T. OnoNagoya Math. J.Vol. 92 (1983), 133-144

ON A GENERALIZATION OF LAPLACE INTEGRALS

TAKASHI ONO

Introduction

Let \mathbf{R}^n be the Euclidean space of dimension $n \geq 1$ with the standard inner product $\langle x, y \rangle = \sum x_i y_i$ and the norm $|x| = \langle x, x \rangle^{1/2}$, S^{n-1} be the unit sphere $\{x \in \mathbf{R}^n; |x| = 1\}$ and $d\omega_{n-1}$ be the volume element of S^{n-1} such that S^{n-1} gets the volume 1. Let Ω be an open set of \mathbf{R}^n containing S^{n-1} and let $f: \Omega \to \mathbf{R}^m$ be a smooth map. With each integer $\nu \geq 0$, we shall associate a form f_{ν} of degree ν on \mathbf{R}^m defined by

$$(0.1) f_{\nu}(\xi) = \int_{S^{n-1}} \langle \xi, f(x) \rangle^{\nu} d\omega_{n-1} , \xi \in \mathbf{R}^{m} .$$

We then consider the number $\sigma_{\nu}(f)$ which is the mean value of the form f_{ν} on the sphere S^{m-1} :

(0.2)
$$\sigma_{\nu}(f) = \int_{S^{m-1}} f_{\nu}(\xi) d\omega_{m-1} , \qquad \nu \in Z_{+} .$$

When f is an affine map: $\mathbb{R}^n \to \mathbb{R}^m$, the function f_{ν} is substantially the Legendre polynomial of order ν and (0.1) is the Laplace integral for it¹⁾. Therefore it is natural to ask questions about forms f_{ν} associated with more general map f.

In this paper, we shall focus our attention on the determination of the number $\sigma_{\nu}(f)$ for any smooth map f. It will turn out that the main ingredient of the number $\sigma_{\nu}(f)$ is the number:

(0.3)
$$N_k(f) = \int_{S^{n-1}} |f(x)|^{2k} d\omega_{n-1}, \qquad \nu = 2k^{2}$$
.

Since $N_{k}(f)=1$ whenever f maps S^{n-1} into S^{m-1} , we see that all these "spherical" maps share the same numbers $\sigma_{\nu}(f)$ for all $\nu \in \mathbb{Z}_{+}$; hence these numbers measure a deviation of f from being spherical. We shall consider

Received August 23, 1982.

¹⁾ See Appendix for a detailed discussion on this matter.

²⁾ As is easily seen, $\sigma_{\nu}(f) = 0$ if ν is odd.

examples of a family of maps $\{f_{\rho}\}_{\rho\in R}$ for which $\{f_{\pm 1}\}$ are Hopf maps and show that the number $N_{k}(f_{\rho})$ can be written as a hypergeometric polynomial.

The author would like to mention here that the idea of associating the number like $\sigma_{\nu}(f)$ with a map f came from his earlier work [7] on functions over finite fields.

Notation and conventions

The symbols Z, Q, R, C denote the set of integers, rational numbers, real numbers and complex numbers. The set of non-negative real numbers is denoted by R_+ . We put $Z_+ = Z \cap R_+$, $Q_+ = Q \cap R_+$. The set of all multi-indices $\alpha = (\alpha_1, \dots, \alpha_n)$ is Z_+^n . We denote by 1_n the multi-index $(1, \dots, 1) \in Z_+^n$. For α , $\beta \in Z_+^n$ and $x = (x_1, \dots, x_n) \in R^n$, $|\alpha| = \alpha_1 + \dots + \alpha_n$, $\alpha! = \alpha_1! \dots \alpha_n!$, $m\alpha = (m\alpha_1, \dots, m\alpha_n)$, $m \in Z_+$, $x^\alpha = x_1^{\alpha_1} \dots x_n^{\alpha_n}$, $\alpha \leq \beta \Leftrightarrow \alpha_i \leq \beta_i$, $1 \leq i \leq n$. For an integer m, $\alpha \equiv 0$ $(m) \Leftrightarrow \alpha = m\beta$ for some β . When $\beta \leq \alpha$ we put

$$\begin{pmatrix} \alpha \\ \beta \end{pmatrix} \stackrel{\text{def}}{=} \frac{\alpha!}{\beta! (\alpha - \beta)!} = \begin{pmatrix} \alpha_1 \\ \beta_1 \end{pmatrix} \cdots \begin{pmatrix} \alpha_n \\ \beta_n \end{pmatrix}.$$

For $a \in C$, $n \in \mathbb{Z}_+$ we use Appell's notation $(a, n) = a(a + 1) \cdots (a + n - 1)$ for $n \ge 1$ and (a, 0) = 1. For $a, b, c \in C$, the hypergeometric series is defined by

$$F(a, b; c; z) = \sum_{n=0}^{\infty} \frac{(a, n)(b, n)}{(c, n)} \frac{z^n}{n!}$$
.

For a smooth function f on an open set of \mathbb{R}^n , we put

$$\operatorname{grad} f = \left(\frac{\partial f}{\partial x_1}, \, \cdots, \, \frac{\partial f}{\partial x_n}\right), \qquad \varDelta f = \frac{\partial^2 f}{\partial x_1^2} \, + \, \cdots \, + \, \frac{\partial^2 f}{\partial x_n^2} \; .$$

We shall use the following formulas freely:

$$(0.4) \qquad (x_1 + \cdots + x_n)^{\nu} = \sum_{|\alpha| = \nu} \frac{\nu!}{\alpha!} x^{\alpha} , \quad \nu \in \mathbb{Z}_+ , \quad x = (x_1, \cdots, x_n) \in \mathbb{R}^n ,$$

(0.5) when
$$|\alpha| = 2\nu$$
, we have $\frac{\Delta^{\nu} x^{\alpha}}{\nu!} = \frac{(2\beta)!}{\beta!}$ if $\alpha = 2\beta$, =0 if $\alpha \not\equiv 0$ (2).

§ 1. Numbers $b_{\nu}(\ell; \lambda)$

Let $\nu \geq 0$, $\ell \geq 1$ be integers and let $\lambda = (\lambda_1, \dots, \lambda_n) \in \mathbb{R}^n$. We assume that $\ell \nu$ is even: $\ell \nu = 2k$, $k \in \mathbb{Z}_+$. We define a number $b_{\nu}(\ell; \lambda) \in \mathbb{Q}_+$ by

$$(1.1) b_{\nu}(\ell;\lambda) = \frac{\nu!}{k!} \sum_{\substack{|\alpha|=k \\ \alpha=0(\ell)}} \frac{(2\alpha)!}{\alpha! (2\alpha/\ell)!} \lambda^{2\alpha/\ell}.$$

If, in particular, $\ell = 2$, then $k = \nu$ and

$$(1.2) b_{\nu}(2;\lambda) = \sum_{|\alpha|=\nu} {2\alpha \choose \alpha} \lambda^{\alpha}.$$

If, $\lambda = 1_n$ in (1.2), we have

$$(1.3) b_{\nu}(2; 1_n) = \sum_{|\alpha|=\nu} {2\alpha \choose \alpha}.$$

Using the equality

$$4^k(\frac{1}{2}, k) = k! \binom{2k}{k}$$

which one verifies easily, we get the following equality as formal power series in t

(1.4)
$$\sum_{k=0}^{\infty} {2k \choose k} t^k = (1-4t)^{-1/2}$$

and, by (1.2), (1.4), we get

(1.5)
$$\sum_{\nu=0}^{\infty} b_{\nu}(2;\lambda)t^{\nu} = \prod_{i=1}^{n} (1 - 4\lambda_{i}t)^{-1/2}.$$

In particular, we have

(1.6)
$$\sum_{n=0}^{\infty} b_{\nu}(2; 1_n) t^{\nu} = (1 - 4t)^{-n/2}$$

and hence, by (1.3), (1.6), we get

(1.7)
$$4^{\nu}\left(\frac{n}{2},\nu\right) = \nu! \ b_{\nu}(2;1_n) \ .$$

§2. Review of a mean value theorem in potential theory

Let $\varphi(x)$ be a complex valued smooth function defined on an open set Ω in \mathbb{R}^n containing S^{n-1} . Assume that either (i) $\Delta^m \varphi = 0$ for some

 $m \ge 1$ or (ii) $\Delta \varphi = \lambda \varphi$ for a constant $\lambda \in C$. In this situation, a mean value theorem in potential theory³⁾ tells us that

(2.1)
$$\int_{S^{n-1}} \varphi(x) d\omega_{n-1} = \sum_{\nu=0}^{\infty} \frac{(\Delta^{\nu} \varphi)(0)}{4^{\nu} \nu! (n/2, \nu)}.$$

Needless to say, when φ is harmonic, then m=1 in (i) or $\lambda=0$ in (ii) and (2.1) is the mean value theorem of Gauss:

$$\int_{S^{n-1}} \varphi(x) d\omega_{n-1} = \varphi(0) .$$

By (1.7), (2.2) can also be written as:

(2.2)
$$\int_{S^{n-1}} \varphi(x) d\omega_{n-1} = \sum_{\nu=0}^{\infty} \frac{(\Delta^{\nu} \varphi)(0)}{(\nu!)^2 b_{\nu}(2; 1_{\nu})}.$$

If φ is a form of even degree $\ell = 2k$, then $\Delta^m \varphi = 0$ for m > k and since $\Delta^m \varphi(0) = 0$ for m < k, we get from (2.2)

(2.3)
$$\int_{S^{n-1}} \varphi(x) d\omega_{n-1} = \frac{\Delta^k \varphi}{(k!)^2 b_k(2; 1_n)}.$$

This shows also that

$$\int_{S^{n-1}} \varphi(x) d\omega_{n-1} = 0$$

if the degree of φ is odd, a fact which can be proved directly. On the other hand, in case (ii), we have

(2.4)
$$\int_{S^{n-1}} \varphi(x) d\omega_{n-1} = \varphi(0) \sum_{\nu=0}^{\infty} \frac{\lambda^{\nu}}{(\nu!)^{2} b_{\nu}(2; 1_{n})} = \varphi(0) \frac{\Gamma(n/2)}{(\sqrt{-\lambda}/2)^{(n-2)/2}} J_{(n-2)/2}(\sqrt{-\lambda})$$

where $J_{\nu}(z)$ is the ν -th Bessel function.

§3. Mean value of quadratic forms

Let

(3.1)
$$\varphi(x) = \sum_{|\beta|=2k} c_{\beta} x^{\beta}$$

be a form of even degree 2k. By (0.5), (2.3), we have

³⁾ See Courant-Hilbert [3], pp. 258-261.

(3.2)
$$\int_{S^{n-1}} \varphi(x) d\omega_{n-1} = \frac{(1/k!) \sum_{|\alpha|=k} c_{2\alpha}((2\alpha)!/\alpha!)}{b_k(2; 1_n)}$$

Consider now a diagonal form $f(x) = \lambda_1 x_1^{\ell} + \cdots + \lambda_n x_n^{\ell}$ and put $\varphi(x) = f(x)^{\nu}$ with $\ell \nu = 2k$. Since we have

$$\varphi(x) = \sum_{|\sigma| = \nu} \frac{\nu!}{\sigma!} \lambda^{\sigma} x^{\ell \sigma} = \sum_{|\beta| = 2k} c_{\beta} x^{\beta}$$

where

(3.3)
$$c_{\beta} = \frac{\nu!}{\sigma!} \lambda^{\sigma} \quad \text{if } \beta = \ell \sigma, =0 \text{ if } \beta \not\equiv 0 \ (\ell) \ ,$$

by (3.2), (3.3), we have

(3.4)
$$\int_{S^{n-1}} f(x)^{\nu} d\omega_{n-1} = \frac{(1/k!) \sum_{\substack{|\alpha|=k \\ 2\alpha=\delta\sigma}} (\nu!/\sigma!) \lambda^{\sigma}((2\alpha)!/\alpha!)}{b_{\nu}(2;1_{n})}$$

From (1.1), (3.4), it follows that

$$(3.5) \qquad \int_{S^{n-1}} (\lambda_1 x_1^{\ell} + \cdots + \lambda_n x_n^{\ell})^{\nu} d\omega_{n-1} = \frac{b_{\nu}(\ell; \lambda)}{b_{\nu}(2; 1_{\nu})}, \qquad \ell \nu = 2k.$$

If, in particular, $\ell = 2$, then $\nu = k$ and we have

(3.6)
$$\int_{S^{n-1}} (\lambda_1 x_1^2 + \cdots + \lambda_n x_n^2)^{\nu} d\omega_{n-1} = \frac{b_{\nu}(2; \lambda)}{b_{\nu}(2; 1_{\nu})}$$

Consider a quadratic form q(x) on \mathbb{R}^n and the integral

$$(3.7) \qquad \int_{S^{n-1}} q(x)^{\nu} d\omega_{n-1}.$$

Since the change of variable $x \mapsto sx$, $s \in O(\mathbb{R}^n)$, the orthogonal group, does not change the integral (3.7) and q(x) can be brought to a diagonal form $\lambda_1 x_1^2 + \cdots + \lambda_n x_n^2$ by such a change of variable, (3.6) implies that

(3.8)
$$\int_{S^{n-1}} q(x)^{\nu} d\omega_{n-1} = \frac{b_{\nu}(2; \lambda)}{b_{\nu}(2; 1_{n})}, \quad \nu \in Z_{+},$$

where $\lambda = (\lambda_1, \dots, \lambda_n)$ denotes arbitrarily ordered eigenvalues of $q(x)^4$. From (3.8), we get the following equality of formal power series

⁴⁾ Note that $b_{\nu}(2; \lambda)$ is a symmetric function of λ_i 's.

(3.9)
$$\int_{S^{n-1}} \sum_{\nu=0}^{\infty} b_{\nu}(2; 1_{n}) q(x)^{\nu} t^{\nu} d\omega_{n-1} = \sum_{\nu=0}^{\infty} b_{\nu}(2; \lambda) t^{\nu}.$$

Replacing t by t/4 in (3.9) and using (1.5), (1.6), we obtain an interesting equality

(3.10)
$$\int_{S^{n-1}} (1 - tq(x))^{-n/2} d\omega_{n-1} = \prod_{i=1}^{n} (1 - \lambda_i t)^{-1/2}$$

which makes sense if |t| is sufficiently small.

§ 4. $f_{\nu}(\xi)$ and $\sigma_{\nu}(f)$

As in Introduction, let Ω be an open set of \mathbb{R}^n containing S^{n-1} and let $f: \Omega \to \mathbb{R}^m$ be a smooth map. With each $\nu \geq 0$, we associate a form $f_{\nu}(\xi)$ on \mathbb{R}^m of degree ν by

$$(4.1) f_{\nu}(\xi) = \int_{S^{n-1}} \langle \xi, f(x) \rangle^{\nu} d\omega_{n-1} .$$

We shall denote by $\sigma_{\nu}(f)$ the mean value of $f_{\nu}(\xi)$:

(4.2)
$$\sigma_{\nu}(f) = \int_{S^{m-1}} f_{\nu}(\xi) d\omega_{m-1}.$$

To study the numbers $\sigma_{\nu}(f)$ simultaneously for all $\nu \in \mathbb{Z}_+$, we introduce the generating function

(4.3)
$$\sigma(f;t) = \sum_{\nu=0}^{\infty} \sigma_{\nu}(f) \frac{t^{\nu}}{\nu!} .$$

As is easily seen, the series (4.3) converges for any $t \in C$ and we have

$$\sigma(f; t) = \sum_{\nu=0}^{\infty} \frac{t^{\nu}}{\nu!} \int_{S^{m-1}} d\omega_{m-1} \int_{S^{m-1}} \langle \xi, f(x) \rangle^{\nu} d\omega_{m-1}
= \sum_{\nu=0}^{\infty} \frac{t^{\nu}}{\nu!} \int_{S^{m-1}} d\omega_{n-1} \int_{S^{m-1}} \langle \xi, f(x) \rangle^{\nu} d\omega_{m-1}
= \int_{S^{m-1}} d\omega_{n-1} \int_{S^{m-1}} \sum_{\nu=0}^{\infty} \frac{1}{\nu!} (t \langle \xi, f(x) \rangle)^{\nu} d\omega_{m-1}
= \int_{S^{m-1}} d\omega_{n-1} \int_{S^{m-1}} \exp(t \langle \xi, f(x) \rangle) d\omega_{m-1}.$$

We are thus reduced to compute the integral

(4.5)
$$\int_{S^{m-1}} \varphi(\xi) d\omega_{m-1} \quad \text{with} \quad \varphi(\xi) = \exp\left(t\langle \xi, f(x) \rangle\right).$$

A simple computation shows that

Since $\varphi(0) = 1$ and $\lambda = t^2 |f(x)|^2$ is a constant for the variable ξ , we have, by (2.4),

$$\int_{S^{m-1}} \varphi(\xi) d\omega_{m-1} = \sum_{k=0}^{\infty} \frac{|f(x)|^{2k}}{(k!)^2 b_k(2; 1_m)} t^{2k}$$

and so

(4.7)
$$\sigma(f;t) = \sum_{k=0}^{\infty} \frac{N_k(f)}{(k!)^2 b_k(2;1_m)} t^{2k}$$

with

$$(4.8) N_k(f) = \int_{S^{n-1}} |f(x)|^{2k} d\omega_{n-1}.$$

Finally, by (4.3), (4.7), we have

$$\sigma_{2k}(f) = \frac{\binom{2k}{k}}{b_k(2;1_m)} N_k(f) , \qquad k \in \mathbf{Z}_+ .$$

We can also write (4.9) as

(4.10)
$$\sigma_{2k}(f) = \frac{b_k(2;1)}{b_k(2;1_m)} N_k(f) , \qquad k \in \mathbb{Z}_+ .$$

We shall call a map $f: \Omega \to \mathbb{R}^m$ spherical if $f(S^{n-1}) \subset S^{m-1}$. Since |f(x)| = 1, $x \in S^{n-1}$, for a spherical map f, we have $N_k(f) = 1$ and so

(4.11)
$$\sigma_{2k}(f) = \frac{b_k(2;1)}{b_k(2;1_m)}$$

when f is spherical.

§5. Examples

EXAMPLE 1. Let f_{ρ} , $\rho \in \mathbf{R}$, be the map $\mathbf{R}^2 \to \mathbf{R}^2$ defined by $f_{\rho}(x) = (x_1^2 - x_2^2, 2\rho x_1 x_2)$. When $\rho = \pm 1$, f_{ρ} sends S^1 onto S^1 ; when $\rho = 1$, f_{ρ} is the map $z \to z^2$ of $C = \mathbf{R}^2$ onto itself and when $\rho = -1$, f_{ρ} is the map $z \to \overline{z}^2$. Now,

$$(f_{\scriptscriptstyle
ho})_{\scriptscriptstyle
ho}(\xi) = \int_{\scriptscriptstyle
m CL} q(x)^{\scriptscriptstyle
ho} d\omega_{\scriptscriptstyle
m L}$$

with $q(x) = \langle \xi, f_{\rho}(x) \rangle = {}^{\iota}xA_{\rho}x$ where

$$A_{\scriptscriptstyle
ho} = egin{pmatrix} \xi_{\scriptscriptstyle 1} &
ho \xi_{\scriptscriptstyle 2} \
ho \xi_{\scriptscriptstyle 2} & - \xi_{\scriptscriptstyle 1} \end{pmatrix}$$

whose eigenvalues are $\lambda_1 = \sqrt{\xi_1^2 + \rho^2 \xi_2^2}$, $\lambda_2 = -\sqrt{\xi_1^2 + \rho^2 \xi_2^2}$. Therefore, by (1.7), (3.8), we have

$$\int_{S^1} q(x)^{\nu} d\omega_1 = \frac{b_{\nu}(2; \lambda)}{b_{\nu}(2; 1_{\nu})} = \frac{b_{\nu}(2; \lambda)}{4^{\nu}}.$$

Since we have

$$\begin{split} \prod_{i=1}^2 (1 - 4\lambda_i t)^{-1/2} &= (1 - 4\lambda_i t)^{-1/2} (1 + 4\lambda_i t)^{-1/2} = (1 - 16\lambda_1^2 t^2)^{-1/2} \\ &= \sum_{k=0}^{\infty} (\frac{1}{2}, k) \frac{16^k \lambda_1^{2k}}{k!} t^{2k} \;, \end{split}$$

we get, by (1.5),

$$b_{2k}(2;\lambda) = rac{(rac{1}{2},k)16^k}{k!} (\xi_1^2 +
ho^2 \xi_2^2)^k \ .$$

Or we have

$$(f_{
ho})_{2k}(\xi) = \int_{S^1} q(x)^{2k} d\omega_1 = rac{(rac{1}{2},k)}{k!} (\xi_1^2 +
ho^2 \xi_2^2)^k$$

and

$$\sigma_{2k}(f_{
ho}) = rac{(rac{1}{2},k)}{k!} \int_{S^1} (\xi_1^2 +
ho^2 \xi_2^2)^k d\omega_1 \ .$$

Since $\xi_1^2 + \xi_2^2 = 1$ on S^1 , we have

$$egin{aligned} \sigma_{2k}(f_
ho) &= rac{(rac{1}{2},k)}{k!} \int_{S^1} (1+(
ho^2-1) \xi_2^2)^k d\,\omega_1 \ &= rac{(rac{1}{2},k)}{k!} \sum_{m=0}^k \left(egin{array}{c} k \ m \end{array}
ight) (
ho^2-1)^m \int_{S^1} \xi_2^{2m} d\,\omega_1 \;. \end{aligned}$$

As for the last integral, since the eigenvalues of the quadratic form ξ_2^2 are $\lambda = (0, 1)$, we have, by (1.7), (3.8),

$$\int_{S^1} \xi_2^{2m} d\omega_1 = rac{b_m(2;(0,1))}{b_m(2;1_o)} = rac{b_m(2;(0,1))}{4^m}$$
 .

Since we have

$$\prod_{i=1}^{2} (1-4\lambda_{i}t)^{-1/2} = (1-4t)^{-1/2} = \sum_{m=0}^{\infty} (\frac{1}{2}, m) \frac{4^{m}t^{m}}{m!},$$

we get, by (1.5),

$$b_m(2;(0,1)) = (\frac{1}{2},m)\frac{4^m}{m!}$$

and

$$egin{aligned} \sigma_{2k}(f_
ho) &= rac{(rac{1}{2},k)}{k!} \sum\limits_{m=0}^k inom{k}{m} rac{(rac{1}{2},m)}{m!} (
ho^2 - 1)^m \ &= rac{(rac{1}{2},k)}{k!} F(-k,rac{1}{2};1;1-
ho^2) \end{aligned}$$

because

$$\binom{k}{m} = (-1)^m \frac{(-k, m)}{m!}.$$

Finally, from (4.10) we get

(5.1)
$$N_k(f_{\varrho}) = F(-k, \frac{1}{2}; 1; 1 - \rho^2).$$

Example 2. Let f_{ρ} , $\rho \in R$, be the map $R^4 \to R^3$ defined by $f_{\rho}(x) = (x_1^2 + x_2^2 - x_3^2 - x_4^2, 2\rho(x_2x_3 - x_1x_4), 2\rho(x_1x_3 + x_2x_4))$. When $\rho = \pm 1$, f is the classical Hopf map sending S^3 onto S^2 (see Hopf [5]). Now,

$$(f_{\scriptscriptstyle
ho})_{\scriptscriptstyle
ho}(\xi) = \int_{S^3} q(x)^{\scriptscriptstyle
ho} d\,\omega_{\scriptscriptstyle 3}$$

with $q(x) = \langle \xi, f_{\rho}(x) \rangle = {}^{t}xA_{\rho}x$ where

$$A_{
ho} = egin{pmatrix} \xi_1 & 0 &
ho \xi_3 & -
ho \xi_2 \ 0 & \xi_1 &
ho \xi_2 &
ho \xi_3 \
ho \xi_3 &
ho \xi_2 & - \xi_1 & 0 \ -
ho \xi_2 &
ho \xi_3 & 0 & - \xi_1 \end{pmatrix}$$

whose eigenvalues are

$$\lambda_1 = \lambda_2 = \sqrt{\xi_1^2 +
ho^2(\xi_2^2 + \xi_3^2)} \;, \qquad \lambda_3 = \lambda_4 = -\sqrt{\xi_1^2 +
ho^2(\xi_2^2 + \xi_3^2)} \;.$$

Therefore, by (1.7), (3.8), we have

$$\int_{S^3} q(x)^{\nu} d\omega_3 = \frac{b_{\nu}(2;\lambda)}{b_{\nu}(2;1_4)} = \frac{b_{\nu}(2;\lambda)}{4^{\nu}(\nu+1)}.$$

Since we have

$$\prod_{i=1}^{4} (1 - 4\lambda_i t)^{-1/2} = (1 - 4\lambda_i t)^{-1} (1 + 4\lambda_i t)^{-1} = (1 - 16\lambda_1^2 t^2)^{-1}$$

$$= 1 + 16\lambda_1^2 t^2 + 16^2\lambda_1^4 t^4 + \cdots,$$

we get, by (1.5),

$$b_{2k}(2;\lambda) = 16^k(\xi_1^2 + \rho^2(\xi_2^2 + \xi_3^2))^k$$
.

Or we have

$$(f_{
ho})_{2k}(\xi) = \int_{S^3} q(x)^{2k} d\omega_3 = rac{1}{2k+1} (\xi_1^2 +
ho^2 (\xi_2^2 + \xi_3^2))^k$$

and

$$\sigma_{2k}(f_
ho) = rac{1}{2k+1} \int_{S^2} (\xi_1^2 +
ho^2(\xi_2^2 + \xi_3^2))^k d\omega_2 \, .$$

Since $\xi_1^2 + \xi_2^2 + \xi_3^2 = 1$ on S^2 , we have

$$egin{aligned} \sigma_{2k}(f_
ho) &= rac{1}{2k+1} \int_{S^2} (1+(
ho^2-1)(\xi_2^2+\,\xi_3^2))^k d\omega_2 \ &= rac{1}{2k+1} \sum_{m=0}^k igg(rac{k}{m}igg) (
ho^2-1)^m \int_{S^2} (\xi_2^2+\,\xi_3^2)^m d\omega_2 \ . \end{aligned}$$

As for the last integral, since the eigenvalues of the quadratic form $\xi_2^2 + \xi_3^2$ are $\lambda = (0, 1, 1)$, we have, by (1.7), (3.8)

$$\int_{S^2} (\xi_2^2 + \xi_3^2)^m d\omega_2 = rac{b_{\,m}(2;\,(0,\,1,\,1))}{b_{\,m}(2;\,1_3)} = rac{m\,!}{4^m(rac{3}{2},\,m)} b_{\,m}(2;\,(0,\,1,\,1)) \;.$$

Since we have

$$\prod_{i=1}^{3} (1-4\lambda_i t)^{-1/2} = (1-4t)^{-1} = 1+4t+4^2t^2+\cdots.$$

we get, by (1.5), $b_m(2; (0, 1, 1)) = 4^m$ and

$$egin{align} \sigma_{2k}(f_{
ho}) &= rac{1}{2k+1} \sum\limits_{m=0}^k \left(rac{k}{m}
ight) rac{m!}{(rac{3}{2},m)} \ &= rac{1}{2k+1} F(-k,1;rac{3}{2};(1-
ho^2)) \;. \end{split}$$

Finally, from (4.10) we get

(5.2)
$$N_k(f_{\rho}) = F(-k, 1; \frac{3}{2}; (1-\rho^2)).$$

Appendix. On Legendre polynomials

Let $f: \mathbb{R}^n \to \mathbb{R}^m$ be an affine map given by f(x) = Ax + b where A is an $(m \times n)$ -matrix with real coefficients a_{ij} , $1 \le i \le m$, $1 \le j \le n$, and b is a real m-vector written vertically. Put $a_i = (a_{i1}, \dots, a_{in})$, the i-th row of $A = (a_{ij})$, $M = (\langle a_i, a_j \rangle)$, an $(m \times m)$ -symmetric real matrix, and $Q(\xi) = {}^t \xi M \xi$, the corresponding quadratic form. We can verify easily that $\lambda = t^2 Q(\xi)$ satisfies $\Delta \varphi = \lambda \varphi$ for $\varphi = \exp(t \langle \xi, f(x) \rangle)$. Hence, by (2.4), we have

(a.1)
$$\begin{split} \sum_{\nu=0}^{\infty} f_{\nu}(\xi) \frac{t^{\nu}}{\nu!} &= \int_{S^{n-1}} \varphi(x) d\omega_{n-1} \\ &= \exp\left(t\langle \xi, b \rangle\right) \sum_{k=0}^{\infty} \frac{Q(\xi)^{k} t^{2k}}{4^{k} k! \left(n/2, k\right)} \end{split}$$

and so

(a.2)
$$f_{\nu}(\xi) = \sum_{k=0}^{\lfloor \nu/2 \rfloor} \frac{\nu!}{4^k k! (n/2, k) (\nu - 2k)!} Q(\xi)^k \langle \xi, b \rangle^{\nu - 2k} .$$

Denote by H the algebraic set in \mathbb{R}^m defined by

$$H = \{ \xi \in \mathbf{R}^m; \ Q(\xi) = \langle \xi, b \rangle^2 - 1 \}$$

and put $z = \langle \xi, b \rangle$. Then, for $\xi \in H$, we have

$$f_{\nu}(\xi) = \sum_{k=0}^{\lfloor \nu/2 \rfloor} \frac{\nu!}{4^{k}k! (n/2, k)(\nu - 2k)!} (z^{2} - 1)^{k} z^{\nu - 2k}$$

$$= z^{\nu} \sum_{k=0}^{\lfloor \nu/2 \rfloor} \frac{(-\nu/2, k)(-\nu/2 + \frac{1}{2}, k)}{(n/2, k)k!} \left(\frac{z^{2} - 1}{z^{2}}\right)^{k}$$

$$= z^{\nu} F\left(-\frac{\nu}{2}, -\frac{\nu}{2} + \frac{1}{2}; \frac{n}{2}; \frac{z^{2} - 1}{z^{2}}\right)$$

$$= F\left(-\nu, \nu + n - 1; \frac{n}{2}; \frac{1 - z}{2}\right)$$

$$= P_{\nu, n+1}(z),$$

where the equality between two hypergeometric series follows from a formula of quadratic transformations⁵⁾ and the last equality is a well-known relation of the Legendre polynomial for R^{n-1} of order ν and the hypergeometric series⁶⁾. On equating the first and the last terms of (a.3), we get

⁵⁾ See Magnus-Oberhettinger-Soni [6], p. 50, line 3 from the bottom.

⁶⁾ See Hochstadt [4], p. 183, line 8.

$$(a.4) P_{\nu,n+1}(\langle \xi,b\rangle) = \int_{S^{n-1}} \langle \xi,Ax+b\rangle^{\nu} d\omega_{n-1} , \xi \in H ,$$

which is substantially the Laplace integral for the Legendre polynomials. If, in particular, m = n + 1,

and $\xi_3 = \cdots = \xi_{n+1} = 0$, then $\xi_1^2 - \xi_2^2 = 1$ for $\xi \in H$ and we get

$$P_{
u,n+1}(\xi_1) = \int_{S^{n-1}} (\xi_1 + \sqrt{\xi_1^2 - 1} \, x_1)^{
u} d\omega_{n-1} \; ,$$

the Laplace integral in its original form⁷.

REFERENCES

- [1] P. Appell and J. Kampé de Fériet, Fonctions hypergéometriques et hypersphériques, polynomes d'Hermite, Gauthier-Villars, Paris 1926.
- [2] B. C. Carlson, Special Functions of Applied Mathematics, Academic Press, New York, 1977.
- [3] R. Courant and D. Hilbert, Methoden der Mathematischen Physik II, 2nd Ed., Springer-Verlag, Berlin Heidelberg New York, 1968.
- [4] H. Hochstadt, The Functions of Mathematical Physics, Wiley-Interscience, New York, 1971.
- [5] H. Hopf, Über die Abbildungen der dreidimensionalen Sphäre auf die Kugelfläche, Math. Ann., 104 (1931), 637-665.
- [6] W. Magnus, F. Oberhettinger and R. P. Soni, Formulas and Theorems for the Special Functions of Mathematical Physics, 3rd Ed., Springer-Verlag, New York, 1966.
- [7] T. Ono, On Certain Numerical Invariants of Mappings over Finite Fields. I., Proc. Japan Acad. Ser. A, 56 (1980), 342-347.

Department of Mathematics The Johns Hopkins University, Baltimore, Maryland, U.S.A.

⁷⁾ See [4], p. 182, Theorem.