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# FLAT AFFINE AND SYMPLECTIC GEOMETRIES ON LIE GROUPS

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**Abstract.** In this paper we exhibit a family of flat left invariant affine structures on the double Lie group of the oscillator Lie group of dimension 4, associated to each solution of classical Yang-Baxter equation given by Boucetta and Medina. On the other hand, using Koszul's method, we prove the existence of an immersion of Lie groups between the group of affine transformations of a flat affine and simply connected manifold and the classical group of affine transformations of  $\mathbb{R}^n$ . In the last section, for each flat left invariant affine symplectic connection on the group of affine transformations of the real line, describe by Medina-Saldarriaga-Giraldo, we determine the affine symplectomorphisms. Finally we exhibit the Hess connection, associated to a Lagrangian bi-foliation, which is flat left invariant affine.

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## 1. Introduction

Let M be a smooth (paracompact) connected real smooth manifold of dimension n. A *linear frame* u at point  $x \in M$  is an ordered basis  $(X_1, \ldots, X_n)$  of the tangent space  $T_x(M)$  it will usually be identified with a linear isomorphism from  $\mathbb{R}^n$  onto  $T_x(M)$ . The set L(M) of all linear frames at all points of M is a smooth principal fiber bundle over M with group  $GL(n, \mathbb{R})$ , the right action of  $GL(n, \mathbb{R})$  on L(M) is given by

$$u \cdot a = u \circ a$$
 with  $u \in L(M)$  and  $a \in GL(n, \mathbb{R})$ .

The diffeomorphism of L(M) associated to a will be denoted  $R_a(u) = u \cdot a$ . The *canonical form*  $\theta$  of L(M) is the  $\mathbb{R}^n$ -valued one-form defined by

$$\theta(X_u) := u^{-1}(\pi_{*,u}(X_u)) \quad \text{for} \quad X_u \in T_u(L(M))$$

where  $\pi: L(M) \to M$  in the natural projection.

In what follows we will privilege the notion of connection due to C. Ehresmann [8]. When necessary we will also use the notions of E. Cartan [2] or J.-L. Koszul [14].

A *Linear connection*  $\Gamma$  over M is a smooth distribution H in L(M) transverse and supplementary to the fiber and invariant under the right action of  $GL(n, \mathbb{R})$ , defined above. The subspace  $H_u$ , where u is a linear frame at  $x \in M$ , is called the horizontal space in u.

It is clear that every vector  $X_u \in T_u(L(M))$  can be uniquely written as  $X_u = X_{u,H} + X_{u,V}$ , with  $X_{u,H} \in H_u$  and  $X_{u,V}$  tangent to the fiber over u.

Each  $\vec{x} \in \mathbb{R}^n$  determines a *standard horizontal vector field*  $B(\vec{x})_u$  defined as the unique element of  $H_u$  such that  $\pi_{*,u}(B(\vec{x})_u) = u(\vec{x})$ .

Moreover,  $\Gamma$  defines a unique one-form  $\omega$  on L(M) with values in  $\mathfrak{g} = \mathfrak{gl}(n, \mathbb{R})$ , called the *connection form*. More precisely  $\omega_u(X_u)$  is the unique  $A \in \mathfrak{g}$  such that  $A_u^* = X_{u,V}$  where

$$A_u^* = \left. \frac{\mathrm{d}}{\mathrm{d}t} \right|_{t=0} u \cdot \exp tA$$

is the fundamental vector field corresponding to A. The connection form  $\omega$  satisfies the following conditions

$$\omega(A^*) = A, \quad A \in \mathfrak{g}, \qquad (R_{\sigma})^* \omega = \mathsf{ad}(\sigma^{-1})\omega, \quad \sigma \in \mathrm{GL}(n, \mathbb{R})$$
(1)

where ad denotes the adjoint representation of  $GL(n, \mathbb{R})$ . Having a connection  $\Gamma$  is equivalent to have a one-form  $\omega$  on L(M) with values in  $\mathfrak{g}$  verifying (1). A smooth vector field  $X \in \mathfrak{X}(L(M))$  such that  $\omega(X) = 0$ , is called *horizontal vector* relative to  $\Gamma$ .

The *curvature form*  $\Omega$  and the torsion form  $\Theta$  of  $(L(M), \Gamma)$  are defined respectively by

$$\Omega_u(X_u, Y_u) = (\mathrm{d}\omega)_u(X_{u,H}, Y_{u,H}), \qquad \Theta_u(X_u, Y_u) = (\mathrm{d}\theta)_u(X_{u,H}, Y_{u,H})$$

with  $X_u, Y_u \in T_u(L(M))$  and  $u \in L(M)$ .

Note that  $\Omega$  takes values in  $\mathfrak{gl}(n,\mathbb{R})$  and  $\Theta$  in  $\mathbb{R}^n$ . The following equations are called *structure equations* (also known as Cartan equations) (see [12])

$$(d\theta)_u(X_u, Y_u) = -\frac{1}{2}(\omega_u(X_u)\theta_u(Y_u) - \omega_u(Y_u)\theta_u(X_u)) + \Theta_u(X_u, Y_u)$$
  
(d\omega)\_u(X\_u, Y\_u) = -\frac{1}{2}[\omega\_u(X\_u), \omega(Y\_u)] + \Omega\_u(X\_u, Y\_u). (2)

Given a global (or local) diffeomorphism f of M we will denote  $\tilde{f}$  to the global (or local) diffeomorphism of L(M) defined by  $\tilde{f}(u) = (f_{*,x}(X_1), \dots, f_{*,x}(X_n))$  with  $u = (X_1, \dots, X_n)$  a linear frame at  $x \in M$ . The diffeomorphism  $\tilde{f}$  is called *the natural lift* of f to L(M).

When f satisfies the condition  $\tilde{f}^*\omega = \omega$ , it is called a global (or local) *affine transformation* of  $(M, \Gamma)$ .

If M is simply connected endowed with a linear connection  $\Gamma$ , it is well known that if  $\Gamma$  is flat, that is, without curvature, the connection over M is isomorphic to the connection  $\Gamma'$  given by the distribution H in  $(M \times \operatorname{GL}(n, \mathbb{R}), \psi)$  where  $H_u$  is the tangent space to  $M \times \{a\}$  in  $u = (x, a) \in M \times \operatorname{GL}(n, \mathbb{R})$ . The connection  $\Gamma'$ is called the *canonical flat (linear) connection* over M (see [12]).

Let  $\pi_2 : M \times GL(n, \mathbb{R}) \to GL(n, \mathbb{R})$  natural projection. Notice that the connection form of the canonical flat linear connection is given by

$$\omega' = \pi_2^* \theta^+$$

where  $\theta^+$  is the left invariant one-form over  $\mathrm{GL}(n,\mathbb{R})$ , with values in  $\mathfrak{gl}(n,\mathbb{R})$  defined by

$$\theta^+(A^+) := A^+$$

and  $A^+ \in \mathfrak{g}$  is the left invariant vector field corresponding to  $A \in T_{\varepsilon} \mathrm{GL}(\mathbf{n}, \mathbb{R})$ . The one-form  $\theta^+$  is called the canonical one-form over  $\mathrm{GL}(\mathbf{n}, \mathbb{R})$ .

If M = G is a Lie group and the left multiplications  $L_{\sigma}$  (respectively right multiplications  $R_{\sigma}$ ) are affine transformations it is said that  $\Gamma$  is *left invariant* (*respectively right invariant*). This means that the lift  $\tilde{L}_{\sigma}$  of  $L_{\sigma}$  (respectively  $\tilde{R}_{\sigma}$ ) preserves the distribution H given by  $\Gamma$ . In this case we denote the connection by  $\Gamma^+$  and we will say that  $(G, \Gamma^+)$  is an affine Lie group. It is known that having a linear connection  $\Gamma$  over M is equivalent to having a covariant derivative  $\nabla$  (see [12]). Let us consider the product

$$XY := \nabla_X Y, \qquad X, Y \in \mathfrak{X}(M). \tag{3}$$

If the curvature tensor relative to  $\nabla$  is identically zero, this product verifies the condition

$$[X,Y]Z = X(YZ) - Y(XZ), \qquad X,Y,Z \in \mathfrak{X}(M).$$

If both torsion and curvature vanish identically we have

$$(XY)Z - X(YZ) = (YX)Z - Y(XZ), \qquad X, Y, Z \in \mathfrak{X}(M).$$
(4)

A vector space endowed with a bilinear product satisfying Equation (4) is called a *left symmetric algebra*. Left symmetric algebras are Lie algebras with the bracket given by  $[X, Y]_1 := XY - YX$ . If  $\nabla$  is torsion free this Lie bracket agrees with the usual Lie bracket of  $\mathfrak{X}(M)$ , that is

$$[X,Y] = XY - YX. (5)$$

Recall that product (3) satisfies (fX)Y = f(XY) and X(gY) = X(g)Y + g(XY), for all  $f, g \in C^{\infty}(M, \mathbb{R})$ , i.e., the product defined above is  $\mathbb{R}$ -bilinear and  $C^{\infty}(M, \mathbb{R})$ -linear in the first component.

If M = G is a Lie group and  $\nabla = \nabla^+$  is left invariant, it is well known that  $(G, \nabla^+)$  is flat affine if and only if Product (3) turn g into a left symmetric algebra (see [16]).

If this is the case, we denote by  $\omega^+$  the corresponding connection form associated to the connection  $\nabla^+$ .

The notion of Poisson-Lie group, as well as the term quantum group, were introduced by Drinfeld [6]. They derived their origins from Faddeev's work and his school on quantum inverse scattering, which inspired those of Sklyanin on the Yang-Baxter equation. The Poisson groups appear in the study of completely integrable Hamiltonian system, r-matrices, the Yang-Baxter equation, quantum inverse scattering, symplectic groupoids, non commutative geometry, etc. Knowing the Lie-Poisson structures on G (see for instance [21]), would allow to study the quantum deformations of G. It is known that every connected Poisson-Lie group arises from a Lie bialgebra and important class of Lie bialgebra is obtained by solving the generalized Yang-Baxter equation.

Diatta and Medina proved in [5] that each solution r of the Classical Yang-Baxter Equation on G, corresponds to a connected Lie subgroup H of  $T^*G$  such that its Lie algebra  $\mathfrak{h}$  is a Lagrangian graph in  $\mathfrak{g} \oplus \mathfrak{g}*$  and H carries a left invariant affine structure. They also prove that, if r is invertible, the Poisson Lie tensor  $\pi$  given by r on G is polynomial of degree at most two and every double Lie group of  $(G, \pi)$  is endowed with an affine and a complex structures  $\nabla^+$  and J, both left invariant and given by r, such that  $\nabla^+ J = 0$ .

## 2. Flat Affine Geometry on Four-Dimensional Oscillator Lie Group

The three-dimensional Heisenberg Lie group is the real manifold  $\mathsf{H}_3 = \mathbb{R} \times \mathbb{C}$  with product

$$(s,w)(s',w') = \left(s+s'+\frac{1}{2}\operatorname{Im}(\overline{w}w'),w+w'\right)$$

**Definition 1.** Let  $\lambda$  be a positive real number and  $\rho : \mathbb{R} \to Aut(H_3)$  the Lie group homomorphism given by

$$\rho(t)(s,z) = (s, \mathrm{e}^{\mathrm{i}t\lambda}z).$$

The real  $\lambda$ -oscillator group of dimension 4, denoted by  $O_{\lambda}$ , is the semidirect product of Heisenberg Lie group  $H_3$  for the group  $(\mathbb{R}, +)$  by the representation  $\rho$ .

The following result is known (see [17]).

**Lemma 2.** The Lie group  $O_{\lambda}$  is independent of  $\lambda$ , for  $\lambda > 0$ .

**Proof:** The statement follows from the fact that the map  $\phi : H_3 \to H_3$  defined by  $\phi(s, z) = \left(s, z^{\frac{\lambda'}{\lambda}}\right)$ , is an automorphism of Lie groups such that for all  $t \in \mathbb{R}$  the following diagram commutes

$$\begin{array}{ccc} \mathsf{H}_3 & \stackrel{\rho(t)}{\longrightarrow} & \mathsf{H}_3 \\ \downarrow \phi & & \downarrow \phi \\ \mathsf{H}_3 & \stackrel{\rho'(t)}{\longrightarrow} & \mathsf{H}_3 \end{array}$$

where  $\rho$  and  $\rho'$  are the representations

$$\rho: \mathbb{R} \to \operatorname{Aut}(\mathsf{H}_3) \qquad and \qquad \begin{array}{c} \rho': \mathbb{R} \to \operatorname{Aut}(\mathsf{H}_3) \\ t \mapsto (s, \mathrm{e}^{\lambda t\mathrm{i}}) \end{array} \qquad and \qquad \begin{array}{c} t \mapsto (s, \mathrm{e}^{\lambda' t\mathrm{i}}). \end{array}$$

In the follows we denote by O the oscillator Lie group of dimension four with  $\lambda = 1$ . Let  $(e_0, e_1, e_2, d)$  be the natural basis for  $\mathbb{R}^4$ . The corresponding left

invariant vector fields are as follows:

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$$e_{0}^{+} = \frac{\partial}{\partial x_{1}}$$

$$e_{1}^{+} = \left(\frac{1}{2}(x_{2}\sin(x_{4}) - x_{3}\cos(x_{4}))\right)\frac{\partial}{\partial x_{1}} + \cos(x_{4})\frac{\partial}{\partial x_{2}} + \sin(x_{4})\frac{\partial}{\partial x_{3}}$$

$$e_{2}^{+} = \left(\frac{1}{2}(x_{2}\cos(x_{4}) + x_{3}\sin(x_{4}))\right)\frac{\partial}{\partial x_{1}} - \sin(x_{4})\frac{\partial}{\partial x_{2}} + \cos(x_{4})\frac{\partial}{\partial x_{3}}$$

$$d^{+} = \frac{\partial}{\partial x_{4}}$$
(6)

It is directly verified that the non zero Lie brackets, except antisymmetry, are given by

$$[e_1^+, e_2^+] = e_0^+, \qquad [d^+, e_1^+] = e_2^+, \qquad [d^+, e_2^+] = -e_1^+.$$
(7)

Consequently the Lie algebra of the oscillator group O is the real vector space with basis  $\{e_0^+, e_1^+, e_2^+, d^+\}$  and bracket defined in (7). This Lie algebra will be denoted by  $\mathfrak{o}$ .

## 2.1. Flat Affine Structures on O

To construct these kind of structures, we start by finding flat left invariant affine structures on the group of isometries of the Euclidean plane E(2), called the Euclidean group. Then we use the cohomology of left symmetric algebras developed in [19] to extend these structures to the oscillator group of dimension 4.

Notice that the oscillator group can be viewed as an extension of the Euclidean group E(2). This extension is given by the exact split sequence of Lie groups

$$\{0\} \longrightarrow (\mathbb{R}, +) \xrightarrow{i} O \xrightarrow{\pi} E(2) \longrightarrow \{I\}$$

where i(s) = (s, 0, 0), for all  $s \in \mathbb{R}$  and

$$\pi(s, x + iy, t) = \begin{pmatrix} \cos(t) & -\sin(t) & x \\ \sin(t) & \cos(t) & y \\ 0 & 0 & 1 \end{pmatrix}, \qquad (s, x + iy, t) \in O$$

and the section is given by the Lie group homomorphism  $\mathbf{s} : E(2) \to O$  defined by  $\mathbf{s}(A, (x, y)) = (0, x + iy, t)$  where  $A = \begin{pmatrix} \cos(t) & -\sin(t) \\ \sin(t) & \cos(t) \end{pmatrix}$ . The Lie algebra of the Euclidean Lie group,  $\mathbf{c}(2) = \operatorname{Vect}_{\mathbb{R}}\{e_1, e_2, d\}$ , has Lie bracket given by  $[d, e_1] = e_2$ ,  $[d, e_2] = -e_1$ , and  $[e_1, e_2] = 0$ . **Proposition 3.** The following left symmetric products determine flat left invariant non-isomorphic affine structures on the Euclidean group E(2).

$$\mathcal{F}_{2}(\alpha) \qquad \frac{\begin{vmatrix} e_{1} & e_{2} & d \\ e_{1} & 0 & 0 & 0 \\ e_{2} & 0 & 0 & 0 \\ d & e_{2} & -e_{1} & \alpha d \end{vmatrix}$$
(9)

$$\mathcal{F}_{3} \qquad \frac{\begin{vmatrix} e_{1} & e_{2} & d \\ e_{1} & d & 0 & 0 \\ e_{2} & 0 & d & 0 \\ d & e_{2} & -e_{1} & 0 \end{vmatrix}$$
(10)

$$F_{4} = \begin{array}{c|c} e_{1} & e_{2} & d \\ \hline e_{1} & e_{1} + 2e_{2} - d & 0 & e_{1} + 2e_{2} - d \\ e_{2} & 0 & e_{1} + 2e_{2} - d & 2e_{1} + 4e_{2} - 2d \\ d & e_{1} + 3e_{2} - d & e_{1} + 4e_{2} - 2d & 3e_{1} + 11e_{2} - 5d \end{array}$$
(11)

**Proof:** Recall that finding flat left invariant affine structures on E(2) is equivalent to finding bilinear products on e(2) satisfying (4) and (5).

Let us denote by  $L_a$  the endomorphism of the vector space  $\mathfrak{e}(2)$  defined by  $L_a(b) = ab$ .

If a bilinear product is compatible with the Lie bracket of  $\mathfrak{e}(2)$ , the endomorphisms  $L_{e_1}$ ,  $L_{e_2}$  and  $L_d$  can be written, in the base  $(e_1, e_2, d)$ , as

$$L_{e_1} = (a_{ij}), \ L_{e_2} = \begin{pmatrix} a_{12} & b_{12} & b_{13} \\ a_{22} & b_{22} & b_{23} \\ a_{32} & b_{32} & b_{33} \end{pmatrix}, \ L_d = \begin{pmatrix} a_{13} & b_{13} - 1 & c_{13} \\ 1 + a_{23} & b_{23} & c_{23} \\ a_{33} & b_{33} & c_{33} \end{pmatrix}$$
(12)

where  $a_{ij}$ ,  $b_{ij}$  and  $c_{ij}$  are real numbers.

Moreover, if the product on  $\mathfrak{e}(2)$  satisfies the identity (4), the endomorphisms (12) satisfy the system of equations

$$[L_{e_1}, L_{e_2}] = 0,$$
  $[L_d, L_{e_1}] = L_{e_2},$   $[L_d, L_{e_2}] = -L_{e_1}.$ 

Solving the system and using the above observations we obtain the products described in the tables.

## 2.2. Left Symmetric Algebra Cohomology

Let A be a left symmetric algebra and M a A-bimodule. We denote by  $H^k_{SG}(A, M)$  the cohomology group of order k of A with coefficients in M. Let us consider  $\mathbb{R}$  as a trivial A-bimodule. A scalar two-cocycle of A is a bilinear map  $f : A \times A \to \mathbb{R}$  verifying

$$f(ab - ba, c) = f(a, bc) - f(b, ac).$$

A scalar two-coboundary of A is a two-cocycle f such that  $f(a, b) = \varphi(ab)$ , for some  $\varphi \in A^*$ . The cohomology group  $H^2_{SG}(A, \mathbb{R}) = Z^2/B^2$ , where  $Z^2$  (respectively  $B^2$ ) is the set of two-cocycles (respectively two-coboundaries), classifies the left symmetric algebras extensions of A for which  $\mathbb{R}$  is an ideal of null square (see [19]). In particular

$$0 \to \mathbb{R} \to \mathbb{R} \times A \to A \to 0$$

is an exact sequence of left symmetric algebras.

**Proposition 4.** The second order cohomology groups of left symmetric algebras of  $\mathfrak{e}(2)$  corresponding to the products  $\mathcal{F}_1(0)$ ,  $\mathcal{F}_1(\alpha)$ ,  $\mathcal{F}_2(\alpha)$ ,  $\mathcal{F}_3$  and  $\mathcal{F}_4$  are respectively given by

$$\begin{split} H^2_{SG,1,0}(\mathfrak{e}(2),\mathbb{R}e_0) &\cong \mathsf{Vect}_{\mathbb{R}}\{E_{12} - E_{21}, E_{22} + E_{11}, E_{33}\}\\ H^2_{SG,1,\alpha}(\mathfrak{e}(2),\mathbb{R}e_0) &\cong \{0\}, \ if \ \alpha \neq 0\\ H^2_{SG,2,\alpha}(\mathfrak{e}(2),\mathbb{R}e_0) &\cong \mathsf{Vect}_{\mathbb{R}}\{E_{11} + E_{22}, E_{12} - E_{21}\}\\ H^2_{SG,3}(\mathfrak{e}(2),\mathbb{R}e_0) &\cong \mathsf{Vect}_{\mathbb{R}}\{E_{12} - E_{21}\}\\ H^2_{SG,4}(\mathfrak{e}(2),\mathbb{R}e_0) &\cong \mathsf{Vect}_{\mathbb{R}}\{E_{12} - E_{21} + 2E_{13} - E_{23}\} \end{split}$$

where  $\{E_{ij}\}$  denote the elementary matrices on  $M_{3\times 3}(\mathbb{R})$ .

**Proof:** Let us consider the A-bimodule  $M = \mathbb{R}e_0$  as a trivial  $\mathfrak{e}(2)$ -bimodule, with  $\mathfrak{e}(2) = \operatorname{Vect}_{\mathbb{R}}(e_1, e_2, d)$ . Each of the following statements can be verified by direct calculations.

In the algebra with product given by  $\mathcal{F}_1(0)$ , the scalar two-cocycles of  $\mathfrak{e}(2)$  with coefficients in M, are of the form

$$g = p(E_{11} + E_{22}) + q(E_{12} - E_{21}) + rE_{31} + sE_{32} + tE_{33}$$

and the two-coboundaries as

$$f = rE_{31} + sE_{32}.$$

In the algebras corresponding to the products  $\mathcal{F}_1(\alpha)$ , with  $\alpha \neq 0$ , the two-cocycles (respectively the two-coboundaries), are given by

$$g = p\left(E_{13} + E_{31} - \frac{1}{\alpha}E_{32}\right) + q\left(E_{23} + \frac{1}{\alpha}E_{31} + E_{32}\right) + rE_{33}$$

and

$$f = p(\alpha E_{13} + \alpha E_{31} - E_{32}) + q(\alpha E_{23} + E_{31} + \alpha E_{32}) + rE_{33}.$$

In the algebras with product  $\mathcal{F}_2(\alpha)$  the two-cocycles (respectively the two-coboundaries) are written

$$g = t(E_{11} + E_{22}) + s(E_{12} - E_{21}) + rE_{31} + pE_{32} + qE_{33}$$

and

$$f = rE_{31} + pE_{32} + qE_{33}.$$

The two-cocycles (respectively the two-coboundaries) associated to the algebra with product  $\mathcal{F}_3$ , are given by

$$g = t(E_{11} + E_{22}) + s(E_{12} - E_{21}) + rE_{31} + pE_{32}$$

and

$$f = t(E_{11} + E_{22}) + rE_{31} + pE_{32}$$

Finally, the scalar two-cocycles (respectively two-coboundaries) of the left symmetric algebra defined by  $\mathcal{F}_4$ , are of the form

$$g = p(E_{11} + E_{22} + E_{13} + 2E_{23}) + q(E_{12} - E_{21} + 2E_{13} - E_{23}) + r(E_{31} + E_{33}) + s(E_{32} + 2E_{33})$$

and

$$f = p(E_{11} + E_{22} + E_{13} + 2E_{23}) + r(E_{31} + E_{33}) + s(E_{32} + 2E_{33})$$

where p, q, r, s, t are real numbers.

The previous statements imply the results enunciated in the proposition.

Each class [g] of the left symmetric two-cohomology of  $\mathfrak{e}(2)$  determines a left symmetric product in  $\mathfrak{o} \cong \mathbb{R}e_0 \rtimes_g \mathfrak{e}(2)$  by the formula  $(e_0, a) \cdot (e_0, b) = (g(a, b), a \cdot b)$ , where g is a class representative [g]. This method and other observations lead to the following result.

**Theorem 5.** The left symmetric products on the Lie algebra  $\mathfrak{o}$  given by the following tables determine non isomorphic flat left invariant affine connections on the four-dimensional oscillator group

$\mathcal{F}_{1}(0,t,s,q) \qquad \begin{array}{ccccccccccccccccccccccccccccccccccc$	2)
	$\mathbf{a}$
$e_0 \mid 0 - ae_0 te_0 = 0$	3)
$c_2 \mid 0  qc_0  cc_0  0$	
$d \mid 0  e_2  -e_1  se_0$	
$e_0 \mid 0  0  0  0$	(14)
$\mathcal{F}_2(\alpha, t, q) \qquad e_1 \mid 0  te_0  qe_0  0 \qquad (1)$	
$e_2 \mid 0  -qe_0  te_0  0$	
$d \mid 0  e_2  -e_1 \; lpha d$	
$e_0 e_1 e_2 d$	
$e_0 \mid 0  0  0  0$	(15)
$\mathcal{F}_3(q) \qquad \qquad e_1 \mid 0  d  qe_0  0 \qquad \qquad (1$	
$e_2 \mid 0  -qe_0  d  0$	
$d \mid 0  e_2  -e_1 \mid 0$	

where  $t, s \in [0, +\infty), \alpha \in \mathbb{R} \setminus \{0\}$  and  $q = \frac{1}{2}$ .

**Corollary 6.** The flat left invariant affine connections on the four-dimensional oscillator group given by the products  $\mathcal{F}_1(0, t, s, \frac{1}{2})$ ,  $\mathcal{F}_3(\frac{1}{2})$  and  $\mathcal{F}_4(\frac{1}{2})$  are geodesically complete.

**Proof:** This results follows from the fact that each left symmetric algebra has right multiplications  $R_a$  of trace zero (see [9] and [16]).

**Remark 7.** It is important to notice that left symmetric products  $\mathcal{F}_1(0,t,s,q), \mathcal{F}_2(\alpha,t,q), \mathcal{F}_3(q)$  and  $\mathcal{F}_4(q)$ , with  $q \neq 0$  are isomorphic respectively to  $\mathcal{F}_1(0,s,t,1/2), \mathcal{F}_2(\alpha,t,1/2), \mathcal{F}_3(1/2)$  and  $\mathcal{F}_4(1/2)$ .

# 2.3. Double Lie Groups of the Oscillator Group *O* Relative to a Solution of Classical Yang-Baxter Equation

The aim of this section is to determine the dual Lie groups of the oscillator group *O* relative to solutions of the classical Yang-Baxter equation. We describe the corresponding double Lie groups and we exhibit flat left invariant affine structures on them, by extending the flat affine structures on the oscillator group given in the previous section.

Let G be a real Lie group with Lie algebra  $\mathfrak{g}$ . A solution of classical Yang-Baxter equation (CYBE) on G, is a bivector  $r \in \wedge^2 \mathfrak{g}$  satisfying

$$[r,r] = 0 \tag{17}$$

where  $[r,r] \in \wedge^3 \mathfrak{g}$  is the Schouten bracket (see [7]). The equation (17) is equivalent to

$$\alpha([r_{\sharp}(\beta), r_{\sharp}(\gamma)]) + \beta([r_{\sharp}(\gamma), r_{\sharp}(\alpha)]) + \gamma([r_{\sharp}(\alpha), r_{\sharp}(\beta)]) = 0$$

where  $r_{\sharp} : \mathfrak{g}^* \to \mathfrak{g}$  is the linear homomorphism given by  $\beta(r_{\sharp}(\alpha)) = r(\alpha, \beta)$ . Every solution r of CYBE on G, determines a Lie algebra structure on the dual vector space  $\mathfrak{g}^*$  (see [13]), with Lie bracket given by

$$[\alpha,\beta]_r = \mathsf{ad}^*_{r_{\mathsf{H}}(\beta)}\alpha - \mathsf{ad}^*_{r_{\mathsf{H}}(\alpha)}\beta.$$

**Definition 8.** The Lie algebra  $\mathfrak{g}^*(r) := (\mathfrak{g}^*, [\cdot, \cdot]_r)$  is called dual Lie algebra of  $\mathfrak{g}$  relative to r. We denote by  $G^*(r)$  a Lie group with Lie algebra  $\mathfrak{g}^*(r)$ .

Notice that, if  $\langle \cdot, \cdot \rangle$  is an orthogonal structure on  $\mathfrak{g}$ , then  $\mathfrak{g}^*(r)$  has a scalar product given by

$$\langle \alpha, \beta \rangle^* = \langle \phi^{-1}(\alpha), \phi^{-1}(\beta) \rangle \tag{18}$$

where  $\phi(x) = \langle x, \cdot \rangle$ . We denote by  $k^*$  the Riemannian pseudo-metric on  $G^*(r)$  determined by  $\langle \cdot, \cdot \rangle^*$  and  $\nabla^*$  the Levi-Civita connection associated to  $k^*$ .

The essential tool to be used in this section is provided by the following result (see [1]).

**Theorem 9 (Boucetta-Medina [1]).** Let (G, k) be a Lie group endowed with a bi-invariant pseudo-Riemannian metric and let r be a solution of the generalized classical Yang-Baxter equation on G. Then, we have the following

1.  $(G^*(r), k^*)$  is a locally symmetric pseudo-Riemannian manifold, i.e.,

 $\nabla^* R = 0$ 

where R is the curvature of  $k^*$ . In particular, R vanishes identically when r is a solution of the CYBE.

2. If  $k^*$  is flat then it is complete if and only if  $G^*(r)$  is unimodular.

If the two conditions are verified then  $G^*(r)$  is solvable.

Let r be a solution of the CYBE on G, then every Lie group of  $G^*(r)$  is endowed with a flat affine left invariant connection given by

$$\alpha \cdot \beta := -\mathsf{ad}^*_{r_{\mathsf{H}}(\alpha)}\beta$$

for  $\alpha, \beta \in \mathfrak{g}^*$  (see for instance [4]). The real vector space  $\mathcal{D}(\mathfrak{g}, r) := \mathfrak{g} \oplus \mathfrak{g}^*(r)$ endowed with the Lie bracket

$$[(x,\alpha),(y,\beta)] = ([x,y] + \mathrm{ad}_\alpha^* y - \mathrm{ad}_\beta^* x, [\alpha,\beta]_r + \mathrm{ad}_x^* \beta - \mathrm{ad}_y^* \alpha)$$

is called the double Lie algebra of Lie algebra  $\mathfrak{g}$  relative to r. Moreover, it is verified that the scalar product

$$\langle (x, \alpha), (y, \beta) \rangle := \beta(x) + \alpha(y)$$

turns  $\mathcal{D}(\mathfrak{g}, r)$  into an orthogonal Lie algebra. It is verified that  $\mathfrak{g}$  and  $\mathfrak{g}^*(r)$  are totally isotropic subalgebras of  $\mathcal{D}(\mathfrak{g}, r)$  (see [6]). Therefore, if dim  $\mathfrak{g} = n$ , it holds that  $\langle \cdot, \cdot \rangle$  has signature (n, n).

Boucetta and Medina [1] determined generic solutions of the CYBE of the  $\lambda$ -oscillator group  $O_{\lambda}$  of dimension n. In the case n = 4 such solutions of the CYBE are written as

$$r = \alpha_1 e_0 \wedge e_1 + \alpha_2 e_0 \wedge e_2 + \alpha_3 e_0 \wedge d$$

for  $\alpha_1, \alpha_2, \alpha_3 \in \mathbb{R}$ . The dual Lie algebra  $\mathfrak{o}^*(r) = \text{Vect}\{e_0^*, e_1^*, e_2^*, d^*\}$  has non zero Lie brackets (except antisymmetry)

$$[e_0^*, e_1^*] = -\alpha_2 d^* + \alpha_3 e_2^*, \qquad [e_0^*, e_2^*] = \alpha_1 d^* - \alpha_3 e_1^*.$$

The left symmetric product compatible with the Lie bracket is given by the following products

$$e_0^* \cdot e_0^* = \alpha_2 e_1^* - \alpha_1 e_2^*, \qquad e_0^* \cdot e_1^* = -\alpha_2 d^* + \alpha_3 e_2^*, \qquad e_0^* \cdot e_2^* = \alpha_1 d^* - \alpha_3 e_1^*.$$
(19)

The following proposition describes the dual Lie groups of O relative to solutions of the CYBE.

**Lemma 10.** The simply connected flat affine Lie groups  $O^*(r)$  corresponding to the Lie algebras  $\mathfrak{o}^*(r)$  where  $r = \alpha_1 e_0 \wedge e_1 + \alpha_2 e_0 \wedge e_2 + \alpha_3 e_0 \wedge d$  is a solution of the CYBE on O, are given by

- 1.  $O^*(r) = \mathbb{R}^4$ , if r = 0.
- 2.  $O^*(r) = \mathsf{H}_3 \times \mathbb{R}$  (direct product), if  $r = \alpha_1 e_0 \wedge e_1 \neq 0$  or  $r = \alpha_2 e_0 \wedge e_2 \neq 0$ .
- 3.  $O^*(r) = \mathbb{R}^4$  with product

$$(x, y, z, w)(x', y', z', w') = (x + x' + (w - z)y' + wz' + zw', y + y', z + z', w + w'),$$
  
if  $\alpha_3 = 0$  and  $\alpha_1 \alpha_2 \neq 0$ .

4.  $O^*(r) = (\mathbb{C} \rtimes \mathbb{R}) \times \mathbb{R}$  with product

$$(z,s,t)(z',s',t') = (z + e^{i\alpha_3 s} z', s + s', t + t'), \text{ if } \alpha_3 \neq 0.$$

**Proof:** When  $\alpha_1, \alpha_2, \alpha_3 = 0$ , we obtain that  $\mathfrak{o}^*(r)$  is abelian, then  $O^*(r) = \mathbb{R}^4$ . In the second case, the Lie algebra  $\mathfrak{o}_4^*(r)$  is isomorphic to direct sum  $\mathfrak{h}_3 \oplus \mathbb{R}d$ , it follows that  $O^*(r) = \mathsf{H}_3 \times \mathbb{R}$ .

If  $\alpha_3 = 0$ , the Lie algebra  $\mathfrak{o}^*(r)$  is isomorphic to Lie algebra

$$\mathbb{R}d^* \times_q \mathfrak{m}$$

central extension of abelian algebra  $\mathfrak{m} = \text{Vect}_{\mathbb{R}}\{e_0^*, e_1^*, e_2^*\}$  by means of the scalar two-cocycle g of  $\mathfrak{m}$  defined by

$$g(e_0^*, e_1^*) = -1, \quad g(e_0^*, e_2^*) = 1, \quad g(e_i^*, e_j^*) = 0, \quad (i, j) \neq (0, 1), (0, 2).$$
 (20)

Then, the Lie group  $O^*(r)$  with Lie algebra  $\mathfrak{o}^*(r)$  is  $\mathbb{R}^4$  with product

$$(x, y, z, w)(x', y', z', w') = (x + x' + (w - z)y' + wz' + zw', y + y', z + z', w + w').$$

For  $\alpha_3 \neq 0$ , we obtain that  $\mathfrak{o}^*(r)$  is isomorphic to Lie algebra  $\operatorname{Vect}_{\mathbb{R}}\{f_0, f_1, f_2, d\}$ 

$$[d, f_1] = \alpha_3 f_2, \qquad [d, f_2] = -\alpha_3 f_1.$$

It follows that

$$\mathfrak{o}^*(r) = \mathsf{Vect}_{\mathbb{R}}\{f_1, f_2\} \rtimes_d \mathbb{R} d \oplus \mathbb{R} f_0$$

where  $d \in \text{End}(\text{Vect}_{\mathbb{R}}\{f_1, f_2\})$ . It verifies that  $O^*(r) = (\mathbb{C} \rtimes \mathbb{R}) \times \mathbb{R}$  with the product

$$(s, z, t)(s', z', t') = (s + s', z + e^{i\alpha_3 s} z', t + t'), \text{ for all } (s, z, t), (s', z', t') \in O^*(r).$$

**Theorem 11.** The Lie groups  $O^*(r)$  defined by each solution of the CYBE are endowed with a flat left invariant Lorentz metric, a symplectic form and a flat affine left invariant connection.

**Proof:** Theorem (9) implies that  $O^*(r)$  admits a flat left invariant Lorentz metric. Moreover, this metric is given by the scalar product on  $\mathfrak{o}^*(r)$  defined by (18) and the left symmetric product is given by (19). It can be verified that the corresponding connection is the Levi-Civita connection.

A symplectic form on Lie group  $O^*(r)$  is given by the non degenerate scalar twococycle of the algebra  $\mathfrak{o}^*$  as follows

$$\omega(e_1^*, e_2^*) = \omega(e_0^*, d^*) = 1, \qquad \mathsf{Vect}_{\mathbb{R}}\{e_1^*, e_2^*\} \bot_{\omega} \mathsf{Vect}_{\mathbb{R}}\{e_0^*, d^*\}$$

for  $x, y, z \in \mathfrak{o}^*(r)$ . The formula

$$\omega(xy,z) = -\omega(y,[x,z]) \tag{21}$$

defines a left symmetric product compatible with the Lie bracket of  $\mathfrak{o}^*(r)$ . The non zero products defined by (21) are given by

 $\begin{aligned} e_0^* e_2^* &= -\alpha_3 e_1^*, \qquad e_1^* e_1^* = -\alpha_3 d^*, \qquad e_2^* e_2^* = -\alpha_3 d^* \\ e_0^* e_0^* &= -\alpha_2 e_2^* - \alpha_1 e_0^*, \qquad e_0^* e_1^* = \alpha_3 e_2^*. \end{aligned}$ 

The theorem follows from the above statements.

**Remark 12.** The left symmetric product on  $\mathfrak{o}^*(r)$  given in (21) defines a complete flat affine left invariant connection on  $O^*(r)$  which is not compatible with the symplectic form  $\omega$ .

**Remark 13 (Double Lie Groups of the Oscillator).** *The double Lie algebra of*  $\mathfrak{o}$  *relative to a solution* r *of the* CYBE *given above, correspond to the vector space* 

$$\mathcal{D}(\mathfrak{o}, r) = \mathfrak{o}^* \oplus \mathfrak{o} = \mathsf{Vect}_{\mathbb{R}} \{ e_0, e_1.e_2, d, e_0^*, e_1^*, e_2^*, d^* \}$$

with Lie brackets (non zero and except antisymmetries)

$$\begin{split} & [e_1, e_0^*] = -\alpha_3 e_2 - e_2^*, \quad [e_2, e_0^*] = \alpha_3 e_1 + e_1^*, \qquad [d, e_0^*] = -\alpha_2 e_1 + \alpha_1 e_2 \\ & [e_2, e_1^*] = -\alpha_3 e_0 - d^*, \quad [d, e_1^*] = \alpha_2 e_0 - e_2^*, \qquad [e_1, e_2^*] = \alpha_3 e_0 + d^* \\ & [d, e_2^*] = -\alpha_1 e_0 - e_1^*, \quad [e_0^*, e_1^*] = -\alpha_2 d^* + \alpha_3 e_2^*, \quad [e_0^*, e_2^*] = \alpha_1 d^* - \alpha_3 e_1^* \\ & [e_1, e_2] = e_0, \qquad [d, e_1] = e_2, \qquad [d, e_2] = -e_1. \end{split}$$

#### 3. Development of Flat Affine Manifold

This section is devoted to the description of the developing map corresponding to each flat affine left invariant structure on  $Aff(\mathbb{R})$  found in [18] and to each flat affine structure on the oscillator group O exhibited in Section 2. The method we will apply has been introduced by Koszul in [15].

**Proposition 14.** Let  $(M, \nabla)$  be a simply connected flat affine manifold of dimension n. Then there exist an affine immersion  $D : (M, \nabla) \to (\mathbb{R}^n, \nabla^0)$  and a Lie group homomorphism  $A : \operatorname{Aff}(M, \nabla) \to \operatorname{Aff}(\mathbb{R}^n, \nabla^0)$  such that the following diagram commutes

$$M \xrightarrow{D} \mathbb{R}^{n}$$

$$\downarrow_{F} \qquad \qquad \downarrow_{A(F)}$$

$$M \xrightarrow{D} \mathbb{R}^{n}$$
(22)

Moreover, A is an immersion.

**Proof:** Let  $p \in M$  and  $v \in T_p(M)$ . For a curve  $\alpha$  of class  $C^1$  on M such that  $\alpha(0) = p$  there exist a unique smooth vector field  $X(\alpha)$  along  $\alpha$  that is parallel relative to  $\nabla$  such that  $X(\alpha)_p = v$ . Since M is connected and it has an affine atlas, i.e., the transition maps are local affine transformations of open sets of  $\mathbb{R}^n$ , we can extend  $X(\alpha)$  to a smooth vector field  $X^v$  of M such that  $X_p^v = v$ , and  $\nabla X^v = 0$ . Let  $\eta$  be a differential one-form on M with values in  $T_p(M)$  such that  $\eta(X^v) = v$ . We shall see that  $\eta$  is closed. Notice that

$$(\nabla \eta)(X^v, Y) = Y(\eta(X^v)) - \eta(\nabla_Y X^v) = 0$$

for all  $Y \in \mathfrak{X}(M)$ . It follows that  $\nabla \eta = 0$  and therefore

$$\begin{aligned} 0 &= (\nabla \eta)(X, Y) - (\nabla \eta)(Y, X) \\ &= Y(\eta(X)) - \eta(\nabla_Y X) - X(\eta(Y)) + \eta(\nabla_X Y) \\ &= Y(\eta(X)) - X(\eta(Y)) - \eta(\nabla_Y X - \nabla_X Y) \\ &= Y(\eta(X)) - X(\eta(Y)) - \eta([Y, X]) \\ &= 2d\eta(Y, X) \end{aligned}$$

for all  $X, Y \in \mathfrak{X}(M)$ .

Since M is simply connected, there exist a unique differential map D from M in  $T_p(M)$  such that  $dD = \eta$ . In fact, taking a smooth curve  $\tau$  in M from p to q, it

follows from Stokes's theorem that  $\int_{\tau} \eta$  does not depend of the curve. Hence the following vector is well defined

$$D(x) = \int_{p}^{x} \eta.$$
(23)

Now we will show that D is an affine immersion from M in  $T_p(M)$ . Let  $\sigma(t)$  be a smooth curve on M, since

$$\dot{\sigma}(t)(\eta(\dot{\sigma}(t))) = \eta(\nabla_{\sigma(t)}\dot{\sigma}(t))$$

it follows that  $\sigma$  is a geodesic if and only if  $\eta(\dot{\sigma}(t))$  is constant. Then the geodesics of M are curves  $\sigma$  such that the map  $\eta \circ \sigma$  is constant.

For a geodesic  $\sigma(t)$  on  $(M, \nabla)$  we have the following

$$D(\sigma(t)) = \int_{p}^{\sigma(0)} \eta + \int_{\sigma(0)}^{\sigma(t)} \eta = \int_{p}^{\sigma(0)} \eta + \int_{\sigma(0)}^{t} \eta(\dot{\sigma}(t)) dt = \eta(\dot{\sigma}(t))t + \int_{p}^{\sigma(0)} \eta.$$

Hence,  $D(\sigma(t))$  is a line on  $T_p(M)$ , i.e.,  $D : (M, \nabla) \to (T_pM, \nabla^0)$  is an affine map. Since  $dD = \eta$ , the rank of D is constant and equal to  $n = \dim M$  in every point, thats means D is an immersion.

Next we will prove that the following map is a Lie group homomorphism

$$A: \mathsf{Aff}(M, \nabla) \to \mathsf{Aff}(\mathbb{R}^n, \nabla^0), \qquad F \mapsto (Q_F, L_F).$$

We will see that this map is well defined. Let  $L_F(v) = u = F_{*,F^{-1}(p)}(X_{F^{-1}(p)}^v)$ , with  $v \in T_p(M)$ , and  $Q_F = D(F(p))$ . Since  $\nabla F_*(X^v) = 0$  we have that  $F_*(X^v) = X^u$ , which means that  $L_F$  is  $\mathbb{R}$ -linear.

Now, we will show that  $A(F \circ G) = A(F) \circ A(G)$  for all  $F, G \in Aff(M, \nabla)$ . It directly follows from definition that  $L_{F \circ G} = L_F \circ L_G$ . Moreover

$$Q_{F \circ G} = D(F(G(p))) = \int_{\tau} \eta, \text{ where } \tau \text{ is a curve from } p \text{ to } F(G(p))$$
  
=  $\int_{\alpha} \eta + \int_{\alpha'} \eta, \text{ with } \alpha \text{ a curve from } p \text{ to } F(p)$   
and  $\alpha' \text{ is a curve from } F(p) \text{ to } F(G(p))$  (24)  
=  $D(F(p)) + \int_{F \circ \beta} \eta, \quad \beta \text{ is a curve from } p \text{ to } G(p))$   
=  $Q_F + \int_{\beta} F^* \eta.$ 

On the other hand

$$(L_F \circ \eta)_p(X_p^v) = L_F(\eta_p(X_p^v)) = L_F(v) = u$$
  
=  $\eta_{F(p)}(X_p^u)$   
=  $\eta_{F(p)}(F_{*,p}(X_p^v)) = (F^*\eta)_p(X_p^v).$  (25)

It follows that  $\int_{\beta} F^* \eta = L_F(Q_G)$ , hence, from (24) we get that A is a group homomorphism.

Now notice that for  $q \in M$  we have

$$D(F(q)) = \int_{\tau} \omega, \quad \text{where } \tau \text{ is a curve from } p \text{ to } F(q)$$
$$= Q_F + L_F(D(q)) = A(F)(D(q))$$

that is, the diagram (22) is a commutative.

Finally we prove that A is an immersion. First notice that the following diagram is commutative

where  $\Phi$  maps an affine transformation F of  $(M, \nabla)$  in the linear frame  $\widetilde{F}(u_0)$ with  $u_0$  a linear frame of M at p fixed and  $\Psi$  maps the affine transformation T of  $(\mathbb{R}^n, \nabla^0)$  in the linear frame  $\widetilde{T}(u_1)$  with  $u_1 = \widetilde{D}(u_0)$  linear frame of  $\mathbb{R}^n$  at 0. Since  $\Phi$  and  $\Psi$  are embeddings (see page 42 [11]), they are, in particular, immersions. Moreover as D is immersion its natural lift is also an immersion. Therefore A is an immersion.

The map  $D: M \to T_p(M)$  defined in the previous proposition is called *developing* map of  $(M, \nabla)$  at p.

For the following proposition we will consider  $M = \text{Aff}(\mathbb{R})_0$  with the flat affine structures  $\nabla_2(\alpha)$  and  $\nabla_1(\alpha)$  respectively given by  $\mathcal{F}_2(\alpha)$  and  $\mathcal{F}_1(\alpha)$  (see [18]).

**Proposition 15.** The corresponding developing maps of  $Aff(\mathbb{R})_0$  with values in  $aff(\mathbb{R})$  relative to  $\nabla_2(\alpha)$  and  $\nabla_1(\alpha)$  are respectively given by

$$D_{2,\alpha}(x,y) = \begin{cases} \left(\frac{1}{\alpha}(x^{\alpha}-1), \ x^{\alpha}(y+1) - \frac{1}{\alpha+1}(\alpha x^{\alpha+1}+1)\right) & \text{if } \alpha \neq 0, -1\\ \left(1 - \frac{1}{x}, \ \frac{1+y}{x} - 1 + \ln x\right) & \text{if } \alpha = -1\\ (\ln(x), \ y) & \text{if } \alpha = 0 \end{cases}$$

$$D_{1,\alpha}(x,y) = \begin{cases} \left(\frac{1}{\alpha}(x^{\alpha}-1), \ y + \frac{1}{\alpha-1}(x^{\alpha}-\alpha x) + 1\right) & \text{if } \alpha \neq 0, 1\\ \\ (x-1, \ y+1 + x(\ln x - 1)) & \text{if } \alpha = 1. \end{cases}$$

### **Proof:**

1. Let X be a smooth vector field on  $Aff(\mathbb{R})_0$  such that  $\nabla_2(\alpha)X = 0$ . The following identities can be verified

$$\nabla_2(\alpha)_{e_1^+} X = 0, \qquad \nabla_2(\alpha)_{e_2^+} X = 0.$$
 (27)

Solving (27) with initial conditions  $X_{(1,0)} = (a, b) \in \mathfrak{g}$ , we have

$$X = a x^{1-\alpha} \frac{\partial}{\partial x} + (b + \alpha a (x - y - 1)) x^{-\alpha} \frac{\partial}{\partial y} \cdot$$

It is easy to verify that the closed one-form  $\eta$  on  $Aff(\mathbb{R})_0$  with values in  $\mathfrak{g}$  such that  $\eta(X) = (a, b)$ , is given by

$$\eta = \left(x^{\alpha - 1} \mathrm{d}x, \ \alpha(y - x + 1)x^{\alpha - 1} \mathrm{d}x + x^{\alpha} \mathrm{d}y\right).$$

Taking the curve  $\tau(t) = (1 + (x - 1)t, yt)$ , with  $t \in [0, 1]$ , the developing map is written as follows

a) if  $\alpha \neq 0, -1$ 

$$D_{2,\alpha}(x,y) = \int_{\tau}^{\eta} \eta$$
  
=  $\int_{0}^{1} \left( (1 + (x - 1)t)^{\alpha - 1}(x - 1)dt, \alpha(x - 1)(y - x + 1) \times t(1 + (x - 1)t)^{\alpha - 1} + (1 + (x - 1)t))^{\alpha}y dt \right)$   
=  $\left( \frac{1}{\alpha} (x^{\alpha} - 1), x^{\alpha}(y + 1) - \frac{1}{\alpha + 1} (\alpha x^{\alpha + 1} + 1) \right).$ 

b) For the case  $\alpha = -1$ 

$$D_{2,-1}(x,y) = \int_{\tau} \eta$$
  
=  $\left(\int_{0}^{1} \frac{x-1}{(1+(x-1)t)^{2}} dt, \int_{0}^{1} \frac{(x-y-1)(x-1)t}{(1+(x-1)t)^{2}} dt + \frac{y}{1+(x-1)t} dt\right)$   
=  $\left(1 - \frac{1}{x}, \frac{1+y}{x} - 1 + \ln x\right).$ 

c) when  $\alpha = 0$ 

$$D_{2,0}(x,y) = (\ln(x), y).$$

2. If 
$$X \in \mathfrak{X}(\mathsf{Aff}(\mathbb{R})_0)$$
 such that  $X_{(1,0)} = (a,b) \in \mathfrak{g}$ , and

$$\nabla_1(\alpha)_{e_1^+} X = 0, \qquad \nabla_1(\alpha)_{e_2^+} X = 0$$

it can be verified that X is as follows

$$X = \begin{cases} ax^{1-\alpha}\frac{\partial}{\partial x} + \left(b + \frac{a\alpha}{\alpha-1}(x^{1-\alpha} - 1)\right)\frac{\partial}{\partial y} & \text{if } \alpha \neq 1\\\\ a\frac{\partial}{\partial x} + (b - a\ln x)\frac{\partial}{\partial y} & \text{if } \alpha = 1. \end{cases}$$

Hence the one-form  $\eta$  on Aff $(\mathbb{R})_0$  verifying  $\eta(X) = (a, b)$ , is given by

$$\eta = \begin{cases} \left( x^{\alpha - 1} \mathrm{d}x , \frac{\alpha}{\alpha - 1} (x^{\alpha - 1} - 1) \mathrm{d}x + \mathrm{d}y \right) & \text{ if } \alpha \neq 1 \\ \\ (\mathrm{d}x , \ln x \mathrm{d}x + \mathrm{d}y) & \text{ if } \alpha = 1. \end{cases}$$

Therefore, the developing map is as follows

$$D_{1,\alpha}(x,y) = \int_{\tau}^{\tau} \eta$$
  
= 
$$\begin{cases} \left(\frac{1}{\alpha}(x^{\alpha} - 1), \ y + \frac{1}{\alpha - 1}(x^{\alpha} - \alpha x) + 1\right) & \text{if } \alpha \neq 0, 1\\ (x - 1, \ y + 1 + x(\ln x - 1)) & \text{if } \alpha = 1. \end{cases}$$

From now on we will denote by  $\nabla_1(t,s)$ ,  $\nabla_2(\alpha,t)$ ,  $\nabla_3$  and  $\nabla_4$  the flat affine connections on O respectively given by the products  $\mathcal{F}_1(0,t,s,1/2)$ ,  $\mathcal{F}_2(\alpha,t,1/2)$ ,  $\mathcal{F}_3(1/2)$  and  $\mathcal{F}_4(1/2)$ .

**Proposition 16.** The corresponding developing maps of oscillator O relative to the flat affine connections  $\nabla_1(t,s)$ ,  $\nabla_2(\alpha,t)$ ,  $\nabla_3$  and  $\nabla_4$  are respectively

$$D_{1,t,s}(x_1, x_2, x_3, x_4) = \left(x_1 + \frac{t}{2}(x_2^2 + x_3^2) + \frac{s}{2}x_4^2, x_2, x_3, x_4\right)$$
  

$$D_{2,\alpha,t}(x_1, x_2, x_3, x_4) = \left(x_1 + \frac{t}{2}(x_2^2 + x_3^2), x_2, x_3, \frac{1}{\alpha}(e^{\alpha x_4} - 1)\right)$$
  

$$D_3(x_1, x_2, x_3, x_4) = \left(x_1, x_2, x_3, x_4 + \frac{1}{2}(x_2^2 + x_3^2)\right)$$
  

$$D_4 = (F_1, F_2, F_3, F_4)$$

where  $F_i$ , i = 1, 2, 3, 4 are detailed at the end of the proof.

### **Proof:**

1. Let  $X \in \mathfrak{X}(O)$  such that  $X_{(0,0,0,0)} = (a_1, a_2, a_3, a_4)$  with

$$\nabla_2(\alpha, t)_{e^+} X = 0, \qquad \nabla_2(\alpha, t)_{d^+} X = 0, \qquad i = 0, 1, 2$$

where  $(e_0^+, e_1^+, e_2^+, d^+)$  is the basis of left invariant vector fields on O defined in (6).

These equations define a system of partial differential equations whose solutions determine the vector field

$$X = (a_1 - a_2 t x_2 - a_3 t x_3) \frac{\partial}{\partial x_1} + a_2 \frac{\partial}{\partial x_2} + a_3 \frac{\partial}{\partial x_3} + a_4 e^{-\alpha x_4} \frac{\partial}{\partial x_4}.$$

It is easy to verify that the one-form  $\eta$  such that  $\eta(X) = a$  with a in  $\mathfrak{o}$ , is written

$$\eta = (\mathrm{d}x_1 + tx_2\mathrm{d}x_2 + tx_3\mathrm{d}x_3, \mathrm{d}x_2, \mathrm{d}x_3, \mathrm{e}^{\alpha x_4}\mathrm{d}x_4).$$

Hence the developing map is as follows

$$D_{2,\alpha,t}(x_1, x_2, x_3, x_4) = \int_{\tau} \eta$$
  
=  $\left( \int_0^1 x_1 dr + tx_2^2 r dr + tx_3^2 r dr, \int_0^1 x_2 dr, \int_0^1 x_3 dr, \int_0^1 x_4 e^{\alpha x_4 r} dr \right)$   
=  $\left( x_1 + \frac{t}{2} (x_2^2 + x_3^2), x_2, x_3, \frac{1}{\alpha} (e^{\alpha x_4} - 1) \right).$ 

2. The smooth vector field X on O such that  $X_{(0,0,0,0)} = (a_1, a_2, a_3, a_4)$  with  $\nabla_1(t, s)X = 0$ , can be written as

$$X = (a_1 - a_2 t x_2 - a_3 t x_3 - a_4 s x_4) \frac{\partial}{\partial x_1} + a_2 \frac{\partial}{\partial x_2} + a_3 \frac{\partial}{\partial x_3} + a_4 \frac{\partial}{\partial x_4}$$

On the other hand, the one-form on O such that  $\eta(X) = (a_1, a_2, a_3, a_4) \in \mathfrak{o}$ , is given by

$$\eta = (\mathrm{d}x_1 + tx_2\mathrm{d}x_2 + tx_3\mathrm{d}x_3 + sx_4\mathrm{d}x_4, \mathrm{d}x_2, \mathrm{d}x_3, \mathrm{d}x_4).$$

Taking a smooth curve  $\tau$  on O like  $\tau(r) = (x_1r, x_2r, x_3r, x_4r)$  with  $r \in [0, 1]$ , we obtain the developing map

$$D_{1,t,s}(x_1, x_2, x_3, x_4) = \int_{\tau} \eta$$
  
=  $\left( \int_0^1 x_1 dr + tx_2^2 r dr + tx_3^2 r dr + sx_4^2 r dr, \int_0^1 x_2 dr, \int_0^1 x_3 dr, \int_0^1 x_4 dr \right)$   
=  $\left( x_1 + \frac{t}{2} (x_2^2 + x_3^2) + \frac{s}{2} x_4^2, x_2, x_3, x_4 \right).$ 

3. Let X be a smooth vector field on O satisfying  $X_{(0,0,0,0)} = (a_1, a_2, a_3, a_4)$ and

$$\nabla_{3,e_i^+} X = 0, \qquad \nabla_{3,d^+} X = 0, \qquad i = 0,1,2$$

whose solutions determine the vector field

$$X = a_1 \frac{\partial}{\partial x_1} + a_2 \frac{\partial}{\partial x_2} + a_3 \frac{\partial}{\partial x_3} + (a_4 - a_2 x_2 - a_3 x_3) \frac{\partial}{\partial x_4}$$

The one-form  $\eta$ , which  $\eta(X) = (a_1, a_2, a_3, a_4) \in \mathfrak{o}$ , is written

$$\eta = (\mathrm{d}x_1, \mathrm{d}x_2, \mathrm{d}x_3, x_2\mathrm{d}x_2 + x_3\mathrm{d}x_3 + \mathrm{d}x_4).$$

As a consequence we get that the developing map of  $(O, \nabla_3)$  is given by

$$D_{3}(x_{1}, x_{2}, x_{3}, x_{4}) = \int_{\tau} \eta$$
  
=  $\left(\int_{0}^{1} x_{1} dr, \int_{0}^{1} x_{2} dr, \int_{0}^{1} x_{3} dr, \int_{0}^{1} x_{2}^{2} r + x_{3}^{2} r + x_{4} dr\right)$   
=  $\left(x_{1}, x_{2}, x_{3}, x_{4} + \frac{1}{2}(x_{2}^{2} + x_{3}^{2})\right).$ 

4. Finally, we calculate the developing map for  $(O, \nabla_4)$ . A smooth vector field X on O that is parallel relative to  $\nabla_4$  with  $X_{(0,0,0,0)} = (a_1, a_2, a_3, a_4)$  has the form

$$X = \left(-\frac{\hat{a}_9}{2}\sin(x_4) + \frac{\hat{a}_1}{2}\cos(x_4) - \frac{\hat{a}_2}{2}x_2 - \frac{\hat{a}_3}{2}x_3 + \hat{a}_4\right)e_0^+ \\ + \left(\hat{a}_5\sin(x_4) + \hat{a}_6\cos(x_4) + \frac{\hat{a}_3}{2}x_2 - \frac{\hat{a}_2}{2}x_3 + \frac{\hat{a}_4}{2}\right)e_1^+ \\ + \left(\hat{a}_7\sin(x_4) + \hat{a}_8\cos(x_4) + \hat{a}_3x_2 - \hat{a}_2x_3 + \hat{a}_4\right)e_2^+ \\ + \left(\hat{a}_1\sin(x_4) + \hat{a}_9\cos(x_4) - \hat{a}_3x_2 + \hat{a}_2x_3 - \hat{a}_4\right)d^+$$

where

$$\begin{aligned} \hat{a}_1 &= \frac{a_2}{2} + a_3 + \frac{5a_4}{4} & \hat{a}_2 &= a_3 + a_4 \\ \hat{a}_3 &= -a_2 - \frac{a_4}{2} & \hat{a}_4 &= a_1 + \frac{a_3}{2} + \frac{a_4}{2} \\ \hat{a}_5 &= \frac{1}{2} \left( -\frac{a_2}{2} + a_3 + \frac{3a_4}{4} \right) & \hat{a}_6 &= \frac{1}{2} \left( a_2 + \frac{a_3}{2} + a_4 \right) \\ \hat{a}_7 &= -\frac{3a_2}{2} - a_3 - \frac{7a_4}{4} & \hat{a}_8 &= -a_2 + \frac{3a_3}{2} + a_4 \\ \hat{a}_9 &= a_2 - \frac{a_3}{2} \cdot \end{aligned}$$

On the other hand, the one-form on O with values in  $\mathfrak{o}$  such that  $\eta(X) = (a_1, a_2, a_3, a_4)$  is  $\eta = (\omega_1, \omega_2, \omega_3, \omega_4)$  where

$$\omega_{1} = \mathrm{d}x_{1} + \frac{1}{2} \Big( (1 - x_{3})\cos(x_{4}) + (1 + x_{2})\sin(x_{4}) - \frac{1}{2} \Big) \mathrm{d}x_{2} \\ + \frac{1}{2} ((1 + x_{2})\cos(x_{4}) + (x_{3} - 1)\sin(x_{4}) - 1)\mathrm{d}x_{3} \\ + \frac{1}{2} \Big( \Big( \frac{3}{2} + x_{2} - \frac{x_{3}}{2} \Big) \cos(x_{4}) + \Big( \frac{x_{2}}{2} + x_{3} - \frac{1}{2} \Big) \sin(x_{4}) - \frac{5}{4} \Big) \mathrm{d}x_{4}$$

$$\omega_{2} = \frac{1}{2} \Big( (1+x_{2})\cos(x_{4}) + \Big(\frac{1}{2}+x_{3}\Big)\sin(x_{4}) + 1 \Big) dx_{2} \\ + \frac{1}{2} \Big( \Big(\frac{1}{2}+x_{3}\Big)\cos(x_{4}) - (x_{2}+1)\sin(x_{4}) - \frac{1}{2} \Big) dx_{3} \\ + \frac{1}{2} \Big( \Big(1+\frac{x_{2}}{2}+x_{3}\Big)\cos(x_{4}) + \Big(x_{3}-x_{2}-\frac{3}{4}\Big)\sin(x_{4}) - 1 \Big) dx_{4} \Big) dx_{4}$$

$$\begin{aligned} \omega_3 &= \left( (x_2 - 1)\cos(x_4) + \left(\frac{3}{2} + x_3\right)\sin(x_4) + 1 \right) \mathrm{d}x_2 \\ &+ \left( \left(\frac{3}{2} + x_3\right)\cos(x_4) + (1 - x_2)\sin(x_4) - \frac{1}{2} \right) \mathrm{d}x_3 \\ &+ \left( \left(1 + \frac{x_2}{2} + x_3\right)\cos(x_4) + \left(\frac{x_3}{2} - x_2 + \frac{7}{4}\right)\sin(x_4) - 1 \right) \mathrm{d}x_4 \end{aligned}$$
$$\omega_4 &= \left( (1 - x_2)\cos(x_4) - \left(\frac{1}{2} + x_3\right)\sin(x_4) - 1 \right) \mathrm{d}x_2 \\ &+ \left(\frac{1}{2} - \left(\frac{1}{2} + x_3\right)\cos(x_4) + (x_2 - 1)\sin(x_4) \right) \mathrm{d}x_3 \\ &+ \left(1 - \left(\frac{x_2}{2} + x_3\right)\cos(x_4) + \left(x_2 - \frac{x_3}{2} - \frac{5}{4}\right)\sin(x_4) \right) \mathrm{d}x_4.\end{aligned}$$

Therefore, the corresponding developing map has the following components

$$F_1(x_1, x_2, x_3, x_4) = x_1 - \frac{1}{4}(x_2 + 5x_4 + 2) - \frac{1}{4x_4}(2x_3x_4 + x_3) + \frac{\sin(x_4)}{4x_4^2}(2x_2^2 + 2x_2x_4(x_4 + 2) + 2x_3^2 - x_3(x_4 - 6)x_4 + 3x_4^2) - \frac{\cos(x_4)}{4x_4}(2x_2^2 + 2x_2x_4 + (2x_3 - 1)(x_3 + 2x_4))$$

$$F_{2}(x_{1}, x_{2}, x_{3}, x_{4}) = -\frac{x_{4}}{2} + \frac{(x_{2}^{2} + x_{3}^{2})}{2x_{4}^{2}}(\cos(x_{4}) - 1) + \frac{1}{8}(4x_{2} - 2x_{3} + 2(x_{2} + 2x_{3} + 2)\sin(x_{4}) + (4x_{2} - 4x_{3} + 3)\cos(x_{4}) - 3) + \frac{1}{4x_{4}}\left((2x_{2}^{2} + x_{3}(2x_{3} + 3))\sin(x_{4}) + 4x_{3}(\cos(x_{4}) - 1)\right)$$

$$F_3(x_1, x_2, x_3, x_4) = \left(\frac{x_2^2 + x_3^2}{4x_4^2} + \frac{4x_2 - 2x_3 - 7}{4} - \frac{x_2}{x_4}\right)\cos(x_4) \\ + \left(2x_2^2 + x_2(x_4 - 4) + 2(x_3(x_3 + x_4 + 2) + x_4)\right)\frac{\sin(x_4)}{2x_4} \\ - \frac{x_2^2 + x_3^2}{x_4^2} + \frac{7 - 2x_3}{4} + \frac{x_2}{x_4} + x_2 - x_4$$

$$F_4(x_1, x_2, x_3, x_4) = \left(\frac{2x_3 + 5}{4} - \frac{x_2^2 + x_3^2}{x_4^2} - x_2\right)\cos(x_4)$$
$$- \left(\frac{x_2^2 + x_3^2}{x_4} - \frac{2x_2 - x_3}{x_4} + \frac{x_2}{2} + x_3\right)\sin(x_4)$$
$$+ \frac{x_2^2 + x_3^2}{x_4^2} + \frac{2x_3 - 5}{4} + x_4 - x_2.$$

# **4.** Symplectic Invariant Geometry on the Lie Group $Aff(\mathbb{R})$

In what follows we consider the group of affine transformations of real line  $G := Aff(\mathbb{R})$  as a symplectic Lie group. On the natural coordinates of G the left invariant symplectic form it is written as  $\omega^+ = -\frac{1}{x^2} dx \wedge dy$ .

The following result (see [18]) will be an important tool in this section.

**Proposition 17.** There exits two left invariant symplectic flat affine, non isomorphic, connections  $\nabla_1$  and  $\nabla_2$  on Aff( $\mathbb{R}$ ). These connections correspond to the left

symmetric products on  $\mathfrak{g} = \mathfrak{aff}(\mathbb{R})$  given by the following tables

Recall that a *symplectomorphism* of  $(M, \omega)$  is a diffeomorphism f of M such that  $f^*\omega = \omega$ , i.e., a diffeomorphism which preserves the symplectic form  $\omega$ . An easy calculation shows the following

**Lemma 18.** A diffeomorphism  $f = (f_1, f_2)$  of the manifold  $Aff(\mathbb{R})_0$  is a symplectomorphism relative to  $\omega^+$  if and only if it verifies the differential equation

$$\frac{\partial f_1}{\partial x}\frac{\partial f_2}{\partial y} - \frac{\partial f_1}{\partial y}\frac{\partial f_2}{\partial x} = \left(\frac{f_1}{x}\right)^2.$$

In particular the following maps are (functionally independent) symplectomorphism

$$f(x,y) = \left(\frac{axe^y}{1 + axe^y p(y)}, ae^y\right), \text{ with } a > 0, p(y) \in C^\infty \text{ and } p > 0.$$
 (29)

$$g(x,y) = \left(e^y, \frac{e^y}{x}\right).$$
(30)

$$h(x,y) = \left(p(x) , \frac{(h(x)^2)y}{h'(x)x^2} + l(x)\right), \text{ with } p(x), l(x) \in C^{\infty} \text{ and } p > 0.$$
 (31)

The previous lemma implies the following.

**Proposition 19.** Let  $\nabla_1$  and  $\nabla_2$  be the connections given by the products (28). *Then the maps* 

$$F(x,y) = (ax, -bx^{-1} + ay + d)$$

and

$$G(x,y) = \left(ax, -2b\sqrt{x} + ay + d\right)$$

where a, b, d are real numbers with a > 0, are symplectomorphisms and affine transformations relatively to  $(G_0, \nabla_1)$  and respectively  $(G_0, \nabla_2)$ .

We finish by computing the Hess connection associated to a Lagrangian bi-foliation of  $Aff(\mathbb{R})$  and we show that each of these connections is flat.

**Proposition 20.** There are on  $(G, \omega^+)$  two transversal Lagrangian foliations given respectively by  $\mathcal{F}_i = \mathbb{R}e_i^+$ , i = 1, 2. The Hess connection associated to this bifolation ( $\mathcal{F}_1$  and  $\mathcal{F}_2$ ) is a flat affine and Lorentz connection.

**Proof:** The formula  $\omega(ab,c) = -\omega(b,[a,c])$  defines on the vector space  $\mathfrak{g} = \mathfrak{aff}(\mathbb{R})$  a left symmetric product compatible with Lie bracket given by following table

$$\begin{array}{c|cccc} \hline \odot & e_1 & e_2 \\ \hline e_1 & -e_1 & 0 \\ e_2 & -e_2 & 0 \end{array}$$
(32)

Clearly  $\mathbb{R}e_1$  and  $\mathbb{R}e_2$  are transversal Lagrangian Lie subalgebras of  $(\mathfrak{g}, \omega)$ . Recall that the Hess connection associated to  $\mathcal{F}_1$  and  $\mathcal{F}_2$  (see [10]) is defined by

$$\nabla^{H}_{(X_1,X_2)}(Y_1,Y_2) = (X_1 \odot Y_1 + [X_2,Y_1]_1, \ X_2 \odot Y_2 + [X_1,Y_2]_2)$$

where  $[, ]_i$  is the projection on the *i*-foliation and  $\odot$  is the product given by (32). A direct computation shows that this connection correspond to bilinear product on g given by

$$\begin{array}{c|cccc} \star & e_1 & e_2 \\ \hline e_1 & -e_1 & e_2 \\ e_2 & 0 & 0 \end{array} \tag{33}$$

From the fact that

$$L_{e_1} \circ L_{e_2} - L_{e_2} \circ L_{e_1} = L_{[e_1, e_2]} = L_{e_2}$$

where  $L_{e_i}$  is the endomorphism  $L_{e_i}(a) = e_i \star a$ , it follows that (33) defines a left invariant flat affine connection on Aff( $\mathbb{R}$ ).

Moreover, since  $[e_1, e_2] = e_1 \star e_2 - e_2 \star e_1$  we get that the connection  $\nabla^H$  is torsion free.

Let  $\langle \cdot, \cdot \rangle$  denote the left invariant Lorentz metric on Aff( $\mathbb{R}$ ) associated to the scalar product on  $\mathfrak{g}$  determined by  $\langle e_1, e_2 \rangle = 1$  and  $\langle e_i, e_i \rangle = 0$ , i = 1, 2, an easy calculation shows that

$$\langle e_i \star e_j, e_k \rangle + \langle e_j, e_i \star e_k \rangle = 0$$

for all i, j, k = 1, 2. Therefore the Hess connection on  $Aff(\mathbb{R})$  is a Lorentzian connection.

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