# THE POWER 3/2 APPEARING IN THE ESTIMATE OF ANALYTIC CAPACITY 

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#### Abstract

We show that the power 3/2 appearing in the estimate of analytic capacity is best possible.


1. Introduction. For a compact set $E$ in the complex plane $\mathbb{C}$, $H^{\infty}\left(E^{c}\right)$ denotes the Banach space of bounded analytic functions in $E^{c}=\mathbb{C} \cup\{\infty\}-E$ with supremum norm $\|\cdot\|_{H^{\infty}}$. The analytic capacity of $E$ is defined by

$$
\gamma(E)=\sup \left\{\left|f^{\prime}(\infty)\right| ;\|f\|_{H^{\infty}} \leq 1, f \in H^{\infty}\left(E^{c}\right)\right\}
$$

where $f^{\prime}(\infty)=\lim _{z \rightarrow \infty} z(f(z)-f(\infty))$, i.e., $f^{\prime}(\infty)$ is the $(1 / z)$-coefficient of the Taylor expansion of $f(z)$ at infinity. It is easily seen that $\gamma(E) \leq|E|$, where $|E|$ is the (generalized) length of $E$; if $E$ is a subset of the real line $\mathbb{R}$, then $|E|$ equals its 1 -dimension Lebesgue measure (cf. Garnett [4, Chap. III]). Vitushkin [12] constructed an example $Q_{\infty}$ such that $\gamma\left(Q_{\infty}\right)=0$ and $\left|Q_{\infty}\right|>0$ (cf. [4, p. 87]). Denjoy [3] showed that $\gamma(E)>0$ if $E$ is a subset of a rectifiable curve such that $|E|>0$. But his proof has a serious gap, and his theorem was, for a while, called the Denjoy conjecture. As is easily seen, we may assume that $E$ is a subset of a rectifiable graph. Let pr $E$ denote the projection of $E$ to $\mathbb{R}$. Since pr is a contraction [6, p. 377], it is natural to try the lower estimate of $\gamma(E)$ by $\gamma(\operatorname{pr} E)$. Pommerenke [11] showed that $\gamma(\operatorname{pr} E)=|\operatorname{pr} E| / 4$. Hence this approach is equivalent to comparing $\gamma(E)$ with $|\operatorname{pr} E|$. To do this, the study of the Cauchy-Hilbert transform on $C^{1}$ graphs is necessary (Davie [2]). In 1977, Calderón [1] succeeded in proving its boundedness, and, using his theorem, Marshall [8] finally settled the Denjoy conjecture in the affirmative. After Marshall's theorem, we are concerned with studying further relations between $\gamma(E)$ and $|\operatorname{pr} E|$. Using an estimate of the Cauchy-Hilbert transform on Lipschitz graphs [10, p. 53], the author [9] showed that

$$
\gamma(E) \geq C_{0}|\operatorname{pr} E|^{3 / 2}
$$

if $E$ is a subset of a rectifiable graph $\Gamma$ satisfying $|\Gamma|=1$, where $C_{0}$ is an absolute constant. The main purpose of this paper is to show
that the power $3 / 2$ is best possible. Our method gives a new approach to the computation of analytic capacity, and suggests that analytic capacity is related to the theory of fractals (Mandelbrot [7]).

For an integer $p \geq 2$, we put

$$
B_{p}(x)=\frac{1}{2 p}\left\{1-(-1)^{k}\right\} \quad\left(\frac{k}{p} \leq x<\frac{k+1}{p}, 0 \leq k \leq p-1\right) .
$$

For an $n$-tuple $\left(p_{1}, \ldots, p_{n}\right)$ of integers larger than or equal to 2 , we put

$$
A\left(x ; p_{1}, \ldots, p_{n}\right)=\sum_{j=1}^{n} B_{p_{1} \cdots p_{j}}(x) .
$$

A set $\Gamma \subset \mathbb{C}$ is called a crank of degree $n$ if it is expressed in the form

$$
\Gamma=\Gamma\left(p_{1}, \ldots, p_{n}\right)=\left\{x+i A\left(x ; p_{1}, \ldots, p_{n}\right) ; 0 \leq x<1\right\}
$$

for some $n$-tuple $\left(p_{1}, \ldots, p_{n}\right)$ of integers larger than or equal to 2 . (The class of cranks in this paper is smaller than a class defined in [10, Chap. III].) We shall show

Theorem. For any $n \geq 1$, there exists a crank $\Gamma_{n}$ of degree $n$ such that

$$
\frac{1}{C_{1}} \frac{1}{\sqrt{n}} \leq \gamma\left(\Gamma_{n}\right) \leq C_{1} \frac{1}{\sqrt{n}},
$$

where $C_{1}$ is an absolute constant.
Once this theorem is established, we can deduce the exactness of the power $3 / 2$ as follows. Adding some segments (perpendicular to the $x$ axis) to $\Gamma_{n}$, we obtain an arc connecting 0 and 1 . Then the length of this arc is less than or equal to $n+1$. Hence we can define a rectifiable graph $\Gamma_{n}^{\prime}$ so that $\left|\Gamma_{n}^{\prime}\right| \leq 3 n,\left|\operatorname{pr} E_{n}^{\prime}\right| \geq 1 / 2$, where $E_{n}^{\prime}=\Gamma_{n} \cap \Gamma_{n}^{\prime}$. Then $\gamma\left(E_{n}^{\prime}\right) \leq \gamma\left(\Gamma_{n}\right) \leq C_{1} / \sqrt{n}$. Contracting $E_{n}^{\prime}$, $\Gamma_{n}^{\prime}$, we define $E_{n}^{\prime \prime}$, $\Gamma_{n}^{\prime \prime}$ so that $\left|\Gamma_{n}^{\prime \prime}\right|=1$. Then

$$
\begin{aligned}
\gamma\left(E_{n}^{\prime \prime}\right) & =\gamma\left(E_{n}^{\prime}\right) /\left|\Gamma_{n}^{\prime}\right| \leq \sqrt{3} C_{1}\left|\Gamma_{n}^{\prime}\right|^{-3 / 2} \\
& \leq 2^{3 / 2} \sqrt{3} C_{1}\left\{\left|\operatorname{pr} E_{n}^{\prime}\right| /\left|\Gamma_{n}^{\prime}\right|\right\}^{3 / 2}=2^{3 / 2} \sqrt{3} C_{1}\left|\operatorname{pr} E_{n}^{\prime \prime}\right|^{3 / 2},
\end{aligned}
$$

which shows that the power $3 / 2$ cannot be replaced by any number less than $3 / 2$.

To prove our theorem, it is necessary to investigate cranks carefully. In $\S 2$, we shall give a formula ((1) in Proposition 1) to compute analytic capacity. Proposition 2 is a generalization of Garnett's example [4, p. 87], and will be used to prove our theorem. Using the
method in the proof of the formula, we shall, in $\S 3$, give the proof of our theorem. In the last section, we shall give a new proof of Pommerenke's theorem [11] as another application of Proposition 1; our method shows how to construct the extremal functions.
2. A formula for the computation of $\gamma(\cdot)$. Let $L^{2}(\Gamma)$ denote the $L^{2}$ space of functions on a finite union $\Gamma$ of smooth arcs with respect to the length element $|d z|$. The norm is denoted by $\|\cdot\|_{L^{2}(\Gamma)}$. The Cauchy-Hilbert transform $\mathscr{H}_{\Gamma}$ from $L^{2}(\Gamma)$ to itself is defined by

$$
\mathscr{H}_{\Gamma} f(z)=\frac{1}{\pi} \mathrm{p} \cdot \mathrm{v} \cdot \int_{\Gamma} \frac{f(w)}{w-z}|d w|
$$

where p.v. is the principal value. This is a bounded operator and the norm is denoted by $\left\|\mathscr{H}_{\Gamma}\right\|_{L^{2}(\Gamma), L^{2}(\Gamma)}$. An operator $\overline{\mathscr{H}}_{\Gamma}$ is defined by $\overline{\mathscr{H}}_{\Gamma} f=\overline{\mathscr{H}}_{\Gamma} \bar{f}$, and $\mathscr{F}_{\Gamma}$ is the identity operator. We show

Proposition 1. Let $\Gamma$ be a finite union of smooth arcs. Then, for any $0<\varepsilon<1 /\left\|\mathscr{H}_{\Gamma}\right\|_{L^{2}(\Gamma), L^{2}(\Gamma)}$,

$$
\begin{equation*}
\gamma(\Gamma)=\frac{1}{\pi}\left\{|\Gamma|+\sum_{m=0}^{\infty}\left(1-\varepsilon^{2}\right)^{m} \sum_{l=0}^{\infty} \varepsilon^{2 l+2} \frac{(l+1)(l+2) \cdots(l+m)}{m!} d_{2 l+2}\left(\mathscr{H}_{\Gamma}\right)\right\} \tag{1}
\end{equation*}
$$

where

$$
d_{2 l}\left(\mathscr{H}_{\Gamma}\right)=\int_{\Gamma}\left(\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma}\right)^{l} 1|d z| \quad\left(l \geq 0,\left(\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma}\right)^{0}=\mathscr{F}_{\Gamma}\right)
$$

and $(l+1) \cdots(l+m) / m!=1$ if $m=0$. (First $\sum_{l=0}^{\infty}$ is taken, and next $\sum_{m=0}^{\infty}$ is taken.) If $\left\|\mathscr{I}_{\Gamma}+\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right\|_{L^{2}(\Gamma), L^{2}(\Gamma)}<2$, then

$$
\begin{equation*}
\gamma(\Gamma)=\frac{1}{\pi}\left\{|\Gamma|+\sum_{m=0}^{\infty} 2^{-m-1} \sum_{l=0}^{\infty}\binom{m}{l} d_{2 l+2}\left(\mathscr{H}_{\Gamma}\right)\right\} \tag{2}
\end{equation*}
$$

where $\binom{0}{0}=1$. If $\lim _{l \rightarrow \infty} d_{2 l}\left(\mathscr{H}_{\Gamma}\right)=0$, then

$$
\begin{equation*}
\gamma(\Gamma)=\frac{1}{\pi} \sum_{l=0}^{\infty} d_{2 l}\left(\mathscr{H}_{\Gamma}\right) \tag{3}
\end{equation*}
$$

This is a version of Garabedian's theorem [4, p. 22] to $\mathscr{H}_{\Gamma}$. Equality (3) is applicable to give a new proof of Pommerenke's theorem. (See $\S 4$.) Notice that $\mathscr{J}_{\mathbb{R}}+\overline{\mathscr{H}}_{\mathbb{R}} \mathscr{H}_{\mathbb{R}}=0$, where $\mathscr{H}_{\mathbb{R}}$ is the Hilbert transform on $\mathbb{R}$. Hence (2) is applicable to compact sets $\Gamma$ on a Lipschitz graph which is a small perturbation of $\mathbb{R}$. For any $M>0$, there exists a
crank $\Gamma^{\prime}$ such that $d_{2}\left(\mathscr{F}_{\Gamma^{\prime}}\right) \geq M[\mathbf{1 0}$, p. 84$]$. Then Cauchy-Schwarz' inequality yields that

$$
d_{2^{\prime}}\left(\mathscr{H}_{\Gamma^{\prime}}\right) \geq d_{2}\left(\mathscr{H}_{\Gamma^{\prime}}\right)^{2^{l-1}} \geq M^{2^{l-1}} \quad(l \geq 1) .
$$

Hence (1) is necessary in this case.
Proof of Proposition 1. Let

$$
\begin{equation*}
\gamma^{*}\left(\mathscr{R}_{\Gamma}\right)=\inf \left\{\|1+\mathscr{R} h\|_{L^{2}(\Gamma)}^{2}+\|h\|_{L^{2}(\Gamma)}^{2} ; h \in L^{2}(\Gamma)\right\} . \tag{4}
\end{equation*}
$$

We begin by showing that

$$
\begin{equation*}
\gamma(\Gamma)=\frac{1}{\pi} \gamma^{*}\left(\mathscr{H}_{\Gamma}\right) . \tag{5}
\end{equation*}
$$

For a compact set $E$ bounded by a finite number of smooth Jordan curves, we have
(6) $\gamma(E)=\frac{1}{2 \pi} \inf \left\{\int_{\partial E}|g(z)|^{2}|d z| ; g(\infty)=1, g\right.$ is analytic in $\left.E^{c}\right\}$
[4, p. 22]. Hence a standard argument yields that (6) holds with $E$ replaced by $\Gamma$; in this case, the boundary $\partial \Gamma$ has two sides. We define a smooth curve $\mathscr{L}$ tending to infinity so that $\Gamma \subset \mathscr{L}$ and that $\mathscr{L}=\mathbb{R}$ outside a large disk. Then $\mathscr{L}$ divides $\mathbb{C}$ into two domains $\Omega_{ \pm}$. For an analytic function $g(z)$ in $\Gamma^{c}$ such that $g(\infty)=1$ and $\int_{\partial \Gamma}|g(z)|^{2}|d z|<\infty$, we can write

$$
g(z)=1+\frac{1}{\pi} \int_{\Gamma} \frac{h(w)}{w-z} d w
$$

where the orientation of $d w$ is chosen so that $\Omega_{+}$lies to the left. Let $g_{ \pm}(z)$ be the nontangential limits of $g$ at $z \in \Gamma$ with respect to $\Omega_{ \pm}$, respectively. Then

$$
\begin{aligned}
g_{+}(z) & =1+\frac{1}{\pi} \mathrm{p} \cdot \mathrm{v} \cdot \int_{\Gamma} \frac{h(w)}{w-z} d w+i h(z) \\
& =1+\mathscr{R}_{\Gamma}(h \psi)(z)+i h(z) \quad(z \in \Gamma),
\end{aligned}
$$

where $\psi(z)=d z /|d z|$. Analogously,

$$
g_{-}(z)=1+\mathscr{H}_{\mathbb{F}}(h \psi)(z)-i h(z) \quad(z \in \Gamma) .
$$

Thus

$$
\begin{aligned}
\int_{\partial \Gamma} & |g(z)|^{2}|d z|=\left\|g_{+}\right\|_{L^{2}(\Gamma)}^{2}+\left\|g_{-}\right\|_{L^{2}(\Gamma)}^{2} \\
& =\left\|1+\mathscr{H}_{\Gamma}(h \psi)+i h\right\|_{L^{2}(\Gamma)}^{2}+\left\|1+\mathscr{H}_{\Gamma}(h \psi)-i h\right\|_{L^{2}(\Gamma)}^{2} \\
& =2\left\{\left\|1+\mathscr{H}_{\Gamma}(h \psi)\right\|_{L^{2}(\Gamma)}^{2}+\|h\|_{L^{2}(\Gamma)}^{2}\right\} \\
& =2\left\{\left\|1+\mathscr{H}_{\Gamma}(h \psi)\right\|_{L^{2}(\Gamma)}^{2}+\|h \psi\|_{L^{2}(\Gamma)}^{2}\right\}
\end{aligned}
$$

because $|\psi(z)|=1 \quad(z \in \Gamma)$. This shows that the quantity in the right-hand side of (6) $(E=\Gamma)$ equals $\frac{1}{\pi} \gamma^{*}\left(\mathscr{H}_{\Gamma}\right)$, i.e., (5) holds.

We next compute $\gamma^{*}\left(\mathscr{H}_{\Gamma}\right)$. Fatou's lemma shows that there exists $h_{\Gamma} \in L^{2}(\Gamma)$ which attains the infimum in (4). A variational method yields that $\left(1+\mathscr{H}_{\Gamma} h_{\Gamma}, \mathscr{H}_{\Gamma} h\right)+\left(h_{\Gamma}, h\right)=0$ for all $h \in L^{2}(\Gamma)$, where $(\cdot, \cdot)$ is the (complex) inner product with respect to $|d z|$. Since the adjoint operator of $\mathscr{H}_{\Gamma}$ is $-\overline{\mathscr{H}}_{\Gamma}$, this shows that

$$
\begin{equation*}
\left(\mathscr{C}_{\Gamma}-\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right) h_{\Gamma}=\overline{\mathscr{H}}_{\Gamma} 1 . \tag{7}
\end{equation*}
$$

Suppose that $h_{\Gamma}^{\prime} \in L^{2}(\Gamma)$ also attains the infimum in (4). Then $h_{\Gamma}^{\prime}$ satisfies (7), and hence

$$
\begin{aligned}
0 & =\left(\left(\mathscr{F}_{\Gamma}-\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right)\left(h_{\Gamma}-h_{\Gamma}^{\prime}\right), h_{\Gamma}-h_{\Gamma}^{\prime}\right) \\
& =\left\|h_{\Gamma}-h_{\Gamma}^{\prime}\right\|_{L^{2}(\Gamma)}^{2}+\left\|\mathscr{H}_{\Gamma}\left(h_{\Gamma}-h_{\Gamma}^{\prime}\right)\right\|_{L^{2}(\Gamma)}^{2} .
\end{aligned}
$$

This shows that $h_{\Gamma}^{\prime}=h_{\Gamma}$. Thus $h_{\Gamma}$ is uniquely determined. By (7), we have

$$
\begin{align*}
\gamma^{*}\left(\mathscr{H}_{\Gamma}\right) & =\left\|1+\mathscr{H}_{\Gamma} h_{\Gamma}\right\|_{L^{2}(\Gamma)}^{2}+\left\|h_{\Gamma}\right\|_{L^{2}(\Gamma)}^{2}  \tag{8}\\
& =\left(1+\mathscr{H}_{\Gamma} h_{\Gamma}, 1\right)+\left(\left(\mathscr{K}_{\Gamma}-\mathscr{\mathscr { H }}_{\Gamma} \mathscr{H}_{\Gamma}\right) h_{\Gamma}-\overline{\mathscr{H}}_{\Gamma} 1, h_{\Gamma}\right) \\
& =\int_{\Gamma}\left\{1+\mathscr{H}_{\Gamma} h_{\Gamma}\right\}|d z| .
\end{align*}
$$

Let

$$
T_{\Gamma}=\left(\mathscr{F}_{\Gamma}-\varepsilon^{2} \overline{\mathscr{H}}_{\Gamma} \mathscr{R}_{\Gamma}\right)^{-1} .
$$

Then we can write

$$
T_{\Gamma}=\sum_{l=0}^{\infty} \varepsilon^{2 l}\left(\overline{\mathscr{H}}_{\Gamma} \mathscr{A}_{\Gamma}\right)^{l}
$$

because $0<\varepsilon<1 /\left\|\mathscr{H}_{\Gamma}\right\|_{L^{2}(\Gamma), L^{2}(\Gamma)}$. We have, for any $h \in L^{2}(\Gamma)$,

$$
\begin{aligned}
& \left\|T_{\Gamma} h\right\|_{L^{2}(\Gamma)}^{2} \leq\left\|T_{\Gamma} h\right\|_{L^{2}(\Gamma)}^{2}+\varepsilon^{2}\left\|\mathscr{R}_{\Gamma} T_{\Gamma} h\right\|_{L^{2}(\Gamma)}^{2} \\
& \quad=\left(\left(\mathscr{C}_{\Gamma}-\varepsilon^{2} \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right) T_{\Gamma} h, T_{\Gamma} h\right)=\left(h, T_{\Gamma} h\right) \leq\|h\|_{L^{2}(\Gamma)}\left\|T_{\Gamma} h\right\|_{L^{2}(\Gamma)},
\end{aligned}
$$

which shows that $\left\|T_{\Gamma}\right\|_{L^{2}(\Gamma), L^{2}(\Gamma)} \leq 1$. Equality (7) can be rewritten as

$$
\begin{equation*}
\left(\mathscr{F}-\varepsilon^{2} \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right) h_{\Gamma}=\left(1-\varepsilon^{2}\right) h_{\Gamma}+\varepsilon^{2} \overline{\mathscr{H}}_{\Gamma} 1 . \tag{9}
\end{equation*}
$$

Observing this equality, we inductively define $\left(h_{m}\right)_{m=0}^{\infty}$ by $h_{0}=0$,

$$
h_{m}=T_{\Gamma}\left\{\left(1-\varepsilon^{2}\right) h_{m-1}+\varepsilon^{2} \overline{\mathscr{H}}_{\Gamma} 1\right\} \quad(m \geq 1) .
$$

Then

$$
\begin{aligned}
\left\|h_{m+1}-h_{m}\right\|_{L^{2}(\Gamma)} & =\left(1-\varepsilon^{2}\right)\left\|T_{\Gamma}\left(h_{m}-h_{m-1}\right)\right\|_{L^{2}(\Gamma)} \\
& \leq\left(1-\varepsilon^{2}\right)\left\|h_{m}-h_{m-1}\right\|_{L^{2}(\Gamma)}
\end{aligned}
$$

Hence $\lim _{m \rightarrow \infty} h_{m}$ exists and satisfies (9), i.e., (7). Thus $h_{\Gamma}=$ $\lim _{m \rightarrow \infty} h_{m}$. Since

$$
\begin{aligned}
h_{m+1}-h_{m} & =\left(1-\varepsilon^{2}\right) T_{\Gamma}\left(h_{m}-h_{m-1}\right)=\cdots=\left(1-\varepsilon^{2}\right)^{m} T_{\Gamma}^{m} h_{1} \\
& =\varepsilon^{2}\left(1-\varepsilon^{2}\right)^{m} T_{\Gamma}^{m+1 \overline{\mathscr{H}}_{\Gamma}} 1
\end{aligned}
$$

we have

$$
h_{\Gamma}=\sum_{m=0}^{\infty}\left(h_{m+1}-h_{m}\right)=\varepsilon^{2} \sum_{m=0}^{\infty}\left(1-\varepsilon^{2}\right)^{m} T_{\Gamma}^{m+1} \overline{\mathscr{H}}_{\Gamma} 1
$$

Consequently, (8) yields that

$$
\begin{aligned}
& \gamma^{*}\left(\mathscr{H}_{\Gamma}\right)=|\Gamma|+\sum_{m=0}^{\infty}\left(1-\varepsilon^{2}\right)^{m} \int_{\Gamma} \varepsilon^{2} \mathscr{\mathscr { H }}_{\Gamma} T_{\Gamma}^{m+1} \overline{\mathscr{H}}_{\Gamma} 1|d z| \\
& =|\Gamma|+\sum_{m=0}^{\infty}\left(1-\varepsilon^{2}\right)^{m} \int_{\Gamma} \varepsilon^{2} \mathscr{H}_{\Gamma}\left\{\sum_{l=0}^{\infty} \varepsilon^{2 l}\left(\overline{\mathscr{H}}_{\Gamma} \mathscr{\mathscr { H }}_{\Gamma}\right)^{l}\right\}^{m+1} \overline{\mathscr{H}}_{\Gamma} 1|d z| \\
& =|\Gamma|+\sum_{m=0}^{\infty}\left(1-\varepsilon^{2}\right)^{m} \sum_{l=0}^{\infty} \varepsilon^{2 l+2} \frac{(l+1) \cdots(l+m)}{m!} \int_{\Gamma}\left(\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma}\right)^{l+1} 1|d z| \\
& =|\Gamma|+\sum_{m=0}^{\infty}\left(1-\varepsilon^{2}\right)^{m} \sum_{l=0}^{\infty} \varepsilon^{2 l+2} \frac{(l+1) \cdots(l+m)}{m!} d_{2 l+2}\left(\mathscr{H}_{\Gamma}\right) .
\end{aligned}
$$

Using (5), we obtain (1).
We can write

$$
\mathscr{\mathscr { I }}_{\Gamma}-\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}=2\left\{\mathscr{F}_{\Gamma}-\frac{1}{2}\left(\mathscr{F}_{\Gamma}+\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right)\right\} .
$$

Hence, if $\left\|\mathscr{C}_{\Gamma}+\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right\|_{L^{2}(\Gamma), L^{2}(\Gamma)}<2$, then

$$
h_{\Gamma}=\sum_{m=0}^{\infty} 2^{-m}\left(\mathscr{C}+\overline{\mathscr{H}}_{\Gamma} \mathscr{\mathscr { H }}_{\Gamma}\right)^{m}\left(\frac{1}{2} \overline{\mathscr{H}}_{\Gamma}\right) 1
$$

Thus (5) and (8) yield (2).
Equality (7) shows that $\mathscr{H}_{\Gamma} h_{\Gamma}=\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} 1+\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma} h_{\Gamma}$, and hence, by (8),

$$
\gamma^{*}\left(\mathscr{H}_{\Gamma}\right)=\int_{\Gamma}\left\{1+\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} 1+\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma} h_{\Gamma}\right\}|d z| .
$$

Repeating this argument, we have

$$
\begin{aligned}
\gamma^{*}\left(\mathscr{H}_{\Gamma}\right) & =\int_{\Gamma}\left\{\sum_{l=0}^{L}\left(\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma}\right)^{l} 1+\mathscr{\mathscr { H }}_{\Gamma}\left(\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}\right)^{L} h_{\Gamma}\right\}|d z| \\
& =\sum_{l=0}^{L} d_{2 l}\left(\mathscr{H}_{\Gamma}\right)-\int_{\Gamma}\left\{\left(\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma}\right)^{L} \mathscr{H}_{\Gamma}\right\} 1(z) h_{\Gamma}(z)|d z| .
\end{aligned}
$$

If $\lim _{L \rightarrow \infty} d_{2 L}\left(\mathscr{F}_{\mathrm{F}}\right)=0$, then

$$
\begin{aligned}
& \lim _{L \rightarrow \infty}\left|\int_{\Gamma}\left\{\left(\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma}\right)^{L} \mathscr{H}_{\mathcal{F}}\right\} 1(z) h_{\Gamma}(z)\right| d z \| \mid \\
& \quad \leq \lim _{L \rightarrow \infty}\left\|\left(\mathscr{F}_{\Gamma} \overline{\mathscr{H}}_{\Gamma}\right)^{L} \mathscr{H}_{\Gamma} 1\right\|_{L^{2}(\Gamma)}\left\|h_{\Gamma}\right\|_{L^{2}(\Gamma)} \\
& \quad=\lim _{L \rightarrow \infty} d_{4 L+2}\left(\mathscr{H}_{\Gamma}\right)^{1 / 2}\left\|h_{\Gamma}\right\|_{L^{2}(\Gamma)}=0 .
\end{aligned}
$$

Hence (5) gives (3). This completes the proof of our proposition.
We now give a remark. There exists an analytic function $g_{\Gamma}(z)$ in $\Gamma^{c}$ such that $g_{\Gamma}(\infty)=1$ and $\gamma(\Gamma)=(1 / 2 \pi) \int_{\partial \Gamma}\left|g_{\Gamma}(z)\right||d z|[4$, p. 19]. This is called the Garabedian function of $\Gamma$. Equality (5) shows that

$$
g_{\Gamma}(z)=\left\{1+\frac{1}{\pi} \int_{\Gamma} \frac{h_{\Gamma}(w)}{w-z}|d w|\right\}^{2}
$$

There exists $f_{\Gamma} \in H^{\infty}\left(\Gamma^{c}\right)$ such that $\left\|f_{\Gamma}\right\|_{H^{\infty}}=1$ and $f_{\Gamma}^{\prime}(\infty)=\gamma(\Gamma)$ [4, p. 18]. This is called the Ahlfors function of $\Gamma$. We have

$$
\begin{equation*}
f_{\Gamma}(z)=\frac{-\frac{1}{\pi}\left\{\int_{\Gamma} \frac{|d w|}{w-z}+\int_{\Gamma} \frac{\overline{\mathscr{H}}_{\Gamma} \bar{h}_{\Gamma}(w)}{w-z}|d w|\right\}}{\left\{1+\frac{1}{\pi} \int_{\Gamma} \frac{h_{\Gamma}(w)}{w-z}|d w|\right\}} \tag{10}
\end{equation*}
$$

To see this, let $f(z)$ denote the function in the right-hand side. Since $g_{\Gamma}(z)$ does not take 0 in $\Gamma^{c}, f(z)$ is analytic in $\Gamma^{c}[4, \mathrm{p} .21]$. We have $f^{\prime}(\infty)=\frac{1}{\pi} \gamma^{*}\left(\mathscr{H}_{\Gamma}\right)=\gamma(\Gamma)$ and

$$
f_{ \pm}(z)=-\frac{\mathscr{H}_{\Gamma} 1(z) \pm i \tilde{\psi}(z)+\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} \bar{h}_{\Gamma}(z) \pm i \overline{\mathscr{H}}_{\Gamma} \bar{h}_{\Gamma}(z) \tilde{\psi}(z)}{1+\mathscr{H}_{\Gamma} h_{\Gamma}(z) \pm i h_{\Gamma}(z) \tilde{\psi}(z)},
$$

where $\tilde{\psi}(z)=|d z| / d z$ and $f_{ \pm}(z)$ are the nontangential limits of $f$ at $z \in \Gamma$ with respect to $\Omega_{ \pm}$, respectively. Equality (7) shows that

$$
\begin{aligned}
\mathscr{H}_{\Gamma} 1+ & i \tilde{\psi}+\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} \bar{h}_{\Gamma}+i\left(\overline{\mathscr{H}}_{\Gamma} \bar{h}_{\Gamma}\right) \tilde{\psi} \\
& =\mathscr{H}_{\Gamma} 1+i \tilde{\psi}+\left(\bar{h}_{\Gamma}-\mathscr{H}_{1} 1\right)+i\left(\overline{\mathscr{H}}_{\Gamma} \bar{h}_{\Gamma}\right) \tilde{\psi} \\
& =i \tilde{\psi}+i\left(\overline{\mathscr{H}}_{\Gamma} \bar{h}_{\Gamma}\right) \tilde{\psi}+\bar{h}_{\Gamma}=i \tilde{\psi}\left\{\overline{1+\mathscr{H}_{\Gamma} h_{\Gamma}+i h_{\Gamma} \tilde{\psi}}\right\},
\end{aligned}
$$

which yields that $\left|f_{+}(z)\right|=1$ on $\Gamma$. Analogously, $\left|f_{-}(z)\right|=1$ on $\Gamma$. Thus $\|f\|_{H^{\infty}}=1$. This shows that $f=f_{\Gamma}$.

For the proof of our theorem, we note

Proposition 2. Let $0<\delta_{0}<1$ and let $\left(q_{n}\right)_{n=1}^{\infty}$ be a sequence of integers larger than or equal to 2 such that

$$
\sum_{n=j}^{\infty}\left(q_{j} \cdots q_{n}\right)^{-1} \leq \delta_{0} \quad(j \geq 1)
$$

Then

$$
\lim _{n \rightarrow \infty} \sup \gamma\left(\Gamma\left(p_{1}, \ldots, p_{n}\right)\right)=0
$$

where the supremum is taken over all n-tuples $\left(p_{1}, \cdots, p_{n}\right)$ satisfying $p_{j} \geq q_{j}(1 \leq j \leq n)$.

This is a generalization of Garnett's example [4, p. 87], and used later. Notice that $\sum_{n=1}^{\infty} 2^{-n}=1$. A sequence $\left(\Gamma\left(\mathbf{2}_{n}\right)\right)_{n=1}^{\infty}\left(\mathbf{2}_{n}\right.$ is the $n$ tuple of 2 ) topologically converges to a segment $\{x+i x ; 0 \leq x<1\}$, and these cranks behave like cranks of degree 1 with respect to this segment. Hence we have $\lim \sup _{n \rightarrow \infty} \gamma\left(\Gamma\left(2_{n}\right)\right)>0$. This shows that our proposition is sharp in a sense. Since a minor change of the argument in [10, p. 81] yields the required equality, we omit the proof (cf. Jones [5]).
3. Proof of Theorem. In this section, we give the proof of our theorem. Let $L^{q}$ denote the $L^{q}$ space of functions on $[0,1)$ with respect to the 1 -dimension Lebesgue measure $|\cdot|(1 \leq q<\infty)$. For a kernel $K=K(x, y)$ on $[0,1) \times[0,1)$, we simply write by the same notation $K$ an operator defined by this kernel, and write by $\bar{K}$ an operator defined by $\overline{K(x, y)} ;\|K\|_{L^{q}, L^{q^{\prime}}}$ denotes the norm of $K$ as an operator from $L^{q}$ to $L^{q^{\prime}}$. The identity operator is denoted by Id. A kernel $K$ is anti-symmetric if $K(x, y)=-K(y, x) \quad(x \neq y)$. A kernel $K$ is of type 0 if

$$
\sup _{x, y \in[0,1)}\left\{|K(x, y)|+\left|\frac{\partial}{\partial x} K(x, y)\right|+\left|\frac{\partial}{\partial y} K(x, y)\right|\right\}<\infty
$$

A kernel $K$ is of type 1 if $\|K\|_{L^{4}, L^{4}}<\infty$ and if there exists a sequence $\left(K_{j}\right)_{j=1}^{\infty}$ of kernels of type 0 such that

$$
\lim _{j \rightarrow \infty}\left\|K_{j}-K\right\|_{L^{4}, L^{2}}=0, \quad \sup _{j \geq 1}\left\|K_{j}\right\|_{L^{4}, L^{4}}<\infty
$$

Kernels used in this section are bounded as operators from $L^{q}$ to itself for all $1<q<\infty$. Let

$$
\begin{aligned}
\gamma^{*}(K) & =\inf \left\{\|1+K h\|_{L^{2}}^{2}+\|h\|_{L^{2}}^{2} ; h \in L^{2}\right\} \\
d_{2 l}(K) & =\int_{0}^{1}(K \bar{K})^{l} 1 d x \quad\left(l \geq 0,(K \bar{K})^{0}=\mathrm{Id}\right) .
\end{aligned}
$$

Recall the function $A\left(x ; p_{1}, \ldots, p_{n}\right)$ in the introduction. Let

$$
H(x, y)=\mathscr{H}_{\mathbb{R}}(x, y)=\frac{1}{\pi} \frac{1}{y-x},
$$

$$
\begin{aligned}
& H\left[p_{1}, \ldots, p_{n}\right](x, y) \\
& \quad=\frac{1}{\pi} \frac{1}{(y-x)+i\left(A\left(y ; p_{1}, \ldots, p_{n}\right)-A\left(x ; p_{1}, \ldots, p_{n}\right)\right)}, \\
& \Delta\left[p_{1}, \ldots, p_{n}\right]=H\left[p_{1}, \ldots, p_{n}\right]-H\left[p_{1}, \ldots, p_{n-1}\right] \quad(n \geq 1),
\end{aligned}
$$

where $H\left[p_{1}, \ldots, p_{n-1}\right]=H$ if $n=1$. Then

$$
H\left[p_{1}, \ldots, p_{n}\right]=H+\sum_{j=1}^{n} \Delta\left[p_{1}, \ldots, p_{j}\right]
$$

Since all components/segments of $\Gamma\left(p_{1}, \cdots, p_{n}\right)$ are parallel to the $x$-axis, we can identify $\mathscr{H}_{\Gamma\left(p_{1}, \ldots, p_{n}\right)}, L^{2}\left(\Gamma\left(p_{1}, \ldots, p_{n}\right)\right)$ with $H\left[p_{1}, \ldots, p_{n}\right], L^{2}$, respectively. We have $\left\|H\left[p_{1}, \ldots, p_{n}\right]\right\|_{L^{2}, L^{2}} \leq$ $C_{2} \sqrt{n}$ for some absolute constant $C_{2}[\mathbf{1 0}$, p. 84]. Hence Proposition 1 shows that

$$
\begin{align*}
& \gamma\left(\Gamma\left(p_{1}, \ldots, p_{n}\right)\right)=\frac{1}{\pi} \gamma^{*}\left(H\left[p_{1}, \ldots, p_{n}\right]\right)  \tag{11}\\
& =\frac{1}{\pi}\left\{1+\sum_{m=0}^{\infty}\left(1-\varepsilon_{n}^{2}\right)^{m}\right. \\
& \left.\quad \times \sum_{l=0}^{\infty} \varepsilon_{n}^{2 l+2} \frac{(l+1) \cdots(l+m)}{m!} d_{2 l+2}\left(H\left[p_{1}, \ldots, p_{n}\right]\right)\right\},
\end{align*}
$$

where $\varepsilon_{n}=\left(2 C_{2} \sqrt{n}\right)^{-1}$. We shall inductively estimate

$$
\lim _{p_{1} \rightarrow \infty} \cdots \lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(H\left[p_{1}, \ldots, p_{n}\right]\right)
$$

where $\lim _{p_{n} \rightarrow \infty}$ is taken first and $\lim _{p_{1} \rightarrow \infty}$ is taken last. For $E \subset \mathbb{R}$, $\chi_{E}$ denotes its characteristic function, and, for $x \in \mathbb{R}, l(x)$ denotes its integral part. Here are some lemmas necessary for the estimate.

Lemma 3. For two kernels $K$ and $K^{\prime}$,

$$
\gamma^{*}\left(K+K^{\prime}\right) \leq 2\left(1+\left\|K^{\prime}\right\|_{L^{2}, L^{2}}^{2}\right) \gamma^{*}(K) .
$$

Proof. We have, for any $h \in L^{2}$,

$$
\left\|1+\left(K+K^{\prime}\right) h\right\|_{L^{2}}^{2}+\|h\|_{L^{2}}^{2} \leq 2\left(1+\left\|K^{\prime}\right\|_{L^{2}, L^{2}}^{2}\right)\left\{\|1+K h\|_{L^{2}}^{2}+\|h\|_{L^{2}}^{2}\right\},
$$ which yields the required inequality.

Lemma 4. Let $K$ be an anti-symmetric kernel such that

$$
\lim _{l \rightarrow \infty} d_{2 l}(K)=0
$$

Then

$$
\begin{equation*}
\gamma^{*}(K)=\sum_{l=0}^{\infty} d_{2 l}(K) . \tag{12}
\end{equation*}
$$

Since this is a version of (3) to $K$, we omit the proof.
Lemma 5. For an anti-symmetric kernel $K, 0<\varepsilon_{0} \leq\left(3\|K\|_{L^{2}, L^{2}}\right)^{-1}$ and $w \in U=\{\zeta \in \mathbb{C} ;|\zeta|<2,|\arg \zeta|<\pi / 4\}$,

$$
\begin{array}{r}
1+\sum_{m=0}^{\infty}\left(1-\varepsilon_{0}^{2}\right)^{m} \sum_{l=0}^{\infty} w^{2 l+2} \varepsilon_{0}^{2 l+2} \frac{(l+1) \cdots(l+m)}{m!} d_{2 l+2}(K)  \tag{13}\\
\quad\left(=\gamma^{*}(w ; K), \text { say }\right)
\end{array}
$$

exists and $\gamma^{*}(w ; K)$ is analytic in $U$.
Proof. Let

$$
T(w ; K)=\left(\operatorname{Id}-w^{2} \varepsilon_{0}^{2} \bar{K} K\right)^{-1} .
$$

Then

$$
\begin{aligned}
& w^{2} \varepsilon_{0}^{2} \int_{0}^{1} K T(w ; K)^{m+1} \bar{K} 1 d x \\
& \quad=\sum_{l=0}^{\infty} w^{2 l+2} \varepsilon_{0}^{2 l+2} \frac{(l+1) \cdots(l+m)}{m!} d_{2 l+2}(K)
\end{aligned}
$$

because $2 \varepsilon_{0}\|K\|_{L^{2}, L^{2}}<1$. Evidently, this is analytic in $U$. Since $K$ is anti-symmetric and $\operatorname{Re} w^{2}>0 \quad\left(w \in U, \operatorname{Re} w^{2}\right.$ is the real part of $w^{2}$ ), we have, in the same manner as in the proof of (1),

$$
\begin{equation*}
\|T(w ; K)\|_{L^{2}, L^{2}} \leq 1 \quad(w \in U) \tag{14}
\end{equation*}
$$

Thus the convergence of $\sum_{m=0}^{\infty}$ in (13) is uniform in $U$, which shows that $\gamma^{*}(w ; K)$ exists and is analytic in $U$.

Lemma 6. For any $l \geq 0$,

$$
\lim _{p \rightarrow \infty} d_{2 l}(\Delta[p]) \quad\left(=d_{2 l}(\Delta[\infty]), \text { say }\right)
$$

exists and

$$
\begin{equation*}
d_{2}(\Delta[\infty]) \leq-\frac{1}{25 \pi^{2}} \tag{15}
\end{equation*}
$$

Proof. We put

$$
\begin{aligned}
R(s, t)= & \frac{1}{\pi}\left\{\frac{1}{t-s+1+i}-\frac{1}{t-s+1}\right\} \\
& -\frac{2}{\pi} \sum_{m=1}^{\infty}\left\{\frac{t-s+1+i}{4 m^{2}-(t-s+1+i)^{2}}-\frac{t-s+1}{4 m^{2}-(t-s+1)^{2}}\right\}
\end{aligned}
$$

and show that

$$
\begin{equation*}
\lim _{p \rightarrow \infty} d_{2 l}(\Delta[p])=\frac{1}{2} \int_{0}^{1}\left\{R^{2 l} 1+\bar{R}^{2 l} 1\right\} d s \tag{16}
\end{equation*}
$$

Let

$$
\begin{gathered}
W_{p}=\bigcup_{m, \mathrm{odd}}\left[\frac{m}{p}, \frac{m+1}{p}\right), \quad W_{p}^{\prime}=\bigcup_{m, \text { even }}\left[\frac{m}{p}, \frac{m+1}{p}\right), \\
X_{p}=\bigcup_{m=l(\log p)}^{p-l(\log p)-1}\left[\frac{m}{p}, \frac{m+1}{p}\right), \\
s_{x}=p x-l(p x) \quad(0 \leq x<1, \quad p \geq 2) .
\end{gathered}
$$

Notice that $\left|[0,1)-X_{p}\right| \leq 2 l(\log p) / p$ and $\|\Delta[p]\|_{L^{4}, L^{4}} \leq 10$. Since $\Delta[p](x, y)=0 \quad\left(x, y \in W_{p} ; x, y \in W_{p}^{\prime}\right)$, we have

$$
\begin{aligned}
d_{2 l}(\Delta[p])= & \int_{0}^{1}(\Delta[p] \overline{\Delta[p]})^{l-1} \Delta[p]\left\{\chi_{W_{p}} \overline{\Delta[p]} \chi_{W_{p}^{\prime}}+\chi_{W_{p}^{\prime}} \overline{\Delta[p]} \chi_{W_{p}}\right\} d x \\
= & \int_{0}^{1}(\Delta[p] \overline{\Delta[p]})^{l-1} \Delta[p]\left\{\chi_{W_{p} \cap X_{p}} \overline{\Delta[p]} \chi_{W_{p}^{\prime}}+\chi_{W_{p}^{\prime} \cap X_{p}} \overline{\Delta[p]} \chi_{W_{p}}\right\} d x \\
& +O\left(\left(\frac{\log p}{p}\right)^{1 / 4}\right) .
\end{aligned}
$$

We now study $\overline{\Delta[p]} \chi_{W_{p}^{\prime}}(x) \quad\left(x \in W_{p} \cap X_{p}\right)$. Without loss of generality, we may assume that $p$ is even. Since $x \in W_{p} \cap X_{p}, l(p x)$ is even and
$l(\log p) \leq l(p x) \leq p-l(\log p)-1$. We may assume that $l(\log p) \leq$ $l(p x) \leq p / 2$. We have

$$
\begin{aligned}
& \overline{\Delta[p]} \chi_{W_{p}^{\prime}}(x)=\frac{1}{\pi} \int_{W_{p}^{\prime}}\left\{\frac{1}{y-x-i / p}-\frac{1}{y-x}\right\} d y \\
& =\frac{1}{\pi} \int_{W_{p}}\left\{\frac{1}{y-x+1 / p-i / p}-\frac{1}{y-x+1 / p}\right\} d y \\
& =\frac{1}{\pi} \sum_{m=0}^{(p / 2)-1} \int_{0}^{1 / p}\left\{\frac{1}{(2 m / p+y)-(l(p x) / p+x-l(p x) / p)+1 / p-i / p}\right. \\
& \left.\quad-\frac{1}{(2 m / p+y)-(l(p x) / p+x-l(p x) / p)+1 / p}\right\} d y \\
& = \\
& =\frac{1}{\pi} \sum_{m=0}^{l(p x)}+\frac{1}{\pi} \sum_{m=l(p x)+1}^{(p / 2)-1}=L_{1}+L_{2},
\end{aligned}
$$

$$
\begin{aligned}
L_{1}= & \frac{1}{\pi} \int_{0}^{1}\left\{\frac{1}{t-s_{x}+1-i}-\frac{1}{t-s_{x}+1}\right\} d t+\frac{1}{\pi} \sum_{0 \leq m \leq l(p x), m \neq l(p x) / 2} \\
= & \frac{1}{\pi} \int_{0}^{1}\left\{\frac{1}{t-s_{x}+1-i}-\frac{1}{t-s_{x}+1}\right\} d t \\
& -\frac{2}{\pi} \sum_{m=1}^{l(p x) / 2} \int_{0}^{1}\left\{\frac{t-s_{x}+1-i}{4 m^{2}-\left(t-s_{x}+1-i\right)^{2}}\right. \\
& \left.-\frac{t-s_{x}+1}{4 m^{2}-\left(t-s_{x}+1\right)^{2}}\right\} d t
\end{aligned}
$$

$$
=\bar{R} 1\left(s_{x}\right)+O\left(\frac{1}{\log p}\right),
$$

$$
L_{2}=-\frac{i}{\pi} \sum_{m=l(p x)+1}^{(p / 2)-1} \int_{0}^{1} \frac{1}{(2 m-l(p x))+\left(t-s_{x}+1-i\right)}
$$

$$
\times \frac{d t}{(2 m-l(p x))+\left(t-s_{x}+1\right)}=O\left(\frac{1}{\log p}\right),
$$

which shows that $\overline{\Delta[p]} \chi_{W_{p}^{\prime}}(x)=\bar{R} 1\left(s_{x}\right)+O(1 / \log p) \quad\left(x \in W_{p} \cap X_{p}\right)$.

In the same manner, $\overline{\Delta[p]} \chi_{W_{p}}(x)=R 1\left(s_{x}\right)+O(1 / \log p) \quad\left(x \in W_{p}^{\prime} \cap\right.$ $X_{p}$ ). Thus

$$
\begin{aligned}
d_{2 l}(\Delta[p])= & \int_{0}^{1}(\Delta[p] \overline{\Delta[p]})^{l-1} \Delta[p]\left\{\chi_{W_{p} \cap X_{p}} \bar{R} 1\left(s_{.}\right)+\chi_{W_{p}^{\prime} \cap X_{p}} R 1\left(s_{.}\right)\right\} d x \\
& +O\left(\frac{1}{\log p}\right) \\
= & \int_{0}^{1}(\Delta[p] \overline{\Delta[p]})^{l-1} \Delta[p]\left\{\chi_{W_{p}} \bar{R} 1\left(s_{.}\right)+\chi_{W_{p}^{\prime}} R 1\left(s_{.}\right)\right\} d x \\
& +O\left(\frac{1}{\log p}\right) \\
= & \int_{0}^{1}(\Delta[p] \overline{\Delta[p]})^{l-1}\left\{\chi_{W_{p} \cap X_{p}} \Delta[p]\left(\chi_{W_{p}^{\prime}} R 1\left(s_{.}\right)\right)\right. \\
& \left.+\chi_{W_{p}^{\prime} \cap X_{p}} \Delta[p]\left(\chi_{W_{p}} \bar{R} 1\left(s_{.}\right)\right)\right\} d x \\
& +O\left(\frac{1}{\log p}\right) .
\end{aligned}
$$

Since $R 1\left(s_{x}\right)$ is a periodic function with period $1 / p$, we have, in the same manner as above,

$$
\begin{array}{ll}
\Delta[p]\left(\chi_{W_{p}^{\prime}} R 1\left(s_{.}\right)\right)(x)=R^{2} 1\left(s_{x}\right)+O\left(\frac{1}{\log p}\right) & \left(x \in W_{p} \cap X_{p}\right), \\
\Delta[p]\left(\chi_{W_{p}} \bar{R} 1\left(s_{.}\right)\right)(x)=\bar{R}^{2} 1\left(s_{x}\right)+O\left(\frac{1}{\log p}\right) & \left(x \in W_{p}^{\prime} \cap X_{p}\right) .
\end{array}
$$

Repeating this argument, we have

$$
\begin{aligned}
d_{2 l}(\Delta[p]) & =\int_{0}^{1}\left\{\chi_{W_{p}}(x) R^{2 l} 1\left(s_{x}\right)+\chi_{W_{p}^{\prime}}(x) \bar{R}^{2 l} 1\left(s_{x}\right)\right\} d x+O\left(\frac{1}{\log p}\right) \\
& =\frac{1}{2} \int_{0}^{1}\left\{R^{2 l} 1+\bar{R}^{2 l} 1\right\} d s+O\left(\frac{1}{\log p}\right),
\end{aligned}
$$

which gives (16).
We have

$$
\begin{aligned}
R(s, t)= & \frac{1}{\pi} \sum_{m=-\infty}^{\infty}\left\{\frac{1}{2 m+1+t-s+i}-\frac{1}{2 m+1+t-s}\right\} \\
= & -\frac{1}{\pi} \sum_{m=-\infty}^{\infty} \frac{1}{(2 m+1+t-s)\left\{(2 m+1+t-s)^{2}+1\right\}} \\
& -\frac{i}{\pi} \sum_{m=-\infty}^{\infty} \frac{1}{(2 m+1+t-s)^{2}+1} \\
= & -R^{\prime}(s, t)-i R^{\prime \prime}(s, t), \text { say. }
\end{aligned}
$$

Then $R^{\prime}$ is anti-symmetric and $R^{\prime \prime}$ is symmetric, i.e., $R^{\prime \prime}(s, t)=$ $R^{\prime \prime}(t, s)$. Thus

$$
\begin{aligned}
d_{2}(\Delta[\infty]) & =\operatorname{Re} \int_{0}^{1} R^{2} 1 d s \\
& =\operatorname{Re} \int_{0}^{1}\left(-R^{\prime} 1+i R^{\prime \prime} 1\right)\left(R^{\prime} 1+i R^{\prime \prime} 1\right) d s \\
& =-\int_{0}^{1}\left\{\left(R^{\prime} 1\right)^{2}+\left(R^{\prime \prime} 1\right)^{2}\right\} d s \\
& \leq-\frac{1}{\pi^{2}} \int_{0}^{1}\left\{\int_{0}^{1} \frac{d t}{(1+t-s)^{2}+1}\right\}^{2} d s \leq-\frac{1}{25 \pi^{2}} .
\end{aligned}
$$

Thus (15) holds.
Lemma 7. Let $K$ be an anti-symmetric kernel of type 1 , and let $\left(g_{p}\right)_{p=2}^{\infty},\left(h_{p}\right)_{p=2}^{\infty}$ be two sequences in $L^{4}$ such that $\left\|g_{p}\right\|_{L^{4}} \leq 1$, $\left\|h_{p}\right\|_{L^{4}} \leq 1$. Then, for any $l \geq 0$,
(17) $\lim _{p \rightarrow \infty}\left\{\int_{0}^{1} K g_{p} \cdot(\Delta[p] \overline{\Delta[p]}) / \bar{K} h_{p} d x\right.$

$$
\left.-d_{2 l}(\Delta[\infty]) \int_{0}^{1} K g_{p} \cdot \bar{K} h_{p} d x\right\}=0
$$

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \int_{0}^{1} K g_{p} \cdot \overline{\Delta[p]}(\Delta[p] \overline{\Delta[p]})^{l} K h_{p} d x=0 \tag{18}
\end{equation*}
$$

Equalities (17) and (18) hold with $K g_{p}$ replaced by 1.
Proof. First we assume that $K$ is of type 0 . Let

$$
\Delta^{\prime}[p](x, y)=\Delta[p](x, y) \chi_{[0, N / p)}(|y-x|) \quad(p \geq 2)
$$

where $N=l(\log p)$. Then $\left\|\Delta[p]-\Delta^{\prime}[p]\right\|_{L^{2}, L^{2}}=O(1 / \log p)$ (cf. Lemma 6), and hence

$$
\begin{aligned}
\int_{0}^{1} K g_{p} \cdot(\Delta[p] \overline{\Delta[p]}) & l \bar{K} h_{p} d x \\
\quad & =\int_{0}^{1} K g_{p} \cdot\left(\Delta^{\prime}[p] \overline{\Delta^{\prime}[p]}\right) l \bar{K} h_{p} d x+O\left(\frac{1}{\log p}\right)
\end{aligned}
$$

Notice that

$$
\left(\Delta^{\prime}[p] \overline{\Delta^{\prime}[p]}\right)^{l}(x, y)=0 \quad(|y-x|>2 l N / p),
$$

and that $\left(\Delta^{\prime}[p] \overline{\Delta^{\prime}[p]}\right)^{l} 1$ is a periodic function on $[2 l N / p, 1-(2 l N / p))$ with period $2 / p$. Let

$$
\begin{aligned}
& \eta_{p}^{(m)}=\frac{p}{2} \int_{2 m / p}^{2(m+1) / p}\left(\Delta^{\prime}[p] \overline{\Delta^{\prime}[p]}\right)^{l} 1 d x-d_{2 l}(\Delta[\infty]) \\
&(0 \leq m \leq l(p / 2)-1)
\end{aligned}
$$

Then $\eta_{p}^{(m)}=\eta_{p}^{(l N)}(l N \leq m \leq l(p / 2)-l N-1)$. Lemma 6 shows that

$$
\begin{aligned}
\limsup _{p \rightarrow \infty}\left|\eta_{p}^{(l N)}\right| & =\limsup _{p \rightarrow \infty}\left|d_{2 l}\left(\Delta^{\prime}[p]\right)-d_{2 l}(\Delta[\infty])\right| \\
& =\underset{p \rightarrow \infty}{\limsup }\left|d_{2 l}(\Delta[p])-d_{2 l}(\Delta[\infty])\right|=0
\end{aligned}
$$

Since $K$ is of type 0 , we have

$$
\sup \frac{|K h(y)-K h(x)|}{|y-x|} \leq \sup _{s, t \in[0,1)}\left|\frac{\partial}{\partial s} K(s, t)\right|<\infty
$$

where the supremum in the left-hand side is taken over all $x, y \in$ $[0,1)$ and all $h \in L^{4}$ satisfying $\|h\|_{L^{4}} \leq 1$. Thus

$$
\begin{aligned}
\int_{0}^{1} K & g_{p} \cdot(\Delta[p] \overline{\Delta[p]}) l \bar{K} h_{p} d x \\
= & \int_{2 l N / p}^{1-(2 l N / p)} K g_{p}(x) \\
& \times \int_{|y-x| \leq 2 l N / p}\left(\Delta^{\prime}[p] \overline{\Delta^{\prime}[p]}\right)^{l}(x, y) \bar{K} h_{p}(y) d y d x+o(1) \\
= & \int_{2 l N / p}^{1-(2 l N / p)} K g_{p} \cdot \bar{K} h_{p} \cdot\left(\Delta^{\prime}[p] \overline{\Delta^{\prime}[p]}\right)^{l} 1 d x+o(1) \\
= & \frac{2}{p} \sum_{m=l N}^{l(p / 2)-l N-1} K g_{p}\left(\frac{2 m}{p}\right) \bar{K} h_{p}\left(\frac{2 m}{p}\right)\left\{\eta_{p}^{(m)}+d_{2 l}(\Delta[\infty])\right\}+o(1) \\
= & d_{2 l}(\Delta[\infty]) \int_{0}^{1} K g_{p} \cdot \bar{K} h_{p} d x+O\left(\eta_{p}^{(l N)}\right)+o(1) \\
= & d_{2 l}(\Delta[\infty]) \int_{0}^{1} K g_{p} \cdot \bar{K} h_{p} d x+o(1)
\end{aligned}
$$

which shows that (17) holds. Let $K$ be of type 1 . Then there exists a
sequence $\left(K_{j}\right)_{j=1}^{\infty}$ of kernels of type 0 such that

$$
\left\|K-K_{j}\right\|_{L^{4}, L^{2}} \leq \frac{1}{j} \quad(j \geq 1), \quad \sup _{j \geq 1}\left\|K_{j}\right\|_{L^{4}, L^{4}}<\infty
$$

Then

$$
\begin{gathered}
\mid \int_{0}^{1} K g_{p} \cdot\left(\Delta[p] \overline{\Delta[p])} l \bar{K} h_{p} d x-\int_{0}^{1} K_{j} g_{p} \cdot(\Delta[p] \overline{\Delta[p]}) / \bar{K}_{j} h_{p} d x \mid \leq C_{3} / j\right. \\
\left|\int_{0}^{1} K g_{p} \cdot \bar{K} h_{p} d x-\int_{0}^{1} K_{j} g_{p} \cdot \bar{K}_{j} h_{p} d x\right| \leq C_{3} / j \quad(j \geq 1)
\end{gathered}
$$

for some constant $C_{3}$ independent of $p$ and $j$. Since (17) holds for all $K_{j}(j \geq 1)$, this shows that (17) holds.

Since $\overline{\Delta[p]}(\Delta[p] \overline{\Delta[p]})^{l}$ is anti-symmetric, we have

$$
\int_{0}^{1} \overline{\Delta[p]}(\Delta[p] \overline{\Delta[p]})^{l} 1 d x=0
$$

Hence, in the same manner as above, we obtain (18). Analogously, we can replace $K g_{p}$ by 1.

Lemma 8. Let $K$ be an anti-symmetric kernel of type 1. Then, for any $l \geq 0$,

$$
\lim _{p \rightarrow \infty} d_{2 l}(\Delta[p]+K) \quad\left(=d_{2 l}(\Delta[\infty]+K), \text { say }\right)
$$

exists, and we can write

$$
\begin{equation*}
d_{2 l}(\Delta[\infty]+K)=\sum_{k=0}^{l} c_{2 k}^{(2 l)} d_{2 l-2 k}(K) \tag{19}
\end{equation*}
$$

so that $c_{2 k}^{(m)}(0 \leq k \leq \imath(m / 2), m \geq 0)$ satisfy

$$
\begin{equation*}
c_{0}^{(m)}=1, \quad c_{2 k}^{(2 k)}=d_{2 k}(\Delta[\infty]) \quad(m \geq 0, k \geq 0) \tag{20}
\end{equation*}
$$

(21) $\quad c_{2 k}^{(m)}=\sum_{j=0}^{k} c_{2 k-2 j}^{(m-2 j-1)} d_{2 j}(\Delta[\infty])$

$$
\left(0 \leq k \leq i\left(\frac{m-1}{2}\right), m \geq 0\right)
$$

Proof. We say that a $2 l$-tuple $\left(\tau_{1}, \ldots, \tau_{2 l}\right), \tau_{j}= \pm 1$ is negligible if there exist two integers $j_{0}, j_{0}^{\prime}\left(1 \leq j_{0}<j_{0}^{\prime} \leq 2 l\right)$ such that $j_{0}^{\prime}-j_{0}-1$ is odd, $\tau_{j}=-1\left(j_{0} \leq j \leq j_{0}^{\prime}\right)$ and $\tau_{j_{0}-1}=\tau_{j_{0}^{\prime}+1}=1$. (We put $\tau_{0}=\tau_{2 l+1}=1$. Hence $\tau_{j_{0}-1}=1$ if $j_{0}=1$, and $\tau_{j_{0}^{\prime}+1}=1$ if $j_{0}^{\prime}=2 l$.) Let $\tau(\Delta[p])=-1(p \geq 2), \tau(K)=1$. Lemmas 6 and 7 show that $d_{2 l}(\Delta[\infty]+K)$ exists and

$$
\begin{aligned}
d_{2 l}(\Delta[\infty]+K) & =\lim _{p \rightarrow \infty} \sum_{\left(K_{1}, \ldots, K_{2 l}\right), K_{j}=\Delta[p], K} \int_{0}^{1} K_{1} \bar{K}_{2} \cdots K_{2 l-1} \bar{K}_{2 l} 1 d x \\
& =\lim _{p \rightarrow \infty} \sum_{(p)} \int_{0}^{1} K_{1} \bar{K}_{2} \cdots K_{2 l-1} \bar{K}_{2 l} 1 d x
\end{aligned}
$$

where $\sum_{(p)}$ is the summation over all $2 l$-tuples $\left(K_{1}, \ldots, K_{2 l}\right), K_{j}=$ $\Delta[p], K$ such that $\left(\tau\left(K_{1}\right), \ldots, \tau\left(K_{2 l}\right)\right)$ is not negligible. If $\left(\tau\left(K_{1}\right), \ldots\right.$, $\tau\left(K_{2 l}\right)$ ) is not negligible, then $K$ appears even times in ( $K_{1}, \ldots, K_{2 l}$ ). We can choose $j_{1}<j_{2}<\cdots<j_{2 \nu}$ so that $K_{j_{\mu}}=K(1 \leq \mu \leq 2 \nu)$, $K_{j}=\Delta[p]\left(j \notin\left\{j_{\mu}\right\}_{\mu=1}^{2 \nu}\right)$. Then $j_{1}-1, j_{\mu+1}-j_{\mu}-1(1 \leq \mu \leq 2 \nu-1)$, $2 l-j_{2 \nu}$ are even. Notice that

$$
d_{2 j}(K)=\int_{0}^{1}(\bar{K} K)^{j} 1 d x \quad(j \geq 0) .
$$

Thus we can write

$$
d_{2 l}(\Delta[\infty]+K)=\sum_{k=0}^{l} c_{2 k}^{(2 l)} d_{2 l-2 k}(K)
$$

Let $\kappa_{0}$ be an operator defined by $h \in L^{2} \rightarrow\left(\int_{0}^{1} h d x\right) \chi_{[0,1)}$. We put $Y_{p,-1}(t)=1$,

$$
Y_{p, m}(t)= \begin{cases}\int_{0}^{1} K_{p, t}^{m / 2} 1 d x & (m \text { is even }) \\ \int_{0}^{1}\left(\kappa_{0}+t \overline{\Delta[p]}\right) K_{p, t}^{(m-1) / 2} 1 d x & (m \text { is odd })\end{cases}
$$

where $K_{p, t}=\left(\kappa_{0}+t \Delta[p]\right)\left(\kappa_{0}+t \overline{\Delta[p]}\right)$. Then $Y_{\infty, m}(t)=\lim _{p \rightarrow \infty} Y_{p, m}(t)$ exists, and $c_{2 k}^{(2 l)}$ equals the $t^{2 k}$-coefficient of $Y_{\infty, 2 l}(t)$. Evidently, (20)
holds. Since $\int_{0}^{1} \Delta[p](\overline{\Delta[p]} \Delta[p])^{j} 1 d x=0 \quad(j \geq 0)$, we have inductively

$$
\begin{aligned}
Y_{p, 2 l}(t)= & Y_{p, 2 l-1}(t)+t \int_{0}^{1} \Delta[p]\left(\kappa_{0}+t \overline{\Delta[p]}\right) K_{p, t}^{l-1} 1 d x \\
= & Y_{p, 2 l-1}(t)+t^{2} \int_{0}^{1} \Delta[p] \overline{\Delta[p]} K_{p, t}^{l-1} 1 d x \\
= & Y_{p, 2 l-1}(t)+t^{2} d_{2}(\Delta[p]) Y_{p, 2 l-3}(t) \\
& +t^{3} \int_{0}^{1} \Delta[p] \overline{\Delta[p]} \Delta[p]\left(\kappa_{0}+t \overline{\Delta[p]}\right) K_{p, t}^{l-2} 1 d x \\
= & Y_{p, 2 l-1}(t)+t^{2} d_{2}(\Delta[p]) Y_{p, 2 l-3}(t) \\
& +t^{4} \int_{0}^{1}(\Delta[p] \overline{\Delta[p]})^{2} K_{p, t}^{l-2} 1 d x \\
= & \cdots=\sum_{j=0}^{l} t^{2 j} d_{2 j}(\Delta[p]) Y_{p, 2 l-2 j-1}(t) .
\end{aligned}
$$

Letting $p$ tend to infinity, we have

$$
Y_{\infty, 2 l}(t)=\sum_{j=0}^{l} t^{2 j} d_{2 \jmath}(\Delta[\infty]) Y_{\infty, 2 l-2 j-1}(t)
$$

In the same manner,

$$
Y_{\infty, 2 l+1}(t)=\sum_{j=0}^{l} t^{2 j} d_{2 j}(\Delta[\infty]) Y_{\infty, 2 l-2 \jmath}(t)
$$

Thus

$$
Y_{\infty, m}(t)=\sum_{j=0}^{l(m / 2)} t^{2 \jmath} d_{2 j}(\Delta[\infty]) Y_{\infty, m-2 J-1}(t)
$$

Comparing the $t^{2 k}$-coefficients of both sides, we obtain (21).
Lemma 9. Let $K$ be an anti-symmetric kernel of type 1. Then, for any $0<\delta \leq 1$,

$$
\lim _{p \rightarrow \infty} \gamma^{*}(\delta \Delta[p]+\delta K) \quad\left(=\gamma^{*}(\delta \Delta[\infty]+\delta K), \text { say }\right)
$$

exists; we write $\gamma^{*}(\delta \Delta[\infty])$ if $K=0$. Moreover,

$$
\begin{equation*}
\gamma^{*}(\delta \Delta[\infty]+\delta K)=\gamma^{*}(\delta \Delta[\infty]) \gamma^{*}\left(\gamma^{*}(\delta \Delta[\infty]) \delta K\right) . \tag{22}
\end{equation*}
$$

Proof. First we show that $\gamma^{*}(\delta \Delta[\infty]+\delta K)$ and $\gamma^{*}(\delta \Delta[\infty])$ exist. Define $\gamma^{*}(w ; \Delta[p]+K), T(w ; \Delta[p]+K)(w \in U)$ for $\varepsilon_{0}=$
$\left(12+3\|K\|_{L^{2}, L^{2}}\right)^{-1}$ in the same manner as in Lemma 5; we have $\varepsilon_{0} \leq\left(3\|\Delta[p]+K\|_{L^{2}, L^{2}}\right)^{-1}$ because $\|\Delta[p]\|_{L^{2}, L^{2}} \leq 4$. Lemma 8 shows that

$$
\begin{aligned}
& \lim _{p \rightarrow \infty} w^{2} \varepsilon_{0}^{2} \int_{0}^{1}(\Delta[p]+K) T(w ; \Delta[p]+K)^{m+1}(\overline{\Delta[p]+K}) 1 d x \\
& \quad=\sum_{l=0}^{\infty} w^{2 l+2} \varepsilon_{0}^{2 l+2} \frac{(l+1) \cdots(l+m)}{m!} d_{2 l+2}(\Delta[\infty]+K) \quad(m \geq 0)
\end{aligned}
$$

Since (14) holds with $K$ replaced by any $\Delta[p]+K(p \geq 2)$, (13) exists with $K$ replaced by $\Delta[\infty]+K$, i.e.,

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \gamma^{*}(w ; \Delta[p]+K) \quad\left(=\gamma^{*}(w ; \Delta[\infty]+K), \text { say }\right) \tag{23}
\end{equation*}
$$

exists. Since

$$
\gamma^{*}(\delta ; \Delta[p]+K)=\gamma^{*}(\delta \Delta[p]+\delta K) \quad(p \geq 2)
$$

$\gamma^{*}(\delta \Delta[\infty]+\delta K)\left(=\gamma^{*}(\delta ; \Delta[\infty]+K)\right)$ exists. Putting $K=0$, we see that $\gamma^{*}(\delta \Delta[\infty])$ exists.

Next we show that $\gamma^{*}(w ; \Delta[\infty]+K)$ and $\gamma^{*}\left(\gamma^{*}(w ; \Delta[\infty]) w ; K\right)$ are analytic in a domain containing ( 0,1 ]. The convergence of (23) is uniform in $U$. By Lemma $5, \gamma^{*}(w ; \Delta[p]+K)$ is analytic in $U$, and hence $\gamma^{*}(w ; \Delta[\infty]+K)$ is analytic in $U$. The definition of $\gamma^{*}(\cdot)$ immediately shows that

$$
\gamma^{*}(\operatorname{Re} w ; \Delta[p])=\gamma^{*}(\operatorname{Re} w \Delta[p]) \leq 1 \quad(w \in U)
$$

Letting $p$ tend to infinity, we have $\gamma^{*}(\operatorname{Re} w ; \Delta[\infty]) \leq 1 \quad(w \in U)$. Since $\gamma^{*}(w ; \Delta[\infty])$ is analytic in $U$, there exists $0<\eta<\pi / 8$ such that

$$
\left|\gamma^{*}(w ; \Delta[\infty])\right| \leq \frac{4}{3}, \quad\left|\arg \gamma^{*}(w ; \Delta[\infty])\right| \leq \frac{\pi}{8}
$$

in $U_{\eta}=\{w \in \mathbb{C} ;|w|<4 / 3,|\arg w|<\eta\}$. Then $\gamma^{*}(w ; \Delta[\infty]) w \in U$ $\left(w \in U_{\eta}\right)$. Thus, by Lemma $5, \gamma^{*}\left(\gamma^{*}(w ; \Delta[\infty]) w ; K\right)$ is analytic in $U_{\eta}$.

By the theorem of identity, it is sufficient to show that (22) holds for $0<\delta<\left(8+2\|K\|_{L^{2}, L^{2}}\right)^{-1}$. Since

$$
\lim _{l \rightarrow \infty} d_{2 l}(\delta \Delta[p])=\lim _{l \rightarrow \infty} d_{2 l}(\delta \Delta[p]+\delta K)=0,
$$

(12) holds for $\delta \Delta[p], \delta \Delta[p]+\delta K(p \geq 2)$. Letting $p$ tend to infinity, we have

$$
\gamma^{*}(\delta \Delta[\infty])=\sum_{l=0}^{\infty} d_{2 l}(\delta \Delta[\infty])=\sum_{l=0}^{\infty} \delta^{2 l} d_{2 l}(\Delta[\infty])
$$

$$
\gamma^{*}(\delta \Delta[\infty]+\delta K)=\sum_{l=0}^{\infty} \delta^{2 l} d_{2 l}(\Delta[\infty]+K) .
$$

Let

$$
\mu_{m}=\sum_{k=0}^{\infty} \delta^{2 k} c_{2 k}^{(m+2 k)} \quad(m \geq 0)
$$

where $c_{2 k}^{(m)}(0 \leq k \leq \imath(m / 2), m \geq 0)$ are numbers in Lemma 8. Then

$$
\mu_{0}=\sum_{k=0}^{\infty} \delta^{2 k} c_{2 k}^{(2 k)}=\gamma^{*}(\delta \Delta[\infty]),
$$

by (20). Equality (21) yields that

$$
\begin{aligned}
\mu_{m} & =\sum_{k=0}^{\infty} \delta^{2 k} \sum_{j=0}^{k} c_{2 k-2 j}^{(m+2 k-2 j-1)} d_{2 j}(\Delta[\infty]) \\
& =\sum_{j=0}^{\infty} \delta^{2 j} d_{2 j}(\Delta[\infty]) \sum_{k=j}^{\infty} \delta^{2(k-j)} c_{2(k-j)}^{(m-1+2(k-j))} \\
& =\mu_{m-1} \mu_{0} \quad(m \geq 1),
\end{aligned}
$$

which gives

$$
\mu_{m}=\mu_{0}^{m+1}=\gamma^{*}(\delta \Delta[\infty])^{m+1} \quad(m \geq 1)
$$

Thus, by (21),

$$
\begin{aligned}
\gamma^{*}(\delta \Delta & {[\infty]+\delta K)=\sum_{l=0}^{\infty} \delta^{2 l} d_{2 l}(\Delta[\infty]+K) } \\
& =\sum_{l=0}^{\infty} \delta^{2 l} \sum_{k=0}^{l} c_{2 k}^{(2 l)} d_{2 l-2 k}(K)=\sum_{k=0}^{\infty} \sum_{l=k}^{\infty} \delta^{2 l} c_{2 k}^{(2 l)} d_{2 l-2 k}(K) \\
& =\sum_{k=0}^{\infty} \sum_{j=0}^{\infty} \delta^{2 j+2 k} c_{2 k}^{(2 j+2 k)} d_{2 j}(K)=\sum_{j=0}^{\infty} \mu_{2 j} \delta^{2 j} d_{2 j}(K) \\
& =\sum_{j=0}^{\infty} \gamma^{*}(\delta \Delta[\infty])^{2 j+1} \delta^{2 j} d_{2 j}(K)=\gamma^{*}(\delta \Delta[\infty]) \gamma^{*}\left(\gamma^{*}(\delta \Delta[\infty]) \delta K\right) .
\end{aligned}
$$

Lemma 10. We inductively define a sequence $\left(\gamma_{n}^{*}\right)_{n=1}^{\infty}$ of positive numbers by

$$
\gamma_{1}^{*}=\gamma^{*}(\Delta[\infty]), \quad \gamma_{n}^{*}=\gamma_{n-1}^{*} \gamma^{*}\left(\gamma_{n-1}^{*} \Delta[\infty]\right) \quad(n \geq 2)
$$

Then

$$
\lim _{p_{1} \rightarrow \infty} \lim _{p_{2} \rightarrow \infty} \cdots \lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(H\left[p_{1}, \ldots, p_{n}\right]-H\right)=\gamma_{n}^{*} \quad(n \geq 1)
$$

where $\lim _{p_{n} \rightarrow \infty}$ is taken first and $\lim _{p_{1} \rightarrow \infty}$ is taken last.
Proof. We define a sequence $\left(\lambda_{n}\right)_{n=1}^{\infty}$ of positive numbers by $\lambda_{1}=$ $\gamma^{*}(\Delta[\infty]), \lambda_{n}=\gamma^{*}\left(\lambda_{1} \ldots \lambda_{n-1} \Delta[\infty]\right)(n \geq 2)$. Then $\gamma_{n}^{*}=\lambda_{1} \cdots \lambda_{n}$ $(n \geq 1)$. Fixing an $(n-1)$-tuple $\left(p_{1}, \ldots, p_{n-1}\right) \quad(n \geq 2)$ of integers larger than or equal to 3 , we study

$$
\begin{aligned}
& \lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(H\left[p_{1}, \ldots, p_{n}\right]-H\right) \\
& \quad=\lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(\Delta\left[p_{1}, \ldots, p_{n}\right]+\left(H\left[p_{1}, \ldots, p_{n-1}\right]-H\right)\right)
\end{aligned}
$$

Put $I_{0}=\left[0,1 /\left(p_{1} \cdots p_{n-1}\right)\right), \quad I_{j}=\left(I_{0}+j /\left(p_{1} \cdots p_{n-1}\right)\right) \quad(0 \leq j \leq$ $\left.\left(p_{1} \ldots p_{n-1}\right)-1\right)$. Then

$$
\begin{aligned}
& \quad\left(H\left[p_{1}, \ldots, p_{n-1}\right]-H\right)(x, y)=0 \quad\left(x, y \in I_{j}\right) \\
& \left|\left(H\left[p_{1}, \ldots, p_{n-1}\right]-H\right)(x, y)\right| \\
& \quad \leq \frac{2}{\pi} p_{1} \cdots p_{n-1}+\frac{1}{\pi} \frac{1}{|y-x|} \quad\left(x \in I_{j}, y \in I_{k}, j \neq k\right)
\end{aligned}
$$

which shows that $H\left[p_{1}, \ldots, p_{n-1}\right]-H$ is of type 1 . Let

$$
\begin{gathered}
\Delta^{\prime}\left[p_{1}, \ldots, p_{n}\right](x, y)=\Delta\left[p_{1}, \ldots, p_{n}\right](x, y) \chi_{\left[0, \frac{N^{\prime}}{p_{1} \cdots p_{n}}\right)}(|y-x|), \\
\Delta^{\prime}\left[\prod_{j=1}^{n} p_{j}\right](x, y)=\Delta\left[\prod_{j=1}^{n} p_{j}\right](x, y) \chi_{\left[0, \frac{N^{\prime}}{p_{1} \cdots p_{n}}\right)}(|y-x|) \\
\left(N^{\prime}=l\left(\log \left(p_{1} \ldots p_{n}\right)\right) .\right. \text { Then } \\
\lim _{p_{n} \rightarrow \infty}\left\|\Delta\left[p_{1}, \ldots, p_{n}\right]-\Delta^{\prime}\left[p_{1}, \ldots, p_{n}\right]\right\|_{L^{4}, L^{4}}=0, \\
\lim _{p_{n} \rightarrow \infty}\left\|\Delta\left[\prod_{j=1}^{n} p_{j}\right]-\Delta^{\prime}\left[\prod_{j=1}^{n} p_{j}\right]\right\|_{L^{4}, L^{4}}=0
\end{gathered}
$$

(cf. Lemmas 6 and 7). Since
$\Delta^{\prime}\left[p_{1}, \ldots, p_{n}\right](x, y)=\Delta^{\prime}\left[\prod_{j=1}^{n} p_{j}\right](x, y)$

$$
\left(x, y \in I_{j}, 0 \leq j \leq\left(p_{1} \cdots p_{n-1}\right)-1\right)
$$

we have

$$
\begin{aligned}
& \lim _{p_{n} \rightarrow \infty}\left\|\Delta\left[p_{1}, \ldots, p_{n}\right]-\Delta\left[\prod_{j=1}^{n} p_{j}\right]\right\|_{L^{4}, L^{2}} \\
& \quad=\lim _{p_{n} \rightarrow \infty}\left\|\Delta^{\prime}\left[p_{1}, \ldots, p_{n}\right]-\Delta^{\prime}\left[\prod_{j=1}^{n} p_{j}\right]\right\|_{L^{4}, L^{2}}=0
\end{aligned}
$$

and hence, in the same manner as in the proof of the existence of (23),

$$
\begin{aligned}
& \lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(\Delta\left[p_{1}, \ldots, p_{n}\right]+\left(H\left[p_{1}, \ldots, p_{n-1}\right]-H\right)\right) \\
& \quad=\lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(\Delta\left[\prod_{j=1}^{n} p_{j}\right]+\left(H\left[p_{1}, \ldots, p_{n-1}\right]-H\right)\right) .
\end{aligned}
$$

Using (22) with $\delta=1, K=H\left[p_{1}, \ldots, p_{n-1}\right]-H$, we have

$$
\begin{aligned}
\lim _{p_{n} \rightarrow \infty} & \gamma^{*}\left(H\left[p_{1}, \ldots, p_{n}\right]-H\right)=\gamma^{*}\left(\Delta[\infty]+\left(H\left[p_{1}, \ldots, p_{n-1}\right]-H\right)\right) \\
& =\lambda_{1} \gamma^{*}\left(\lambda_{1}\left(H\left[p_{1}, \ldots, p_{n-1}\right]-H\right)\right)
\end{aligned}
$$

In the same manner, using (22) with $\delta=\lambda_{1}, K=H\left[p_{1}, \ldots, p_{n-2}\right]-$ $H$, we have

$$
\begin{aligned}
& \lim _{p_{n-1} \rightarrow \infty} \lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(H\left[p_{1}, \ldots, p_{n}\right]-H\right) \\
& \quad=\lambda_{1} \gamma^{*}\left(\lambda_{1} \Delta[\infty]\right) \gamma^{*}\left(\gamma^{*}\left(\lambda_{1} \Delta[\infty]\right) \lambda_{1}\left(H\left[p_{1}, \ldots, p_{n-2}\right]-H\right)\right) \\
& \quad=\lambda_{1} \lambda_{2} \gamma^{*}\left(\lambda_{1} \lambda_{2}\left(H\left[p_{1}, \ldots, p_{n-2}\right]-H\right)\right)
\end{aligned}
$$

Repeating this argument,

$$
\begin{aligned}
\lim _{p_{1} \rightarrow \infty} & \cdots \lim _{p_{n} \rightarrow \infty} \gamma^{*}\left(H\left[p_{1}, \ldots, p_{n}\right]-H\right) \\
& =\lambda_{1} \cdots \lambda_{n-1} \lim _{p_{1} \rightarrow \infty} \gamma^{*}\left(\lambda_{1} \cdots \lambda_{n-1} \Delta\left[p_{1}\right]\right)=\lambda_{1} \ldots \lambda_{n}=\gamma_{n}^{*}
\end{aligned}
$$

This completes the proof of our lemma.
We now give the proof of our theorem. By Proposition 2, there exists a positive integer $n_{0}$ such that

$$
\begin{equation*}
\sup \gamma\left(\Gamma\left(p_{1}, \ldots, p_{n}\right)\right) \leq 10^{-5} \quad\left(n \geq n_{0}\right) \tag{24}
\end{equation*}
$$

where the supremum is taken over all $n$-tuples $\left(p_{1}, \ldots, p_{n}\right)$ of integers larger than or equal to 3 . By Lemma 10 , we can inductively choose a sequence $\left(p_{n}^{0}\right)_{n=1}^{\infty}$ of integers larger than or equal to 3 so that

$$
\frac{1}{2} \gamma_{n}^{*} \leq \gamma^{*}\left(H\left[p_{1}^{0}, \ldots, p_{n}^{0}\right]-H\right) \leq 2 \gamma_{n}^{*} \quad(n \geq 1)
$$

where $\left(\gamma_{n}^{*}\right)_{n=1}^{\infty}$ is the sequence in Lemma 10 . We show that $\Gamma_{n}=$ $\Gamma\left(p_{1}^{0}, \ldots, p_{n}^{0}\right)(n \geq 1)$ are required cranks. We may assume that $n \geq n_{0}$. Lemma 3 shows that

$$
\begin{aligned}
\frac{1}{4} \gamma^{*}\left(H\left[p_{1}^{0}, \ldots, p_{n}^{0}\right]-H\right) & \leq \gamma^{*}\left(H\left[p_{1}^{0}, \ldots, p_{n}^{0}\right]\right) \\
& \leq 4 \gamma^{*}\left(H\left[p_{1}^{0}, \ldots, p_{n}^{0}\right]-H\right)
\end{aligned}
$$

and hence

$$
\frac{1}{8} \gamma_{n}^{*} \leq \gamma^{*}\left(H\left[p_{1}^{0}, \ldots, p_{n}^{0}\right]\right) \leq 8 \gamma_{n}^{*}
$$

Thus, by (11),

$$
\begin{equation*}
\frac{1}{8 \pi} \gamma_{n}^{*} \leq \gamma\left(\Gamma_{n}\right) \leq \frac{8}{\pi} \gamma_{n}^{*} \tag{25}
\end{equation*}
$$

Using (24) and (25), we have $\gamma_{n}^{*} \leq 8 \pi \cdot 10^{-5}$. Recall (15), and notice that $d_{2 l}(\Delta[\infty]) \leq 4^{l} \quad(l \geq 1)$. Since $\lim _{l \rightarrow \infty} d_{2 l}\left(\gamma_{n}^{*} \Delta[p]\right)=0$, (12) holds for $\gamma_{n}^{*} \Delta[p]$. Letting $p$ tend to infinity, we have

$$
\begin{aligned}
\gamma_{n+1}^{*} & =\gamma_{n}^{*} \gamma^{*}\left(\gamma_{n}^{*} \Delta[\infty]\right)=\gamma_{n}^{*} \sum_{l=0}^{\infty} d_{2 l}\left(\gamma_{n}^{*} \Delta[\infty]\right) \\
& =\gamma_{n}^{*} \sum_{l=0}^{\infty} \gamma_{n}^{*^{*^{l}}} d_{2 l}(\Delta[\infty]) \leq \gamma_{n}^{*}-\frac{1}{25 \pi^{2}} \gamma_{n}^{*^{3}}+\sum_{l=2}^{\infty} 4^{l} \gamma_{n}^{*^{l+1}} \\
& \leq \gamma_{n}^{*}-10^{-3} \gamma_{n}^{*^{3}}, \\
& \gamma_{n+1}^{*} \geq \gamma_{n}^{*}-\sum_{l=1}^{\infty} 4^{l} \gamma_{n}^{*^{l+1}} \geq \gamma_{n}^{*}-10 \gamma_{n}^{*^{3}}, \quad \text { i.e. }, \\
& \gamma_{n}^{*}-10 \gamma_{n}^{*^{*}} \leq \gamma_{n+1}^{*} \leq \gamma_{n}^{*}-10^{-3} \gamma_{n}^{*^{3}} .
\end{aligned}
$$

Since this holds for all $n \geq n_{0}$, a simple induction yields that

$$
\frac{1}{C_{4}} \frac{1}{\sqrt{n}} \leq \gamma_{n}^{*} \leq C_{4} \frac{1}{\sqrt{n}} \quad\left(n \geq n_{0}\right)
$$

for some absolute constant $C_{4}$. Using (25) again,

$$
\frac{1}{8 \pi C_{4}} \frac{1}{\sqrt{n}} \leq \gamma\left(\Gamma_{n}\right) \leq \frac{8}{\pi} C_{4} \frac{1}{\sqrt{n}} \quad\left(n \geq n_{0}\right) .
$$

This completes the proof of our theorem.
Remark 11. It is not known whether $\gamma(\cdot)$ is semi-additive [4, p. 11]. For $0<\eta \leq 1$, we define $B_{p}^{\eta}(x)$ replacing $1 / 2 p$ by $\eta / 2 p$ in the definition of $B_{p}(x)$. Then cranks $\Gamma^{\eta}\left(p_{1}, \ldots, p_{n}\right)$ of degree $n$ are
analogously defined. We see that there exists a crank $\Gamma_{n}^{\eta}$ of degree $n$ such that $\gamma\left(\Gamma_{n}^{\eta}\right) \leq C_{\eta} / \sqrt{n}$, where $C_{\eta}$ is a constant depending only on $\eta$. Adding some segments (perpendicular to the $x$-axis) to $\Gamma_{n}^{\eta}$, we obtain an arc $\widetilde{\Gamma}_{n}^{\eta}$ connecting 0 and 1 . Then the diameter of $\widetilde{\Gamma}_{n}^{\eta}$ is larger than or equal to 1 . Since $\widetilde{\Gamma}_{n}^{\eta}$ is connected, $\gamma\left(\widetilde{\Gamma}_{n}^{\eta}\right) \geq 1 / 4[4, \mathrm{p}$. 9]. Hence, from the point of view of the above semi-additive problem, it seems interesting to compute $\gamma\left(\widetilde{\Gamma}_{n}^{\eta}-\Gamma_{n}^{\eta}\right)$.
4. Another application of Proposition 1. In this section, we show another application of our method. Let $E$ be a compact set on $\mathbb{R}$. Pommerenke [11] showed that

$$
\begin{equation*}
\gamma(E)=|E| / 4, \tag{26}
\end{equation*}
$$

$$
\begin{equation*}
f_{E}(z)=\left\{1-\exp \left(\frac{1}{2} \int_{E} \frac{d t}{t-z}\right)\right\} /\left\{1+\exp \left(\frac{1}{2} \int_{E} \frac{d t}{t-z}\right)\right\} \tag{27}
\end{equation*}
$$

We deduce (26), (27) from (3), (10); our method explains a quarter and (27). Let $L^{2}(\mathbb{R})$ denote the $L^{2}$ space of functions on $\mathbb{R}$, and let $M_{E}$ denote the multiplier: $h \in L^{2}(\mathbb{R}) \rightarrow \chi h \in L^{2}(\mathbb{R})$, where $\chi=\chi_{E}$. We inductively define a sequence $\left(H_{E}^{(m)}\right)_{m=0}^{\infty}$ of operators from $L^{2}(\mathbb{R})$ to itself by $H_{E}^{(0)}=M_{E}, H_{E}^{(m)}=H M_{E} H_{E}^{(m-1)}(m \geq 1)$. Notice that

$$
\begin{aligned}
\gamma(E) & =\frac{1}{\pi} \gamma^{*}\left(M_{E} H M_{E}\right), \\
d_{2 l}\left(M_{E} H M_{E}\right) & =\int_{E} H_{E}^{(2 l)} \chi d x \quad\left(l \geq 0, \chi=\chi_{E}\right) .
\end{aligned}
$$

We also remark that

$$
\begin{equation*}
H(g \cdot H h)+H(H g \cdot h)=H g \cdot H h-g h \quad\left(g, h \in L^{2}(\mathbb{R})\right) . \tag{28}
\end{equation*}
$$

We first show that, for any $m \geq 1$,

$$
\begin{equation*}
\chi H \chi \cdot H_{E}^{(m)} \chi=(m+1) \chi H_{E}^{(m+1)} \chi+m \chi H_{E}^{(m-1)} \chi . \tag{29}
\end{equation*}
$$

Equality (28) shows that $2 H(\chi H \chi)=(H \chi)^{2}-\chi \chi$, which gives $\chi H \chi$. $H_{E}^{(1)} \chi=2 \chi H_{E}^{(2)} \chi+\chi H_{E}^{(0)} \chi$. Suppose that (29) holds for $m$. Using (28) with $g=\chi, h=\chi H_{E}^{(m)} \chi$, we have

$$
\begin{aligned}
& \chi H \chi \cdot H_{E}^{(m+1)} \chi=\chi H \chi \cdot H\left(\chi H_{E}^{(m)} \chi\right) \\
& \quad= \chi H\left\{\chi H\left(\chi H_{E}^{(m)} \chi\right)+H \chi \cdot \chi H_{E}^{(m)} \chi\right\}+\chi\left\{\chi \cdot \chi H_{E}^{(m)} \chi\right\} \\
& \quad= \chi H_{E}^{(m+2)} \chi+\chi H\left\{(m+1) \chi H_{E}^{(m+1)} \chi+m \chi H_{E}^{(m-1)} \chi\right\}+\chi H_{E}^{(m)} \chi \\
& \quad=(m+2) \chi H_{E}^{(m+2)} \chi+(m+1) \chi H_{E}^{(m)} \chi, \quad \text { i.e., }
\end{aligned}
$$

(29) holds for $m+1$. Thus (29) holds for all $m \geq 1$.

We next show that

$$
\begin{equation*}
\int_{E} H_{E}^{(2 l)} \chi d x=\frac{(-1)^{l}}{2 l+1}|E| \quad(l \geq 0) \tag{30}
\end{equation*}
$$

We put $a_{2 l}=\int_{E} H_{E}^{(2 l)} \chi d x \quad(l \geq 0)$. Evidently, $a_{0}=|E|$. Suppose that $a_{2 l-2}=\left\{(-1)^{l-1} /(2 l-1)\right\}|E|$. Equality (29) $(m=2 l-1)$ shows that

$$
\begin{aligned}
\int_{E} H \chi \cdot H_{E}^{(2 l-1)} \chi d x & =2 l \int_{E} H_{E}^{(2 l)} \chi d x+(2 l-1) \int_{E} H_{E}^{(2 l-2)} \chi d x \\
& =2 l a_{2 l}+(2 l-1) a_{2 l-2}
\end{aligned}
$$

Since the adjoint operator of $H$ equals $-H$, we have

$$
\int_{E} H \chi \cdot H_{E}^{(2 l-1)} \chi d x=-\int_{E} H\left\{\chi H_{E}^{(2 l-1)} \chi\right\} d x=-a_{2 l}
$$

Thus $-a_{2 l}=2 l a_{2 l}+(2 l-1) a_{2 l-2}$, which yields that

$$
a_{2 l}=-\frac{2 l-1}{2 l+1} a_{2 l-2}=\frac{(-1)^{l}}{2 l+1}|E|
$$

Now the deduction of (26) is immediate. By (30),

$$
\lim _{l \rightarrow \infty} d_{2 l}\left(M_{E} H M_{E}\right)=\lim _{l \rightarrow \infty} \int_{E} H_{E}^{(2 l)} \chi d x=0
$$

Hence we can apply (3). Leibniz's formula and (30) yield that

$$
\begin{aligned}
\gamma(E) & =\frac{1}{\pi} \gamma^{*}\left(M_{E} H M_{E}\right)=\frac{1}{\pi} \sum_{l=0}^{\infty} d_{2 l}\left(M_{E} H M_{E}\right) \\
& =\frac{1}{\pi} \sum_{l=0}^{\infty} \frac{(-1)^{l}}{2 l+1}|E|=\frac{1}{4}|E|
\end{aligned}
$$

Last, we deduce (27) from (10). Equality (10) gives that

$$
f_{E}(z)=-\frac{1}{\pi}\left\{\int_{E} \frac{d s}{s-z}+\int_{E} \frac{H_{E}^{(1)} h_{E}(s)}{s-z} d s\right\} /\left\{1+\frac{1}{\pi} \int_{E} \frac{h_{E}(s)}{s-z} d s\right\}
$$

where $h_{E}(s)$ is the function which attains $\gamma^{*}\left(M_{E} H M_{E}\right)$. We show that this equals the function in the right-hand side of (27). Let

$$
\begin{aligned}
& u_{0}(z)=1, \quad u_{m}(z)=\frac{1}{\pi} \int_{E} \frac{H_{E}^{(m-1)} \chi(s)}{s-z} d s \\
& v_{m}(z)=\frac{1}{\pi} \int_{E} \frac{H \chi(s) H_{E}^{(m-2)} \chi(s)}{s-z} d s \quad(m \geq 1)
\end{aligned}
$$

where $H_{E}^{(-1)} \chi=\chi$. Let

$$
P_{t}(z)=\sum_{m=0}^{\infty} t^{m} u_{m}(z) \quad(t \in \mathbb{C},|t|<1)
$$

We begin by showing that

$$
\begin{equation*}
\left(1+t^{2}\right) \frac{\partial}{\partial t} P_{t}(z)=u_{1}(z) P_{t}(z) \quad(0<t<1) \tag{31}
\end{equation*}
$$

Let $m \geq 1$. We have, on $\mathbb{R}$,

$$
\begin{aligned}
& \lim _{\eta \downarrow 0}\left\{u_{m+1}(\cdot+i \eta)+v_{m+1}(\cdot+i \eta)\right\} \\
&= H\left(\chi H_{E}^{(m)} \chi\right)+i \chi H_{E}^{(m)} \chi \\
&+H\left\{\chi H \chi \cdot H_{E}^{(m-1)} \chi\right\}+i \chi H \chi \cdot H_{E}^{(m-1)} \chi \\
&= H\left\{\chi H\left(\chi H_{E}^{(m-1)} \chi\right)+H \chi \cdot \chi H_{E}^{(m-1)} \chi\right\} \\
&+i\left\{\chi H\left(\chi H_{E}^{(m-1)} \chi\right)+H \chi \cdot \chi H_{E}^{(m-1)} \chi\right\}, \\
& \lim _{\eta \downarrow 0} u_{1}(\cdot+i \eta) u_{m}(\cdot+i \eta) \\
&=\{H \chi+i \chi\}\left\{H\left(\chi H_{E}^{(m-1)} \chi\right)+i \chi H_{E}^{(m-1)} \chi\right\} \\
&= H \chi \cdot H\left(\chi H_{E}^{(m-1)} \chi\right)-\chi \cdot \chi H_{E}^{(m-1)} \chi \\
&+i\left\{\chi H\left(\chi H_{E}^{(m-1)} \chi\right)+H \chi \cdot \chi H_{E}^{(m-1)} \chi\right\}
\end{aligned}
$$

Hence (28) ( $g=\chi, \quad h=\chi H_{E}^{(m-1)} \chi$ ) shows that

$$
\lim _{\eta \downarrow 0}\left\{u_{m+1}(\cdot+i \eta)+v_{m+1}(\cdot+i \eta)-u_{1}(\cdot+i \eta) u_{m}(\cdot+i \eta)\right\}=0
$$

on $\mathbb{R}$. In particular, this holds on $\mathbb{R}-E$. Hence, by the theorem of identity, $u_{m+1}(z)+v_{m+1}(z)-u_{1}(z) u_{m}(z)=0$. Equality (29) shows that $v_{m+1}(z)=m u_{m+1}(z)+(m-1) u_{m-1}(z)$. Thus

$$
(m+1) u_{m+1}(z)+(m-1) u_{m-1}(z)-u_{1}(z) u_{m}(z)=0 \quad(m \geq 1)
$$

which yields that

$$
\begin{gathered}
\sum_{m=0}^{\infty} m t^{m} u_{m}(z)+t^{2} \sum_{m=0}^{\infty} m t^{m} u_{m}(z)=t u_{1}(z) \sum_{m=0}^{\infty} t^{m} u_{m}(z), \quad \text { i.e. } \\
t \frac{\partial}{\partial t} P_{t}(z)+t^{3} \frac{\partial}{\partial t} P_{t}(z)=t u_{1}(z) P_{t}(z)
\end{gathered}
$$

This is the required equality (31).

We can choose $x_{0} \in \mathbb{R}-E, \eta>0$ so that $P_{t}(x)>0, u_{1}(x)>0$ for all $x \in\left(x_{0}-\eta, x_{0}+\eta\right), 0<t<1$. Equality (31) shows that

$$
\frac{1}{1+t^{2}} u_{1}(x)=\frac{\partial}{\partial t} P_{t}(x) / P_{t}(x) \quad\left(x \in\left(x_{0}-\eta, x_{0}+\eta\right), 0<t<1\right)
$$

which gives that

$$
P_{t}(x)=\exp \left\{\int_{0}^{t} \frac{d s}{1+s^{2}} u_{1}(x)\right\} \quad\left(x \in\left(x_{0}-\eta, x_{0}+\eta\right), 0<t<1\right)
$$

because $P_{0}=1$. By the theorem of identity,

$$
P_{t}(z)=\exp \left\{\int_{0}^{t} \frac{d s}{1+s^{2}} u_{1}(z)\right\} \quad(0<t<1)
$$

Since $P_{t}(z)$ and $\exp \left\{\left(\int_{0}^{t}\left(d s /\left(1+s^{2}\right)\right) u_{1}(z)\right\}\right.$ are analytic in the unit disk as functions of $t$, this equality holds for $-1<t<0$ also. Thus

$$
\begin{aligned}
1+ & \frac{1}{\pi} \int_{E} \frac{h_{E}(s)}{s-z} d s=1+\frac{1}{\pi} \int_{E} \frac{1}{s-z} \sum_{l=1}^{\infty} H_{E}^{(2 l-1)} \chi(s) d s \\
& =\lim _{t \uparrow 1} \sum_{l=0}^{\infty} t^{2 l} u_{2 l}(z)=\frac{1}{2} \lim _{t \uparrow 1}\left\{P_{-t}(z)+P_{t}(z)\right\} \\
& =\frac{1}{2}\left\{\exp \left(-\frac{\pi}{4} u_{1}(z)\right)+\exp \left(\frac{\pi}{4} u_{1}(z)\right)\right\} \\
- & \frac{1}{\pi} \int_{E} \frac{d s}{s-z}-\frac{1}{\pi} \int_{E} \frac{H_{E}^{(1)} h_{E}(s)}{s-z} d s=\frac{1}{2} \lim _{t \uparrow 1}\left\{P_{-t}(z)-P_{t}(z)\right\} \\
& =\frac{1}{2}\left\{\exp \left(-\frac{\pi}{4} u_{1}(z)\right)-\exp \left(\frac{\pi}{4} u_{1}(z)\right)\right\}
\end{aligned}
$$

which gives (27).

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