

# ON THE NON-MINIMAL MARTIN BOUNDARY POINTS

TERUO IKEGAMI

Dedicated to Professor KIVOSHI NOSHIRO on his 60th birthday

1. In a Green space<sup>1)</sup>  $\Omega$  we can introduce Martin's topology and make it the Martin space<sup>2)</sup>  $\hat{\Omega}$ .  $\Omega$  is a dense open subset of  $\hat{\Omega}$  and the kernel

$$K(p, x) = \begin{cases} \frac{G(p, x)}{G(p, y_0)} & p \neq y_0 \\ 0 & p = y_0 \neq x \\ 1 & p = y_0 = x \end{cases}$$

can be extended continuously to  $(p, x) \in \hat{\Omega} \times \Omega$ , where  $G(p, x)$  is a Green function in  $\Omega$  and  $y_0$  the fixed point of  $\Omega$ .  $\hat{\Omega}$  is a metric space.  $\Delta = \hat{\Omega} - \Omega$  is divided into two disjoint subsets  $\Delta_0, \Delta_1$  and  $s \in \Delta_1$  is characterized by the fact that  $K(s, x)$  is a minimal positive harmonic function<sup>3)</sup> in  $x \in \Omega$ .

2. We shall show the following theorem:

**THEOREM.** *No point of  $\Delta_0$  is an isolated point.*

*Proof.* Let  $\omega$  be an open subset of  $\Omega$ ,  $\{x_n\}$  ( $n = 1, 2, \dots$ ) be a sequence of points in  $\omega$  such that  $x_n \rightarrow x_0 \in \Delta$ . If we denote by  $\mathcal{H}$  the family of positive superharmonic functions in  $\Omega$ , each of which dominates  $K(x_n, y)$  on  $\Omega - \omega$ , then  $\inf_{v \in \mathcal{H}} v(y)$  is equal to the positive superharmonic function except a polar set. We shall write this superharmonic function  $\mathcal{E}_{x_n}^\omega(y)$ .

In this case

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<sup>1)</sup> M. Brelot, G. Choquet, Espaces et lignes de Green. *Annales Inst. Fourier* 3 (1951), pp. 199-263.

<sup>2)</sup> M. Brelot, Le problème de Dirichlet. *Axiomatique et frontière de Martin*. *Journal de Math.* 35 (1956), pp. 297-335 (pp. 329-330). Cf. also R. S. Martin, Minimal positive harmonic functions, *Trans. Amer. Math. Soc.*, 49 (1941), pp. 137-172. M. Parreau, Sur les moyennes des fonctions harmoniques et analytiques et la classification des surfaces de Riemann, *Annales Inst. Fourier* 3 (1952), pp. 103-197. L. Naïm, Sur le rôle de la frontière de R. S. Martin dans la théorie du potentiel, *Annales Inst. Fourier* 7 (1957), pp. 183-281.

<sup>3)</sup> R. S. Martin, loc. cit., p. 137.

$$\mathcal{G}_{K_n}^{\omega}(y) = \int K(x, y) d\mu_n(x)$$

where  $\mu_n$  is a positive mass-distribution on  $\hat{\omega} \cap \Omega$  and the total mass of  $\mu_n$  does not exceed 1,  $\hat{\omega}$  being the boundary of  $\omega$  in  $\hat{\Omega}$ . By the theorem of choice, we can extract from  $\{\mu_n\}$  the subsequence  $\{\mu'_n\}$  such that  $\mu'_n$  converges vaguely to  $\mu$  and the carrier of  $\mu$  is contained in  $\overline{\hat{\omega} \cap \Omega}$ .

$$v(y) = \int K(x, y) d\mu(x)$$

is a positive superharmonic function in  $\Omega$ , and we have

$$\mathcal{G}_{K_{\omega_0}}^{\omega}(y) \leq v(y).$$

In fact, for fixed  $y \in \Omega$  and  $r > 0$  we shall denote by  $\varepsilon_y^{!r}$  the mass-distribution which can be obtained after sweeping out the unit mass on  $y$  to the exterior of the sphere (circle) of radius  $r$  and with center  $y$ . Then

$$U^r(x) = \int K(x, z) d\varepsilon_y^{!r}(z)$$

is bounded and continuous on  $\hat{\Omega}$ . Therefore

$$\lim_{n \rightarrow \infty} \int U^r(x) d\mu'_n(x) = \int U^r(x) d\mu(x).$$

By reciprocal law

$$\lim_{n \rightarrow \infty} \int \mathcal{G}_{K_{\omega_n}}^{\omega}(z) d\varepsilon_y^{!r}(z) = \int v(z) d\varepsilon_y^{!r}(z)$$

and by Fatou's lemma

$$\int \mathcal{G}_{K_{\omega_0}}^{\omega}(z) d\varepsilon_y^{!r}(z) \leq \liminf_{n \rightarrow \infty} \int \mathcal{G}_{K_{\omega_n}}^{\omega}(z) d\varepsilon_y^{!r}(z) = \int v(z) d\varepsilon_y^{!r}(z).$$

By making  $r \rightarrow 0$  we can get for each  $y \in \Omega$

$$\mathcal{G}_{K_{\omega_0}}^{\omega}(y) \leq v(y).$$

If  $\mu_1$  denotes the restriction of  $\mu$  to  $\mathcal{A}$  and  $\mu_2$  the restriction of  $\mu$  to  $\Omega$ , then

$$\begin{aligned} v(y) &= \int K(x, y) d\mu_1(x) + \int K(x, y) d\mu_2(x) \\ &= u(y) + w(y) \end{aligned}$$

where  $u$  is harmonic and  $w$  is a potential, and this is just the Riesz decomposi-

tion.

From now on let  $x_0$  be a point of  $A_0$  and  $x_0$  be isolated. Let

$$K(x_0, y) = \int_{A_1} K(x, y) d\nu(x)$$

be the canonical representation<sup>4)</sup> of  $K(x_0, y)$ . Then we can find a neighbourhood  $\hat{\delta}$  of  $x_0$  such that

$$(1) \quad \overline{\hat{\delta}} \cap A_0 = \{x_0\}$$

and

$$(2) \quad \nu(A_1 - \overline{\hat{\delta}}) > 0.$$

If we set  $\omega = \hat{\delta} \cap \Omega$ , then

$$\begin{aligned} \mathcal{E}_{Kx_0}^\omega(y) &= \int_{A_1} \mathcal{E}_{Kx}^\omega(y) d\nu(x) \\ &= \int_{A_1 - \overline{\hat{\delta}}} \mathcal{E}_{Kx}^\omega(y) d\nu(x) + \int_{\overline{\hat{\delta}} \cap A_1} \mathcal{E}_{Kx}^\omega(y) d\nu(x) \end{aligned}$$

The first term of the last side is harmonic, because  $\omega$  is thin at each point of  $A_1 - \overline{\hat{\delta}}$ <sup>5)</sup> and therefore we can get  $\mathcal{E}_{Kx}^\omega(y) \equiv K(x, y)$ .

We note that  $\mu_1$  is the restriction of the mass-distribution  $\mu$  to  $(\overline{\hat{\delta}} \cap \Omega) \cap A$ , which is contained in  $\overline{\hat{\delta}} \cap A$  and does not contain the point  $x_0$ . By (1) we can get  $\mu_1(A_0) = 0$ , that is,  $\mu_1$  is the canonical mass-distribution of  $u$ , and by (2)

$$u_1(y) = \int_{A_1 - \overline{\hat{\delta}}} K(x, y) d\nu(x) > 0.$$

Since

$$\begin{aligned} v(y) &= u(y) + w(y) \\ &\geq \mathcal{E}_{Kx_0}^\omega(y) \geq \int_{A_1 - \overline{\hat{\delta}}} \mathcal{E}_{Kx}^\omega(y) d\nu(x) = \int_{A_1 - \overline{\hat{\delta}}} K(x, y) d\nu(x) = u_1(y) \end{aligned}$$

and  $u$  is the greatest harmonic minorant of  $v$ , we have

$$u \geq u_1,$$

but the canonical mass-distribution of  $u$  has the carrier in  $A_1 \cap \overline{\hat{\delta}}$ , whereas the canonical mass-distribution of  $u_1$  has positive mass only in  $A_1 - \overline{\hat{\delta}}$ . As  $u_1 > 0$

<sup>4)</sup> R. S. Martin, loc. cit., p. 157.

<sup>5)</sup> L. Naïm, loc. cit., p. 203 (théorème 3) and p. 205 (théorème 5).

the canonical mass-distribution of  $\mu$  has positive mass in  $\Delta_1 - \bar{\delta}$ ; this is a contradiction. Q.E.D.

**COROLLARY.** *If  $\Delta_0 \neq \emptyset$  then  $\Delta_0$  contains at least countable points.*

*Remark.* In the above consideration we rely upon the following argument: we have always  $\mu \geq \mu_1$ , and, if  $\mu_1 > 0$ , then  $\mu_1$  is not the canonical mass-distribution of  $\mu$ .

Mr. K. Matsumoto has kindly pointed out the following result:

*Let  $x_0$  be a point of  $\Delta_0$  and  $\nu$  be the canonical mass-distribution of  $K_{x_0}(y)$ , then the common part of the carrier of  $\nu$  with  $\Delta_1$  is contained in  $\bar{\Delta}_0$ .*

The proof follows from the preceding remark; let  $E$  be the carrier of  $\nu$ . If  $E \cap \Delta_1 \not\subset \bar{\Delta}_0$ , then there exist a point  $z_0$  and a set  $A$  satisfying the following conditions:

- 1)  $A$  is an open neighbourhood of  $z_0$  in  $\Delta$ ,
- 2)  $\nu(A) > 0$ ,
- 3)  $A \subset \Delta_1$ .

We can construct an open set (in  $\hat{\mathcal{Q}}$ )  $G : G \cap \Delta = A$ . In this case, there exist two positive numbers  $0 < \rho_1 < \rho$  such that:

- i)  $\text{dist}(z_0, x_0) > \rho$ ,
- ii)  $\overline{U_\rho(z_0)} \subset G^{\delta_1}$ ,
- iii)  $\nu(U_{\rho_1}(z_0) \cap \Delta) > 0$ .

If we set  $\omega = \mathcal{Q} - \overline{U_\rho(z_0)}$ , under the same notations as in the proof of the above theorem, we see from i),  $x_0 \in \bar{\omega}$  and, as  $(\bar{\delta} \cap \mathcal{Q}) \cap \Delta \subset \overline{U_\rho(z_0)} \cap \Delta \subset G \cap \Delta = A \subset \Delta_1$ ,  $\mu_1$  is canonical and from  $\Delta - \bar{\omega} \supset U_{\rho_1}(z_0) \cap \Delta$ ,  $\mu_1 > 0$ , this is a contradiction.

*Osaka City University*

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<sup>6)</sup> We denote  $U_\rho(z_0) = \{x \in \hat{\mathcal{Q}}; \text{dist}(x, z_0) < \rho\}$ , where the metric  $\text{dist}(x, z_0)$  is the Martin's metric.