E-ENTROPY OF THE BROWNIAN MOTION WITH THE MULTI-DIMENSIONAL SPHERICAL PARAMETER

YOSHIKAZU BABA

§ 1. Introduction

M.S. Pinsker [3] has given a general method of calculating the ε -entropy of a Gaussian process and obtained, for example, an exact proof of the estimate for the ε -entropy of the ordinary Brownian motion B(t), $0 \le t \le 1$, which was presented without proof by A.N. Kolmogorov [1].

In this article, we estimate the ε -entropy of the Brownian motion with the multidimensional spherical parameter, by using the expansion of the Brownian motion with a multidimensional parameter by H.P. McKean [4] and by generalizing the Pinsker's method of calculating the ε -entropy.

Let $X(A, \omega)$, $A \in E^d$ (d-dimensional Euclidean space), $\omega \in \Omega(P)$, be a Brownian motion with a parameter space E^d , that is, $\{X(A), A \in E^d\}$ forms a Gaussian system and

- 1) E[X(A)] = 0 for every A,
- 2) X(O) = 0, where O is the origin of E^d ,
- 3) $E[(X(A) X(B))^2] = \operatorname{dis}(A, B)$, where E(X) and $\operatorname{dis}(A, B)$ denote the expectation of a random variable X and the Euclidean distance between A and B, respectively.

We shall call X(A) when the parameter A is restricted to the unit sphere¹⁾ S^{d-1} in E^d the Brownian motion with the d-dimensional spherical parameter and denote it, as in the preceding case, by X(A), $A \in S^{d-1}$.

The ε -entropy $H_{\varepsilon}(X)$ of the process X(A) is defined as follows: Let $\varepsilon > 0$ be arbitrarily fixed, and consider an approximating process X'(A) for the process X(A) on S^{d-1} satisfying the condition of reproducing accuracy,

(1)
$$\int_{S^{d-1}} E[(X'(A) - X(A))^2] d\sigma(A) \leq \varepsilon^2$$

Received March 8, 1967.

¹⁾ Without loss of generality we may consider the unit sphere only.

where $d\sigma$ is the uniform probability measure on S^{d-1} . Then, the ε -entropy of the process X(A) is defined as

$$(2) H_{\varepsilon}(X) = \inf I(X', X),$$

where I(X', X) is the amount of information contained in a process X' with respect to the process X and the infimum is taken for all processes X' satisfying the condition (1).

Our aim is to prove that the ε -entropy of the Brownian motion on S^{d-1} is of order $\varepsilon^{-2(d-1)}$ (Theorem 2);

$$(3) H_{\varepsilon}(X) = O(\varepsilon^{-2(d-1)}).$$

It seems to be interesting to note that the ε -entropy (in Kolmogorov-Tihomirov's sense, cf. Kolmogorov-Tihomirov [2]) of the space of $\frac{1}{2}$ -Hölder continuous functions of (d-1)-variables with the sup-norm has the same order $O(\varepsilon^{-2(d-1)})$.

The author is greatly indebted to Professors T. Hida and N. Ikeda for their kind suggestions and constant encouragement.

§ 2. The generalization of Pinsker's method

Pinsker's method of calculating the ε -entropy of a Gaussian process with one dimensional parameter is as follows: Let X(t), $0 \le t \le T$, be a Gaussian process with mean 0 whose covariance function r(s,t) = E[X(s)X(t)] is continuous in (s,t). Then the ε -entropy $H_{\varepsilon}(X)$ of the process X(t) is given by the formula

(4)
$$H_{\epsilon}(X) = \frac{1}{2} \sum_{\lambda_i > \theta^2} \log \frac{\lambda_i}{\theta^2} ,$$

where λ_i $(i=1,2,\cdots)$ are the eigen-values of the integral operator with the kernel r(s,t) in $L^2[0,T]$, $\lambda_1 \ge \lambda_2 \ge \cdots \ge 0$, and θ is determined (uniquely) by the equation

(5)
$$\sum_{i=1}^{\infty} \min \left(\theta^2, \lambda_i\right) = \varepsilon^2 \cdot 2$$

$$\sum_{i=1}^{\infty} \lambda_i = \sum_{i=1}^{\infty} \lambda_i \int_0^T [\varphi_i(t)]^2 dt = \int_0^T \sum_{i=1}^{\infty} \lambda_i [\varphi_i(t)]^2 dt = \int_0^T r(t,t) dt < \infty.$$

²⁾ By Mercer's theorem

The right-hand side of the relation (4) also equals to the ε -entropy of the infinite dimensional Gaussian random variable $X^* = (X_1^*, X_2^*, \cdots)^3$:

(6)
$$X_i^* = \int_0^T \varphi_i(t)X(t)dt^4$$
 $(i = 1, 2, \cdots)$

where $\varphi_i(t)$ is the eigen-function of the integral operator corresponding to the eigenvalue λ_i and $E[X_i^* X_j^*] = \lambda_i \delta_{ij}$.

As an example, if in particular the sequence $\lambda_1 \ge \lambda_2 \ge \cdots \ge 0$ of the eigen-values of the integral operator with the kernel corresponding to a Gaussian process takes the form: $\lambda_k = ck^{-s}(s > 1; k = 1, 2, \cdots)$, then, the ε -entropy of the process is

(7)
$$H_{\varepsilon}(X) = O(\varepsilon^{-\frac{2}{s-1}}).$$

Now, we proceed to a Gaussian process X(A), $A \in S^{d-1}$, with mean 0. Assume the continuity of the covariance function r(A, B) = E[X(A) X(B)] in $S^{d-1} \times S^{d-1}$, so $\sum_{i=1}^{\infty} \lambda_i$ is finite (see the discussion in the footnote 2)) where λ_i , $i = 1, 2, \cdots$, are the eigenvalues of the integral operator with the kernel r(A, B) in $L^2(S^{d-1}, d\sigma)$. Then, the following entirely analogous result holds, and we state it as a theorem.

Theorem 1. The ε -entropy $H_{\varepsilon}(X)$ of the above Gaussian process X(A), $A \in S^{d-1}$ is

(4')
$$H_{\varepsilon}(X) = \frac{1}{2} \sum_{i \geq \theta^2} \log \frac{\lambda_i}{\theta^2}$$

where $\lambda_i (i = 1, 2, \cdots)$ with $\lambda_1 \ge \lambda_2 \ge \cdots \ge 0$ are eigen-values of the integral operator and θ is determined by the equation (5). The right-hand side of the relation (4') equals also to the ε -entropy of the infinite dimensional Gaussian random variable $X^* = (X_1^*, X_2^*, \cdots)$:

(6')
$$X_i^* = \int_{S^{d-1}} \varphi_i(A) X(A) d\sigma(A) \qquad (i = 1, 2, \cdots)$$

³⁾ The ε -entropy of X^* is defined as $H_{\varepsilon}(X^*) = \inf I(\widetilde{X}^*, X^*)$ where the infimum is taken for all infinite dimensional approximating random variables $\widetilde{X}^* = (\widetilde{X}_1^*, \widetilde{X}_2^*, \cdots)$ satisfying the condition: $\sum_{i=1}^{\infty} E[(\widetilde{X}_i^* - X_i^*)^2] \leq \varepsilon^2$.

⁴⁾ This (Bochner) integral is determined as an element of $L^2(\Omega)$.

where $\varphi_i(A)$ is the eigen-function of the integral operator corresponding to the eigenvalue λ_i , and $E[X_i^*X_j^*] = \lambda_i \delta_{ij}$.

Proof. The proof is quite similar to the proof for one dimensional parameter case dealt by M.S. Pinsker [3], except for the construction of the process $\dot{\xi}$ ([3], formula (132)). The proof, however, can be carried out by using the extension theorem of Urysohn, so that we shall not continue the proof further.

§ 3. The main result

We are now in a position to prove our main result.

THEOREM 2. The ε -entropy of the Brownian motion with the d-dimensional spherical parameter is of order $\varepsilon^{-2(d-1)}$;

(8)
$$H_{\varepsilon}(X) = O(\varepsilon^{-2(d-1)}).$$

Proof. According to H.P. McKean [4] the Brownian motion with the d-dimensional parameter can be expanded as a sum of mutually independent Gaussian processes associated with spherical harmonics. We state this expansion and some related results with the Gaussian process X(A), $A \in S^{d-1}$.

(9)
$$X(A) = \sum_{n \ge 0} \sum_{l=1}^{D(n)} x_n^l(1) h_n^l(A), A \in S^{d-1}$$

where $h_n^t(A)$ is a spherical harmonics of degree n satisfying

(10)
$$\int_{S^{d-1}} h_n^l(A) h_m^k(A) d\sigma(A) = \begin{cases} 1, & \text{if } l = k, \ n = m \\ 0, & \text{otherwise,} \end{cases}$$

D(n) is the dimension of the vector space spanned by all the spherical harmonics of degree n,

(11)
$$D(n) = (2n - 2 + d) \frac{(n - 3 + d)!}{(d - 2)! \, n!} \qquad (d \ge 2, \ n \ge 0)^{5}$$

and $x_n^l(1)$ $(n \ge 0, 1 \le l \le D(n))$ are mutually independent Gaussian random variables which can be expressed in the form

(12)
$$x_n^l(1) \equiv x_n^l = C(d) \int_0^1 C_n(u) dB_n^l(u) .$$

⁵⁾ For d=2 and n=0, D(n)=1.

The processes $B_n^l(u)$ $(n \ge 0, 1 \le l \le D(n))$ appeared in the above expression are mutually independent standard Brownian motions and

(13)
$$C_n(u) = \frac{\int_0^{\cos^{-1}u} p_n(\cos\theta) \sin^{d-2}\theta d\theta}{\int_0^{\pi} \sin^{d-2}\theta d\theta} , \quad n \ge 0$$

with $p_n(\cos\theta) = C_n^{\frac{d-2}{2}}(\cos\theta) \bigg/ C_n^{\frac{d-2}{2}}(1)$, where $C_n^{\nu}(\cdot)$ is the Gegenbauer polynomial and C(d) is a constant depending only on d.

By the expansion (9) and by the independence of the random variables x_n^l with $E[x_n^l] = 0$ $(n \ge 0, 1 \le l \le D(n))$ we easily see that the covariance function of the process X(A) is expressed in the form

(14)
$$r(A,B) = \sum_{n\geq 0} \sum_{l=1}^{D(n)} E[(x_n^l)^2] h_n^l(A) h_n^l(B) .$$

Using this, Mercer's expansion theorem shows us that the eigen-values $\lambda_n^l(n \ge 0, 1 \le l \le D(n))$ of the integral operator with the kernel r(A, B) are equal to $E[(x_n^l)^2]$. Therefore, if we know the amount $E[(x_n^l)^2]$ we can obtain the ε -entropy of the Brownian motion with the parameter space S^{d-1} by the formula (4'). In fact, we can prove in the following that for large n, $E[(x_n^l)^2] = O(n^{-d})$, $1 \le l \le D(n)$, holds. Once the result is shown, then just by renumbering the double sequence of random variables x_0^1 , x_1^1 , x_1^2 , \cdots , $x_n^{D(1)}$, x_2^1 , \cdots into the ordinary sequence x_1' , x_2' , \cdots , while keeping the original order, we can easily apply Theorem 1 in §2. If x_k' , for large k, corresponds to the original random variable $x_n^M(1 \le M \le D(N))$, then by the relation $\sum_{n=0}^{N} n^{d-2} = O(N^{d-1})$ (this nearly equals to k) and by the formula (11) $D(n) = O(n^{d-2})$ for large n), we obtain $N = O(n^{d-1})$ so that

formula (11) $(D(n) = O(n^{d-2})$ for large n), we obtain $N = O(k^{\frac{1}{d-1}})$, so that $E[(x_k')^2] = O((k^{\frac{1}{d-1}})^{-d}) = O(k^{-\frac{d}{d-1}})$. Then, by this and the formula (7), follows the desired result $H_{\epsilon}(X) = O(\epsilon^{-\frac{2}{d-1}-1}) = O(\epsilon^{-2(d-1)})$.

Therefore, in the following, we are to prove that

(15)
$$E[(x_n^i)^2] = O(n^{-d}), \quad 1 \le l \le D(n)$$

holds for large n.

First of all, we show the formula (15) in case the dimension d=2 and 3, and then, generalizing it, we proceed to prove the formula (15) for $d \ge 4$, that is, (I) in case d is an even integer and (II) when d is odd.

In case d=2, $p_n(\cos\theta)$ in the expression (13) turns out to be $\cos n\theta$, so that $C_n(u) = \frac{1}{n\pi} \sin(n\cos^{-1}u)$. From this we have,

$$E[(x_n^l)^2] = \frac{1}{n^2\pi} \int_0^1 \sin^2(n \cos^{-1}u) du$$
$$= \frac{1}{n^2\pi} \int_0^{\frac{\pi}{2}} \sin^2 n\theta \sin \theta d\theta$$
$$= O(n^{-2}).$$

While in case d=3, $p_n(\cos\theta)=P_n(\cos\theta)$, hence we have $C_n(u)=\frac{1}{2}\frac{P_{n-1}(u)-P_{n+1}(u)}{2n+1}$ where $P_n(\cdot)$ is the n-th Legendre polynomial. Then, by the orthogonality of the Legendre polynomials, we obtain

$$\begin{split} E\left[(x_n^l)^2\right] &= \frac{1}{(2n+1)^2} \left\{ \int_0^1 (P_{n+1}(u))^2 \ du + \int_0^1 (P_{n-1}(u))^2 \ du \right\} \\ &= O(n^{-3}) \,. \end{split}$$

In case $d \ge 4$, by the formula (12), we have

$$\begin{split} E\left[(x_n^1)^2\right] &= (C(d))^2 \int_0^1 (C_n(u))^2 du \\ &= (\text{a constant depending on } d \text{ only}) \times \left\{ C_n^{\frac{d-2}{2}}(1) \right\}^{-2} \\ &\times \int_0^1 \left\{ \int_0^{\cos^{-1}u} C_n^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta \right\}^2 du \end{split}$$

and this expression becomes,

$$O(n^{-2d+6}) \times \int_0^1 \left\{ \int_0^{\cos^{-1}u} C_n^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta \right\}^2 du$$

for large
$$n$$
, since $C_n^{\frac{d-2}{2}}(1) = \frac{\Gamma(n+d-2)}{n! \Gamma(d-2)} = O(n^{d-3})$.

To prove $E[(x_n^{l})^2] = O(n^{-d})$, we must show that the above integral (we denote it by I_a) is of order $O(n^{d-6})$.

(I) The proof of the fact that $I_d = O(n^{d-6})$ for d = 2p + 2 $(p \ge 1$, integer).

First we estimate the integrand of the above integral. Let the following integral be denoted by $I_p(u)$,

$$I_p(u) = \int_0^{\cos^{-1}u} C_n^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta = \int_0^{\cos^{-1}u} C_n^p(\cos\theta) \sin^{2p}\theta d\theta.$$

The integrand $C_n^p(\cos \theta) \sin^{2p} \theta$ of the above integral becomes, by using the recurrence formula for the Gegenbauer polynomials

(16)
$$\sin^2 \theta C_n^{\nu+1} (\cos \theta) = \frac{1}{2\nu} \left\{ (n+2\nu) C_n^{\nu} (\cos \theta) - (n+1) \cos \theta C_{n+1}^{\nu} (\cos \theta) \right\}$$

and the formula $\sin \theta C_n^1(\cos \theta) = \sin (n+1)\theta$,

$$C_{n}^{p}(\cos\theta)\sin^{2p}\theta = \sin^{2}\theta C_{n}^{p}(\cos\theta)\sin^{2(p-1)}\theta$$

$$= \frac{1}{2(p-1)} \left\{ (n+2(p-1))C_{n}^{p-1}(\cos\theta)\sin^{2(p-1)}\theta - (n+1)\cos\theta C_{n+1}^{p-1}(\cos\theta)\sin^{2(p-1)}\theta \right\}$$

$$= \frac{1}{2^{p-1}(p-1)!} \left\{ A_{1}^{p}(n)\sin\theta\sin(n+1)\theta + A_{2}^{p}(n)\cos\theta\sin\theta\sin(n+2)\theta + A_{3}^{p}(n)\cos^{2}\theta\sin\theta\sin(n+3)\theta + \dots + A_{p}^{p}(n)\cos^{2}\theta\sin\theta\sin(n+p)\theta \right\}$$

where $A_1^p(n), A_2^p(n), \dots, A_p^p(n)$ are polynomials of n of order (p-1). Noticing that $\sin \theta \sin (n+1)\theta, \cos \theta \sin \theta \sin (n+2)\theta, \dots$ and $\cos^{p-1}\theta \sin \theta \sin (n+p)\theta$ are all expressed as the linear combinations of $\cos n\theta$, $\cos (n+2)\theta$, \dots , $\cos (n+2p)\theta$, we can show that the integral becomes

(17)
$$I_p(u) = \sum_{k=0}^{p} \frac{B_k^p(n)}{n+2k} \sin{(n+2k)\alpha}, \quad \alpha = \cos^{-1}u$$

where B_k^p , $k = 0, 1, \dots, p$, are polynomials of n of order at most (p-1). Therefore, changing the variable of integration into α , and making use of the fact

$$\int_0^{\frac{\pi}{2}} \sin{(n+2k)\alpha} \sin{(n+2l)\alpha} \sin{\alpha} d\alpha = \frac{1}{2} \left\{ \frac{1}{1-4(k-l)^2} + O(n^{-2}) \right\} ,$$

we have

$$I_{d} = \int_{0}^{1} \{I_{p}(u)\}^{2} du = \int_{0}^{\frac{\pi}{2}} \left\{ \sum_{k=0}^{p} \frac{B_{k}^{p}(n)}{n+2k} \sin{(n+2k)\alpha} \right\}^{2} \sin{\alpha} d\alpha$$

$$= \sum_{k,l=0}^{p} \frac{B_{k}^{p}(n)B_{l}^{p}(n)}{(n+2k)(n+2l)} \int_{0}^{\frac{\pi}{2}} \sin(n+2k)\alpha \sin(n+2l)\alpha \sin\alpha d\alpha$$
$$= O(n^{2p-4}) (=O(n^{d-6})).$$

The last estimation is valid if the coefficient of the term n^{2p-4} never vanishes, that is, if at least one of the coefficients of the term n^{p-1} of the polynomials $B_k^p(n)$ $(k=0,1,\dots,p)$ does not vanish. But this is true, for example, $B_0^p(n)$ has non zero coefficient of n^{p-1} .

(II) The proof of the fact that $I_d = O(n^{d-6})$ for d = 2p + 3 ($p \ge 1$, integer).

Similarly to (I), we denote the following integral by $I_p(u)$,

$$I_{p}(u) = \int_{0}^{\cos^{-1}u} C_{n}^{\frac{d-2}{2}}(\cos\theta) \sin^{d-2}\theta d\theta = \int_{0}^{\cos^{-1}u} C_{n}^{p+\frac{1}{2}}(\cos\theta) \sin^{2p+1}\theta d\theta$$

then, by the relation

(18)
$$C_n^{p+\frac{1}{2}}(\cos\theta) = \frac{2^p p!}{(2p)! \sin^p \theta} P_{n+p}^p(\cos\theta)$$

for the half-integer Gegenbauer polynomial $C_n^{p+\frac{1}{2}}$ and the associated Legendre polynomial P_{n+p}^p , we have

$$I_p(u) = c(d) \int_0^{\cos^{-1} u} P_{n+p}^p(\cos \theta) \sin^{p+1} \theta d\theta$$

where c(d) is a constant depending on d. By definition,

$$P_{n+p}^{p}(x) = (1-x^{2})^{\frac{p}{2}} \frac{d^{p}}{dx^{p}} P_{n+p}(x)$$

and by changing the variable of integration into $x = \cos \theta$, we get

$$\begin{split} \frac{1}{c(d)} I_p(u) &= \int_u^1 \frac{d^p}{dx^p} P_{n+p}(x) (1-x^2)^p dx \\ &= - (1-u^2)^p \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) + 2p \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \end{split}$$

From this, the desired integral I_d is

$$\begin{split} & [c(d)]^2 \cdot I_d = [c(d)]^2 \int_0^1 \{I_p(u)\}^2 \, du \, = \int_0^1 \left\{ (1-u^2)^p \, \frac{d^{p-1}}{du^{p-1}} \, P_{n+p}(u) \, \right\}^2 du \\ & (19) \qquad -4p \int_0^1 (1-u^2)^p \, \frac{d^{p-1}}{du^{p-1}} \, P_{n+p}(u) \, \Big\{ \int_u^1 x (1-x^2)^{p-1} \, \frac{d^{p-1}}{dx^{p-1}} \, P_{n+p}(x) dx \, \Big\} du \end{split}$$

$$+4p^2\int_0^1\!\!\left\{\int_u^1x(1-x^2)^{p-1}\,\frac{d^{p-1}}{dx^{p-1}}\!\!-\!P_{n+p}(x)dx\right\}^2\!du.$$

To estimate these integrals, we first express $(1-u^2)^p \frac{d^p}{du^p} P_n(u)$ in terms of $P_n(u)$ and $P_{n-1}(u)$. For this purpose, we make use of the recurrence formula of the Legendre polynomials $(1-x^2)P_n'(x) = n(P_{n-1}(x)-xP_n(x))$ and the differential equation derived from the Legendre's differential equation

$$(20) (1-x^2) \frac{d^k}{dx^k} P_n(x) - 2(k-1)x \frac{d^{k-1}}{dx^{k-1}} P_n(x)$$

$$+ (n+(k-1)) (n-(k-2)) \frac{d^{k-2}}{dx^{k-2}} P_n(x) = 0, (k \ge 2).$$

For any $p \ge 1$, we have

$$(21) (1-u^2)^p - \frac{d^p}{du^p} P_n(u) = P_{n-1}(u)Q_{n-1,p}(u) + P_n(u)Q_{n,p}(u)$$

where $Q_{n-1,p}(u)$ and $Q_{n,p}(u)$ are polynomials of u of the form

(22)
$$Q_{n-1,p}(u) = \sum_{k=0}^{p-1} C_k(n) u^k, \quad Q_{n,p}(u) = \sum_{k=0}^{p} D_k(n) u^k.$$

The coefficients $C_0(n)$, $C_1(n)$, \cdots , $C_{p-1}(n)$, $D_0(n)$, $D_1(n)$, \cdots , $D_p(n)$ have the following properties: (i) $C_{p-1}(n) \neq 0$, $D_p(n) \neq 0$ (ii) they are the polynomials of n with the order at most p (iii) if p is an even integer, then $D_0(n)$ is the polynomial of order p and if p is odd, $C_0(n)$ is the polynomial of order p. By these facts and by the property of the Legendre polynomial: $\int_0^1 \{P_n(x)\}^2 dx = O(n^{-1})$ for large n, we can easily show that the first integral of the right-hand side of the equality (19) becomes,

$$\int_0^1 \left\{ (1-u^2)^p \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \right\}^2 du = \int_0^1 (1-u^2)^2 \left\{ (1-u^2)^{p-1} \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \right\}^2 du = O(n^{2(p-1)}) \cdot O(n^{-1}) = O(n^{d-6}).$$

For the second integral of the right-hand side of (19), we have

$$\begin{split} & \left| \int_0^1 (1-u^2)^p \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\} du \right| \\ & \leq & \left\{ \int_0^1 \left\{ (1-u^2)^p \frac{d^{p-1}}{du^{p-1}} P_{n+p}(u) \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_0^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_0^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_0^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_0^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_0^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_0^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_0^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 du \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^{1/2} \cdot \left\{ \int_u^1 \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^{1/2} \cdot \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^{1/2} \cdot \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^{1/2} \cdot \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^{1/2} \cdot \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^2 dx \right\}^2 dx + \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^2 dx + \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{p-1}} P_{n+p}(x) dx \right\}^2 dx \right\}^2 dx + \left\{ \int_u^1 x (1-x^2)^{p-1} \frac{d^{p-1}}{dx^{$$

The first term of the product on the right-hand side of the inequality, by the above result, has the order $O(n^{\frac{d-6}{2}})$ and the integrand of the second term can be evaluated as follows:

$$\begin{split} \left\{ \int_{u}^{1} x (1-x^{2})^{p-1} \, \frac{d^{p-1}}{dx^{p-1}} \, P_{n+p}(x) dx \right\}^{2} & \leqq \int_{u}^{1} x^{2} dx \cdot \int_{u}^{1} \left\{ (1-x^{2})^{p-1} \, \frac{d^{p-1}}{dx^{p-1}} \, P_{n+p}(x) \, \right\}^{2} dx \\ & < \int_{0}^{1} x^{2} dx \cdot \int_{0}^{1} \left\{ (1-x^{2})^{p-1} \, \frac{d^{p-1}}{dx^{p-1}} \, P_{n+p}(x) \, \right\}^{2} dx = O(n^{d-6}) \, . \end{split}$$

Hence the second integral is at most of order $O(n^{d-6})$. As for the last integral of the equality (19), by a similar approach, we estimate it to be at most of order $O(n^{d-6})$. This proves the desired result for $d=2p+3(p\ge 1)$, and thus we have proved the theorem completely.

REFERENCES

- [1] A.N. Kolmogorov: Theory of the transmission of information, Amer. Math. Soc. Translations Ser. 2, 33 (1963), 291–321.
- [2] A.N. Kolmogorov and V.M. Tihomirov: ε-entropy and ε-capacity of sets in functional spaces, Amer. Math. Soc. Translations Ser. 2, 17 (1961), 277-364.
- [3] М.С. Пинскер, Гауссовские источники, Проблемы Перебачи Информации, 14 (1963) 59–100.
- [4] H.P. McKean: Brownian motion with a several-dimensional time, Theory Prob. Appl., 8-4 (1963), 335-354.
- [5] A. Erdélyi and others: Higher transcendental functions I-III, McGraw-Hill Publ. New York (1953–1955).

Kobe University of Commerce