# RESOLUTION OF THE CANONICAL FIBER METRICS FOR A LEFSCHETZ FIBRATION 

Richard Melrose \& Xuwen Zhu


#### Abstract

We consider the family of constant curvature fiber metrics for a Lefschetz fibration with regular fibers of genus greater than one. A result of Obitsu and Wolpert is refined by showing that on an appropriate resolution of the total space, constructed by iterated blow-up, this family is log-smooth, i.e., polyhomogeneous with integral powers but possible multiplicities, at the preimage of the singular fibers in terms of parameters of size comparable to the logarithm of the length of the shrinking geodesic.


## Introduction

In the setting of complex surfaces, a Lefschetz fibration is a holomorphic map to a curve, generalizing an elliptic fibration in that it has only a finite number of singular points near which it is holomorphically reducible to normal crossing. Donaldson [1] showed that a four-dimensional simply-connected compact symplectic manifold, possibly after stabilization by a finite number of blow-ups, admits a Lefschetz fibration, in an appropriately generalized sense, over the sphere; Gompf [3] showed the converse. The reader is referred to the book of Gompf and Stipsicz [4] for a description of the important role played by Lefschetz fibrations in the general theory of 4-manifolds.

To cover these cases we go beyond the holomorphic case and consider a compact connected almost-complex 4-manifold $M$ and a smooth map, with complex fibers, to a Riemann surface $Z$

$$
\begin{equation*}
M \xrightarrow{\psi} Z . \tag{1}
\end{equation*}
$$

We then require that this map be pseudo-holomorphic, have surjective differential outside a finite set $F \subset M$, on which $\psi$ is injective, so $\psi: F \longleftrightarrow S \subset Z$. Near each of these singular points we require that the map be reducible to the normal crossing, or plumbing variety, model (2) below.

A curve of genus $g$ with $b$ punctures is stable if its automorphism group is finite, which is the case when $3 g-3+b>0$. In this paper we discuss Lefschetz fibrations with regular fibers having genus $g>1$ and,
hence, stable; it is straightforward to extend the results below to the punctured case when $3 g-3+b>0$ provided the singular points are away from the punctures. All fibers carry a unique metric of curvature -1 , for the singular fibers with cusp points replacing the nodes. In view of uniqueness and stability, these metrics necessarily vary smoothly near a regular fiber. We discuss here the precise uniform behavior of this family of metrics near the singular fibers, showing that in terms of appropriate (logarithmic) resolutions, of both the total and parameter spaces, to manifolds with corners the resulting fiber metric is log-smooth, i.e., essentially smooth except for the appearance of logarithmic terms in the expansions at boundary surfaces, as a Hermitian metric on an appropriate complex line bundle (reducing to the fiber tangent bundle away from the singular variety).

This refines a result of Obitsu and Wolpert [10] who gave a two-term expansion with remainder estimate. A weaker form (i.e., identification of the uniform quasi-isometry class) of this result was already implicit in the work of Masur [6] on the asymptotic behavior of the Weil-Petersson metric on the moduli space of Riemann surfaces. We refer the reader to $[\mathbf{1 0}]$ for a discussion of the earlier results in this area. In $[\mathbf{8}]$ the universal case of the Knudsen-Deligne-Mumford compactification of the moduli spaces of pointed Riemann surfaces is discussed. Again Obitsu and Wolpert have given a similar second order expansion in this general case; we use the full expansion to give a corresponding regularity result for both the Weil-Petersson and Takhtajan-Zograf metrics. The essential analytic difficulties are already present in the case of a Lefschetz fibration (in the general case more singularities occur but still essentially of the same type). Moreover, as indicated above, the Lefschetz case is of independent interest in view of the possible extensions of results from holomorphic fibrations (see, e.g., Fine [2]) to these 'minimally singular' fibrations.

The local model for degeneration for the complex structure on a Riemann surface to a surface with a node is the 'plumbing variety' with its projection to the parameter space. We add boundaries, away from the singularity at the origin, to make this into a manifold with corners:

$$
\begin{gather*}
P=\left\{(z, w) \in \mathbb{C}^{2} ; \exists t \in \mathbb{C}, z w=t,|z| \leq \frac{3}{4},|w| \leq \frac{3}{4},|t| \leq \frac{1}{2}\right\}  \tag{2}\\
P \xrightarrow{\phi} \mathbb{D}_{\frac{1}{2}}=\left\{t \in \mathbb{C} ;|t| \leq \frac{1}{2}\right\}
\end{gather*}
$$

Thus, near each point of $F$ we require that $\psi$ can be reduced to $\phi$ in (almost) holomorphic coordinates in $M$ and $Z$.

The fibers of the model Lefschetz fibration $P$ carry an explicit family of metrics of constant curvature -1 , the 'plumbing metric':

$$
\begin{equation*}
g_{P}=\left(\frac{\pi \log |z|}{\log |t|} \csc \frac{\pi \log |z|}{\log |t|}\right)^{2} g_{0}, g_{0}=\left(\frac{|d z|}{|z| \log |z|}\right)^{2} . \tag{3}
\end{equation*}
$$

We follow Obitsu and Wolpert in using this as a model for the degeneration of the global fiber metrics; indeed, our main result is that the global metric is conformal to this model with a log-smooth strictly positive conformal factor.

A (real) manifold with corners $M$ has a principal ideal $\mathcal{I}_{F} \subset \mathcal{C}^{\infty}(M)$ corresponding to each boundary hypersurface (by assumption embedded and connected) generated by a boundary defining function $\rho_{F} \geq 0$ with $F=\left\{\rho_{F}=0\right\}$ and $d \rho_{F} \neq 0$ on $F$. A smooth map between manifolds with corners $f: M \longrightarrow Y$ is an interior b-map if each of these ideals on $Y$ pulls back to non-trivial finite products of the corresponding ideals on $M$, it is b-normal if there is no common factor in these product decompositions - this is always the case here since the range space is a manifold with boundary. Such a map is a b-fibration if, in addition, every smooth vector field tangent to all boundaries on $Y$ is locally (and, hence, globally) $f$-related to such a vector field on $M$; it is then surjective. There is a slightly weaker notion than a manifold with corners, a tied manifold, which has the same local structure but in which the boundary hypersurfaces need not be embedded, meaning that transversal self-intersection is allowed. This arises below, although not in any essential way. There is still a principal ideal associated to each boundary hypersurface and the notions above carry over.

The assumptions above mean that each singular fiber of $\psi$ has one singular point at which it has a normal crossing in the (almost) complex sense as a subvariety of $M$. The real resolution of $M$ that supports the analysis below is constructed in three steps. In the first step the singular fibers of $\psi$ are blown up in the real sense; i.e., this divisor is locally replaced by a boundary which is its normal circle bundle. This is well-defined in view of the transversality of the self-intersection but results in a tied manifold since the boundary faces are not globally embedded. The effect of this step is to introduce $|z|,|w|, \theta_{z}$ and $\theta_{w}$ as local coordinates where

$$
\begin{equation*}
z=|z| e^{i \theta_{z}}, w=|w| e^{i \theta_{w}} \tag{4}
\end{equation*}
$$

The second step in the resolution is to replace this $\mathcal{C}^{\infty}$ structure by its logarithmic weakening, i.e., replacing each (local) boundary defining function $x=|z|$, or $|w|$ by

$$
\operatorname{ilog} x=\left(\log x^{-1}\right)^{-1}, x=e^{-1 / \operatorname{ilog} x}
$$

where we introduce this notation in view of the frequent appearance of this function. This gives a new tied manifold mapping smoothly to the previous one by a homeomorphism. These two steps can be thought of in combination as the 'logarithmic blow up' of the singular fibers.

The third and final step in the resolution is to blow up the corners, of codimension two, in the preimages of the singular fibers, i.e., ilog $|z|=$ ilog $|w|$. This results in a manifold with corners, $M_{\mathrm{mr}}$, with the two


Figure 1. The metric resolution.
boundary hypersurfaces denoted $B_{\mathrm{I}}$, resolving the singular fiber, and $B_{\text {II }}$ arising at the final stage of the resolution. The 'angular' variable ilog $|z| / \operatorname{ilog}|w|$ is smooth on the interior of $B_{\text {II }}$ and this function is precisely what this third blow-up resolves. The parameter space $Z$ is similarly resolved to a manifold with corners by the logarithmic blow up of each of the singular points.

It is shown below that the Lefschetz fibration lifts to a smooth map

$$
\begin{equation*}
M_{\mathrm{mr}} \xrightarrow{\psi_{\mathrm{mr}}} Z_{\mathrm{mr}}, \tag{5}
\end{equation*}
$$

which is a b-fibration. For the local model, this reduces to the statement that the function

$$
\begin{equation*}
\frac{s_{z} s_{w}}{s_{z}+s_{w}}, s_{z}=\operatorname{ilog}|z|, s_{w}=\operatorname{ilog}|w| \tag{6}
\end{equation*}
$$

which is not smooth on $[0,1) \times[0,1)$ lifts to be smooth; this is the reason for the third step in the construction.

This construction is perhaps best justified by the main result below; its relevance can be seen by lifting the model fibre metrics. This is carried out in Lemma 2 which shows that the plumbing metric lifts to be smooth and non-degenerate as a fiber metric on a rescaling (at the face $B_{\mathrm{II}}$ ) of the null bundle of (5) which is to say the fiber tangent bundle. This 'metric resolution' is rather clearly the simplest space on which the model, plumbing, metric is a smooth family of fiber metrics on a line bundle.

More precisely, the smooth vector fields on $M_{\mathrm{mr}}$ which are tangent to all boundaries and to the fibers of $\psi_{\mathrm{mr}}$ form the sections of a smooth vector subbundle of ${ }^{\mathrm{b}} T M_{\mathrm{mr}}$ of rank two. The boundary hypersurface $B_{\mathrm{II}}$ has a preferred class of boundary defining functions, an element of which
is denoted $\rho_{\mathrm{II}}$, arising from the logarithmic nature of the resolution, and this allows a Lie algebra of vector fields to be defined by

$$
\begin{equation*}
V \in \mathcal{C}^{\infty}\left(M_{\mathrm{mr}} ;{ }^{\mathrm{b}} T M_{\mathrm{mr}}\right), V \psi^{*} \mathcal{C}^{\infty}\left(Z_{\mathrm{mr}}\right)=0, V \rho_{\mathrm{II}} \in \rho_{\mathrm{II}}^{2} \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right) \tag{7}
\end{equation*}
$$

The possibly singular vector fields of the form $\rho_{\mathrm{II}}^{-1} V$, with $V$ as in (7), also form all the sections of a smooth vector bundle, denoted ${ }^{L} T M_{\mathrm{mr}}$. This vector bundle inherits a complex structure and, hence, has a smooth Hermitian metric, which is unique up to a positive smooth conformal factor on $M_{\mathrm{mr}}$. The main result of this paper is:

Theorem. The fiber metrics of fixed constant curvature on a Lefschetz fibration, in the sense discussed above, extend to a continuous Hermitian metric on ${ }^{L} T M_{\mathrm{mr}}$ which is related to a smooth Hermitian metric on this complex line bundle by a log-smooth conformal factor.

The notion of log-smoothness here, for a function, is the same as polyhomogeneous conormality with non-negative integral powers and linear multiplicity of slope one. Conormality in this context for $f$ : $M_{\mathrm{mr}} \longrightarrow \mathbb{R}$ can be interpreted as the 'symbol estimates' that

$$
\begin{equation*}
f \in \mathcal{A}\left(M_{\mathrm{mr}}\right) \Longleftrightarrow \operatorname{Diff}_{\mathrm{b}}^{*}\left(M_{\mathrm{mr}}\right) f \subset L^{\infty}\left(M_{\mathrm{mr}}\right) \tag{8}
\end{equation*}
$$

which, in fact, implies that the space of these functions is stable under the action, $\operatorname{Diff}_{\mathrm{b}}^{*}\left(M_{\mathrm{mr}}\right) \mathcal{A}\left(M_{\mathrm{mr}}\right) \subset \mathcal{A}\left(M_{\mathrm{mr}}\right)$. On a manifold with boundary, $M$, log-smoothness of a conormal function $f \in \mathcal{A}(M)$ means the existence of an expansion at the boundary, generalizing the Taylor series of a smooth function, so for any product decomposition near the boundary with boundary defining function $x$, there exist coefficients $a_{j, k} \in \mathcal{C}^{\infty}(\partial M), j \geq 0, j \geq k \geq 0$ such that for any finite $N$,

$$
\begin{equation*}
f-\sum_{j \leq N, 0 \leq k \leq j} a_{j, k} x^{j}(\log x)^{k} \in x^{N} \mathcal{A}([0,1) \times \partial M), \forall N \tag{9}
\end{equation*}
$$

We denote the linear space of such functions $\mathcal{C}_{\log }^{\infty}(M)$, it is independent of choices.

In the case of a manifold with corners the definition may be extended by iteration of boundary codimension. Thus, $f \in \mathcal{C}_{\log }^{\infty}\left(M_{\mathrm{mr}}\right)$ if for any product decompositions of $M_{\mathrm{mr}}$ near the two boundaries there are corresponding coefficients $a_{j, k, b} \in \mathcal{C}_{\log }^{\infty}\left(B_{b}\right), b=\mathrm{I}$, II, such that

$$
\begin{equation*}
f-\sum_{j \leq N, 0 \leq k \leq j} a_{j, k, b} x_{b}^{j}\left(\log x_{b}\right)^{k} \in x_{b}^{N} \mathcal{A}\left([0,1) \times B_{b}\right), b=\mathrm{I}, \mathrm{II}, \quad \forall N \tag{10}
\end{equation*}
$$

There are necessarily compatibility conditions between the two expansions at the corners, $B_{\mathrm{I}} \cap B_{\mathrm{II}}$, and together they determine $f$ up to a smooth function on $M_{\mathrm{mr}}$ vanishing to infinite order on both boundaries. In this sense the conformal factor in the main result above is 'essentially smooth'.

For any $z \neq 0$, the metric $g_{P}$ tends to $g_{0}$ as $t$ tends to 0 . This metric can be extended ('grafted' as in $[\mathbf{1 0}]$ ) to give an Hermitian metric on ${ }^{L} T M_{\mathrm{mr}}$ which has curvature $R$ equal to -1 near $B_{\text {II }}$ and to second order at $B_{\mathrm{I}}$. We prove the Theorem above by constructing the conformal factor $e^{2 f}$ for this metric which satisfies the curvature equation, ensuring that the new metric has curvature -1 :

$$
\begin{equation*}
(\Delta+2) f+(R+1)=-e^{2 f}+1+2 f=O\left(f^{2}\right) \tag{11}
\end{equation*}
$$

This equation is first solved in the sense of formal log-power series at both boundaries, $B_{\mathrm{I}}$ and $B_{\mathrm{II}}$, which gives us an approximate solution $f_{0}$ with

$$
-\Delta f_{0}=R+e^{2 f_{0}}+g, g \in s_{t}^{\infty} \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)
$$

Then a solution $f=f_{0}+\tilde{f}$ to (11) amounts to solving

$$
\tilde{f}=-(\Delta+2)^{-1}\left(2 \tilde{f}\left(e^{2 f_{0}}-1\right)+e^{2 f_{0}}\left(e^{2 \tilde{f}}-1-2 \tilde{f}\right)-g\right)=K(\tilde{f})
$$

Here the non-linear operator $K$ is at least quadratic in $\tilde{f}$ and the boundedness of $(\Delta+2)^{-1}$ on $\rho_{\mathrm{II}}^{-\frac{1}{2}} H_{\mathrm{b}}^{M}\left(M_{\mathrm{mr}}\right)$ for all $M$ allow the Inverse Function Theorem to be applied to show that $\tilde{f} \in s_{t}^{\infty} \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)$ and, hence, that $f$ itself is log-smooth.

In $\S 1$ the model space and metric are analyzed and in $\S 2$ the global resolution is described and the proof of the Theorem above is outlined. The linearized model involves the inverse of $\Delta+2$ for the Laplacian on the fibers and the uniform behavior, at the singular fibers, of this operator is explained in $\S 3$. The solution of the curvature problem in formal log-power series is discussed in $\S 4$ and using this the regularity of the fiber metric is shown in $\S 5$.

In [7] Mazzeo and Swoboda discussed the expansion of the WeilPetersson metric on the moduli space of genus $g$ surfaces which is closely related to the expansion for the metric discussed here. Our interest in the behavior of the fiber metrics was stimulated by the possibility, arising in discussion with Michael Singer, of extending the work of Fine [2], to the Lefschetz case.

## 1. The plumbing model

We start with a description of the real resolution of the plumbing variety, given in (2), and the properties of the fiber metric, (3), on the resolved space. As noted above there are three steps in this resolution, first the fiber complex structure is resolved, in a real sense, then two further steps are required to resolve the fiber metric.

The plumbing variety itself is smooth with $z$ and $w$ global complex coordinates - it is the model singular fibration $\phi$ which is to be 'resolved'
in the real sense. The fibers above each $t \neq 0$ are annuli

$$
\begin{equation*}
\left\{|t| \leq|z| \leq \frac{3}{4}\right\} \xrightarrow{w=t / z}\left\{|t| \leq|w| \leq \frac{3}{4}\right\}, \tag{1.1}
\end{equation*}
$$

whereas the singular fiber above $t=0$ is the union of the two discs at $z=0$ and $w=0$ identified at their origins

$$
\begin{equation*}
\phi^{-1}(0)=\left\{|z| \leq \frac{3}{4}\right\} \cup\left\{|w| \leq \frac{3}{4}\right\} /(\{z=0\} \sim\{w=0\}) \tag{1.2}
\end{equation*}
$$

Note that the differential of $\phi$ vanishes at the singular point $z=w=0$ so any smooth vector field on the range which lifts under it, i.e., is $\phi$ related to a smooth vector field on $P$, vanishes at $t=0$. Conversely, $t \partial_{t}$ is $\phi$-related to both $z \partial_{z}$ and $w \partial_{w}$ whereas the vector field

$$
\begin{equation*}
V=z \partial_{z}-w \partial_{w} \tag{1.3}
\end{equation*}
$$

annihilates $\phi^{*} t$ and so is everywhere tangent to the fibers of $\phi$.
The first step in the resolution of $\phi: P \longrightarrow \mathbb{D}_{\frac{1}{2}}$ consists in passing to the commutative square


Here $\left[\mathbb{D}_{\frac{1}{2}}, 0\right]$ is the space obtained by real blow up of the origin in the disk, which can be realized globally as

$$
\begin{equation*}
\left[\mathbb{D}_{\frac{1}{2}}, 0\right] \simeq\left[0, \frac{1}{2}\right] \times \mathbb{S} \ni(r, \theta) \longmapsto t=r e^{i \theta} \in \mathbb{D}_{\frac{1}{2}} \tag{1.5}
\end{equation*}
$$

if $\mathbb{S}=\mathbb{R} / 2 \pi \mathbb{Z}$. As a real blow-up $\left[\mathbb{D}_{\frac{1}{2}}, 0\right]$ is a well-defined manifold with boundary and any diffeomorphism of $\mathbb{D}_{\frac{1}{2}}$ fixing the origin lifts to a global diffeomorphism. The complex structure on $\mathbb{D}_{\frac{1}{2}}$ lifts to a complex structure on ${ }^{\mathrm{b}} T\left[\mathbb{D}_{\frac{1}{2}}, 0\right]$ generated by $t \partial_{t}=r \partial_{r}+i \partial_{\theta}$ in terms of (1.5).

Proposition 1. The space

$$
\begin{equation*}
P_{\bar{\partial}}=[P ;\{z=0\} \cup\{w=0\}], \tag{1.6}
\end{equation*}
$$

obtained by the real blow-up of the two normally-intersecting divisors forming the singular fiber of $\phi$, gives a commutative diagram (1.4) in which $\phi_{\bar{\partial}}$ is a b-fibration with

$$
\begin{equation*}
\phi_{\bar{\partial}}^{\frac{*}{\mathcal{I}} \mathcal{I}_{\bar{\partial}}=\mathcal{I}_{I, L} \mathcal{I}_{I, R}, ~} \tag{1.7}
\end{equation*}
$$

where $\mathcal{I}_{I, L}$ and $\mathcal{I}_{I, R}$ correspond to the two boundary components introduced by the blow-up, forming the proper transforms of $z=0$ and $w=0$, respectively.

Proof. The two divisors forming the singular fiber $\phi^{-1}(0)$ are each contained in a product neighborhood $\mathbb{D}_{\frac{1}{2}} \times \mathbb{D}_{\frac{3}{4}} \subset P$ and $\mathbb{D}_{\frac{3}{4}} \times \mathbb{D}_{\frac{1}{2}} \subset P$. The transversality of their intersection is clear and it follows that the blow-up is well-defined independently of order with the new front faces being

$$
\begin{equation*}
B_{\mathrm{I}, \mathrm{~L}}=\mathbb{S} \times\left[\mathbb{D}_{\frac{3}{4}},\{0\}\right] \subset P_{\bar{\partial}}, B_{\mathrm{I}, \mathrm{R}}=\left[\mathbb{D}_{\frac{3}{4}},\{0\}\right] \times \mathbb{S} \subset P_{\bar{\partial}} \tag{1.8}
\end{equation*}
$$

Here each of the blown up disks corresponds to the introduction of polar coordinates, so $r_{z}=|z|$ is a defining function (globally) for $B_{\mathrm{I}, \mathrm{L}}$ and $r_{w}=|w|$ for $B_{\mathrm{I}, \mathrm{R}}$. Since $r_{t}=|t|$ is a defining function for the blown-up disk in the range and

$$
\begin{equation*}
r_{t}=r_{z} r_{w} \tag{1.9}
\end{equation*}
$$

the b-fibration condition follows from the behavior of the corresponding angular variables

$$
\begin{equation*}
e^{i \theta_{t}}=e^{i \theta_{z}} e^{i \theta_{w}} \tag{1.10}
\end{equation*}
$$

As a compact manifold with corners, $P_{\bar{\partial}}$ is globally the product of an embedded manifold in $\mathbb{R}^{2}$ and a 2 -torus

$$
\begin{equation*}
P_{\bar{\partial}}=\left\{\left(r_{z}, r_{w}\right) ; 0 \leq r_{z}, r_{w} \leq \frac{3}{4}, r_{z} r_{w} \leq \frac{1}{2}\right\} \times \mathbb{S}_{z} \times \mathbb{S}_{w} \tag{1.11}
\end{equation*}
$$

This first step in the resolution resolves the complex structure in a real sense. In particular, the vector fields tangent to the fibers of $\phi_{\bar{\partial}}$ and to the boundaries form all the sections of a subbundle of ${ }^{\mathrm{b}} T P_{\overline{\bar{\partial}}}$ which has a complex structure, spanned by the lift of the single vector field (1.3).

Although the complex structure is effectively resolved, the plumbing metric in (3) is not. That $g_{P}$ has curvature -1 on the fibers, away from the singular point, can be seen by changing variables to $s=\log r, r=r_{z}$ and $\theta=\theta_{z}$ in terms of which

$$
g_{P}=\left(\frac{\pi / \log |t|}{\sin (\pi s / \log |t|)}\right)^{2}\left(d s^{2}+d \theta^{2}\right)
$$

It then follows from the standard formula for the Gauss curvature that

$$
R=-\frac{1}{2 \sqrt{f g}}\left(\partial_{s}\left(\frac{\partial_{s} g}{\sqrt{f g}}\right)+\partial_{\theta}\left(\frac{\partial_{\theta} f}{\sqrt{f g}}\right)\right)=-1
$$

In view of the coefficients in $g_{P}$ it is natural to introduce the inverted logarithms of the new boundary defining functions, so replacing the radial by the logarithmic blow-up. Thus,

$$
\begin{equation*}
s_{z}=\mathrm{i} \log r_{z}=\frac{1}{\log \frac{1}{r_{z}}}, s_{w}=\mathrm{i} \log r_{w} \tag{1.12}
\end{equation*}
$$

become new boundary defining functions in place of $r_{z}$ and $r_{w}$. The space with this new $\mathcal{C}^{\infty}$ structure can be written

$$
\begin{equation*}
\left[P ;\{z=0\}_{\log } \cup\{w=0\}_{\log }\right] \tag{1.13}
\end{equation*}
$$

However, even after this second step, the fiber metric does not have smooth coefficients:

$$
g_{P}=\frac{\pi^{2} s_{t}^{2}}{\sin ^{2}\left(\frac{\pi s_{t}}{s_{w}}\right)}\left(\frac{d s_{w}^{2}}{s_{w}^{4}}+d \theta_{w}^{2}\right)
$$

Indeed, $s_{t}=\frac{s_{z} s_{w}}{s_{z}+s_{w}}$ is not a smooth function on the space (1.13).
The final part of the metric resolution is to blow up, radially, the corner formed by the intersection of the two logarithmic boundary faces

$$
\begin{equation*}
P_{\mathrm{mr}}=\left[\left[P ;\{z=0\}_{\log } \cup\{w=0\}_{\log }\right] ;\left\{s_{z}=s_{w}=0\right\}\right] . \tag{1.14}
\end{equation*}
$$

In terms of the presentation (1.11) this preserves the torus factor and replaces the 2-manifold with corners by a new one with more smooth functions and an extra boundary hypersurface.

Proposition 2. The model Lefschetz fibration $\phi$ lifts to a b-fibration $\phi_{\mathrm{mr}}$ giving a commutative diagram


Proof. The radial variables on the spaces $P_{\bar{\partial}}$ and $\left[\mathbb{D}_{\frac{1}{2}},\{0\}\right]$ are related by

$$
\begin{equation*}
|t|=|z||w| \Longrightarrow s_{t}=\frac{s_{z} s_{w}}{s_{z}+s_{w}}, s_{t}=\mathrm{i} \log |t| \tag{1.16}
\end{equation*}
$$

so $\phi$ does not lift to be smooth. However, consider the further introduction of the radial variable $R=\left(s_{z}^{2}+s_{w}^{2}\right)^{\frac{1}{2}}$ and the smooth defining functions $R_{z}=s_{z} / R, R_{w}=s_{w} / R$ for the lifts of the two boundary hypersurfaces. Then

$$
\begin{equation*}
s_{t}=\frac{R_{z} R R_{w}}{R_{z}+R_{w}} \tag{1.17}
\end{equation*}
$$

which is smooth since $R_{z}$ and $R_{w}$ have disjoint zero sets. It follows that $\phi$ lifts to a b-fibration as in (1.15) under which the boundary ideal lifts to the product of the three ideals

$$
\begin{equation*}
\phi_{\mathrm{mr}}^{*} \mathcal{I}_{s_{t}}=\mathcal{I}_{R_{z}} \mathcal{I}_{R} \mathcal{I}_{R_{w}} \tag{1.18}
\end{equation*}
$$

The generator $V$, in (1.3), of the fiber tangent space of $\phi$ lifts to $P_{\bar{\partial}}$ as

$$
V=r_{z} \partial_{r_{z}}-r_{w} \partial_{r_{w}}-i \partial_{\theta_{z}}+i \partial_{\theta_{w}}
$$

in terms of the coordinates in (1.10) and (1.9). Under the introduction of the logarithmic variables in (1.12) it further lifts to

$$
V=s_{z}^{2} \partial_{s_{z}}-s_{w}^{2} \partial_{s_{w}}-i \partial_{\theta_{z}}+i \partial_{\theta_{w}} .
$$

In a neighborhood of the lift of the face $s_{z}=0$ to $P_{\mathrm{mr}}$ the variables $s_{w}$ (defining the new front face) and $\rho_{z}=s_{z} / s_{w} \in[0, \infty)$ (defining the lift of $s_{z}=0$ ) are valid and

$$
\begin{equation*}
V=-s_{w}\left(s_{w} \partial_{s_{w}}-\rho_{z} \partial_{\rho_{z}}-\rho_{z}^{2} \partial_{\rho_{z}}\right)-i \partial_{\theta_{z}}+i \partial_{\theta_{w}} . \tag{1.19}
\end{equation*}
$$

Reviewing the three steps in the construction of $P_{\mathrm{mr}}$, notice that the two holomorphic defining functions $z$ and $w$ are well-defined up to constant multiples and addition of (holomorphic) terms $O\left(|z|^{2}\right)$ and $O\left(|w|^{2}\right)$, respectively. Under these two changes, the logarithmic variables $s_{z}$ change to $s_{z}+s_{z}^{2} G$ with $G \in \mathcal{C}^{\infty}\left(P_{\mathrm{mr}}\right)$ smooth. The same is true of $s_{w}$ so it follows that the radial variable

$$
\begin{equation*}
R=\left(s_{z}^{2}+s_{w}^{2}\right)^{1 / 2} \in \mathcal{C}^{\infty}\left(P_{\mathrm{mr}}\right) \tag{1.20}
\end{equation*}
$$

which defines the front face, is also uniquely defined up to an additive term vanishing quadratically there. This determines a 'cusp' structure at $B_{\text {II }}$ and from (1.19) we conclude that

Lemma 1. The vector field $R^{-1} V$ on $P_{\mathrm{mr}}$ spans a smooth complex line bundle, ${ }^{L} T P_{\mathrm{mr}}$ over $P_{\mathrm{mr}}$ with underlying real plane bundle having smooth sections precisely of the form $R^{-1} W$ where $W$ is a smooth vector field tangent to the boundaries, to the fibers of $\phi_{\mathrm{mr}}$ and satisfying $W R=$ $O\left(R^{2}\right)$ at $R=0$.

It is natural to consider this bundle, precisely because
Lemma 2. The plumbing metric defines an Hermitian metric on ${ }^{L} T P_{\mathrm{mr}}$.

Proof. On $P_{\mathrm{mr}}$, in a neighborhood of the lift of $\left\{s_{z}=0\right\}$ as discussed above,

$$
s_{t}=\mathrm{i} \log |t|=\frac{s_{z} s_{w}}{s_{z}+s_{w}}=\frac{\rho_{z} s_{w}}{1+\rho_{z}}, \frac{\log |z|}{\log |t|}=\frac{1}{1+\rho_{z}},
$$

so the fiber metric lifts to

$$
\begin{equation*}
g=\frac{\pi^{2} s_{t}^{2}}{\sin ^{2}\left(\frac{\pi s_{t}}{s_{w}}\right)}\left(\frac{d s_{w}^{2}}{s_{w}^{4}}+d \theta_{w}^{2}\right)=\frac{\pi^{2} s_{t}^{2}}{\sin ^{2}\left(\frac{\pi}{1+\rho_{z}}\right)}\left(\frac{d \rho_{z}^{2}}{s_{t}^{2}\left(1+\rho_{z}\right)^{4}}+d \theta_{z}^{2}\right) \tag{1.21}
\end{equation*}
$$

This is Hermitian and the length of $V$ relative to it is a smooth positive multiple of $R^{2}$.

## 2. Global resolution and outline

It is now straightforward to extend the resolution of the plumbing variety to a global resolution of any Lefschetz fibration as outlined in the Introduction. By hypothesis, the singular fibers of a Lefschetz fibration $\psi$, as in (1), are isolated and each contains precisely one singular point. Near the singular point the map $\psi$ is reduced to $\phi$ by local complex diffeomorphisms. Thus, each singular fiber is a connected compact real manifold of dimension two with a transversal self-intersection. The real blow-up of such a submanifold is well-defined, since it is locally well-defined away from the self-intersection and well-defined near the intersection in view of the transversality. Thus,

$$
\begin{equation*}
M_{\bar{\partial}}=\left[M, \phi^{-1}(S)\right] \xrightarrow{\psi_{\bar{\jmath}}}[Z, S] \tag{2.1}
\end{equation*}
$$

reduces to $\phi_{\bar{\partial}}$ near the preimage of the finite singular set $F \subset M$. Similarly, the logarithmic step can be extended globally since away from the singular set it corresponds to replacing $|z|$, by ilog $|z|$. Here $z$ is a local complex defining function with holomorphic differential along the singular fiber. Finally, the third step is within the preimage of the set of the singular points and so is precisely the same as for the plumbing variety.

Thus, the resolved space $M_{\mathrm{mr}}$ with its global b-fibration (5) is welldefined as is the Hermitian bundle ${ }^{L} T M_{\mathrm{mr}}$ which reduces to ${ }^{L} T P_{\mathrm{mr}}$ near the singular points and is otherwise the bundle of fiber tangents to $M_{\mathrm{mr}}$ with its inherited complex structure.

To arrive at the description of the constant curvature fiber metric, as an Hermitian metric on ${ }^{L} T M_{\mathrm{mr}}$ we start with the "grafting" construction of Obitsu and Wolpert which we interpret as giving a good initial choice of Hermitian metric. Namely choose any smooth Hermitian metric $h_{0}$ on ${ }^{L} T M_{\mathrm{mr}}$; from Lemma 2

$$
\begin{equation*}
g_{\mathrm{Pl}}=e^{f_{\mathrm{Pl}}} h_{0} \text { near } B_{\mathrm{II}}, f_{\mathrm{Pl}} \text { smooth. } \tag{2.2}
\end{equation*}
$$

Away from the singular set, near the singular fiber, $\psi$ is a fibration in the real sense. Thus, it has a product decomposition, with the fibration $\psi$ the projection, and this can be chosen to be consistent with the product structure on $P$ away from the singular point. Then the complex structure on the fibers is given by a smoothly varying tensor $J$. The constant curvature metric $g_{0}$ on the resolved singular fiber may, therefore, be extended trivially to a metric on the fibers nearby, away from the singular points. This has non-Hermitian part vanishing at the singular fiber, so removing this gives a smooth family of Hermitian metrics reducing to $g_{0}$ and so with curvature equal to -1 at the singular fiber. After blow-up this remains true since the regular part of the singular fiber is replaced by a trivial circle bundle over it. On the introduction of the logarithmic variables in the base and total space, the
curvature of this smooth family, $g_{\mathrm{I}}$, is constant to infinite order at the singular fiber since it is equal to the limiting metric $g_{0}$ to infinite order. Comparing $g_{\mathrm{I}}$ to the chosen Hermitian metric gives a conformal factor $g_{\mathrm{I}}=e^{f_{\mathrm{I}}} h, f_{\mathrm{I}} \in \mathcal{C}^{\infty}(N)$ where $N$ is a neighborhood of $B_{\mathrm{I}}$ excluding a neighborhood of $B_{\mathrm{II}}$. Moreover, $g_{\mathrm{PI}}$ is also equal to the trivial extension of $g_{0}$ to second order in a compatible trivialization so the two conformal factors

$$
\begin{equation*}
f_{\mathrm{I}}=f_{\mathrm{Pl}} \text { to second order, } \tag{2.3}
\end{equation*}
$$

in their common domain of definition.
The grafting construction of Obitsu and Wolpert interpreted in this setting is then to choose a cutoff $\chi \in \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)$ equal to 1 in a neighborhood of $B_{\mathrm{II}}$ and supported near it and to set

$$
\begin{equation*}
h=e^{\chi f_{\mathrm{PI}}+(1-\chi) f_{\mathrm{I}}} h_{0} . \tag{2.4}
\end{equation*}
$$

It follows from the discussion above that $h$ is a smooth Hermitian metric on ${ }^{L} T M_{\mathrm{mr}}$ near the preimage of the singular fibers and that its curvature

$$
R(h)=\left\{\begin{array}{l}
-1 \text { near } B_{\mathrm{II}}  \tag{2.5}\\
-1+O\left(s_{t}^{2}\right) \text { near } B_{\mathrm{I}}
\end{array}\right.
$$

We, therefore, use this in place of the initial choice of Hermitian metric.
Let $g$ be the unique Hermitian constant curvature metric on the regular fibers of $\psi$, so $g=e^{2 f} h$. The curvatures are related by

$$
R(g) e^{2 f}=\Delta_{h} f+R(h)
$$

which reduces to the curvature equation

$$
\begin{equation*}
\Delta f+R(h)=-e^{2 f}, \Delta=\Delta_{h} \tag{2.6}
\end{equation*}
$$

The linearization of this equation is

$$
\begin{equation*}
(\Delta+2) f=-(R(h)+1) . \tag{2.7}
\end{equation*}
$$

The uniform invertibility of $\Delta+2$ with respect to the metric $L^{2}$ norm, shown below, implies that (2.6) has a unique small solution for small values of the parameter. The proof of the Theorem in the Introduction, therefore, reduces to the statement that (2.6) has a log-smooth solution vanishing at the boundary.

## 3. Bounds on $(\Delta+2)^{-1}$

In the linearization of the curvature equation (2.7), the operator $\Delta+2$, for the fixed initial choice of smooth fiber hermitian metric, appears. For the Laplacian on a compact manifold, $\Delta+2$ is an isomorphism of any Sobolev space $H^{k+1}$ to $H^{k-1}$, in particular, this is the case for the map from the Dirichlet space to its dual, corresponding to the case $k=0$. For a smooth family of metrics on a fibration the family of Dirichlet spaces forms the fiber $H^{1}$ space and its dual the fiber $H^{-1}$ space and $\Delta+2$
is again an isomorphism between them. These spaces are modules over the $\mathcal{C}^{\infty}$ functions of the total space and this, plus a simple commutation argument, shows that in this case of a fibration $\Delta+2$ is an isomorphism for any $k \geq 1$ between the space of functions with up to $k$ derivatives, in all directions, in the Dirichlet domain to the space with up to $k$ derivatives in the dual to the Dirichlet space. In particular, it follows from this that $\Delta+2$ is an isomorphism on functions supported away from the boundary:

$$
\begin{equation*}
\Delta+2: \mathcal{C}_{\mathrm{c}}^{\infty}\left(M_{\mathrm{reg}}\right) \longleftrightarrow \mathcal{C}_{\mathrm{c}}^{\infty}\left(M_{\mathrm{reg}}\right), M_{\mathrm{reg}}=M_{\mathrm{mr}} \backslash \partial M_{\mathrm{mr}} \tag{3.1}
\end{equation*}
$$

We extend this result up to the boundary of the resolved space for the Lefschetz fibration in terms of tangential regularity.

Proposition 3. For the Laplacian of the grafted metric

$$
\begin{equation*}
(\Delta+2)^{-1}: \rho_{I I}^{-\frac{1}{2}} H_{\mathrm{b}}^{k}\left(M_{\mathrm{mr}}\right) \longrightarrow \rho_{I I}^{-\frac{1}{2}} H_{\mathrm{b}}^{k}\left(M_{\mathrm{mr}}\right) \forall k \in \mathbb{N} \tag{3.2}
\end{equation*}
$$

The main complication in the proof arises from the fact that the Dirichlet space is not a $\mathcal{C}^{\infty}$ module.

First consider the following analog of Fubini's theorem.
Lemma 3. For the fiber metrics corresponding to an Hermitian metric on ${ }^{L} T M_{\mathrm{mr}}$, the metric density is of the form

$$
\begin{equation*}
|d g|=\rho_{I I} \nu_{\mathrm{b}, \mathrm{fib}} \tag{3.3}
\end{equation*}
$$

and the space of weighted $L^{2}$ functions with values in the $L^{2}$ spaces of the fibers can be realized as

$$
\begin{equation*}
L^{2}\left(M_{\mathrm{mr}} ;|d g| \phi_{\mathrm{mr}}^{*} \nu_{\mathrm{b}}\left(Z_{\mathrm{mr}}\right)\right)=L_{\mathrm{b}}^{2}\left(Z_{\mathrm{mr}} ; L^{2}(|d g|)\right)=\rho_{I I}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(M_{\mathrm{mr}}\right) \tag{3.4}
\end{equation*}
$$

Proof. Away from $B_{\mathrm{II}} \subset M_{\mathrm{mr}}$ the resolved map $\psi_{\mathrm{mr}}$ is a fibration, ${ }^{L} T M_{\mathrm{mr}}$ is the fiber tangent bundle and the boundary is in the base. Thus, (3.3) and (3.4) reduce to the local product decomposition for a fibration and Fubini's Theorem.

It, therefore, suffices to localize near $B_{\mathrm{II}}$ and to consider the plumbing metric since all hermitian metrics on ${ }^{L} T M_{\mathrm{mr}}$ are quasi-conformal. The symmetry in $z$ and $w$ means that it suffices to consider the region in which $\rho_{z}=s_{z} / s_{w}$ and $s_{w}$ are defining functions for the two boundary hypersurfaces $B_{\mathrm{I}}$ and $B_{\mathrm{II}}$, respectively. The plumbing metric may then be written

$$
g=\frac{\pi^{2} s_{t}^{2}}{\sin ^{2}\left(\frac{\pi s_{t}}{s_{w}}\right)}\left(\frac{d s_{w}^{2}}{s_{w}^{4}}+d \theta_{w}^{2}\right)=\frac{\pi^{2} s_{t}^{2}}{\sin ^{2}\left(\frac{\pi}{1+\rho_{z}}\right)}\left(\frac{d \rho_{z}^{2}}{s_{t}^{2}\left(1+\rho_{z}\right)^{4}}+d \theta_{z}^{2}\right)
$$

Thus, the fiber area form,

$$
|d g|=\frac{\pi^{2} s_{t}^{2}}{\sin ^{2}\left(\frac{\pi}{1+\rho_{z}}\right)} \frac{d \rho_{z}}{s_{t}\left(1+\rho_{z}\right)^{2} d \theta_{z}}=f\left(\rho_{z}\right) \frac{s_{t}}{\rho_{z}} \frac{d \rho_{z}}{\rho_{z}} d \theta_{z}=\tilde{f}\left(\rho_{z}\right) s_{w} \frac{d \rho_{z}}{\rho_{z}} d \theta_{z}
$$

is a positive multiple of $s_{w} \frac{d \rho_{z}}{\rho_{z}} d \theta_{z}$ from which (3.3) follows.

The identification (3.4) holds after localization away from $B_{\text {II }}$ and locally near it

$$
\|f\|_{\left.L_{b}^{2}\left(Z_{\mathrm{mr}}\right) ; L^{2}(d g)\right)}^{2}=\iint|f|^{2}|d g| \frac{d s_{t}}{s_{t}} d \theta_{t}=\int_{Z_{\mathrm{mr}}}|f|^{2} \rho_{\mathrm{II}} \nu_{\mathrm{b}}
$$

Since $(\Delta+2)^{-1}$ is a well-defined bounded operator on the metric $L^{2}$ space which depends continuously on the parameter in $Z \backslash S$ with norm bounded by $1 / 2$, it follows from (3.4) that

$$
\begin{equation*}
(\Delta+2)^{-1} \text { is bounded on } \rho_{\mathrm{II}}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(M_{\mathrm{mr}}\right) \tag{3.5}
\end{equation*}
$$

We consider the 'total' Dirichlet space based on this $L^{2}$ space - we are free to choose the weighting in the parameter space. Thus, let $D$ be the completion of the smooth functions on $M_{\mathrm{mr}}$ supported in the interior with respect to

$$
\begin{equation*}
\|u\|_{D}^{2}=\int\left(\left|d_{\mathrm{fib}} u\right|_{g}^{2}+2|u|^{2}\right)|d g| \phi^{*} \nu_{\mathrm{b}}\left(Z_{\mathrm{mr}}\right) \tag{3.6}
\end{equation*}
$$

Note that $D$ depends only on the quasi-isometry class of the fiber Hermitian metric but does depend on the induced fibration of the boundary $B_{\mathrm{II}}$.

The dual space, $D^{\prime}$, to $D$ as an abstract Hilbert space, may be embedded in the extendible distributions on $M_{\mathrm{mr}}$ using the volume form $\phi_{\mathrm{mr}}^{*} \nu_{\mathrm{b}}|d g|$. As is clear from the discussion below, the image is independent of the choice of, $\nu_{\mathrm{b}}$, of a logarithmic area form on $Z_{\mathrm{mr}}$ but the embedding itself depends on this choice. Thus, $\tilde{v} \in D^{\prime}$ is identified as a map $v: \dot{\mathcal{C}}_{c}^{\infty}\left(M_{\mathrm{mr}}\right) \longrightarrow \mathbb{C}$ by

$$
\begin{equation*}
\int v \phi|d g| \phi^{*} \nu_{\mathrm{b}}\left(Z_{\mathrm{mr}}\right)=\tilde{v}(\phi) \tag{3.7}
\end{equation*}
$$

We consider the space of vector fields $\mathcal{W} \subset \rho_{\text {II }}^{-1} \mathcal{V}_{\mathrm{b}}\left(M_{\mathrm{mr}}\right)$ which are tangent to the fibers of $\psi_{\mathrm{mr}}$ and to the fibers of $B_{\mathrm{II}}$ and which commute with $\partial_{\theta_{z}}$ and $\partial_{\theta_{w}}$ near $B_{\mathrm{II}}$.

Proposition 4. For the grafted metric

$$
\Delta+2: D \rightarrow D^{\prime} \subset \mathcal{C}^{-\infty}\left(M_{\mathrm{mr}}\right)
$$

is an isomorphism, where the elements of $D^{\prime}$ are precisely those extendible distributions which may be written as finite sums

$$
\begin{equation*}
v=\sum_{j} W_{j} u_{j}, W_{j} \in \mathcal{W}, u_{j} \in \rho_{I I}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(Z_{\mathrm{mr}}\right) \tag{3.8}
\end{equation*}
$$

and has the injectivity property that

$$
\begin{equation*}
u \in \mathcal{C}^{-\infty}\left(M_{\mathrm{mr}}\right),(\Delta+2) u \in D^{\prime} \Longrightarrow u \in D \tag{3.9}
\end{equation*}
$$

This result remains true for any Hermitian metric on ${ }^{L} T M_{\mathrm{mr}}$ but is only needed here for the grafted metric which is equal to the plumbing metric near $B_{\mathrm{II}}$.

Proof. Although defined above by completion of the space of smooth functions supported away from the boundary of $M_{\mathrm{mr}}$ with respect to the norm (3.6) the space $D$ can be identified in the usual way with the subspace of $\mathcal{C}^{-\infty}\left(M_{\mathrm{mr}}\right)$ consisting of those

$$
\begin{equation*}
u \in \rho_{\mathrm{II}}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(M_{\mathrm{mr}}\right) \text { s.t. } \mathcal{W} \cdot u \subset \rho_{\mathrm{II}}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(M_{\mathrm{mr}}\right) \tag{3.10}
\end{equation*}
$$

with the derivatives taken in the sense of extendible distributions. Indeed, choosing a cutoff $\mu \in \mathcal{C}_{c}^{\infty}(\mathbb{R})$ which is equal to 1 near 0 the sequence of multiplication operators $1-\mu\left(n \rho_{\text {II }}\right)$ tends strongly to the identity on $\rho_{\mathrm{II}}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(M_{\mathrm{mr}}\right)$. By assumption this commutes with the elements of $\mathcal{W}$ and it follows that elements with support in the interior of $M_{\mathrm{mr}}$, where $\psi_{\mathrm{mr}}$ is a fibration, are dense in $D$; for these approximation by smooth elements is standard.

That $\Delta+2: D \longrightarrow D^{\prime} \subset \mathcal{C}^{-\infty}\left(M_{\mathrm{mr}}\right)$ is the explicit form of the Riesz representation theorem in this setting. Then the identification, (3.8), of elements of $D^{\prime}$ follows from the form of $\Delta$. Away from $B_{\mathrm{II}}, D$ is a $\mathcal{C}^{\infty}$ module (since the elements of $\mathcal{W}$ are smooth there) and then (3.8) is the identification of the fiber $H^{-1}$ space. Near $B_{\text {II }}$ we may use the explicit form of the Laplacian for the plumbing metric.

Indeed, the local version of the Dirichlet form is

$$
\begin{equation*}
D(\phi, \psi)=\int\left(V_{\operatorname{Re}} \phi \overline{V_{\operatorname{Re}} \phi}+V_{\operatorname{Im}} \phi \overline{V_{\operatorname{Im}} \phi}\right) \frac{d s_{w} d \theta_{w}}{s_{w}^{2}} \tag{3.11}
\end{equation*}
$$

where $V$ is given by (1.19) and it follows that the Laplacian acting on functions on the fibers can be written

$$
\begin{equation*}
\Delta=-\frac{\sin ^{2}\left(\frac{\pi}{1+\rho_{z}}\right)}{\pi^{2} s_{t}^{2}}\left(V_{\mathbb{R}}^{2}+\left(\partial_{\theta_{z}}-\partial_{\theta_{w}}\right)^{2}\right) \tag{3.12}
\end{equation*}
$$

in the coordinates $s_{w}, \rho_{z}, \theta_{w}$ and $\theta_{z}$.
The vector fields $V_{\mathbb{R}}$ and $\rho_{\mathrm{II}}^{-1}\left(\partial_{\theta_{z}}-\partial_{\theta_{w}}\right)$ generate $\mathcal{W}$ near $B_{\mathrm{II}}$ over the functions which are constant in $\theta_{w}$ and $\theta_{z}$. If we write $\operatorname{Diff}{ }_{\mathcal{W}}^{k}\left(M_{\mathrm{mr}}\right)$ for the differential operators which can be written as sums of products of elements of at most $k$ elements of $\mathcal{W}$ with smooth coefficients which are independent of the angular variables near $B_{\text {II }}$ then

$$
\begin{equation*}
\Delta \in \operatorname{Diff}_{\mathcal{W}}^{2}\left(M_{\mathrm{mr}}\right) \tag{3.13}
\end{equation*}
$$

Moreover,

$$
\begin{gather*}
\operatorname{Diff}_{\mathcal{W}}^{1}\left(M_{\mathrm{mr}}\right): D \longrightarrow \rho_{\mathrm{II}}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(M_{\mathrm{mr}}\right) \text { and } \\
\operatorname{Diff}_{\mathcal{W}}^{1}\left(M_{\mathrm{mr}}\right): \rho_{\mathrm{II}}^{-\frac{1}{2}} L_{\mathrm{b}}^{2}\left(M_{\mathrm{mr}}\right) \longrightarrow D^{\prime} \tag{3.14}
\end{gather*}
$$

where the second statement follows by duality from the first. Together (3.13) and (3.14) imply (3.8). q.e.d.

Consider the space $\mathcal{U} \subset \mathcal{V}_{\mathrm{b}}\left(M_{\mathrm{mr}}\right)$, defined analogously to $\mathcal{W}$, as consisting of the vector fields which commute with $\partial_{\theta_{z}}$ and $\partial_{\theta_{w}}$ near $B_{\mathrm{II}}$. Then let $\operatorname{Diff}_{\mathcal{U}}^{k}\left(M_{\mathrm{mr}}\right)$ be the part of the enveloping algebra of $\mathcal{U}$ up to order $k$, this just consists of the elements of $\operatorname{Diff}_{\mathrm{b}}^{k}\left(M_{\mathrm{mr}}\right)$ which commute with $\partial_{\theta_{z}}$ and $\partial_{\theta_{w}}$ near $B_{\mathrm{II}}$. We may define 'higher order' versions of the spaces $D$ and $D^{\prime}$ :

$$
\begin{align*}
& D_{k}=\left\{u \in D ; \operatorname{Diff}_{\mathcal{U}}^{k}\left(M_{\mathrm{mr}}\right) \cdot u \subset D\right\}  \tag{3.15}\\
& \quad D_{k}^{\prime}=\left\{u \in D^{\prime} ; \operatorname{Diff}_{\mathcal{U}}^{k}\left(M_{\mathrm{mr}}\right) \cdot u \subset D^{\prime}\right\}, k \in \mathbb{N}
\end{align*}
$$

Since $\mathcal{U}$ spans $\mathcal{V}_{\mathrm{b}}\left(M_{\mathrm{mr}}\right)$ over $\mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)$ it follows that

$$
\begin{equation*}
D_{k} \subset \rho_{\mathrm{II}}^{-\frac{1}{2}} H_{\mathrm{b}}^{k}\left(M_{\mathrm{mr}}\right) \subset D_{k}^{\prime} \forall k \tag{3.16}
\end{equation*}
$$

Proposition 5. For any $k, \dot{\mathcal{C}}^{\infty}\left(M_{\mathrm{mr}}\right)$ is dense in $D_{k}$ and $D_{k}^{\prime}$ and

$$
\begin{equation*}
\Delta+2: D_{k} \longrightarrow D_{k}^{\prime} \tag{3.17}
\end{equation*}
$$

is an isomorphism.
Proof. The density statement follows from the same argument as for $D$ and $D^{\prime}$.

Consider the commutator relation which follows directly from the definitions
$[\mathcal{U}, \mathcal{W}] \subset \mathcal{W} \Longrightarrow\left[\operatorname{Diff}_{\mathcal{U}}^{k}\left(M_{\mathrm{mr}}\right), \Delta\right] \subset \operatorname{Diff}_{\mathcal{W}}^{2}\left(M_{\mathrm{mr}}\right) \cdot \operatorname{Diff}_{\mathcal{U}}^{k-1}\left(M_{\mathrm{mr}}\right), k \in \mathbb{N}$.
To prove (3.17) we need to show that if $u \in D, Q \in \operatorname{Diff}_{\mathcal{U}}^{k}\left(M_{\mathrm{mr}}\right)$ and $f=(\Delta+2) u \in D_{k}^{\prime}$ then $Q u \in D$. Assuming the result for $Q \in$ $\operatorname{Diff}_{\mathcal{U}}^{k-1}\left(M_{\mathrm{mr}}\right)$ it follows from (3.18) that

$$
\begin{align*}
& \Delta Q u=Q \Delta u+\sum_{p} L_{p} Q_{p} u  \tag{3.19}\\
& \quad \begin{aligned}
\text { with } L_{p} \in \operatorname{Diff}_{\mathcal{W}}^{2}\left(M_{\mathrm{mr}}\right), \quad & Q_{p}
\end{aligned} \quad \in \operatorname{Diff}_{\mathcal{U}}^{k-1}\left(M_{\mathrm{mr}}\right) \\
& \Longrightarrow \Delta Q u \in D^{\prime} \Longrightarrow Q u \in D
\end{align*}
$$

by distributional uniqueness. q.e.d.

Proof of Proposition 3. The boundedness in (3.2) follows directly from (3.17) and (3.16). q.e.d.

## 4. Formal solution of $(\Delta+2) u=f$

In the previous section the uniform invertibility of $\Delta+2$ for the grafted metric was established. In particular, the case $k=\infty$ in (3.2) shows the invertibility on conormal functions. In this section, we solve
the same equation, $(\Delta+2) u=f$ in formal power series with logarithmic terms.

Let $\mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right) \subset \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)$ denote the subspace annihilated to infinite order at $B_{\text {II }}$ by the angular operators $D_{\theta_{z}}$ and $D_{\theta_{w}}$.

Lemma 4. The restriction, $\Delta_{I}$, of the Laplacian to $B_{I}$ satisfies

$$
\begin{align*}
& \left(\Delta_{I}+2\right)^{-1}\left(\rho_{I I}\left(\log \rho_{I I}\right)^{k} g_{k}\right)  \tag{4.1}\\
& \quad=\rho_{I I} \sum_{0 \leq p \leq k+1}\left(\log \rho_{I I}\right)^{p} u_{p}, u_{p} \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right) \forall g_{k} \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)
\end{align*}
$$

Proof. The fiber metric on $B_{\mathrm{I}}$ is a trivial family with respect to the product decomposition $B_{\mathrm{I}}=A \times \mathbb{S}$ where $A$ has the complete metric on the Riemann surface with cusps arising from the 'removal' of the nodal points. The Laplacian is, therefore, essentially self-adjoint and nonnegative, so $\Delta+2$ is invertible. Either from the form of a parametrix or by Fourier expansion near the cusps it follows that rapid decay in the non-zero Fourier modes (in both angular variables) is preserved by $\left(\Delta_{I}+2\right)^{-1}$. Near the boundary the zero Fourier mode satisfies a reduced, ordinary differential, equation with regular singular points and having indicial roots 1 and -2 in terms of a defining function for the (resolved) cusps. Then (4.1) follows directly.
q.e.d.

Lemma 5. If $u \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)$ then $\Delta u \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)$ restricts to $B_{\text {II }}$ to $\widetilde{\Delta}_{I I} v, v=\left.u\right|_{B_{I I}}$ where $\widetilde{\Delta}_{I I}$ is an ordinary differential operator of order 2 elliptic in the interior with regular singular endpoints, with indicial roots $-1,2$ such that

$$
\begin{equation*}
\operatorname{Nul}\left(\widetilde{\Delta}_{I I}+2\right) \subset \rho_{I}^{-1} \mathcal{C}^{\infty}\left(B_{I I}\right) \tag{4.2}
\end{equation*}
$$

has no smooth elements and for $h_{j} \in \mathcal{C}_{F}^{\infty}\left(B_{I I}\right)$

$$
\begin{gather*}
\left(\widetilde{\Delta}_{I I}+2\right)^{-1}\left(\log \rho_{I}\right)^{j} h_{j}=\sum_{0 \leq q \leq j}\left(\log \rho_{I}\right)^{q} v_{q, j}+\rho_{I}^{2}\left(\log \rho_{I}\right)^{j+1} w_{j}  \tag{4.3}\\
\text { with } v_{q, j}, w_{j} \in \mathcal{C}_{F}^{\infty}\left(B_{I I}\right)
\end{gather*}
$$

Proof. The form of the Laplacian in (3.12) shows that the reduced operator $\widetilde{\Delta}_{\text {II }}$ exists and after the change coordinates on $B_{\text {II }}$ to

$$
\begin{equation*}
\rho=\frac{1}{1+\rho_{\mathrm{II}}} \tag{4.4}
\end{equation*}
$$

becomes

$$
\begin{equation*}
\Delta+2=2-\left(\frac{\sin (\pi \rho)}{\pi \rho}\right)^{2}\left[\left(\rho \partial_{\rho}\right)^{2}-\rho \partial_{\rho}\right] \tag{4.5}
\end{equation*}
$$

The indicial roots of this operator are 2 and -1 and its homogeneity shows that the null space has no logarithmic terms. The absence of smooth elements in the null space follows by integration by parts and positivity.

The problem that we need to solve at $B_{\text {II }}$ is

$$
\begin{equation*}
(\Delta+2)\left(\rho_{\mathrm{II}} w\right)=\rho_{\mathrm{II}} g+O\left(\rho_{\mathrm{II}}^{2}\right) \Longrightarrow\left(\widetilde{\Delta}_{\mathrm{II}}^{(1)}+2\right)\left(\left.w\right|_{B_{\mathrm{II}}}\right)=\left.g\right|_{B_{\mathrm{II}}} . \tag{4.6}
\end{equation*}
$$

Since the parameter, $s_{t}$, is the product of defining functions for $B_{\mathrm{I}}$ and $B_{\text {II }}$ and commutes through the problem this can be solved by dividing by it. Thus, $\widetilde{\Delta}_{\text {II }}^{(1)}$ is obtained from $\widetilde{\Delta}_{\text {II }}$ by conjugating by a boundary defining function on $B_{\text {II }}$ so the preceding Lemma can be applied after noting the shift of the indicial roots.

Lemma 6. For the conjugated operator on $B_{I I}$,

$$
\begin{equation*}
\operatorname{Nul}\left(\widetilde{\Delta}_{I I}^{(1)}+2\right) \subset \mathcal{C}^{\infty}\left(B_{I I}\right) \tag{4.7}
\end{equation*}
$$

with the Dirichlet problem uniquely solvable and

$$
\begin{gather*}
\left(\widetilde{\Delta}_{I I}+2\right)^{-1}\left(\log \rho_{I}\right)^{j} h_{j}=\sum_{0 \leq q \leq j}\left(\log \rho_{I}\right)^{q} v_{q, j}+\rho_{I}^{3}\left(\log \rho_{I}\right)^{j+1} w_{j}  \tag{4.8}\\
\text { with } v_{q, j}, w_{j} \in \mathcal{C}_{F}^{\infty}\left(B_{I I}\right) .
\end{gather*}
$$

To express the form of the expansion which occur below, consider the space of polynomials in $\log \rho_{\mathrm{I}}$ and $\log \rho_{\mathrm{II}}$ with coefficients in $\mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)$

$$
\begin{equation*}
\mathcal{P}^{k}=\left\{u=\sum_{0 \leq l+p \leq k}\left(\log \rho_{\mathrm{I}}\right)^{l}\left(\log \rho_{\mathrm{II}}\right)^{p} u_{l, p}, u_{l, p} \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)\right\} \tag{4.9}
\end{equation*}
$$

We also consider the filtration of these spaces by the maximal order in each of the variables:
(4.10)

$$
\begin{aligned}
\mathcal{P}_{\mathrm{I}}^{k, j} & = \\
\{u & \left.=\sum_{0 \leq l+p \leq k,}\left(\log \rho_{\mathrm{I}}\right)^{l}\left(\log \rho_{\mathrm{II}}\right)^{p} u_{l, p}, u_{l, p} \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)\right\}, j \leq k ; \\
\mathcal{P}_{\mathrm{II}}^{k, m} & = \\
\{u & \left.=\sum_{0 \leq l+p \leq k, p \leq m}\left(\log \rho_{\mathrm{I}}\right)^{l}\left(\log \rho_{\mathrm{II}}\right)^{p} u_{l, p}, u_{l, p} \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)\right\}, m \leq k .
\end{aligned}
$$

Since the coefficients are in $\mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right), \Delta$ acts as a smooth b-differential operator on all of these spaces. If $u \in \mathcal{P}_{\mathrm{I}}^{k, p}$, then $u=u_{p}+u^{\prime}$ with $u^{\prime} \in \mathcal{P}_{\mathrm{I}}^{k, p-1}$ and $u_{p}=v\left(\log \rho_{\mathrm{I}}\right)^{p}$ where $v \in \mathcal{P}_{\mathrm{I}}^{k-p, 0}$. Then $\Delta u=$ $\left(\Delta_{\mathrm{I}} v\right)\left(\log \rho_{\mathrm{I}}\right)^{p}+f^{\prime}, f^{\prime} \in \mathcal{P}_{\mathrm{I}}^{k-1, p-1}+\rho_{\mathrm{I}} \mathcal{P}_{\mathrm{I}}^{k, p-1}$ where the first error term corresponds to at least one derivation of $\left(\log \rho_{\mathrm{I}}\right)^{p}$. Similar statements apply to $B_{\text {II }}$ and $\tilde{\Delta}_{\text {II }}$.

As a basis for iteration, to capture the somewhat complicated behavior of the logarithmic terms, we first consider a partial result.

Proposition 6. For each $k$

$$
\begin{equation*}
f \in \rho_{I I} \mathcal{P}^{k}+\rho_{I} \rho_{I I} \mathcal{P}^{k+1} \Longrightarrow \exists u \in \rho_{I I} \mathcal{P}^{k+1}+\rho_{I}^{2} \rho_{I I} \mathcal{P}_{I I}^{k+2, k+1} \tag{4.11}
\end{equation*}
$$

such that

$$
\begin{equation*}
(\Delta+2) u-f \in s_{t}\left(\rho_{I I} \mathcal{P}^{k+1}+\rho_{I} \rho_{I I} \mathcal{P}^{k+2}\right) \tag{4.12}
\end{equation*}
$$

Proof. We first solve on $B_{\mathrm{I}}$, then on $B_{\mathrm{II}}$. The second term in $f$ in (4.11) vanishes on $B_{\mathrm{I}}$ so the restriction $\left.f_{\mathrm{I}} \in \rho_{\mathrm{II}} \mathcal{P}^{k}\right|_{B_{\mathrm{I}}}$. Proceeding iteratively, suppose

$$
f \in \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{I}}^{k, j}+\rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}^{k+1}
$$

with $j \leq k$ and consider the term of order $j$ in $\log \rho_{\mathrm{I}}$; this is a polynomial in $\log \rho_{\mathrm{II}}$ of degree at most $k-j$ with coefficients in $\rho_{\mathrm{II}} \mathcal{C}_{F}^{\infty}\left(B_{\mathrm{I}}\right)$. Applying Lemma 4 to the restriction to $B_{\mathrm{I}}$ gives a polynomial in $\log \rho_{\mathrm{II}}$ of degree at most $k-j+1$ with coefficients in $\rho_{\mathrm{II}} \mathcal{C}_{F}^{\infty}\left(B_{\mathrm{I}}\right)$. Extending these coefficients off $B_{\mathrm{I}}$ and restoring the coefficient of $\left(\log \rho_{\mathrm{I}}\right)^{j}$ gives $v_{j} \in \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{I}}^{k+1, j}$ such that

$$
(\Delta+2) v_{j}-f=-f^{\prime}, f^{\prime} \in \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{I}}^{k, j-1}+\rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}^{k+1}
$$

Here the first part of the error arises from differentiation of the factor $\left(\log \rho_{\mathrm{I}}\right)^{j}$ in $v_{j}$ at least once. If we start with $j=k$ and proceed iteratively over decreasing $j$ this allows us to find $v \in \rho_{\text {II }} \mathcal{P}^{k+1}$ such that

$$
\begin{equation*}
(\Delta+2) v-f=-g \in \rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}^{k+1} \tag{4.13}
\end{equation*}
$$

Now we proceed similarly by solving on $B_{\text {II }}$ using Lemma 6 . So, suppose $h \in \rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{k+1, p}$, for $p \leq k+1$. Then the coefficient $h_{p}$ of $\left(\log \rho_{\mathrm{II}}\right)^{p}$ is a polynomial of degree at most $k+1-p$ in $\log \rho_{\mathrm{I}}$ with coefficients in $\rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)$. Conjugating away the factor of $\rho_{\mathrm{II}}$ and applying Lemma 6 to the restriction to $B_{\text {II }}$ and then extending the coefficients off $B_{\mathrm{II}}$ allows us to find $w_{p} \in \rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{k+1, p}+\rho_{\mathrm{I}}^{2} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{k+2, p}$, where the second term arises from the possible increase in multiplicity of the logarithmic coefficient of $\rho_{\mathrm{I}}^{2}$ in the solution, satisfying
$(\Delta+2) w_{p}-g=-g^{\prime}+e, g^{\prime} \in \rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{k+1, p-1}, e \in \rho_{\mathrm{I}} \rho_{\mathrm{II}}^{2} \mathcal{P}_{\mathrm{II}}^{k+1, p}+\rho_{\mathrm{I}}^{2} \rho_{\mathrm{II}}^{2} \mathcal{P}_{\mathrm{II}}^{k+2, p}$, where the first part of the error arises from differentiation of $\left(\log \rho_{\mathrm{II}}\right)^{p}$ at least once. Starting with $p=k+1$ and iterating over decreasing $p$ allows us to find $w \in \rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}^{k+1}+\rho_{\mathrm{I}}^{2} \rho_{\mathrm{II}} \mathcal{P}^{k+2, k+1}$ such that

$$
\begin{equation*}
(\Delta+2) w-g \in \rho_{\mathrm{I}} \rho_{\mathrm{II}}^{2} \mathcal{P}^{k+1}+\rho_{\mathrm{I}}^{2} \rho_{\mathrm{II}}^{2} \mathcal{P}^{k+2} \tag{4.15}
\end{equation*}
$$

Combining (4.13) and (4.15) gives (4.12) since $\rho_{\mathrm{I}} \rho_{\mathrm{II}}$ is a smooth multiple of $s_{t}$. q.e.d.

Proposition 6 allows iteration since $s_{t}$ commutes through $\Delta+2$.
Proposition 7. If $f \in \rho_{I I} \mathcal{P}^{k}+\rho_{I} \rho_{I I} \mathcal{P}^{k+1}$ then $u=(\Delta+2)^{-1} f \in$ $s_{t}^{-\epsilon} H_{\mathrm{b}}^{\infty}\left(M_{\mathrm{mr}}\right)$ for any $\epsilon>0$, has a complete asymptotic expansion of the form

$$
\begin{equation*}
u \simeq \sum_{j \geq 0} s_{t}^{j} u_{j}, u_{j} \in \rho_{I I} \mathcal{P}^{k+j}+\rho_{I} \rho_{I I} \mathcal{P}_{I I}^{k+j+1, k+j} \tag{4.16}
\end{equation*}
$$

Proof. For any $\epsilon>0, g=s_{t}^{\epsilon} f \in \rho_{\mathrm{II}}^{-\frac{1}{2}} H_{\mathrm{b}}^{\infty}\left(M_{\mathrm{mr}}\right)$ so $u=s_{t}^{-\epsilon}(\Delta+2)^{-1} g$ exists by (3.2). Comparing $u$ to the expansion cut off at a finite point gives (4.16).
q.e.d.

This result can itself be iterated, asymptotically summed and then the rapidly decaying remainder term again removed to show the logsmoothness of the solution for an asymptotically convergent sum over terms on the right in (4.16).

For the solution of the curvature equation the leading term is smooth because of the special structure of the forcing term.

Lemma 7. If $f \in \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)$ has support disjoint from $B_{I I}$ then $u=$ $(\Delta+2)^{-1} f$ is log-smooth and has an asymptotic expansion of the form

$$
\begin{equation*}
u \simeq \rho_{I I} v_{0}+\sum_{k \geq 1} s_{t}^{k} v_{k}, v_{k} \in \rho_{I I} \mathcal{P}^{k}+\rho_{I} \rho_{I I} \mathcal{P}_{I I}^{k+1, k} \tag{4.17}
\end{equation*}
$$

Note that log-smoothness follows from the fact that $s_{t}=a \rho_{\mathrm{I}} \rho_{\mathrm{II}}, a \in$ $\mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right)$ so each term in the expansion can be written as a polynomial in $\rho_{\mathrm{I}}, \rho_{\mathrm{I}} \log \rho_{\mathrm{I}}, \rho_{\mathrm{II}}$ and $\rho_{\mathrm{II}} \log \rho_{\mathrm{II}}$ of degree at least $2 k$.

## 5. Log-smoothness for the curvature equation

Under a conformal change from the grafted metric $h$ with curvature $R$ to $e^{2 f} h$ the condition for the curvature of the new metric to be -1 is given by (11). To construct the canonical metrics on the fibers we proceed, as in the linear case discussed above, to solve (11) in the sense of formal power series at the two boundaries above $s_{t}=0$ and then, using the Implicit Function Theorem deduce that the actual solution has this asymptotic expansion.

Lemma 8. For the grafted metric there is a formal log-power series

$$
\begin{equation*}
\sum_{k \geq 2} s_{t}^{k} f_{k}, f_{2} \in \mathcal{C}_{F}^{\infty}\left(M_{\mathrm{mr}}\right), f_{k} \in \rho_{I I} \mathcal{P}^{k-2}+\rho_{I} \rho_{I I} \mathcal{P}_{I I}^{k-1, k-2}, k \geq 3 \tag{5.1}
\end{equation*}
$$

## solving (11).

The $\mathcal{P}^{k}$ are defined in (4.9); in the last term there is no factor of $\left(\log \rho_{\mathrm{II}}\right)^{k-1}$.

Proof. Since $R+1 \in s_{t}^{2} \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)$ is supported away from $B_{\mathrm{II}}$, Lemma 7 shows that $g_{1}=-(\Delta+2)^{-1}(R+1)$ is of the form (5.1). We look for the formal power series solution of the non-linear problem as

$$
\begin{equation*}
f \simeq \sum_{k \geq 1} g_{k} \tag{5.2}
\end{equation*}
$$

Inserting this sum into the equation gives

$$
\begin{equation*}
-(\Delta+2)\left(\sum_{i \geq 1} g_{i}\right)=\sum_{j \geq 2} \frac{2^{j}}{j!}\left(g_{1}+\sum_{k \geq 2} g_{k}\right)^{j}+1+R . \tag{5.3}
\end{equation*}
$$

For each $i \geq 2$ we fix $g_{i}$ by

$$
\begin{array}{r}
-(\Delta+2) g_{i}=\sum_{j \geq 1} \frac{2^{j}}{j!}\left(g_{1}+\sum_{i-1 \geq k \geq 2} g_{k}\right)^{j}-\frac{2^{j}}{j!}\left(g_{1}+\sum_{i-2 \geq k \geq 2} g_{k}\right)^{j}  \tag{5.4}\\
=g_{i-1} P_{i}\left(g_{1}, g_{2}, \ldots g_{i-1}\right)
\end{array}
$$

where $P_{i}$ is a formal power series in $g_{1}, \ldots g_{i-1}$ without constant term.
Proceeding by induction we claim that

$$
\begin{equation*}
g_{i} \simeq \sum_{j \geq 2 i} s_{t}^{j} g_{i, j}, g_{i, j} \in \rho_{\mathrm{II}} \mathcal{P}^{j-2 i}+\rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{j-2 i+1, j-2 i} \tag{5.5}
\end{equation*}
$$

We have already seen that this holds for $i=1$ and using the obvious multiplicitivity properties

$$
\mathcal{P}^{k} \cdot \mathcal{P}^{j} \subset \mathcal{P}^{j+k}, \mathcal{P}^{k} \cdot \mathcal{P}_{\mathrm{II}}^{j, j-1} \subset \mathcal{P}_{\mathrm{II}}^{j+k, j+k-1}
$$

it follows from the inductive assumption, that (5.5) holds for all smaller indices, that

$$
\begin{align*}
& g_{i-1} P_{i}\left(g_{1}, g_{2}, \ldots g_{i-1}\right)  \tag{5.6}\\
& \simeq s_{t}^{2 i} \sum_{k \geq 2, j \geq 2 i-2}\left(\rho_{\mathrm{II}} \mathcal{P}^{j-2 i+2}+\rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{j-2 i+3, j-2 i+2}\right) \\
& \quad\left(\rho_{\mathrm{II}} \mathcal{P}^{k-2}+\rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{k-1, k-2}\right) \\
& \simeq \sum_{k \geq 2 i} s_{k}^{j} F_{k}, F_{k} \in \rho_{\mathrm{II}} \mathcal{P}^{k-2 i}+\rho_{\mathrm{I}} \rho_{\mathrm{II}} \mathcal{P}_{\mathrm{II}}^{k-2 i+1, k-2 i}
\end{align*}
$$

Applying Proposition 7 we recover the inductive hypothesis at the next step. Then (5.1) follows from (5.2) and (5.5).
q.e.d.

Summing the formal power series solution gives a log-smooth function with

$$
\begin{equation*}
-\Delta f_{0}=R+e^{2 f_{0}}+g, g \in O\left(s_{t}^{\infty}\right) \tag{5.7}
\end{equation*}
$$

Now we look for the solution as a perturbation $f=f_{0}+\tilde{f}$, so $\tilde{f}$ satisfies

$$
\begin{equation*}
-\Delta \tilde{f}=-g+e^{2 f_{0}}\left(e^{2 \tilde{f}}-1\right) \tag{5.8}
\end{equation*}
$$

which can be rewritten as

$$
\tilde{f}=-(\Delta+2)^{-1}\left(2 \tilde{f}\left(e^{2 f_{0}}-1\right)+e^{2 f_{0}}\left(e^{2 \tilde{f}}-1-2 \tilde{f}\right)-g\right)
$$

So consider the nonlinear operator

$$
\begin{equation*}
K: \tilde{f} \mapsto(\Delta+2)^{-1}\left(2 \tilde{f}\left(e^{2 f_{0}}-1\right)+e^{2 f_{0}}\left(e^{2 \tilde{f}}-1-2 \tilde{f}\right)-g\right) \tag{5.9}
\end{equation*}
$$

which acts on $s_{t}^{N} H_{b}^{M}\left(M_{\mathrm{mr}}\right)$ for all $N \geq 1$ and $M>2$. Note that for $M>2$, the b-space $H_{b}^{M}\left(M_{\mathrm{mr}}\right)$ is closed under multiplication, therefore,
this weighted Sobolev space is also an algebra. Since the nonlinear terms are at least quadratic, $K$ is well-defined on this domain. The solution to (5.8) satisfies $\tilde{f}=K(\tilde{f})$.

Proposition 8. For any $M>1$ and $N \geq 1$ there is a unique solution $\tilde{f} \in s_{t}^{N} H_{b}^{M}\left(M_{\mathrm{mr}}\right)$ to equation (5.8).

Proof. We construct the solution $\tilde{f}$ by iteration. Let

$$
\tilde{f}=s_{t}^{N} \sum_{i \geq 2} s_{t}{ }^{i} f_{i}
$$

put it into equation (5.8), divide by the common factor $s_{t}^{N}$ on both sides and then we get

$$
\begin{align*}
& \sum_{i \geq 2} s_{t}^{i} f_{i}=K\left(\sum s_{t}^{i} f_{i}\right)  \tag{5.10}\\
& \quad=(\Delta+2)^{-1}\left(\left(e^{2 f_{0}}-1\right) \sum s_{t}^{i} f_{i}+s_{t}^{N}\left(\sum s_{t}^{i} f_{i}\right)^{2}+s_{t}^{-N} g\right)
\end{align*}
$$

The right hand side belongs to $(\Delta+2)^{-1}\left(O\left(s_{t}^{2}\right)\right)$ because of the quadratic structure and the fact that $e^{2 f_{0}}-1 \in O\left(s_{t}^{2}\right)$. Therefore, the right hand side is the form $(\Delta+2)^{-1}\left(s_{t} h\right)$ where $s_{t} h \in \rho_{\mathrm{II}}^{-\frac{1}{2}} H_{\mathrm{b}}^{M}\left(M_{\mathrm{mr}}\right)$ so this quantity is well-defined using Proposition 3.

Now we proceed by induction. Assume that the first k terms in the expansion have been solved, then the equation for the next term $f_{k}$ is given by

$$
f_{k}=(\Delta+2)^{-1}\left(\left(e^{2 f_{0}}-1\right) f_{k-2}+s_{t}^{N} Q\left(f_{0}, \ldots f_{k-1}\right)\right)
$$

where the polynomial $Q$ on the right hand side is a quadratic polynomial of order $k-N$. By using the invertibility property in Proposition 3, we can now solve $f_{k}$. Therefore, the induction gives us the total expansion for $\tilde{f}$. q.e.d.

Proof of Theorem. From Proposition 8 we obtain the solution, $f=$ $f_{0}+\tilde{f}$, to the curvature equation $R\left(e^{2 f} h\right)=-1$. Since $f_{0}$ is the formal power series and $\tilde{f} \in s_{t}^{\infty} \mathcal{C}^{\infty}\left(M_{\mathrm{mr}}\right)$, we get the solution with required regularity. q.e.d.

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> Department of Mathematics
> Massachusetts Institute of Technology Cambridge, MA 02139
> USA
> E-mail address: rbm@math.mit.edu

Department of Mathematics
Stanford University
Stanford, CA 94305
USA
E-mail address: xuwenzhu@stanford.edu

