > 1/2$. They will be used in §4 to show the limiting absorption principle for the operator H .

Let us start with some notation. For $z \in \mathbb{C} \setminus \mathbb{R}$ and $x \in \mathbb{R}^N$ we set

$$(3.1) \quad k = k(x) = k(x, z) = [z\mu(x)]^{1/2},$$

$$(3.2) \quad a = a(x) = a(x, z) = \operatorname{Re} k(x, z),$$

$$(3.3) \quad b = b(x) = b(x, z) = \operatorname{Im} k(x, z),$$

where the branch of $[z\mu(x)]^{1/2}$ is taken so that $b(x, z) > 0$. With fixed z the functions $k(x)$, $a(x)$ and $b(x)$ are step functions on \mathbb{R}^N which are constant in each Ω_1 and Ω_2 . Let us next introduce some differential expressions of the first order;

$$(3.4) \quad \mathcal{D}_j u = \partial_j u + \{(N-1)/(2r)\} \tilde{x}_j u - ik(x) \tilde{x}_j u \\ (r = |x|, \tilde{x}_j = x_j/r, j = 1, 2, \dots, N),$$

$$(3.5) \quad \mathcal{D}u = \nabla u + \{(N-1)/(2r)\} \tilde{x}u - ik(x) \tilde{x}u \\ (\tilde{x} = (\tilde{x}_1, \tilde{x}_2, \dots, \tilde{x}_N)),$$

$$(3.6) \quad |\mathcal{D}u| = \left[\sum_{j=1}^N |\mathcal{D}_j u|^2 \right]^{1/2},$$

$$(3.7) \quad \mathcal{D}_r u = \mathcal{D}u \cdot \tilde{x} = \frac{\partial u}{\partial r} + \{(N-1)/(2r)\} u - ik(x)u,$$

$$(3.8) \quad \mathcal{D}_n u = \mathcal{D}u \cdot n = \frac{\partial u}{\partial n} + \{(N-1)/(2r)\} (\tilde{x} \cdot n)u - ik(x)(\tilde{x} \cdot n)u,$$

where n is a unit vector in \mathbb{R}^N .

LEMMA 3.1. *Let $u \in H^2(\mathbb{R}^N)_{\text{loc}}$ and set $f = \mu(x)^{-1}(-\Delta - k^2)u$, where $k = [z\mu(x)]^{1/2}$ with $z \in \mathbb{C} \setminus \mathbb{R}$. Let $\xi \in C^1([0, \infty))$ with $\xi(r) = 0$*

in a neighborhood of $r = 0$ and set $\varphi(x) = \xi(|x|)$. Then we have

$$\begin{aligned}
 (3.9) \quad & 2^{-1} \int_{B_R} \frac{\partial \varphi}{\partial r} |\mathcal{D}u|^2 dx + \sum_{l=1}^2 \int_{\partial \Omega_l \cap B_R} \varphi \cdot \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\
 & + \int_{B_R} b\varphi |\mathcal{D}u|^2 dx \\
 & + \int_{B_R} \left(\frac{\varphi}{r} - \frac{\partial \varphi}{\partial r} \right) (|\mathcal{D}u|^2 - |\mathcal{D}_r u|^2) dx \\
 & + c_N \int_{B_R} r^{-2} \left(\frac{\varphi}{r} - 2^{-1} \frac{\partial \varphi}{\partial r} + b\varphi \right) |u|^2 dx \\
 & = \operatorname{Re} \int_{B_R} \varphi \mu(x) \overline{f(\mathcal{D}_r u)} dx \\
 & + 2^{-1} \sum_{l=1}^2 \int_{\partial \Omega_l \cap B_R} \varphi \left\{ \frac{(N-1)b}{r} + |k|^2 \right\} (\tilde{x} \cdot n) |u|^2 dx \\
 & + 2^{-1} \int_{S_R} \varphi (2|\mathcal{D}_r u|^2 - |\mathcal{D}u|^2 - c_N r^{-2} |u|^2) dx
 \end{aligned}$$

for each $R > 0$, where $\partial/\partial n$ in the integrand of the surface integral over $\partial \Omega_l \cap B_R$ means the directional derivative in the direction of the outward normal of $\partial \Omega_l$, B_R is the open ball with center $x = 0$ and radius R , S_R is the sphere with center $x = 0$ and radius R and

$$(3.10) \quad c_N = (N-1)(N-3)/4.$$

Proof. Since $k(x)$ is a constant $[z\mu_1]^{1/2}$ or $[z\mu_2]^{1/2}$ in each region Ω_1 or Ω_2 , the equation $-\Delta u - k^2 u = \mu(x)f$ is rewritten as

$$(3.11) \quad - \sum_{j=1}^N \partial_j \mathcal{D}_j u + \left\{ \frac{N-1}{2r} - ik \right\} \mathcal{D}_r u + V_0(x)u = \mu(x)f$$

in Ω_1 or Ω_2 , where $\mathcal{D}_j u$ and $\mathcal{D}_r u$ are as above, $\partial_j = \partial/\partial x_j$ and $V_0(x)$ is given by

$$(3.12) \quad V_0(x) = c_N r^{-2}$$

(cf. Ikebe & Saitō [5], (2.10)). Multiply the both sides of (3.11) by $\varphi(\mathcal{D}_r u)$, integrate on each $B_R \cap \Omega_l$ ($l = 1, 2$) and take the real part. Then, using integration by parts or the Green formula and making the sum of these two integrals, we obtain the relation (3.9) (cf. [5], Lemma 2.2). \square

In order to get our first a priori estimate we are going to introduce the following weight functions.

Notation 3.2. (i) Let $\rho(r)$ be a C^2 function on $[0, \infty)$ such that

$$(3.13) \quad \rho(r) = 0 \quad (0 \leq r \leq R_0), \quad = 1 \quad (r \geq R_0 + 1)$$

with $R_0 > 0$, $0 \leq \rho(r) \leq 1$ and $\rho'(r) \geq 0$.

(ii) For each $\varepsilon > 0$ the function $\sigma_\varepsilon(r)$ is defined by

$$(3.14) \quad \sigma_\varepsilon(r) = \exp\{-\varepsilon^{-1}(1+r)^{-\varepsilon}\} \quad (r \geq 0).$$

PROPOSITION 3.3. *Let us assume Assumption 1.1. Let $u = R(z)f$ with $z = \lambda + i\eta \in \mathbf{C} \setminus \mathbf{R}$ and $f \in L_2(\mathbf{R}^N)$. Here $z = \lambda + i\eta$ moves in a bounded set K of $\mathbf{C} \setminus \mathbf{R}$ such that*

$$(3.15) \quad |\lambda| \geq \lambda_0 > 0 \quad (z = \lambda + i\eta \in K)$$

with a positive number λ_0 . Then there exists a constant $C = C(K)$ depending only on K such that

$$(3.16) \quad 2^{-1} \int_{\mathbf{R}^N} \frac{\partial}{\partial r} \{\rho^2 \sigma_\varepsilon\} |\mathcal{D}u|^2 dx \\ + \sum_{l=1}^2 \int_{\partial\Omega_l} \rho^2 \sigma_\varepsilon \cdot \text{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\ \leq C \left\{ (\sigma_\varepsilon |f|, |\nabla u| + |u|)_{0, E_{R_0}} \right. \\ \left. + \int_{E_{R_0}} (1 + |x|)^{-2} \sigma_\varepsilon (|\nabla u|^2 + |u|^2) dx \right\}$$

holds for each $0 < \varepsilon$, where

$$(3.17) \quad E_{R_0} = \{x/|x| > R_0\},$$

the functions ρ and σ_ε are as in Notation 3.2 with $r = |x|$.

Proof. Set $\varphi(x) = \rho(|x|)^2 \sigma_\varepsilon(|x|)$ in (3.9). We are going to evaluate each term of (3.9). Noting that

$$(3.18) \quad \frac{\varphi}{r} - \frac{\partial \varphi}{\partial r} = \rho^2 \sigma_\varepsilon \{r^{-1} - (1+r)^{-1-\varepsilon}\} - 2\rho\rho'\sigma_\varepsilon \geq -2\rho\rho'\sigma_\varepsilon$$

and

$$(3.19) \quad |\mathcal{D}u|^2 - |\mathcal{D}_r u|^2 = |\nabla u|^2 - \left| \frac{\partial u}{\partial r} \right|^2 \geq 0,$$

we have for $R > R_0$

$$(3.20) \quad \int_{B_R} \left\{ \frac{\varphi}{r} - \frac{\partial \varphi}{\partial r} \right\} \{ |\mathcal{D}u|^2 - |\mathcal{D}_r u|^2 \} dx \\ \geq -2 \int_{B_R} \rho \rho' \sigma_\varepsilon \{ |\nabla u|^2 - |u|^2 \} dx \geq -2 \int_{B_R} \rho \rho' \sigma_\varepsilon |\nabla u|^2 dx.$$

Therefore it follows from (3.9) that

$$(3.21) \quad 2^{-1} \int_{B_R} \frac{\partial \varphi}{\partial r} |\mathcal{D}u|^2 dx + \sum_{l=1}^2 \int_{\partial \Omega_l \cap B_R} \varphi \cdot \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\ \leq \operatorname{Re} \int_{B_R} \varphi \cdot \mu \cdot f \cdot \overline{(\mathcal{D}_r u)} dx + 2 \int_{B_R} \rho \rho' \sigma_\varepsilon |\nabla u|^2 dx \\ + 2^{-1} \sum_{l=1}^2 \int_{\partial \Omega_l \cap B_R} \varphi \left\{ \frac{(N-1)b}{r} + |k|^2 \right\} (\tilde{x} \cdot n) |u|^2 dS \\ - c_N \int_{B_R} r^{-2} \left(\frac{\varphi}{r} - 2^{-1} \frac{\partial \varphi}{\partial r} + b\varphi \right) |u|^2 dx \\ + 2^{-1} \int_{S_R} \varphi (2|\mathcal{D}_r u|^2 - |\mathcal{D}u|^2 - c_N r^{-2} |u|^2) dx.$$

Since $u = R(z)f \in H^2(\mathbf{R}^N)$, the fifth term of the right-hand side of (3.21) will converge to zero when $R \rightarrow \infty$ along a suitable sequence $\{R_m\}$. Thus we obtain from (3.21)

$$(3.22) \quad 2^{-1} \int_{\mathbf{R}^N} \frac{\partial \varphi}{\partial r} |\mathcal{D}u|^2 dx + \sum_{l=1}^2 \int_{\partial \Omega_l} \varphi \cdot \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\ \leq \operatorname{Re} \int_{\mathbf{R}^N} \varphi \cdot \mu \cdot f \cdot \overline{(\mathcal{D}_r u)} dx + 2 \int_{\mathbf{R}^N} \rho \rho' \sigma_\varepsilon |\nabla u|^2 dx \\ + 2^{-1} \sum_{l=1}^2 \int_{\partial \Omega_l} \varphi \left\{ \frac{(N-1)b}{r} + |k|^2 \right\} (\tilde{x} \cdot n) |u|^2 dS \\ + c_N \int_{\mathbf{R}^N} r^{-2} \left(\frac{\varphi}{r} - 2^{-1} \frac{\partial \varphi}{\partial r} + b\varphi \right) |u|^2 dx,$$

where we should note that the surface integrals over $\partial \Omega_l$ in (3.22) is absolutely convergent, because $u, \nabla u \in L_2(S)$ by Lemmas 2.2 and 2.4.

Let us denote by J the third term of the right-hand side of (3.22) and let us estimate the surface integral J . It follows from (1.7) in Assumption 1.1 that

$$(3.23) \quad |J| \leq C_1 \int_S \varphi (1 + |x|)^{-1} |u|^2 dS \quad (z \in K),$$

where the constant $C_1 = C_1(K)$ depends only on the bounded set K (and the constant c_2 in (1.7)). Set

$$(3.24) \quad v(x) = \varphi^{1/2}(1+r)^{-1/2}u = \rho\sigma_\varepsilon^{1/2}(1+r)^{-1/2}u \equiv \psi u.$$

Since v satisfies the equation

$$(3.25) \quad -\Delta v - k^2 v = \psi \mu f - 2(\nabla \psi) \cdot (\nabla u) - (\Delta \psi)u \equiv F$$

it follows that Lemma 2.2 that

$$(3.26) \quad |J| \leq C_1 \int_S |v|^2 dS \leq C_2(|F|, |\psi u + \partial_N(\psi u)|)_0$$

with a constant $C_2 = C_2(K)$. By a straightforward computation we get

$$(3.27) \quad |\partial^\alpha \psi(x)| \leq C_3 \sigma_\varepsilon(|x|)^{1/2}(1+|x|)^{-(1/2)-|\alpha|} \\ (|x| \geq R_0, |\alpha| = 0, 1, 2),$$

with $\partial^\alpha \psi(x) = 0$ for $|x| \leq R_0$ ($|\alpha| = 0, 1, 2$), where C_3 is a positive constant depending only on R_0 though the constant R_0 is fixed throughout this work. Since it follows from (3.27) that

$$(3.28) \quad |F| \cdot |\psi u + \partial_N(\psi u)| \leq C_4 \sigma_\varepsilon \{ |f| \cdot (|\nabla u| + |u|) \\ + (1+|x|)^{-2}(|\nabla u|^2 + |u|^2) \},$$

with a constant $C_4 = C_4(K)$, we have

$$(3.29) \quad |J| \leq C_5 \left\{ (\sigma_\varepsilon |f|, |\nabla u| + |u|)_{0, E_{R_0}} \right. \\ \left. + \int_{E_{R_0}} (1+|x|)^{-2} \sigma_\varepsilon (|\nabla u|^2 + |u|^2) dx \right\}$$

with $C_5 = C_5(K)$.

The other three terms in the right-hand side of (3.22) can be easily evaluated. Thus we get (3.16). \square

COROLLARY 3.4. *Let p, σ_ε, K and $\varphi = \rho^2 \sigma_\varepsilon$ be as above. Then there exists a positive constant $\tilde{C} = \tilde{C}(K)$ such that*

$$(3.30) \quad \int_{E_{R_0}} \frac{\partial \varphi}{\partial r} |\nabla u - ik\tilde{x}u|^2 dx + 2 \sum_{l=1}^2 \int_{\partial \Omega_l} \varphi \cdot \text{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\ \leq \tilde{C} \left\{ (\sigma_\varepsilon |f|, |\nabla u| + |u|)_{0, E_{R_0}} \right. \\ \left. + \int_{R_0} (1+|x|)^{-2} \sigma_\varepsilon (|\nabla u|^2 + |u|^2) dx \right\}$$

holds for $u = R(z)f$ with $f \in L_2(\mathbf{R}^N)$, $\varepsilon > 0$ and $z \in K$.

Proof. The estimate (3.30) directly follows from (3.16) by noting that

$$(3.31) \quad \begin{aligned} & |\mathcal{D}u|^2 - |\nabla u - ik\tilde{x}u|^2 \\ &= \frac{N-1}{2r} \left(\frac{\partial |u|^2}{\partial r} \right) - \frac{(N-1)b|u|^2}{r} + \frac{(N-1)^2|u|^2}{4r^2}. \quad \square \end{aligned}$$

REMARK 3.5. (i) It should be noted that the constants C in (3.16) and \tilde{C} in (3.30) do not depend on ε in the function σ_ε .

(ii) The estimate (3.16) or (3.30) can be regarded as a (weaker) generalization of the estimate

$$(3.32) \quad \int_{\mathbf{R}^N} (1 + |x|)^{2\delta-2} |\nabla v - ik\tilde{x}v|^2 dx \leq C \{ \|g\|_\delta^2 + \|v\|_{-\delta}^2 \}$$

with $\delta > 1/2$ for solutions v of the Schrödinger equation

$$(-\Delta + Q(x) - k^2)v = g$$

with a short-range or long-range potential (cf. e.g., Ikebe & Saitō [5], Lemma 1.7).

The estimate obtained above will be used to prove an a priori estimate of the $H^1_{-\delta}(\mathbf{R}^N)$ -norm of $u = R(z)f$. First we need the following lemma.

LEMMA 3.6. *Let S be the separating surface as above. Then there exists $R_1 > 0$ such that*

$$(3.33) \quad \begin{aligned} & \int_{S \cap E_R} F(x) dS \\ &= \int_R^\infty \left\{ \int_{S^{N-2}} F(t\omega, \varphi(t\omega)) \frac{(1 + |\nabla \varphi|^2)^{1/2} t^{N-2} r}{t + \varphi(t\omega) \cdot (\nabla \varphi \cdot \omega)} d\omega \right\} dr \end{aligned}$$

holds for any $R \geq R_1$ and any integrable function $F(x)$ over $S \cap E_R$, where $E_R = \{x \in \mathbf{R}^N / |x| > R\}$, s^{N-2} is the unit sphere in \mathbf{R}^{N-1} , $\omega \in S^{N-2}$, $t = t(r, \omega)$ is defined by $r = (t^2 + \varphi(t\omega)^2)^{1/2}$ and $\nabla \varphi = (\nabla \varphi)(t\omega)$.

Proof. By use of the relation

$$(3.34) \quad \begin{aligned} dS &= (1 + |\nabla \varphi|^2)^{1/2} dx' \\ &= (1 + |\nabla \varphi|^2)^{1/2} t^{N-2} dt d\omega \quad (x' = t\omega) \end{aligned}$$

it follows that

$$(3.35) \quad \begin{aligned} & \int_{S \cap E_R} F dS \\ &= \int_{S^{N-2}} d\omega \int_{(t^2 + \varphi(t\omega)^2)^{1/2} > R} F(t\omega, \varphi(t\omega)) (1 + |\nabla \varphi|^2)^{1/2} t^{N-2} dt. \end{aligned}$$

For fixed $\omega \in S^{N-2}$ let us consider the relation

$$(3.36) \quad r = (t^2 + \varphi(t\omega))^2)^{1/2}.$$

Then

$$(3.37) \quad \frac{dr}{dt} = \left(t + \varphi(t\omega) \cdot \frac{\partial \varphi(t\omega)}{\partial t} \right) (t^2 + \varphi(t\omega)^2)^{-1/2}.$$

Since the outward normals $n_l(x)$ of $\partial\Omega_l$ are expressed as

$$(3.38) \quad n_l(x) = \pm (1 + |\nabla \varphi(x')|^2)^{-1/2} \left(\frac{\partial \varphi}{\partial x_1}, \frac{\partial \varphi}{\partial x_2}, \dots, \frac{\partial \varphi}{\partial x_{N-1}}, -1 \right) \\ (x = (x', x_N) \in S),$$

the inner product $n_l(x) \cdot x$ ($x \in S$) in \mathbf{R}^N has the form

$$(3.39) \quad n_l(x) \cdot x = \pm (1 + |\nabla \varphi(x')|^2)^{-1/2} \left\{ r' \frac{\partial \varphi(x')}{\partial r'} - \varphi(x') \right\}$$

with $r' = |x'|$. By setting $x' = t\omega$ in (3.39) and using (1.6) and (1.7) in Assumption 1.1, we obtain

$$(3.40) \quad \frac{\partial \varphi(t\omega)}{\partial t} = t^{-1} \varphi(t\omega) + O(t^{-1}) \quad (t \rightarrow \infty),$$

where the term $O(t^{-1})$ is uniform for $\omega \in S^{N-2}$. Here we should note that the boundedness of $|\nabla \varphi(x')|$ on $\mathbf{R}^{N-1} \setminus \{0\}$ follows from (1.6) since the N th component $n_{lN}(x)$ of the outward normal $n_l(x)$ has the form $\pm(1 + |\nabla \varphi(x')|)^{-1/2}$. It follows from (3.40) that

$$(3.41) \quad \frac{dr}{dt} = \{t + t^{-1} \varphi(t\omega)^2 + \varphi(t\omega) \cdot O(t^{-1})\} \{t^2 + \varphi(t\omega)^2\}^{-1/2} \\ = \{1 + t^{-2} \varphi(t\omega)^2 + t^{-1} \varphi(t\omega) \cdot O(t^{-1})\} \{1 + t^{-2} \varphi(t\omega)^2\}^{-1/2} \\ = \{1 + t^{-2} \varphi(t\omega)^2\}^{1/2} + t^{-1} \varphi(t\omega) \{1 + t^{-2} \varphi(t\omega)^2\}^{-1/2} \cdot O(t^{-1}) \\ = \{1 + t^{-2} \varphi(t\omega)^2\}^{1/2} + O(t^{-1})$$

as $t \rightarrow \infty$. Therefore there exists a positive number R_1 ($\geq R_0$) such that $dr/dt > 0$ for all $r \geq R_1$. Therefore the inverse function $t = t(r) = t(r, \omega)$ is well-defined and is a C^1 function. Thus we have (3.33). \square

In the following proposition we shall study some integrals which are closely related to the left-hand side of (3.30).

PROPOSITION 3.7. *Let us assume Assumption 1.1. Let $u = R(z)f$ with $z \in \mathbf{C} \setminus \mathbf{R}$ and $f \in L_2(\mathbf{R}^N)$.*

(i) Then we have

$$(3.42) \quad \begin{aligned} & \sum_{l=1}^2 \int_{\partial\Omega_l} \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\ &= - \int_{\mathbf{R}^N} b(|\nabla u|^2 + |k|^2|u|^2) dx \\ & \quad - \int_{\mathbf{R}^N} \mu(x) \cdot \operatorname{Im}(\bar{k} \cdot f \cdot \bar{u}) dx, \end{aligned}$$

where $k = [z\mu]^{1/2}$, $b = \operatorname{Im} k$ and n in the surface integral over $\partial\Omega_l$ means the outward unit normal of $\partial\Omega_l$.

(ii) We have

$$(3.43) \quad \begin{aligned} & \int_{S_r} |\nabla u - ik\tilde{x}u|^2 dS + 2 \sum_{l=1}^2 \int_{\partial\Omega_l \cap E_r} \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\ &= \int_{S_r} (|\nabla u|^2 + |k|^2|u|^2) dS - 2 \int_{E_r} b(|\nabla u|^2 + |k|^2|u|^2) dx \\ & \quad - 2 \int_{E_r} \mu(x) \cdot \operatorname{Im}\{\bar{k} \cdot f \cdot \bar{u}\} dx, \end{aligned}$$

for $r > 0$, where $S_r = \{x \in \mathbf{R}^N / |x| = r\}$, $E_r = \{x \in \mathbf{R}^N / |x| > r\}$ and n in the surface integral over $\partial\Omega_l \cap E_r$ means the outward unit normal of $\partial\Omega_l$.

Proof. Multiply the equation $-\Delta u - k^2u = \mu f$ by $\bar{k}u$ and integrate on E_r . Then we have

$$(3.44) \quad \begin{aligned} & - \sum_{l=1}^2 \int_{\partial\Omega_l \cap E_r} \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} dS + \int_{S_r} \bar{k} \cdot \frac{\partial u}{\partial r} \cdot \bar{u} dS \\ & \quad + \int_{E_r} \bar{k} |\nabla u|^2 dx - \int_{E_r} k |k|^2 |u|^2 dx = \int_{E_r} \bar{k} \mu(x) f \bar{u} dx. \end{aligned}$$

The relation (3.42) is obtained by taking the imaginary part of (3.44) and letting $r \rightarrow 0$ along a suitable sequence $\{r_m\}$. The second relation (3.43) follows from (3.44) and

$$(3.45) \quad \begin{aligned} & \int_{S_r} |\nabla u - ik\tilde{x}u|^2 dS \\ &= \int_{S_r} (|\nabla u|^2 + |k|^2|u|^2) dS - 2 \int_{S_r} \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial r} \cdot \bar{u} \right\} dS. \quad \square \end{aligned}$$

Let us now evaluate the norm $\|u\|_{-\delta, E_r}$.

PROPOSITION 3.8. *Suppose that Assumption 1.1 is satisfied. Let R_1 be as in Lemma 3.6. Let $\sigma_\varepsilon(r)$ be as in Notation 3.2. Let $u = R(z)f$ with $z \in K$ and $f \in L_2(\mathbf{R}^N)$, where K is as in Proposition 3.3. Then there exists a positive constant $C = C(K)$ such that*

$$(3.46) \quad \int_{E_R} (1 + |x|)^{-2\tau} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \leq C(1 + R)^{-(2\tau-1-\varepsilon)} \left\{ (|f|, |\nabla u| + |u|)_0 + \int_{E_{R_1}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \right\}$$

holds for all $R \geq R_1 + 1$ and all pairs (τ, ε) satisfying

$$(3.47) \quad 0 < \varepsilon \leq 2\tau - 1,$$

where $(\cdot, \cdot)_0$ means the inner product of $L_2(\mathbf{R}^N)$.

Proof. The proof will be divided into several steps.

(I) Set

$$(3.48) \quad \mathcal{A}(r) = \mathcal{A}(r, z, f) = \int_{S_r} (|\nabla u|^2 + |k|^2 |u|^2) dS.$$

Let $\rho(r)$ be as in Notation 3.2 with R_0 replaced by R_1 and let $R \geq R_1 + 1$. Then we have

$$(3.49) \quad \begin{aligned} & \int_{E_R} (1 + |x|)^{-2\tau} \sigma_\varepsilon(|\nabla u|^2 + |k|^2 |u|^2) dx \\ &= \int_R^\infty (1 + r)^{-2\tau} \sigma_\varepsilon(r) \mathcal{A}(r) dr \\ &\leq (1 + R)^{-(2\tau-1-\varepsilon)} \int_R^\infty (1 + r)^{-(1+\varepsilon)} \sigma_\varepsilon(r) \mathcal{A}(r) dr \\ &\leq (1 + R)^{-(2\tau-1-\varepsilon)} \int_{R_1}^\infty \rho(r)^2 (1 + r)^{-(1+\varepsilon)} \sigma_\varepsilon(r) \mathcal{A}(r) dr \\ &\leq (1 + R)^{-(2\tau-1-\varepsilon)} \int_{R_1}^\infty \frac{\partial}{\partial r} \{ \rho^2 \sigma_\varepsilon \} \mathcal{A}(r) dr, \end{aligned}$$

where we should note that $2\tau - 1 - \varepsilon \geq 0$, $\rho(r) = 1$ for $r \geq R_1 + 1$ and

$$(3.50) \quad \frac{\partial}{\partial r} \{ \rho(r)^2 \sigma_\varepsilon(r) \} \geq \rho(r)^2 (1 + r)^{-(1+\varepsilon)} \sigma_\varepsilon(r) \geq 0$$

for $r \geq R_1$.

(II) It follows from (3.43) that

$$\begin{aligned}
 (3.51) \quad & \int_{R_1}^{\infty} \frac{\partial \varphi}{\partial r} \mathcal{A}(r) dr \\
 &= \int_{R_1}^{\infty} \frac{\partial \varphi}{\partial r} \left[\int_{S_r} |\nabla u - ik\tilde{x}u|^2 dS \right. \\
 &\quad + 2 \sum_{l=1}^2 \int_{\partial\Omega_l \cap E_r} \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\
 &\quad + 2 \int_{E_r} b(|\nabla u|^2 + |k|^2|u|^2) dx \\
 &\quad \left. + 2 \int_{E_r} \mu(x) \cdot \operatorname{Im} \{ \bar{k} \cdot f \cdot \bar{u} \} dx \right] dr,
 \end{aligned}$$

where we set $\varphi = \rho^2 \sigma_\varepsilon$. Let us first look at the term

$$(3.52) \quad J_l = \int_{R_1}^{\infty} \frac{\partial \varphi}{\partial r} \left[\int_{\partial\Omega_l \cap E_r} \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \right] dr.$$

It follows from Lemma 3.6 that the term J_l is expressed in the form

$$(3.53) \quad J_l = \int_{R_1}^{\infty} \frac{\partial \varphi}{\partial r} \left[\int_r^{\infty} F_l(s) ds \right] dr,$$

whence we get

$$(3.54) \quad J_l = \int_{R_1}^{\infty} \varphi(r) F_l(r) dr = \int_{\partial\Omega_l} \varphi \cdot \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS.$$

Here we should note that $\varphi(r) = \rho(r)^2 \sigma_\varepsilon(r) = 0$ for $r \leq R_1$. Thus, together with Corollary 3.4, we obtain

$$\begin{aligned}
 (3.55) \quad & \int_{R_1}^{\infty} \frac{\partial \varphi}{\partial r} \left[\int_{S_r} |\nabla u - ik\tilde{x}u|^2 dS \right. \\
 &\quad \left. + 2 \sum_{l=1}^2 \int_{\partial\Omega_l \cap E_r} \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \right] dr \\
 &= \int_{E_{R_1}} \frac{\partial \varphi}{\partial r} |\nabla u - ik\tilde{x}u|^2 dx \\
 &\quad + 2 \sum_{l=1}^2 \int_{\partial\Omega_l \cap E_{R_1}} \varphi \cdot \operatorname{Im} \left\{ \bar{k} \cdot \frac{\partial u}{\partial n} \cdot \bar{u} \right\} dS \\
 &\leq C_1 \left\{ (\sigma_\varepsilon |f|, |\nabla u| + |u|)_{0, E_{R_1}} \right. \\
 &\quad \left. + \int_{E_{R_1}} (1 + |x|)^{-2} \sigma_\varepsilon (|\nabla u|^2 + |u|^2) dx \right\}
 \end{aligned}$$

with a constant $C_1 = C_1(K)$.

(III) It is seen from Lemma 2.1 that we have

$$(3.56) \quad \int_{E_t} b(|\nabla u|^2 + |k|^2|u|^2) dx \leq C_2(|f|, |u|)_0$$

with a constant $C_2 = C_2(K)$, where we should note that $2ab = \mu(x)\eta$ and $|a| \geq \sqrt{\lambda_0}$. Therefore we obtain

$$(3.57) \quad \int_{R_1}^{\infty} \frac{\partial \varphi}{\partial t} \left[\int_{E_t} b(|\nabla u|^2 + |k|^2|u|^2) dx + \int_{E_t} \mu \cdot \text{Im}\{\bar{k} \cdot f \cdot \bar{u}\} dS \right] dt \leq C_3(|f|, |u|)_0$$

with $C_3 = C_3(K)$, where we have used the fact that $\varphi(r) \rightarrow 1$ as $r \rightarrow \infty$. The estimate (3.46) follows from (3.49), (3.51), (3.55) and (3.57). \square

The estimate (3.46) can be improved in the following way.

COROLLARY 3.9. *Let K, R_1 be as above. Then there exist positive constants $\tilde{C} = \tilde{C}(K)$ and $R_2 = R_2(K)$ such that*

$$(3.58) \quad \int_{E_R} (1 + |x|)^{-2\tau} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \leq \tilde{C}(1 + R)^{-(2\tau-1-\varepsilon)} \left\{ (|f|, |\nabla u| + |u|)_0 + \int_{B_{R_1, R_2}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \right\} \\ (R \geq R_1 + 1, u = R(z)f, f \in L_2(\mathbf{R}^N), z \in K)$$

holds for τ and ε satisfying $0 < \varepsilon \leq 2\tau - 1$ and $\varepsilon < 1$. Here we set

$$(3.59) \quad B_{R_1, R_2} = \{x \in \mathbf{R}^N / R_1 < |x| < R_2\}$$

and σ_ε is as in Notation 3.2.

Proof. It follows from (3.46) that

$$(3.60) \quad \int_{E_R} (1 + |x|)^{-2\tau} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \leq C(1 + R)^{-(2\tau-1-\varepsilon)} \left\{ (|f|, |\nabla u| + |u|)_0 + \int_{E_{R_1}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \right\}$$

for $R \geq R_1 + 1$ and $u = R(z)f$ with $z \in K$ and $f \in L_2(\mathbf{R}^N)$, where $C = C(K)$. Take $R_2 = R_2(K) (\geq R_1 + 1)$ so that

$$(3.61) \quad C(1 + R_2)^{-(1-\varepsilon)} \leq \frac{1}{2},$$

where we should note that $1 - \varepsilon > 0$. Setting $\tau = 1$ and $R = R_2$ in (3.60), we get

$$(3.62) \quad \int_{E_{R_2}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \\ \leq (|f|, |\nabla u| + |u|)_0 + \int_{B_{R_1, R_2}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx.$$

The estimate (3.58) directly follows from (3.62) and (3.46) with $\tilde{C} = 2C$. \square

By the use of Proposition 2.3 the second term in the right-hand side of (3.58) can be eliminated.

PROPOSITION 3.10. *Let Assumption 1.1 be satisfied. Let $\sigma_\varepsilon(r)$, τ , R_1 , R_2 and ε be as above. Let K be a bounded set in $\mathbf{C} \setminus \mathbf{R}$ such that $\operatorname{Re} z \geq \lambda_0$ for $z \in K$ with a positive constant λ_0 . Then there exists a positive constant $\bar{C} = \bar{C}(K)$ such that*

$$(3.63) \quad \int_{E_R} (1 + |x|)^{-2\tau} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \\ \leq \bar{C}(1 + R)^{-(2\tau-1-\varepsilon)} (|f|, |\nabla u|^2 + |u|^2)_0 \\ (R \geq R_1 + 1, u = R(z)f, f \in L_2(\mathbf{R}^N), z \in K).$$

Proof. By the use of Proposition 2.3 the integral over B_{R_1, R_2} in (3.58) is estimated as

$$(3.64) \quad \int_{B_{R_1, R_2}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u|^2 + |u|^2) dx \\ \leq (1 + R_2)^{2\gamma-2} \int_{R_1, R_2} (1 + |x|)^{-2\gamma} (|\nabla u|^2 + |u|^2) dx \\ \leq (1 + R_2)^{2\gamma-2} C (|f|, |\nabla u| + |u|)_0,$$

where the constants γ and C are as given in Proposition 2.3. The estimate (3.63) is obtained from (3.58) and (3.64) by noting that $R_2 = R_2(K)$ depend only on K . \square

4. The limiting absorption principle. The results obtained in the preceding section will be used to show the limiting absorption principle for the operator H .

THEOREM 4.1. *Suppose that Assumptions 1.1 are satisfied. Let K be a bounded set of $\mathbf{C} \setminus \mathbf{R}$ such that we have for $z = \lambda + i\eta \in K$*

$$(4.1) \quad \lambda \geq \lambda_0 > 0$$

with a positive number λ_0 . Let δ be a real number such that

$$(4.2) \quad \delta > 1/2.$$

Then there exists a positive constant $C = C(K, \delta)$ such that

$$(4.3) \quad \|R(z)f\|_{2,-\delta} \leq C\|f\|_\delta$$

holds for any $f \in L_{2,\delta}(\mathbf{R}^N)$ and any $z \in K$, where $R(z)$ is the resolvent of H .

Proof. (I) Let us first show that

$$(4.4) \quad \|R(z)f\|_{1,-\delta} \leq C\|f\|_\delta \quad (z \in K, f \in L_{2,\delta}(\mathbf{R}^N)).$$

We can assume with no loss of generality that

$$(4.5) \quad 1/2 < \delta \leq 1.$$

Let us assume (4.4) is false. Then there exist sequences $\{z_m\} \subset K$ and $\{f_m\} \subset L_{2,\delta}(\mathbf{R}^N)$ such that

$$(4.6) \quad \|R(z_m)f_m\|_{1,-\delta} = 1 \quad \text{and} \quad \|f_m\|_\delta < 1/m \quad (m = 1, 2, \dots).$$

We shall set

$$(4.7) \quad u_m = R(z_m)f_m.$$

Since the sequence $\{z_m\}$ is contained in a bounded subset K of \mathbf{C} , there exists a subsequence of $\{z_m\}$ which converges to an element z_0 of \mathbf{C} . For the sake of avoiding complication of notations we shall express the subsequence by $\{z_m\}$ again. Since u_m satisfies the equation

$$(4.8) \quad -\Delta u_m - z_m \mu(x)u_m = \mu(x)f_m,$$

it follows from the Rellich Lemma and the interior estimate that there exists a subsequence of $\{u_m\}$ which converges to u_0 in $H^1(\mathbf{R}^N)_{\text{loc}}$. The subsequence will be denoted by $\{u_m\}$ again.

Set $\tau = \delta$, $\varepsilon = (2\delta - 1)/2$, $u = u_m$, and $f = f_m$ in (3.63). Here we should note that the conditions $2\tau - 1 \geq \varepsilon > 0$ and $\varepsilon < 1$ are satisfied. Then we have

$$(4.9) \quad \int_{E_R} (1 + |x|)^{-2\delta} \sigma_\varepsilon (|\nabla u_m|^2 + |u_m|^2) dx \leq \bar{C}(1 + R)^{-(2\delta-1)/2} (|f_m|, |\nabla u_m| + |u_m|)_0 \quad (m = 1, 2, \dots)$$

for any $R \geq R_1 + 1$, where R_1 is given in Lemma 3.6. Note that

$$(4.10) \quad 0 < \exp\{-\varepsilon^{-1}(1 + R_1)^{-\varepsilon}\} \leq \sigma_\varepsilon(r) \leq 1 \quad (r \geq R_1, \varepsilon > 0),$$

and the term $(|f_m|, |\nabla u_m| + |u_m|)$ is uniformly bounded for $n = 1, 2, \dots$ by (4.6) and the Schwarz inequality. Thus it follows from (4.9) and (4.10) that

$$(4.11) \quad \|u_m\|_{1, -\delta, E_R} \leq C_1(1 + R)^{-(2\tau-1)/2} \quad (m = 1, 2, \dots)$$

with a constant C_1 , where $\|\cdot\|_{1, -\delta, E_R}$ is the norm of $H^1_{-\delta}(E_R)$ with $E_R = \{x \in \mathbf{R}^N / |x| > R\}$. Thus the (sub)sequence $\{u_m\}$ not only converges to u_0 in $H^1(\mathbf{R}^N)_{\text{loc}}$ but also has a uniformly small $H^1_{-\delta}(E_R)$ -norm when $R \rightarrow \infty$. Therefore we have shown that the sequence $\{u_m\}$ converges to u_0 in $H^1_{-\delta}(\mathbf{R}^N)$. Especially we have

$$(4.12) \quad \|u_0\|_{1, -\delta} = 1.$$

Let $m \rightarrow \infty$ and $R = R_1 + 1$ in (4.9). Then, since $\{u_m\}$ is proved to converge to u_0 in $H^1_{-\delta}(\mathbf{R}^N)$, in the left-hand side of (4.9) we have

$$(4.13) \quad \begin{aligned} \lim_{m \rightarrow \infty} \int_{E_{R_1+1}} (1 + |x|)^{-2\delta} \sigma_\varepsilon(|\nabla u_m|^2 + |u_m|^2) dx \\ = \int_{E_{R_1+1}} (1 + |x|)^{-2\delta} \sigma_\varepsilon(|\nabla u_0|^2 + |u_0|^2) dx. \end{aligned}$$

On the other hand, by (4.6) and the Schwarz inequality, the right-hand side of (4.9) converges to zero as $m \rightarrow \infty$. Therefore we arrive at

$$(4.14) \quad u_0 = 0 \quad \text{on } E_{R_1+1},$$

which implies, by noticing that u_0 is a solution of the elliptic equation $-\Delta u - z_0\mu(x)u = 0$ and by the use of the unique continuation theorem, that u_0 is identically zero on \mathbf{R}^N . This contradicts (4.12). Thus we have proved (4.4).

(II) It is easily seen that (4.3) follows from (4.4) and the fact that u is a solution of the equation $-\Delta u - k^2u = \mu(x)f$. □

REMARK 4.2 As we have seen, Proposition 3.10 plays a crucial role in the above proof of Theorem 4.1, and Proposition 3.10 is an improvement of Corollary 3.9 by Proposition 2.3. It might be interesting to discuss what we can get if we use only Corollary 3.9 without using Proposition 2.3. Let us suppose that (4.6) holds with $\{z_m\} \in K$, $\{f_m\} \subset L_{2,\delta}(\mathbf{R}^N)$ and $\{u_m\} = \{R(z_m)f_m\}$. We can assume, with no loss of generality, that the sequence $\{z_m\}$ converges to a positive number λ_1 (if the limit is a nonreal number, a contradiction follows much

more easily). It follows from the estimate (3.58) in Corollary 3.9 that there exists a subsequence of $\{u_m\}$, which will be denoted by $\{u_m\}$ again, such that

$$(4.15) \quad s\text{-}\lim_{m \rightarrow \infty} u_m = u_0 \quad \text{in } H^1_{-\tau}(\mathbf{R}^N)$$

for any $\tau > 1/2$. As has been shown, u_0 is a solution of the homogeneous equation $-\Delta u - \lambda_1 \mu(x)u = 0$.

Let ε satisfy $0 < \varepsilon \leq 1$ and set $\tau = (1 + \varepsilon)/2$, $u = u_m$, $f = f_m$ and $R = R_* = \text{Max}\{R_1 + 1, R_2\}$ and (3.58). Then we have

$$(4.16) \quad \int_{E_{R_*}} (1 + |x|)^{-(1+\varepsilon)} \sigma_\varepsilon(|\nabla u_m|^2 + |u_m|^2) dx \\ \leq C' \left\{ (|f_m|, |\nabla u_m|^2 + |u_m|^2)_0 \right. \\ \left. + \int_{B_{R_1, R_2}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u_m|^2 + |u_m|^2) dx \right\}$$

with a constant C' which is independent of ε in $(0,1)$. Let $m \rightarrow \infty$ in (4.16). Then we get

$$(4.17) \quad \int_{E_{R_*}} (1 + |x|)^{-(1+\varepsilon)} \sigma_\varepsilon(|\nabla u_0|^2 + |u_0|^2) dx \\ \leq C' \int_{B_{R_1, R_2}} (1 + |x|)^{-2} \sigma_\varepsilon(|\nabla u_0|^2 + |u_0|^2) dx.$$

By noting that

$$(4.18) \quad \sigma_\varepsilon(r) \geq \sigma_\varepsilon(R_*) \geq \sigma_\varepsilon(R_2) = \exp\{-\varepsilon^{-1}(1 + R_2)^{-1}\} \quad (r \geq R_*)$$

and

$$(4.19) \quad \sigma_\varepsilon(R_2)^{-1} \cdot \sigma_\varepsilon(r) \leq 1 \quad (R_1 \leq r \leq R_2),$$

it follows from (4.17) that

$$(4.20) \quad \int_{E_{R_*}} (1 + |x|)^{-(1+\varepsilon)} (|\nabla u_0|^2 + |u_0|^2) dx \\ \leq C' \int_{B_{R_1, R_2}} (1 + |x|)^{-2} (|\nabla u_0|^2 + |u_0|^2) dx$$

for any ε in $(0,1)$, whence we get, by letting ε to zero,

$$(4.21) \quad \int_{E_{R_*}} (1 + |x|)^{-1} (|\nabla u_0|^2 + |u_0|^2) dx < \infty,$$

i.e., we have $u_0 \in H^1_{-1/2}(\mathbf{R}^N)$.

If u_0 is a solution of the homogeneous Schrödinger equation

$$(4.22) \quad -\Delta u + Vu - \lambda_1 u = 0 \quad (\lambda_1 > 0)$$

with, e.g., a short-range or long-range potential $V(x)$, then, by the theorems on the asymptotic behavior of a solution of the equation (4.22) (e.g., Kato [7], Eastham & Kalf [3]), the estimate (4.21) is sufficient to guarantee that u_0 is identically zero. However, it seems to be an open problem whether the above theorems on the asymptotic behavior of the solution can be extended to our case.

Now that the estimate (4.3) has been shown, we can make use of Theorem 1.2 to show that the limiting absorption principle for the operator H holds between $L_{2,\delta}(\mathbf{R}^N)$ and $L_{2,-\delta}(\mathbf{R}^N)$.

THEOREM 4.3. *Suppose that Assumptions 1.1 are satisfied. Let $R(z)$ be the resolvent of the selfadjoint operator H and let $\delta > 1/2$. Then for each $\lambda > 0$ there exist the operators $R^\pm(\lambda)$ in $B(L_{2,\delta}(\mathbf{R}^N), L_{2,-\delta}(\mathbf{R}^N))$ such that we have*

$$(4.23) \quad \text{s-}\lim_{\eta \rightarrow \pm 0} R(\lambda + i\eta)f = R^\pm(\lambda)f \quad \text{in } H_{-\delta}^2(\mathbf{R}^N)$$

for all $f \in L_{2,\delta}(\mathbf{R}^N)$. Furthermore, $R^\pm(\lambda)$ is an $H_{-\delta}^2(\mathbf{R}^N)$ -valued, strongly continuous function on $(0, \infty)$ for each $f \in L_{2,\delta}(\mathbf{R}^N)$.

Proof. It is sufficient to assume that $1/2 < \delta \leq 1$. Let us show the existence of the limit (4.23) only when $\eta \rightarrow +0$. The case where $\eta \rightarrow -0$ can be treated in the same way. Let us denote by $R^+(\lambda)$ the operator in $B(L_{2,1}(\mathbf{R}^N), L_{2,-1}(\mathbf{R}^N))$ whose existence has been proved by Theorem 1.2. Thus we have

$$(4.24) \quad \text{s-}\lim_{\eta \rightarrow +0} R(\lambda + i\eta)f = R^+(\lambda)f \quad \text{in } L_{2,-1}(\mathbf{R}^N)$$

for all $f \in L_{2,1}(\mathbf{R}^N)$.

Let us first prove (4.23) for $f \in L_{2,1}(\mathbf{R}^N)$. Let $\{\eta_m\}$ be an arbitrary positive sequence such that $\eta_m \rightarrow +0$ and let $u_m = R(z_m)f$ with $z_m = \lambda + i\eta_m$. Then, proceeding as in the proof of Theorem 4.1, we can find a subsequence $\{u_{m_j}\}$ of $\{u_m\}$ which converges strongly to an element u_0 in $H_{-\delta}^2(\mathbf{R}^N)$ with the estimate

$$(4.25) \quad \|u_0\|_{2,-\delta} \leq C\|f\|_\delta$$

with a positive constant C which remains bounded when λ moves in a compact set in $(0, \infty)$. On the other hand it follows from (4.24) that

the sequence $\{u_m\}$ itself converges to $R^+(\lambda)f$ in $L_{2,-1}(\mathbf{R}^N)$. Therefore we have $u_0 = R^+(\lambda)f$,

$$(4.26) \quad \text{s-lim}_{j \rightarrow \infty} u_{m_j} = R^+(\lambda)f \quad \text{in } H_{-\delta}^2(\mathbf{R}^N),$$

and

$$(4.27) \quad \|R^+(\lambda)f\|_{2,-\delta} \leq C\|f\|_{\delta}.$$

Since the sequence $\{\eta_m\}$ was taken arbitrarily, we can conclude that (4.23) is true for $f \in L_{2,1}(\mathbf{R}^N)$. It follows from (4.27) and the denseness of $L_{2,1}(\mathbf{R}^N)$ in $L_{2,\delta}(\mathbf{R}^N)$ that $R^+(\lambda)$ can be uniquely extended to a bounded linear operator from $L_{2,\delta}(\mathbf{R}^N)$ into $H_{-\delta}^2(\mathbf{R}^N)$. The extension will be denoted by $R^+(\lambda)$ again. Then, by use of the denseness of $L_{2,1}(\mathbf{R}^N)$ in $L_{2,\delta}(\mathbf{R}^N)$ and the estimates (4.3) and (4.27), it is easy to see that (4.23) is true for $f \in L_{2,\delta}(\mathbf{R}^N)$. Noting that $u = R^+(\lambda)f$ satisfies the estimate (3.63) as well as (4.27), we can almost repeat the above arguments to show the continuity of $R^+(\lambda)f$ in $H_{-\delta}^2(\mathbf{R}^N)$ with respect to λ for any $f \in L_2(\mathbf{R}^N)$. \square

REFERENCES

- [1] S. Agmon, *Spectral properties of Schrödinger operators and scattering theory*, Ann. Scuola Nor. Sup. Pisa, (4) 2 (1975), 151–218.
- [2] A. Devinatz and P. Rejto, *A limiting absorption principle for Schrödinger operators with oscillating potentials*. I, II, J. Diff. Equations, 49 (1983), 29–84, 85–104.
- [3] M. S. P. Eastham and H. Kalf, *Schrödinger-Type Operators With Continuous Spectra*, Pitman, London, 1982.
- [4] D. Eidus, *The limiting absorption and amplitude principles for the diffraction problem with two unbounded media*, Comm. Math. Phys., 107 (1986), 29–38.
- [5] T. Ikebe and Y. Saitō, *Limiting absorption method and absolute continuity for the Schrödinger operator*, J. Kyoto Univ., 7 (1972), 513–542.
- [6] W. Jäger, *Ein gewöhnlicher Differentialoperator zweiter Ordnung für Funktionen mit Werten in einem Hilbertraum*, Math. Z., 113 (1970), 68–98.
- [7] T. Kato, *Growth properties of solutions of the reduced wave equation with a variable coefficient*, Comm. Pure Appl. Math., 12 (1959), 403–425.
- [8] R. Lavine, *Absolute continuity of positive spectrum for Schrödinger operators with long-range potentials*, J. Funct. Anal., 12 (1973), 30–54.
- [9] K. Mochizuki and J. Uchiyama, *Radiation conditions and spectral theory for 2-body Schrödinger operators with “oscillatory” long-range potentials*, I. J. Math. Kyoto Univ., 18 (1978), 377–408.
- [10] Y. Saitō, *The principle of limiting absorption for second-order differential equations with operator-valued coefficients*, Publ. Res. Inst. Math. Sci. Kyoto Univ., 7 (1972/72), 581–619.
- [11] ———, *Spectral representations for Schrödinger operators with long-range potentials*, Lecture Notes in Mathematics 727, Springer, Berlin-Heidelberg-New York, 1979.

- [12] ———, *Schrödinger operators with nonspherical radiation condition*, *Pacific J. Math.*, **126** (1987), 331–359.

Received August 24, 1987 and in revised form March 18, 1988. Partially supported by NSF EPSCoR grant number RII-8610669.

UNIVERSITY OF ALABAMA AT BIRMINGHAM
BIRMINGHAM, AL 35294