

ON STATIC THREE-MANIFOLDS WITH POSITIVE SCALAR CURVATURE

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Abstract

We compute a Bochner type formula for static three-manifolds and deduce some applications in the case of positive scalar curvature. We also explain in details the known general construction of the (Riemannian) Einstein $(n + 1)$ -manifold associated to a maximal domain of a static n -manifold where the static potential is positive. There are examples where this construction inevitably produces an Einstein metric with conical singularities along a codimension-two submanifold. By proving versions of classical results for Einstein four-manifolds for the singular spaces thus obtained, we deduce some classification results for compact static three-manifolds with positive scalar curvature.

1. Introduction and statements of the main results

Let (M^n, g) be a complete Riemannian manifold. A *static potential* is a non-trivial solution $V \in C^\infty(M)$ to the equation

$$(1) \quad \text{Hess}_g V - \Delta_g V g - V \text{Ric}_g = 0.$$

Riemannian manifolds admitting static potentials are very interesting geometric objects that arise in different contexts. For instance, the left hand side of (1) defines the formal adjoint of the linearization of the scalar curvature and the existence of static potentials plays an important role in problems related to prescribing the scalar curvature function (see [23], [14] and also [21]). On the other hand, given a solution to (1) on a three-manifold it is possible to construct a space-time satisfying the vacuum Einstein equations (with cosmological constant), whose properties, physically interpreted, justify the name static (see, for example, [59]).

The existence of a static potential imposes many restrictions on the geometry of the underlying manifold. For example, its scalar curvature must be constant and the set of zeroes of V is a totally geodesic regular hypersurface [14]. It might seem reasonable to believe that a description of all Riemannian manifolds admitting a static potential is possible. In dimension $n = 3$, a solution to this classification

problem seems to be more likely (because the Ricci tensor determines completely the Riemann curvature tensor) and has also physical interest (since $n = 3$ is the relevant dimension in the General Theory of Relativity).

Definition. A *static triple* is a triple (M^n, g, V) consisting of a connected n -dimensional smooth manifold M with boundary ∂M (possibly empty), a complete Riemannian metric g on M and a static potential $V \in C^\infty(M)$ that is *non-negative and vanishes precisely on ∂M* . Two static triples (M_i, g_i, V_i) , $i = 1, 2$, are said to be *equivalent* when there exists a diffeomorphism $\phi : M_1 \rightarrow M_2$ such that $\phi^*g_2 = cg_1$ for some constant $c > 0$ and $V_2 \circ \phi = \lambda V_1$ for some constant $\lambda > 0$.

In other words, a static triple is a maximal connected domain where a static potential is positive, and two static triples are equivalent when they are isometric (up to scaling) with proportional static potentials. The classification problem is to describe all equivalence classes of static triples.

The literature on the subject is extensive and many classification results are known particularly in the zero and negative scalar curvature cases (for example, see [35], [49], [54], [18], [4], [47], [25] and [12], [60], [52], [5], [26], [20], [41], [27]).

In this work we focus on three-dimensional static triples with positive scalar curvature.

Fischer and Marsden [23] conjectured that the standard unit round spheres (S^n, g_{can}) are the only closed Riemannian manifolds with scalar curvature $n(n-1)$ admitting static potentials. In fact, the linear combinations of the coordinate functions are all solutions to (1) in (S^n, g_{can}) . This conjecture was soon realized to be too optimistic. Kobayashi [36] and Lafontaine [39] proved more generally a classification result for locally conformally flat static triples in all dimensions. We state their result in dimension $n = 3$ as follows:

Theorem 1 (Kobayashi [36], Lafontaine [39]). *Let (M^3, g, V) be a static triple with positive scalar curvature. If (M^3, g) is locally conformally flat, then (M^3, g, V) is covered by a static triple that is equivalent to one of the following static triples:*

i) *The standard round hemisphere,*

$$(S_+^3, g_{can}, V = x_{n+1}).$$

ii) *The standard cylinder over S^2 with the product metric,*

$$\left(\left[0, \frac{\pi}{\sqrt{3}} \right] \times S^2, g_{prod} = dt^2 + \frac{1}{3}g_{can}, V = \frac{1}{\sqrt{3}} \sin(\sqrt{3}t) \right).$$

iii) For some $m \in (0, \frac{1}{3\sqrt{3}})$, the triple

$$\left([r_h, r_c] \times S^2, g_m = \frac{dr^2}{1 - r^2 - \frac{2m}{r}} + r^2 g_{can}, V_m = \sqrt{1 - r^2 - \frac{2m}{r}} \right),$$

where $r_h = r_h(m) < r_c = r_c(m)$ are the positive zeroes of V_m .

REMARK 1. A static triple $(\tilde{M}^n, \tilde{g}, \tilde{V})$ covers a static triple (M^n, g, V) when there exists a covering map $\pi : \tilde{M} \rightarrow M$ such that $\tilde{g} = \pi^*g$ and $\tilde{V} = V \circ \pi$.

REMARK 2. We have normalized the scalar curvature of the examples above to be 6. They are sometimes referred as (time-symmetric slices of) the *de Sitter space*, the *Nariai space* and the *Schwarzschild–de Sitter spaces of positive mass m* , respectively (see [12]).

REMARK 3. The three types of locally conformally flat manifolds described above can be distinguished from the behavior of their Ricci tensor. The round hemisphere is an Einstein manifold, the standard cylinder has parallel Ricci tensor but is not Einstein, and the Schwarzschild–de Sitter spaces of positive mass are locally conformally flat but do not have parallel Ricci tensor.

In dimension $n = 3$, the above list contains all known examples of compact simply connected static triples with positive scalar curvature. Qing and Yuan [53] proved uniqueness of these examples assuming a weaker hypothesis on the Cotton tensor. Lafontaine and Rozoy [40] gave an argument to show the uniqueness of these examples under an additional hypothesis on the regularity of the critical level sets of the static potential. A speculation on the possibility of constructing examples with more than two boundary components was presented in [12].

It is interesting to observe that precisely two more examples can be obtained as quotients of the standard cylinder, one of them *non-orientable* with *two* boundary components, the other one *orientable* with *connected* boundary (see their description in the end of Section 7). One should think about this last example in the context of the “cosmic no-hair conjecture” suggested in [12] and correctly stated in [11]. In any case, the question remains whether the standard round hemisphere is the only compact *simply connected* static triple (M^3, g, V) with positive scalar curvature and *connected* boundary.

An important result towards the answer to this question states that the standard hemisphere has the maximum possible boundary area among static triples with positive scalar curvature and connected boundary. More precisely, the following theorem holds:

Theorem 2 (Boucher–Gibbons–Horowitz [12], Shen [57]). *Let (M^3, g, V) be a compact oriented static triple with connected boundary and*

scalar curvature 6. Then ∂M is a two-sphere whose area satisfies the inequality

$$|\partial M| \leq 4\pi.$$

Moreover, equality holds if and only if (M^3, g, V) is equivalent to the standard hemisphere.

This result has also a version allowing more boundary components (see Section 2, Proposition 6). Hijazi, Montiel and Raulot [34] established a similar result that involves an inequality for the first eigenvalue of the induced Dirac operator of each boundary component.

Theorem 2 has some interesting interpretations. The inequality for the area of the boundary can be thought as a Penrose type inequality in the context of positive scalar curvature (as suggested in [12]). The rigidity phenomenon associated to the equality case is related to Min-Oo's Conjecture, which is known to be false after the work of Brendle, Marques and Neves [17]. In view of that, Theorem 2 gives an interesting rigidity result for the standard round hemisphere under a nice extra geometric assumption (for related results, see for example [33], [22] and [44]).

In this paper we prove some classification results for static three-manifolds with positive scalar curvature. In particular, we obtain a stronger version of Theorem 2 that includes the oriented static quotient of the standard cylinder in the picture. Our results may also be used to rule out some possible new examples of compact static triples with positive scalar curvature.

The main ingredient is a Bochner type formula that holds for any three-dimensional static triple. Let R denote the scalar curvature, $\mathring{Ric} = Ric - (R/3)g$ denote the trace-free part of the Ricci tensor and C denote the Cotton tensor of (M^3, g) . C is the 3-tensor defined by

$$\begin{aligned} C(X, Y, Z) &= (\nabla_Z Ric)(X, Y) - (\nabla_Y Ric)(X, Z) \\ &\quad - \frac{1}{4}(dR(Z)g(X, Y) - dR(Y)g(X, Z)), \end{aligned}$$

for all $X, Y, Z \in \mathcal{X}(M)$. A classical result asserts that (M^3, g) is locally conformally flat if and only if its Cotton tensor vanishes (see [9], Chapter 4, Section 3.2).

If (M^3, g) admits a static potential V , we prove the following identity:

$$(2) \quad \frac{1}{2}div(V\nabla|\mathring{Ric}|^2) = \left(|\nabla\mathring{Ric}|^2 + \frac{|C|^2}{2} \right) V + \left(R|\mathring{Ric}|^2 + 18det(\mathring{Ric}) \right) V.$$

Assuming a pointwise inequality for the traceless Ricci tensor we deduce the first consequence of formula (2).

Theorem A. *Let (M^3, g, V) be a compact oriented static triple with positive scalar curvature. If*

$$|\mathring{Ric}|^2 \leq \frac{R^2}{6},$$

then one of the following alternatives holds:

- i) $Ric = 0$ and (M^3, g, V) is equivalent to the standard hemisphere;*
- or*
- ii) $|\mathring{Ric}|^2 = R^2/6$ and (M^3, g, V) is covered by a static triple that is equivalent to the standard cylinder.*

REMARK 4. One can verify that the above pinching condition on the traceless Ricci tensor implies that (M^3, g) has non-negative Ricci curvature. As we will see later, the above result can also be deduced as a corollary of more general results (see Theorem E below).

For the next applications we work with the *singular* Einstein manifold that can be constructed from a static triple. The general construction has already been discussed in the literature (for instance, in [28] and [12]). When not singular, the obtained Riemannian manifolds are sometimes called *gravitational instantons*, although this terminology might be confusing.

As mentioned before, as a general fact, given a static triple (M^n, g, V) with scalar curvature $R = \epsilon n(n-1)$, $\epsilon \in \{-1, 0, 1\}$, the metrics

$$h_{\pm} = \pm V^2 dt^2 + g$$

satisfy the Einstein equation $Ric_{h_{\pm}} = \epsilon n h_{\pm}$. While the Lorentzian metric h_- has no singularities on $\mathbb{R} \times \partial M$ (as the vanishing of V only means that ∂_t is a light-like vector on this set), the Riemannian metric h_+ becomes singular on $\mathbb{R} \times \partial M$. One way to overcome this problem could be to identify the variable t with period 2π (as done in [12] and [28]). When this procedure removes the metric singularity, one obtains precisely the type of Einstein manifolds studied by Seshadri [56]. However, although metric singularities might be inevitable, the singular space obtained is well-behaved: it has a cone-like singularity along a codimension-two submanifold.

This type of singular space, sometimes called an *edge space*, has been long studied in the literature and its theory presents many interesting geometric and analytic phenomena. For instance, there has been very recent progress in understanding the Yamabe problem in these spaces, see the work of Akutagawa, Carron and Mazzeo [1] and Mondello [48]. On the other hand, classical formulas like the Gauss–Bonnet–Chern formula for closed 4-manifolds have generalizations to these spaces, see the work of Liu and Shen [43] (and also the work of Atiyah and LeBrun [8]). For a comprehensive exposition on singular Riemannian spaces of stratified type, see [51].

In the second part of this paper, after describing carefully the above construction and the properties of the singular metric h_+ (see Section 6), we argue that part of the classical Bonnet–Myers Theorem and a deep theorem of Gursky [31] on Einstein four-manifolds with positive scalar curvature hold for the singular Einstein manifolds obtained from compact static triples with positive scalar curvature. As a consequence, we prove two results.

The first one gives a description of the topology of compact static triples (M^3, g, V) with positive scalar curvature. It is inspired by the work of Galloway [24] on the topology of black holes and generalizes results of [56]. Before stating it, we recall that a closed minimal surface Σ^2 in (M^3, g) is called *stable* when the second variation of the area is non-negative for all variations, and *unstable* otherwise.

Theorem B. *Let (M^3, g, V) be a compact oriented static triple with positive scalar curvature. Then*

- i) The universal cover of M is compact.*
- ii) If ∂M contains unstable connected components, then ∂M contains exactly one unstable connected component. In this case, M is simply connected.*
- iii) Each connected component of ∂M is diffeomorphic to a sphere.*

REMARK 5. Regarding the classification problem of compact static triples with positive scalar curvature, item *i)* above allows to assume simply connectedness without loss of generality.

The second result can be seen as a gap result for compact static triples (M^3, g, V) with positive scalar curvature.

Theorem C. *Let (M^3, g, V) be a compact simply connected static triple with scalar curvature 6. Then, one of the following alternatives holds:*

- i) (M^3, g, V) is equivalent to the standard hemisphere; or*
- ii) (M^3, g, V) is equivalent to the standard cylinder; or*
- iii) Denoting by $\partial_i M$, $i = 1, \dots, r$, the connected components of ∂M and by k_i the (constant) value of $|\nabla V|$ on $\partial_i M$, the following inequality holds:*

$$\sum_{i=1}^r k_i |\partial_i M| < \frac{4\pi}{3} \sum_{i=1}^r k_i.$$

REMARK 6. This theorem should be compared with Theorem B in [31]. In fact, when the associated Einstein manifold to a static triple satisfying the hypotheses of Theorem C has no singularities the result is a direct corollary of that theorem (see Section 6 for more details).

REMARK 7. The inequality in *iii*) is invariant under rescaling of the static potential. Theorem C is sharp in the sense that one can explicitly verify that:

- i*) For the standard cylinder, $k_1 = k_2$ and $|\partial_1 M| = |\partial_2 M| = 4\pi/3$.
- ii*) For the Schwarzschild–de Sitter spaces of mass $m \in (0, 1/3\sqrt{3})$,

$$\frac{k_1|\partial_1 M| + k_2|\partial_2 M|}{k_1 + k_2}$$

is an increasing function of m that converges to 0 as $m \rightarrow 0$ and to $4\pi/3$ as $m \rightarrow 1/3\sqrt{3}$. Moreover, $|\partial_1 M| < 4\pi/3 < |\partial_2 M|$ and $k_1 > k_2$ for all $m \in (0, 1/3\sqrt{3})$.

As an immediate application, we can state

Theorem D. *Let (M^3, g, V) be a compact simply connected static triple with connected boundary and scalar curvature 6. If*

$$|\partial M| \geq \frac{4\pi}{3},$$

then (M^3, g, V) is equivalent to the standard hemisphere.

Another corollary of the previous results is the following:

Theorem E. *Let (M^3, g, V) be a compact static triple with scalar curvature 6 and non-negative Ricci curvature. Then (M^3, g, V) is equivalent to the standard hemisphere or is covered by the standard cylinder.*

In fact, in this case, one can see using the Gauss equation that each component of ∂M is a totally geodesic two-sphere with Gaussian curvature $K \leq 3$, hence each component has area greater than or equal to $4\pi/3$.

The paper is organized as follows. After reviewing some basic material on static triples (Section 2), we deduce some formulas for the Cotton tensor of such triples (Section 3) and prove the Bochner type formula (2) (Section 4). Theorem A is proven in Section 5. In Section 6 we describe the associated singular Einstein manifold and briefly discuss some of its geometric and topological properties. In Section 7 we prove the topological classification Theorem B. The remaining two sections are devoted to the proof of Theorem C. The proof mimics the arguments of Gursky in [31], which involve obtaining the solution to a Yamabe type problem and the Gauss–Bonnet–Chern formula for closed oriented four-manifolds. Technical work is needed to justify why these arguments also work for the singular Einstein four-manifold associated to a compact three-dimensional static triple. The paper ends after three appendixes. The first one contains the proof of a key inequality used in Section 8. Then we classify certain solutions to an equation of type $HessV = (\Delta V/2)g$ on a compact surface, a result needed in Section 7.

The last appendix discusses the regularity properties of the solution to a degenerate elliptic problem that appeared in the proof of Proposition 21.

Acknowledgments. I am grateful to Fernando Codá Marques, Harold Rosenberg and especially André Neves for their encouragement and support, for interesting conversations and for their kind interest in this work. I would also like to thank Rafe Mazzeo for correspondence regarding the regularity properties of solutions to the singular Yamabe problem. I was supported by CNPq-Brazil.

2. General properties of static triples

This section aims to present some of the basic properties of static manifolds that will be used frequently throughout the paper (see also [21]).

Let (M^n, g) be a complete Riemannian manifold. A static potential is a non-trivial solution $V \in C^\infty(M)$ to the second order overdetermined elliptic equation

$$(3) \quad \text{Hess}_g V - \Delta_g V g - V \text{Ric}_g = 0.$$

Equation (3) is equivalent to a useful system of equations. Moreover, the set of zeroes of V has a very special geometry.

Lemma 3 ([23], [14]). *The static equation (3) is equivalent to the equations*

$$(4) \quad \text{Hess}_g V = V(\text{Ric}_g - \Lambda g),$$

$$(5) \quad \Delta_g V + \Lambda V = 0,$$

where the scalar curvature $R_g = (n-1)\Lambda$ is constant. Moreover, if (M^n, g) admits a static potential $V \in C^\infty(M)$, then

- i) 0 is a regular value of V ;
- ii) $\{V = 0\}$ is totally geodesic; and
- iii) $|\nabla V|$ is locally constant and positive on $\{V = 0\}$.

Proof. (see [21]) Taking the trace of (3), we obtain

$$(6) \quad \Delta_g V + \frac{R_g}{n-1} V = 0.$$

The equivalence between (3) and the system (4), (5) follows immediately. To prove that R_g is constant, we observe that the divergence of (3) gives

$$(7) \quad V dR_g = 0,$$

as can be seen by the Ricci formula for commuting derivatives and the contracted second Bianchi identity $2\text{div}_g \text{Ric}_g = dR_g$. Considering the homogeneous ODE satisfied by the restriction of V to geodesics (which follows from (4)), we conclude that if $\nabla V(p) = 0$ at $p \in \{V = 0\}$, then V

must vanishes identically near p . By analytical continuation of solutions to the elliptic equation (5), V must vanish identically, a contradiction. Thus, 0 is a regular value of V and hence, by (7), R must be constant. The remaining assertions follow because the second fundamental form of the level set $\{V = 0\}$ and the derivative of $|\nabla V|^2$ on $\{V = 0\}$ depends on $HessV$, which vanishes on this set by (4). q.e.d.

The above proposition implies that static triples (M^n, g, V) as defined in the Introduction have constant scalar curvature and totally geodesic boundary. When dealing with the classification problem, one can, therefore, assume (M^n, g, V) has scalar curvature $R = \epsilon n(n - 1)$ for some $\epsilon \in \{-1, 0, 1\}$. The constant value of $|\nabla V|$ on a connected component of ∂M is sometimes called the *surface gravity* of the component (as in [28]).

As a consequence of the Maximum Principle, the following result holds for compact static triples:

Proposition 4 ([23], [14]). *Let (M^n, g, V) be a static triple with scalar curvature $R = \epsilon n(n - 1)$, $\epsilon \in \{-1, 0, 1\}$.*

- i) If M is closed, then V is a positive constant and (M, g) is Ricci-flat.*
- ii) If M is compact and ∂M is non-empty, then $R > 0$ and n is the first Dirichlet eigenvalue of the Laplacian of (M, g) .*

Proof. If M is closed, Hopf's maximum principle implies that R is zero and V is constant. Thus, the static equation (3) becomes $Ric = 0$. If M is compact and $\partial M \neq \emptyset$, $R/(n-1)$ is an eigenvalue of the Laplacian for the Dirichlet problem in (M, g) , by equation (6). Hence, it is positive. Since V does not change its sign, $R/(n-1)$ is the first eigenvalue. q.e.d.

Regarding the classification problem of static triples, item *ii*) of Proposition 4 has a useful consequence. In order to verify that two compact static triples with the same scalar curvature are equivalent it is enough to verify that they are isometric because, as first eigenfunctions, the static potentials must be proportional to each other in that case.

Given a static triple (M^n, g, V) with scalar curvature $R = \epsilon n(n - 1)$, the function $|\nabla V|^2 + \epsilon V^2$ plays an important role. If (M^n, g) is Einstein, in which case the static equation reduces to Obata's equation $HessV = -\epsilon Vg$ [50], it is possible to verify that this function is constant. The following proposition shows the converse. More generally, we have

Proposition 5. *Let (M^n, g, V) be a static triple with scalar curvature $R = \epsilon n(n - 1)$, $\epsilon \in \{-1, 0, 1\}$. Then the function*

$$|\nabla V|^2 + \epsilon V^2$$

does not achieve a non-negative maximum on $M \setminus \partial M$ unless it is constant and (M^n, g) is Einstein.

Proof. Using the Bochner formula, one proves Shen's identity [57],

$$(8) \quad \operatorname{div} \left(\frac{1}{V} \nabla (|\nabla V|^2 + \epsilon V^2) \right) = 2V |\mathring{Ric}|^2,$$

and applies the strong maximum principle. q.e.d.

Integrating (8) on a compact three-dimensional static triple, one obtains a fundamental formula that can be used to prove Theorem 2.

Proposition 6 ([57]). *Let (M^3, g, V) be a compact oriented static triple with scalar curvature $R = 6$. Denoting by $\partial_1 M, \dots, \partial_r M$ the connected components of ∂M and by k_i the value of $|\nabla V|$ on $\partial_i M$, the following formula holds:*

$$(9) \quad \sum_{i=1}^r k_i |\partial_i M| + \int_M |\mathring{Ric}|^2 V d\mu = 2\pi \sum_{i=1}^r k_i \chi(\partial_i M),$$

where $\chi(\partial_i M)$ denotes the Euler characteristic of $\partial_i M$.

Proof. By Shen's identity (8), the vector field

$$X = \frac{1}{V} \nabla (|\nabla V|^2 + V^2)$$

is such that $\operatorname{div} X = 2V |\mathring{Ric}|^2$ on M . Moreover, as can be seen using the static equation (4), $X = 2(\operatorname{Ric}(\nabla V, -) - 2\nabla V)$. The Gauss equation for the totally geodesic boundary ∂M gives

$$K = \frac{R}{2} - \operatorname{Ric} \left(\frac{\nabla V}{|\nabla V|}, \frac{\nabla V}{|\nabla V|} \right) = 3 - \operatorname{Ric} \left(\frac{\nabla V}{|\nabla V|}, \frac{\nabla V}{|\nabla V|} \right),$$

where K denotes the Gaussian curvature of ∂M . Observing that the outward-pointing normal to ∂M is $-\nabla V/|\nabla V|$, we integrate (8) to obtain

$$\begin{aligned} \int_M V |\mathring{Ric}|^2 d\mu &= \frac{1}{2} \int_M \operatorname{div} X d\mu = \frac{1}{2} \int_{\partial M} g \left(X, -\frac{\nabla V}{|\nabla V|} \right) d\sigma = \\ &= \int_{\partial M} \left(2 - \operatorname{Ric} \left(\frac{\nabla V}{|\nabla V|}, \frac{\nabla V}{|\nabla V|} \right) \right) |\nabla V| d\sigma = \int_{\partial M} (K - 1) |\nabla V| d\sigma. \end{aligned}$$

Formula (9) follows now by the Gauss–Bonnet theorem. q.e.d.

In order to understand the behavior of the static potential near ∂M , we compute the expansions of V and $|\nabla V|$ along a normalized geodesic issuing from ∂M orthogonally.

Proposition 7. *Let (M^n, g, V) be a static triple with scalar curvature $R = \epsilon n(n-1)$, $\epsilon \in \{-1, 0, 1\}$. Given $p \in \partial M$, let $\gamma : [0, \epsilon) \rightarrow M$ be the normalized geodesic such that $\gamma(0) = p$ and $\gamma'(0)$ is the unit normal to*

∂M pointing inside M . Then the following expansions hold as s goes to zero:

$$\begin{aligned} V^2 \circ \gamma(s) &= |\nabla V|^2(p)(s^2 + \frac{1}{3}(Ric(\gamma'(0), \gamma'(0)) - \epsilon n)s^4 + O(s^6)), \\ |\nabla V|^2 \circ \gamma(s) &= |\nabla V|^2(p)(1 + (Ric(\gamma'(0), \gamma'(0)) - \epsilon n)s^2 + O(s^4)). \end{aligned}$$

Proof. Without loss of generality, we can assume that $|\nabla V|(p) = 1$, so that $\gamma'(0) = \nabla V(p)$. We need the formulas obtained by taking successive derivatives of the static equation (4):

$$\begin{aligned} HessV &= V(Ric - \epsilon ng), \\ \nabla HessV &= (Ric - \epsilon ng) \oplus dV + V\nabla Ric, \\ \nabla^2 HessV &= V(Ric - \epsilon ng) \oplus (Ric - \epsilon ng) + V\nabla^2 Ric + T, \end{aligned}$$

where T depends on ∇Ric and dV .

For $p \in \partial M$, we claim that $\nabla Ric(\nabla V, \nabla V, \nabla V)(p) = 0$. In fact, given an orthonormal basis $\{e_1, \dots, e_{n-1}, \nabla V(p)\}$ of $T_p \partial M$, since (M, g) has constant scalar curvature, the contracted second Bianchi inequality implies

$$0 = \frac{dR}{2} = div Ric(\nabla V) = \sum_{i=1}^{n-1} (\nabla_{e_i} Ric)(\nabla V, e_i) + (\nabla_{\nabla V} Ric)(\nabla V, \nabla V).$$

On the other hand, since ∂M is totally geodesic, Codazzi equation implies $Ric(\nabla V, e_i) = 0$ for all $i = 1, \dots, n-1$, and, therefore, $\sum (\nabla_{e_i} Ric)(\nabla V, e_i) = 0$. This proves the claim.

Hence, at the point $p \in \partial M$,

$$\begin{aligned} HessV(\nabla V, \nabla V)(p) &= 0, \\ (\nabla HessV)(\nabla V, \nabla V, \nabla V)(p) &= Ric(\nabla V, \nabla V)(p) - \epsilon n, \\ (\nabla^2 HessV)(\nabla V, \nabla V, \nabla V, \nabla V)(p) &= 0. \end{aligned}$$

The result follows now by a direct computation using the above formulas. q.e.d.

As a consequence of Propositions 5 and 7, we obtain some control on the behavior of the curvature of the boundary component where the function $|\nabla V|^2 + \epsilon V^2$ attains a maximum.

Proposition 8. *Let (M^n, g, V) be a static triple with scalar curvature $R^M = \epsilon n(n-1)$, $\epsilon \in \{-1, 0, 1\}$. If $|\nabla V|^2 + \epsilon V^2$ attains a maximum at a boundary component $\partial_1 M$, then $\partial_1 M$ has scalar curvature greater than or equal to $\epsilon(n-1)(n-2)$.*

Proof. Since $|\nabla V|$ is constant on each connected component of $\partial M = V^{-1}(0)$, the maximum of $|\nabla V|^2 + \epsilon V^2$ is attained at every point of $\partial_1 M$. Let γ be a geodesic issuing from $p \in \partial_1 M$ as in Proposition 7. Then $|\nabla V|^2(p) \geq (|\nabla V|^2 + \epsilon V^2) \circ \gamma(s) = |\nabla V|^2(p)(1 + (Ric(\gamma'(0), \gamma'(0)) -$

$\epsilon(n-1)s^2 + O(s^4)$) implies that $Ric(\gamma'(0), \gamma'(0)) \leq \epsilon(n-1)$. Since $\partial_1 M$ is totally geodesic, the result follows from the Gauss equation $R^{\partial_1 M} = R^M - 2Ric(\gamma'(0), \gamma'(0))$. q.e.d.

We finish this section with a result about the Jacobi operator of a minimal hypersurface Σ^{n-1} in (M^n, g, V) . We denote by N a unit vector field normal to Σ and by A the second fundamental form of Σ given by $A = g(\nabla_X N, Y)$ for all X, Y tangent to Σ .

Proposition 9. *Let (M^n, g, V) be a static triple. Let Σ^{n-1} be an immersed compact hypersurface in M . Then the Jacobi operator of Σ is such that*

$$\mathcal{L}_\Sigma V := \Delta_\Sigma V + (Ric(N, N) + |A|^2)V = g(\nabla V, \vec{H}) + |A|^2 V.$$

Proof. Since $\Delta_M f = \Delta_\Sigma f - g(\nabla f, \vec{H}) + Hess f(N, N)$ for every function $f \in C^\infty(M)$, the formula for $\mathcal{L}_\Sigma V$ follows from the static equation (3). q.e.d.

3. The Cotton tensor of a three-dimensional static triple

We deduce some identities satisfied by the Cotton tensor of three-dimensional static triples. They will be used in the proof of the Bochner type formula (2) in the next section. It may be interesting to compare the computations carried out in this section and in the following one with the more general computations related to conformal flatness of [53].

Let (M^3, g) be a Riemannian three-manifold. The Riemann curvature tensor is completely determined by the Ricci tensor. In fact,

$$(10) \quad Rm = \left(Ric - \frac{R}{4}g \right) \odot g = \left(\mathring{Ric} + \frac{R}{12}g \right) \odot g,$$

where $A \odot B$ denotes the Kulkarni–Nomizu product of symmetric two-tensors given in coordinates by $(A \odot B)_{ijkl} = A_{ik}B_{jl} - A_{il}B_{jk} - A_{jk}B_{il} + A_{jl}B_{ik}$ (see [10]).

The Cotton tensor of (M^3, g) is defined by

$$C(X, Y, Z) = (\nabla_Z Ric)(X, Y) - (\nabla_Y Ric)(X, Z) - \frac{1}{4}(dR(Z)g(X, Y) - dR(Y)g(X, Z)),$$

for all $X, Y, Z \in \mathcal{X}(M)$. If $V \in C^\infty(M)$ is a static potential on (M^3, g) , the Cotton tensor has a particularly simple expression.

Lemma 10. *Assume (M^3, g) admits a static potential $V \in C^\infty(M)$. Let*

$$Hess^\circ V = Hess V - \frac{\Delta V}{3}g$$

denote the traceless part of the Hessian of V and let T be the three-tensor defined by

$$T(X, Y, Z) = (\nabla_Z \mathring{Hess}V)(X, Y) - (\nabla_Y \mathring{Hess}V)(X, Z),$$

for all $X, Y, Z \in \mathcal{X}(M)$. Then

$$(11) \quad T(X, Y, Z) = -(\mathring{Ric}(X, Y)dV(Z) - \mathring{Ric}(X, Z)dV(Y)) \\ - (g(X, Y)\mathring{Ric}(Z, \nabla V) - g(X, Z)\mathring{Ric}(Y, \nabla V)),$$

and

$$(12) \quad C(X, Y, Z) = (\nabla_Z \mathring{Ric})(X, Y) - (\nabla_Y \mathring{Ric})(X, Z) \\ = \frac{1}{V} \left(T(X, Y, Z) - \mathring{Ric}(X, Y)dV(Z) + \mathring{Ric}(X, Z)dV(Y) \right).$$

Proof. Without loss of generality we assume that $R = 6\epsilon$ for $\epsilon \in \{-1, 0, 1\}$. The static equations (4), (5) can be rewritten in coordinates as

$$(\mathring{Hess}V)_{ij} = V(\mathring{Ric})_{ij}, \\ (\mathring{Hess}V)_{ij} = V_{;ij} = (\mathring{Hess}V)_{ij} - \epsilon V g_{ij}.$$

By the Ricci identity for commuting covariant derivatives and (10), we obtain formula (11):

$$T_{ijk} = (V_{;ijk} + \epsilon V_k g_{ij}) - (V_{;ikj} + \epsilon V_j g_{kj}) \\ = -R_{pikj}V_{;^p} + \epsilon(g_{ij}V_{;k} - g_{kj}V_{;j}) \\ = -((\mathring{Ric})_{ij}V_{;k} - (\mathring{Ric})_{ik}V_{;j}) - (g_{ij}(\mathring{Ric})_{kp}V_{;^p} - g_{ik}(\mathring{Ric})_{jp}V_{;^p}).$$

On the other hand, $C_{ijk} = (\mathring{Ric})_{ij;k} - (\mathring{Ric})_{ik;j}$, because the scalar curvature is constant. Formula (12) follows since $\mathring{Hess}V = V\mathring{Ric}$. q.e.d.

The norm of the Cotton tensor C can be computed in different ways, each way allowing to deduce interesting consequences. For instance, on a regular level set $\Sigma = V^{-1}(c)$, $c > 0$, $|C|$ depends on the norm of the traceless part of the second fundamental form of Σ and on the gradient of $|\nabla V|^2$ on Σ (see [42], Formula 14). That computation is useful to prove the classification results of [36] and [39]. In the next lemma we show other two formulas that will be used in our applications.

Lemma 11. *Assume (M^3, g) admits a static potential $V \in C^\infty(M)$. Then*

$$|C|^2 = \frac{1}{V^2} \left(8|\mathring{Ric}|^2|\nabla V|^2 - 12|\mathring{Ric}(\nabla V, -)|^2 \right) = -\frac{4}{V} C_{ijk}(\mathring{Ric})^{ij}V^k.$$

In particular, C vanishes at a point $p \in M$ such that $V(p) \neq 0$ only if $\nabla V(p) = 0$ or \mathring{Ric} has at most two different eigenvalues at p .

Proof. By Lemma 10, $VC_{ijk} = T_{ijk} + U_{ijk}$, where

$$\begin{aligned} T_{ijk} &= -((\mathring{Ric})_{ij}V_{;k} - (\mathring{Ric})_{ik}V_{;j}) \\ &\quad - (g_{ij}(\mathring{Ric})_{kp}V_{;p} - g_{ik}(\mathring{Ric})_{jp}V_{;p}), \\ U_{ijk} &= -((\mathring{Ric})_{ij}V_{;k} - (\mathring{Ric})_{ik}V_{;j}). \end{aligned}$$

Thus, since

$$(13) \quad |U|^2 = 2(\mathring{Ric})_{ij}(\mathring{Ric})^{ij}V_{;p}V_{;p} - 2(\mathring{Ric})_{jp}V_{;p}(\mathring{Ric})^{jq}V_{;q},$$

$$(14) \quad T_{ijk}U^{ijk} = 2(\mathring{Ric})_{ij}(\mathring{Ric})^{ij}V_{;p}V_{;p} - 4(\mathring{Ric})_{jp}V_{;p}(\mathring{Ric})^{jq}V_{;q},$$

$$(15) \quad |T|^2 = 2(\mathring{Ric})_{ij}(\mathring{Ric})^{ij}V_{;p}V_{;p} - 2(\mathring{Ric})_{jp}V_{;p}(\mathring{Ric})^{jq}V_{;q},$$

we obtain

$$V^2|C|^2 = |T|^2 + 2T_{ijk}U^{ijk} + |U|^2 = (8|\mathring{Ric}|^2|\nabla V|^2 - 12|\mathring{Ric}(\nabla V, -)|^2).$$

Moreover, since $C_{ijk} + C_{ikj} = 0$, from (13) and (14) we have

$$\begin{aligned} -VC_{ijk}(\mathring{Ric})^{ij}V^k &= \frac{V}{2}C_{ijk}U^{ijk} = \frac{1}{2}(T_{ijk} + U_{ijk})U^{ijk} \\ &= (2|\mathring{Ric}||\nabla V|^2 - 3|\mathring{Ric}(\nabla V, -)|^2) = \frac{V^2}{4}|C|^2. \end{aligned}$$

To finish the proof of the lemma, given $p \in M$ with $V(p) \neq 0$, we choose $\{e_1, e_2, e_3\} \subset T_pM$ an orthonormal basis diagonalizing \mathring{Ric} . Writing $(\mathring{Ric})_{ij}(p) = \lambda_i\delta_{ij}$ and $\nabla V(p) = \sum_{i=1}^3 a_i e_i$, we have

$$\begin{aligned} V^2|C|^2(p) &= 8|\mathring{Ric}|^2(p)|\nabla V|^2(p) - 12|(\mathring{Ric})(\nabla V, -)|^2(p) \\ &= 4a_1^2(2\lambda_2^2 + 2\lambda_3^2 - \lambda_1^2) + 4a_2^2(2\lambda_1^2 + 2\lambda_3^2 - \lambda_2^2) \\ &\quad + 4a_3^2(2\lambda_1^2 + 2\lambda_2^2 - \lambda_3^2). \end{aligned}$$

Since $tr(\mathring{Ric}) = \lambda_1 + \lambda_2 + \lambda_3 = 0$, we obtain

$$V^2|C|^2(p) = 4a_1^2(\lambda_1 + 2\lambda_2)^2 + 4a_2^2(\lambda_2 + 2\lambda_1)^2 + 4a_3^2(\lambda_1 - \lambda_2)^2.$$

The last statement of the lemma follows. q.e.d.

4. The Bochner type formula

Proposition 12. *Assume (M^3, g) admits a static potential $V \in C^\infty(M)$. Then*

$$(16) \quad \frac{1}{2}div(V\nabla|\mathring{Ric}|^2) = (|\mathring{Ric}|^2 + \frac{|C|^2}{2})V + (R|\mathring{Ric}|^2 + 18det(\mathring{Ric}))V.$$

Proof. Without loss of generality, we rescale the metric in such way that $R = 6\epsilon$ for some $\epsilon \in \{-1, 0, 1\}$. During the computation, we will use the static equations written as

$$(17) \quad (HessV)_{ij} = V(\mathring{Ric})_{ij},$$

$$(18) \quad V_{;ij} = (HessV)_{ij} - \epsilon Vg_{ij},$$

and the fact that \mathring{Ric} is divergence free, which is a consequence of the contracted second Bianchi identity $2\mathit{div}(\mathring{Ric}) = dR = 0$.

We have

(19)

$$\begin{aligned} \frac{1}{2}\mathit{div}(V|\mathring{Ric}|^2) &= \frac{1}{2}g(\nabla V, \nabla|\mathring{Ric}|^2) + V\Delta\frac{|\mathring{Ric}|^2}{2} \\ &= (\mathring{Ric})_{ij;p}(\mathring{Ric})^{ij}V_{;^p} + V|\nabla\mathring{Ric}|^2 + Vg(\Delta\mathring{Ric}, \mathring{Ric}). \end{aligned}$$

By the static equations (17) and (18),

$$(\Delta(HessV))_{ij} = (\Delta V)\mathring{Ric}_{ij} + V(\Delta\mathring{Ric})_{ij} + 2(\mathring{Ric})_{ij;p}V_{;^p},$$

hence

$$(20) \quad V(\Delta\mathring{Ric})_{ij} = (\Delta(HessV))_{ij} + 3\epsilon V(\mathring{Ric})_{ij} - 2(\mathring{Ric})_{ij;p}V_{;^p}.$$

Now we compute $\Delta(HessV)$. Following the notations of Lemma 10, let $T_{ijk} = (Hess)_{ij;k} - (Hess)_{ik;j}$. Taking two derivatives of $HessV$, we obtain

$$\begin{aligned} (HessV)_{ij;kl} &= (HessV)_{ik;jl} + T_{ijk;l} = (HessV)_{ki;jl} + T_{ijk;l} \\ &= (HessV)_{ki;l j} + R^p_{kjl}(HessV)_{pi} \\ &\quad + R^p_{ijl}(HessV)_{kp} + T_{ijk;l} \\ &= (HessV)_{kl;ij} + R^k_{plj}(HessV)^p_i - R_{pilj}(HessV)_k^p \\ &\quad + T_{kil;j} + T_{ijk;l}. \end{aligned}$$

Contracting the two last indexes and using (10), we have

(21)

$$\begin{aligned} (\Delta(HessV))_{ij} &= (HessV)^q_{q;ij} + R_{pj}(HessV)^p_i - R_{piqj}(HessV)^{qp} \\ &\quad + T^q_{iq;j} + T_{ijq;}^q \\ &= (\mathring{Ric})_{pj}(HessV)^p_i + 2\epsilon g_{pj}(HessV)^p_i \\ &\quad - \epsilon(g_{pq}g_{ij} - g_{pj}g_{iq})(HessV)^{qp} \\ &\quad - (\mathring{Ric})_{pq}g_{ij}(HessV)^{qp} + (\mathring{Ric})_{pj}g_{iq}(HessV)^{qp} \\ &\quad + (\mathring{Ric})_{iq}g_{pj}(HessV)^{qp} - (\mathring{Ric})_{ij}g_{pq}(HessV)^{qp} \\ &\quad + T^q_{iq;j} + T_{ijq;}^q \\ &= 3\epsilon V(\mathring{Ric})_{ij} - Vg_{ij}(\mathring{Ric})_{pq}(\mathring{Ric})^{qp} \\ &\quad + 2V(\mathring{Ric})_{pj}(\mathring{Ric})^p_i + V(\mathring{Ric})_{iq}(\mathring{Ric})^q_j \\ &\quad + T^q_{iq;j} + T_{ijq;}^q. \end{aligned}$$

By Lemma 10,

$$T_{ijk} = -((\mathring{Ric})_{ij}V_k - (\mathring{Ric})_{ik}V_j) - (g_{ij}(\mathring{Ric})_{kp}V_{;^p} - g_{ik}(\mathring{Ric})_{jp}V_{;^p}).$$

Hence,

$$\begin{aligned}
(22) \quad T_{iq;j}^q &= -(((\mathring{Ric})^q_i V_{;q} - (\mathring{Ric})^q_q V_{;i}) + (g^q_i (\mathring{Ric})_{qp} V_{;j}^p - g^q_q (\mathring{Ric})_{ip} V_{;j}^p))_{;j} \\
&= ((\mathring{Ric})_{ip} V_{;j}^p)_{;j} = (\mathring{Ric})_{ip;j} V_{;j}^p + (\mathring{Ric})_{ip} V_{;j}^p \\
&= (\mathring{Ric})_{ip;j} V_{;j}^p + (\mathring{Ric})_{ip} ((HessV)^p_j - \epsilon V g^p_j) \\
&= (\mathring{Ric})_{ip;j} V_{;j}^p + V (\mathring{Ric})_{ip} (\mathring{Ric})^p_j - \epsilon V (\mathring{Ric})_{ij}.
\end{aligned}$$

Moreover, recalling that $(\mathring{Ric})_{iq}^q = 0$, we also have

$$\begin{aligned}
(23) \quad T_{ijq}^q &= -((\mathring{Ric})_{ij}^q V_{;q} - (\mathring{Ric})_{iq}^q V_{;j}) - ((\mathring{Ric})_{ij} V_{;q}^q - (\mathring{Ric})_{iq} V_{;j}^q) \\
&\quad - (g_{ij} (\mathring{Ric})_{qp}^q - g_{iq} (\mathring{Ric})_{jp}^q) V_{;j}^p - (g_{ij} (\mathring{Ric})_{qp} - g_{iq} (\mathring{Ric})_{jp}) V_{;j}^{pq} \\
&= -(\mathring{Ric})_{ij;q} V_{;j}^q - (-3\epsilon V) (\mathring{Ric})_{ij} + (\mathring{Ric})_{iq} ((HessV)^q_j - \epsilon V g^q_j) \\
&\quad + (\mathring{Ric})_{jp;i} V_{;j}^p - (g_{ij} (\mathring{Ric})_{qp} - g_{iq} (\mathring{Ric})_{jp}) ((HessV)^{pq}_j - \epsilon V g^{pq}_j) \\
&= \epsilon V (\mathring{Ric})_{ij} + V (\mathring{Ric})_{iq} (\mathring{Ric})^q_j + V (\mathring{Ric})_{jp} (\mathring{Ric})^p_i \\
&\quad - V g_{ij} (\mathring{Ric})_{pq} (\mathring{Ric})^{qp} + (\mathring{Ric})_{jp;i} V_{;j}^p - (\mathring{Ric})_{ij;q} V_{;j}^q.
\end{aligned}$$

Putting together (21), (22) and (23) we conclude that

$$\begin{aligned}
(24) \quad (\Delta(HessV))_{ij} &= 3\epsilon V (\mathring{Ric})_{ij} - 2V (\mathring{Ric})_{pq} (\mathring{Ric})^{pq} g_{ij} \\
&\quad + 3V ((\mathring{Ric})_{jp} (\mathring{Ric})^p_i + (\mathring{Ric})_{ip} (\mathring{Ric})^p_j) \\
&\quad + ((\mathring{Ric})_{ip;j} + (\mathring{Ric})_{jp;i} - (\mathring{Ric})_{ij;p}) V_{;j}^p.
\end{aligned}$$

Combining (19), (20) and (24) we finally obtain

$$\begin{aligned}
\frac{1}{2} div(V \nabla |\mathring{Ric}|^2) &= V |\nabla \mathring{Ric}|^2 + 6\epsilon V |\mathring{Ric}|^2 \\
&\quad + 3V ((\mathring{Ric})_{ip} (\mathring{Ric})^p_j + (\mathring{Ric})_{jp} (\mathring{Ric})^p_i) (\mathring{Ric})^{ij} \\
&\quad - 2((\mathring{Ric})_{ij;p} - (\mathring{Ric})_{ip;j}) (\mathring{Ric})^{ij} V_{;j}^p \\
&= V |\nabla \mathring{Ric}|^2 - 2C_{ijp} (\mathring{Ric})^{ij} V_{;j}^p \\
&\quad + 6\epsilon V |\mathring{Ric}|^2 + 6V tr(\mathring{Ric} \circ \mathring{Ric} \circ \mathring{Ric}).
\end{aligned}$$

To finish the proof, we use Lemma 11 and observe that if $\lambda_1, \lambda_2, \lambda_3$ are the eigenvalues of \mathring{Ric} at some point of M , then $\lambda_1 + \lambda_2 + \lambda_3 = 0$ implies

$$\begin{aligned}
tr(\mathring{Ric} \circ \mathring{Ric} \circ \mathring{Ric}) &= \lambda_1^3 + \lambda_2^3 + \lambda_3^3 \\
&= -3\lambda_1^2 \lambda_2 - 3\lambda_1 \lambda_2^2 = 3\lambda_1 \lambda_2 \lambda_3 = 3det(\mathring{Ric}).
\end{aligned}$$

q.e.d.

5. Proof of Theorem A

Theorem 13. *Let (M^3, g, V) be a compact oriented static triple with positive scalar curvature. If $|\mathring{Ric}|^2 \leq R^2/6$, then one of the following alternatives holds:*

- i) $\mathring{Ric} = 0$ and (M^3, g, V) is equivalent to the standard hemisphere;
- or
- ii) $|\mathring{Ric}|^2 = R^2/6$ and (M^3, g, V) is covered by a static triple that is equivalent to the standard cylinder.

Proof. Using Lagrange's multipliers method, one proves that every symmetric traceless linear endomorphism $A : \mathbb{R}^3 \rightarrow \mathbb{R}^3$ satisfies

$$54(\det(A))^2 \leq |A|^6,$$

with equality if and only if A has at most two distinct eigenvalues. Hence, since V is non-negative, the last two terms of the Bochner type formula (16) satisfy the inequality

$$(25) \quad (R|\mathring{Ric}|^2 + 18\det(\mathring{Ric}))V \geq (R - \sqrt{6}|\mathring{Ric}|)|\mathring{Ric}|^2V,$$

with equality at $p \in M \setminus \partial M$ if and only if $\mathring{Ric}(p)$ has at most two distinct eigenvalues. Since V vanishes on ∂M , integrating (16) on M and using (25) we obtain

$$0 \geq \int_M \left(|\nabla \mathring{Ric}|^2 + \frac{|C|^2}{2} \right) V d\mu + \int_M \left(R - \sqrt{6}|\mathring{Ric}| \right) |\mathring{Ric}|^2 V d\mu.$$

Our assumptions, therefore, imply that (M^3, g, V) is a locally conformally flat static triple with parallel Ricci tensor that either satisfies $\mathring{Ric} = 0$ or $R = \sqrt{6}|\mathring{Ric}|$. The result follows then by the classification results of [36] and [39], see Theorem 1 in the Introduction. q.e.d.

6. The associated singular Einstein manifold

Let (M^n, g, V) be a static triple. As usual, we assume it has constant scalar curvature $\epsilon n(n-1)$ for some $\epsilon \in \{1, 0, -1\}$.

Let \mathcal{U} be a small tubular neighborhood of ∂M diffeomorphic to $[0, 1) \times \partial M$. Given $p \in \mathcal{U}$ and the corresponding point $(r, x) \in [0, 1) \times \partial M$, we write $p = (r, x)$, for simplicity. Let B_1^2 denote the open unit ball centered at the origin of \mathbb{R}^2 and let $S^1 = \mathbb{R}/2\pi\mathbb{Z}$ denote the unit circle.

Let \mathcal{N}^{n+1} be the quotient set

$$(S^1 \times (M \setminus \partial M)) \sqcup (B_1^2 \times \partial M) / \sim,$$

where the equivalence relation \sim identifies $(\theta, p) \in (S^1 \times (\mathcal{U} \setminus \partial M))$ with $(r \cos \theta, r \sin \theta, x) \in (B_1^2 \setminus \{0\}) \times \partial M$ if $p = (r, x)$.

\mathcal{N}^{n+1} canonically inherits a structure of smooth manifold. $\{0\} \times \partial M$ is a codimension-two submanifold of \mathcal{N} that we can identify with $\partial M \subset M$.

One can intuitively imagine \mathcal{N}^{n+1} as the manifold obtained by rotating M around ∂M .

The map that associates to each $\theta_0 \in S^1$ the diffeomorphism of \mathcal{N} determined by $(\theta, p) \in S^1 \times (M \setminus \partial M) \mapsto (\theta + \theta_0, p) \in S^1 \times (M \setminus \partial M)$, $p \in \partial M \mapsto p \in \partial M$ is a smooth action of S^1 on \mathcal{N} whose set of fixed points is $\partial M \subset \mathcal{N}$. We denote by $X \in \mathcal{X}(\mathcal{N})$ the vector field generating this action.

For every $\theta \in S^1$ the map $\phi_\theta : M \rightarrow \mathcal{N}$ defined by $p \in M \setminus \partial M \mapsto (\theta, p) \in S^1 \times (M \setminus \partial M)$, $p \in \partial M \mapsto (0, p) \in B_1^2 \times \partial M$ is a smooth embedding. It gives a natural identification between M and the space of orbits of the S^1 action on \mathcal{N} .

We denote by Ω the open dense subset $\mathcal{N} \setminus \partial M \subset \mathcal{N}$. Ω is naturally identified with $S^1 \times (M \setminus \partial M)$. Since V is positive on $M \setminus \partial M$, we can define a smooth Riemannian metric h on Ω by the formula

$$h = V^2 d\theta^2 + g.$$

The maps $\phi_\theta : (M \setminus \partial M, g) \rightarrow (\Omega, h)$ are then isometric embeddings. Moreover, the vector field X is orthogonal to every hypersurface $\phi_\theta(M \setminus \partial M)$ and such that $|X| \circ \phi_\theta = V$ on $M \setminus \partial M$. These conditions uniquely determine the metric h .

A computation in coordinates shows that the Ricci tensor of h is given by

$$\begin{aligned} Ric_h(X, X) &= - \left(\frac{1}{V} \Delta_g V \right) h(X, X), \\ Ric_h(X, Z) &= 0, \\ Ric_h(Y, Z) &= Ric_g(Y, Z) - \frac{1}{V} (Hess_g V)(Y, Z), \end{aligned}$$

where Y, Z are tangent to M . Since V is a static potential and (M^n, g) has scalar curvature $\epsilon n(n-1)$, the metric h is Einstein,

$$Ric_h = \epsilon n h.$$

It is also a straightforward computation to verify that $X = \partial_\theta$ is a Killing vector field, i.e., that the S^1 action described before is an action by isometries.

To understand the behavior of the metric h near a connected component $\partial_i M$ of ∂M , we introduce Fermi coordinates. Let s be the distance function to ∂M in (M, g) and let (x^2, x^3) be a local coordinate system on a neighborhood W of a point in $\partial_i M$. Using $(x^0 = \theta, x^1 = s, x^2, x^3)$ as local coordinates in $\mathcal{N} \setminus \partial M$, we can write

$$(26) \quad h = (ds^2 + k_i^2 s^2 d\theta^2 + \sum_{i,j=2}^3 g_{ij}(0, x^2, x^3) dx^i dx^j) + E,$$

where k_i denote the constant value of $|\nabla V|$ on $\partial_i M$ and $|E| \leq Cs^2$ for some constant $C > 0$. Indeed, using the fact that ∂M is totally geodesic and Proposition 7, we obtain the expansions

$$g_{ij}(s, x^1, x^2) = g_{ij}(0, x^1, x^2) + O(s^2) \quad \text{and} \quad V(s, x^1, x^2) = k_i s + O(s^3).$$

We remark that $E_{1A} = E(\partial_s, \partial_{x^A}) = 0$ for all $A = 0, 1, 2, 3$. Moreover, the coefficients of E in this coordinate system depend only on the coordinates $x^1 = s, x^2, x^3$ and have a Taylor expansion in s as s goes to zero with coefficients that are smooth functions of the coordinates x^2, x^3 .

Equation (26) has the interpretation that near a connected component $\partial_i M$ of the singular set $\partial M \subset \mathcal{N}$, the metric h has the structure of cone metric with link $S^1(k_i)$ (the circle with length $2\pi k_i$, a number sometimes called the *cone angle*) along the codimension-two smooth submanifold $(\partial_i M, g_{\partial M})$, up to a perturbation. Metrics with this singular behavior have been already described and studied in the literature. We refer the reader to [43] (see their definition of singular Riemannian metric), [8] (where the singularities are called of “edge-cone” type) and to the more general notion of iterated edge metrics on iterated edge spaces used in [1] (see also [2]).

Summarizing up, the manifold \mathcal{N}^{n+1} endowed with the Einstein edge metric h and the S^1 action by isometries generated by the vector field X , which is everywhere orthogonal to the orbit space (M^n, g) , vanishes on the singular set ∂M and has norm V , will be called the *associated singular Einstein manifold* to the static triple (M^n, g, V) .

We proceed with a description of some topological and geometrical properties of the associated singular Einstein manifold.

\mathcal{N}^{n+1} is compact if and only if M^n is compact, and it is simply connected if M and ∂M are simply connected (by Van Kampen’s Theorem). When \mathcal{N} is compact, since the vector field $X \in \mathcal{X}(\mathcal{N})$ does not vanish on Ω , \mathcal{N} and $\partial M = \{X = 0\}$ have the same Euler characteristic (one can argue as in [37]). When \mathcal{N} is oriented, $X \mapsto -X$ gives rise to an orientation inverting diffeomorphism.

The metric $h = V^2 d\theta^2 + g$ extends smoothly to a component $\partial_i M$ of $\partial M \subset \mathcal{N}$ if and only if there exists a smooth function $U : [0, \sqrt{\eta}) \times \partial_i M \rightarrow [0, +\infty)$ such that $U(0, p) = 1$ for all $p \in \partial M$ and $V(s, p) = sU(s^2, p)$ in the Fermi coordinate system $[0, \eta) \times \partial_i M$ (see [10], 9.114). A necessary condition is that $k_i = |\nabla V| = 1$ on $\partial_i M$. This condition is usually stated to be also sufficient (see [12] or [28]). When h extends smoothly to all \mathcal{N} , (\mathcal{N}, h) is an example of the type of Einstein manifolds studied in [56].

We emphasize that, in general, equivalent static triples with the same scalar curvature give rise to *geometrically distinct* associated Einstein manifolds. In fact, the orbits of the S^1 action will have different lengths if the static potentials are proportional by a factor $\lambda \neq 1$.

When ∂M is compact it is possible to normalize the static potential V in such a way that $|\nabla V| \leq 1$ on ∂M . This means that in this situation, to study (M^n, g, V) one can work with an edge space (\mathcal{N}^{n+1}, h) such that all cone angles are less than or equal to 2π .

Since (M, g) is complete, \mathcal{N} endowed with the distance function induced by h is a complete length space and the Hopf–Rinow Theorem holds (see [51], Section 2.4). In particular, \mathcal{N} is compact if and only if (\mathcal{N}, h) has finite diameter.

The structure of the metric h near the singular set clearly implies that geodesics realizing the distance between a point in \mathcal{N} and a component of ∂M meet ∂M orthogonally. The proof of this fact is essentially the same as the proof of the Gauss’ Lemma.

For the computation of the volume of an associated singular Einstein manifold (\mathcal{N}^{n+1}, h) to a compact oriented static triple with scalar curvature $R = n(n-1)$, the singular codimension-two set ∂M is negligible. The relation between the volume element of h and the volume element of g is $d\mu_h = V d\theta \wedge d\mu_g$. Hence, Fubini’s theorem and the static equations yield

$$\begin{aligned} |N| &= \int_0^{2\pi} \int_M d\mu_h = 2\pi \int_M V d\mu_g = -\frac{2\pi}{n} \int_M \Delta V d\mu_g = \\ &= \frac{2\pi}{n} \int_{\partial M} g \left(\nabla V, \frac{\nabla V}{|\nabla V|} \right) d\sigma_g = \frac{2\pi}{n} \sum_{i=1}^r k_i |\partial_i M|. \end{aligned}$$

Each function ϕ on M^n corresponds to a function on \mathcal{N}^{n+1} invariant under the S^1 action. For the sake of simplicity, we make no notational distinction between them. If $\psi \in C^\infty(M)$, a direct computation in coordinates shows that $V \Delta_h \psi = V \Delta_g \psi + g(\nabla_g V, \psi) = \operatorname{div}_g(V \nabla_g \psi)$ on Ω .

In the four-dimensional case, the Riemann curvature tensor of the Einstein metric h is determined by the Weyl tensor W of h . An explicit calculus in coordinates shows that W depends only on the Ricci tensor of the metric g . In fact, if $\{e^0 = V^{-1} \partial_\theta, e^1, e^2, e^3\}$ is an orthonormal basis of $T_p \mathcal{N}$, $p \in \Omega$, all the components W_{ABCD} vanish except

$$W_{ijkl} = (\mathring{Ric}_g \odot g)_{ijkl}, \quad W_{i0j0} = -(\mathring{Ric}_g)_{ij} \quad \text{for } i, j, k, l = 1, 2, 3.$$

When \mathcal{N}^4 is oriented, considering W as a symmetric endomorphism of the space of two-forms on M and its usual self-dual/anti-self-dual decomposition $W = W^+ + W^-$ (see [10]), one can also show by a direct computation that

$$|W^+|_h^2 = |W^-|_h^2 = |\mathring{Ric}_g|_g^2.$$

In fact, W^+ and $-\mathring{Ric}_g$ have the same set of eigenvalues. We will use these facts about W in the last sections.

As a final remark concerning the four-dimensional case, we can now unveil the origin of the Bochner type formula (16) for a static triple (M^3, g, V) : it is just the Bochner formula for the self-dual part of the Weyl tensor of the Einstein metric $h = V^2 d\theta^2 + g$ (see [31], Formula 1.3).

We finish our description of the associated singular Einstein manifolds to static triples by discussing some examples.

1) Given the unit hemisphere $(S_+^n, g_{can}, V = x_{n+1})$, the associated Einstein manifold is the unit round sphere (S^{n+1}, g_{can}) . The S^1 action is generated by rotations in the plane $\{x_1 = \dots = x_n = 0\} \subset \mathbb{R}^{n+2}$.

2) Given the standard three-dimensional cylinder over S^2 with scalar curvature 6, the associated Einstein manifold is the product $(S^2 \times S^2, \frac{1}{3}(g_{can} + g_{can}))$. The S^1 action is generated by rotations of one of the S^2 factors around a fixed axis.

3) Any Schwarzschild–de Sitter space of positive mass m has the property that the norm of the gradient of its static potential attains different values on different boundary components (see Remark 7). No rescaling of V can remove the singularities of the associated Einstein manifold, which is, topologically, $S^2 \times S^2$.

4) The Euler characteristic of the associated Einstein four-manifold (\mathcal{N}^4, h) to a compact orientable static triple (M^3, g, V) with positive scalar curvature is equal to the Euler characteristic of the two-dimensional orientable closed surface ∂M . In particular, $\chi(\mathcal{N})$ is an even integer. Therefore, no Einstein metric on $\mathbb{C}\mathbb{P}^2$ (e.g., the Fubini–Study metric) or on $\mathbb{C}\mathbb{P}^2 \# 2\overline{\mathbb{C}\mathbb{P}^2}$ (e.g., the Shen–LeBrun–Weber metric [19]) arises from a compact static three-manifold.

5) If a compact static triple with positive scalar curvature (M^3, g, V) generates an associated Einstein manifold $(\mathcal{N}^4, h = V^2 d\theta^2 + g)$ that is Kähler, then it must be isometric to the standard cylinder. In fact, it is known that the self-dual part W^+ of the Weyl tensor of a Kähler four-manifold is such that $|W_h^+|_h^2 = R_h^2/24$ (see [10], 16.62). Since $|Ric_g|_g^2 = |W_h^+|_h^2$ and $R_h = 2R_g$, the claim follows from Theorem A.

7. Topology of compact static three-manifolds with positive scalar curvature

The topology of three-dimensional static triples (M^3, g, V) can be studied using area-minimizing surfaces that can be produced in mean-convex manifolds by direct variational methods if the topology of M is sufficiently complicated. For, as we will see, locally area-minimizing surfaces exist in $M \setminus \partial M$ only in exceptional cases.

Let (M^3, g, V) be a compact static triple. Let Σ^2 be a compact surface and $\phi : \Sigma \rightarrow M$ a proper embedding, i.e., an embedding of Σ in M such that $\phi(\Sigma) \cap \partial M = \phi(\partial\Sigma)$. Sometimes we identify Σ with its image in

M . We will frequently assume that Σ is two-sided, i.e., that there exists a smooth unit normal vector field N on Σ .

Let $(\overline{M}, \overline{g}) := (M \setminus \partial M, V^{-2}g)$. Since M is compact, V vanishes on ∂M and $|\nabla V|$ is a positive constant on each boundary component of M , $(\overline{M}, \overline{g})$ is a conformally compact Riemannian manifold. These Riemannian manifolds have bounded geometry [3], i.e., they have positive injectivity radius and the Riemann curvature tensor and all of its covariant derivatives are bounded. This allows one to consider, for some small $\delta > 0$, the smooth flow $\Phi : [0, \delta) \times \Sigma \rightarrow M$ of Σ with normal velocity V ,

$$(27) \quad \frac{d}{dt}\Phi_t(p) = V(\Phi_t(p))N_t(p) \quad \text{for all } p \in \Sigma, \quad \Phi_0 = \phi.$$

In fact, (27) is just the flow of Σ in $(\overline{M}, \overline{g})$ by equidistant surfaces. This flow is smooth if δ is less than the injectivity radius of $(\overline{M}, \overline{g})$.

The surfaces $\Phi_t(\Sigma)$, $t \in [0, \delta)$, are compact properly embedded surfaces in (M, g) with the same boundary as Σ , since V vanishes on ∂M .

We show that the flow (27) starting at a minimal surface typically decreases the area (compare with [16], Proposition 3.3, and [41], Lemma 2.4). More precisely, we prove the following:

Proposition 14. *Let (M^3, g, V) be a static triple with scalar curvature $R = 6\epsilon$, $\epsilon \in \{-1, 0, 1\}$. Let Σ_0^2 be a compact properly embedded two-sided minimal surface in (M^3, g) and choose N_0 a smooth unit normal vector field on Σ_0 .*

Let Φ_t , $t \in [0, \delta)$, be the smooth normal flow with speed V starting at Σ_0 defined by (27). Write $\Sigma_t = \Phi_t(\Sigma_0)$. Then

$$t \in [0, \delta) \mapsto |\Sigma_t|$$

is monotone non-increasing. When $t \mapsto |\Sigma_t|$ is constant,

i) If Σ is closed, then

$$\Phi : ([0, \delta) \times \Sigma_0, (V \circ \Phi)^2 dt^2 + g_{\Sigma_0}) \rightarrow (M, g)$$

is an isometry onto its image $W \subset M$. Moreover, (Σ_0, g_{Σ_0}) has constant Gaussian curvature $K = 3\epsilon$ and V is constant on each Σ_t , so that g on W is isometric to the product metric $ds^2 + g_{\Sigma_0}$.

ii) If $\partial\Sigma$ is not empty, then $\epsilon = 1$ and

$$\Phi : ([0, \delta) \times (\Sigma_0 \setminus \partial\Sigma_0), (V \circ \Phi)^2 dt^2 + g_{\Sigma_0}) \rightarrow (M, g)$$

is an isometry onto its image $W \subset M \setminus \partial M$. Moreover, (Σ_0, g_{Σ_0}) is isometric to the round hemisphere (S_+^2, g_{can}) with Gaussian curvature $K = 1$ and V restricted to Σ_t is a static potential, so that g on W is isometric to the constant sectional curvature metric $V_{|\Sigma_0}^2 d\theta^2 + g_{\Sigma_0}$.

Proof. Let $\vec{H}_t = -H_t N_t$ define the mean curvature H_t of Σ_t . The variation of H_t is given by $\partial_t H_t = -\mathcal{L}_{\Sigma_t} V$. By Proposition 9,

$$\partial_t H_t = -g(\nabla V, \vec{H}_t) - |A_t|^2 V \leq g(\nabla V, N_t) H_t.$$

Since $H_0 = 0$, Gronwall's inequality implies that $H_t \leq 0$ for every $t \in [0, \delta)$. By the first variation formula of area for variations that fix the boundary $\partial \Sigma_t = \partial \Sigma_0$, this implies that $t \in [0, \delta) \mapsto |\Sigma_t|$ is non-increasing.

If $t \in [0, \delta) \mapsto |\Sigma_t|$ is constant, the previous analysis implies that each Σ_t is totally geodesic. It follows that each (Σ_t, g_{Σ_t}) is isometric to (Σ_0, g_{Σ_0}) .

Denote by V_t the restriction of V to Σ_t . Since each Σ_t is totally geodesic, using the general formula for the variation of the second fundamental form of a surface Σ under a normal variation with speed ψ ,

$$\partial_t A_{ij} = -\psi R_{iNjN} - (\text{Hess}_{\Sigma} \psi)_{ij} + \psi |A|^2,$$

we deduce that on each Σ_t we have

$$\text{Hess}_{\Sigma_t} V_t(X, Y) = -g(R(X, N_t)N_t, Y) V_t,$$

for every pair of vectors $X, Y \in T\Sigma_t$. Given an orthonormal basis $\{X, Y, N_t\}$ of tangent vectors at a point in the totally geodesic surface Σ_t , we have

$$\begin{aligned} \frac{1}{V_t} \text{Hess}_{\Sigma_t} V_t(X, X) &= \frac{1}{V_t} \text{Hess} V_t(X, X) \\ &= \text{Ric}(X, X) - 3\epsilon \\ &= \text{sec}(X \wedge Y) + \text{sec}(X \wedge N_t) - 3\epsilon \\ &= K_t + g(R(X, N_t)N_t, X) - 3\epsilon \\ &= (K_t - 3\epsilon) - \frac{1}{V_t} \text{Hess} V_t(X, X). \end{aligned}$$

Since $X \in T\Sigma_t$ is arbitrary, we conclude that V_t is a non-negative smooth function on Σ_t that satisfies

$$(28) \quad \text{Hess}_{\Sigma} V_t = \frac{1}{2}(K_t - 3\epsilon)V_t g \quad \text{and} \quad V_t = 0 \text{ on } \partial \Sigma_t.$$

If $\partial \Sigma_t$ is empty, the existence of a positive solution to (28) implies $K_t = 3\epsilon$ and V_t constant. If $\partial \Sigma_t$ is non-empty, one must have $K = \epsilon = 1$ and V_t is a solution to the static equation on (Σ_t, g_{Σ_t}) . See Appendix B for a proof of these facts. The last statements of *i*) and *ii*) follows by a reparametrization. q.e.d.

The next proposition exemplifies how one can combine Proposition 14 and the existence of area-minimizing surfaces to study the topology of a three-dimensional static triple in a quite effective way.

Proposition 15. *Let (M^3, g, V) be a compact static triple with non-empty boundary. The inclusion map $i : \partial M \rightarrow M$ induces an injective map $i_* : \pi_1(\partial M) \rightarrow \pi_1(M)$ between the fundamental groups.*

Proof. Let $[\Gamma] \in \pi_1(\partial M)$ be such that $i_*[\Gamma] = 0$ in $\pi_1(M)$. We may assume that Γ is a smooth embedded closed curve. We want to show that Γ bounds a disk in ∂M .

Since Γ bounds a disk in M and (M^3, g) is a compact mean convex three-manifold, there exists a solution Σ_0 to the Plateau problem for this given Γ [46], i.e., Σ_0 is a disk in M with $\partial\Sigma = \Gamma$ such that

$$|\Sigma_0| = \inf\{|\Sigma|; \Sigma \text{ immersed disk in } (M^3, g) \text{ with } \partial\Sigma = \Gamma\}.$$

In [46] it is also shown that any solution is an embedded minimal disk Σ , with boundary $\partial\Sigma = \Gamma$, which is either contained in ∂M or properly embedded in M . In the first case $[\Gamma] = 0$ in $\pi_1(\partial M)$. Therefore, it remains to be analyzed the case where Σ_0 is a properly embedded disk in M .

Since V vanishes on $\Gamma \subset \partial M$, the flow (27), $\{\Sigma_t\}_{t \in [0, \delta)}$, with normal speed V starting at Σ_0 , is such that $\partial\Sigma_t = \Gamma$ for all $t \in [0, \delta)$. Thus, $|\Sigma_t| \geq |\Sigma_0|$, as Σ_0 is a solution to the Plateau problem. By Proposition 14, the opposite inequality also holds. Thus, each Σ_t is also a solution to the minimization problem we are considering and the statement *ii*) in Proposition 14 holds true.

To finish the proof, one uses a continuation argument. Let $T > 0$ be the maximal time in which the flow (27) exists and is smooth. By the above argument, we are in the situation described by Proposition 14, item *ii*). We claim that $T = +\infty$. In fact, the surfaces $\Sigma_t \setminus \partial\Sigma_t$ never touch the boundary of M in finite time (for it would imply that $(M \setminus \partial M, V^{-2}g)$ is incomplete) and, if $T < +\infty$, a sequence Σ_{t_i} with $t_i \rightarrow T$ would be a sequence of Plateau solutions converging to another solution to the Plateau problem for Γ (which is smooth and properly embedded by [46]) and it would be possible to continue the flow beyond T .

Now we do the same argument flowing in the direction of the opposite normal. In the end, we have obtained an isometric embedding of $(S_+^3 \setminus \partial S_+^3, g_{can})$ in $(M \setminus \partial M, g)$. In particular, M is diffeomorphic to S_+^3 . And in this case, obviously $[\Gamma] = 0$ in $\pi_1(\partial M)$, as we wanted to prove. q.e.d.

REMARK 8. There are oriented static triples (M^3, g, V) with negative scalar curvature and non-spherical boundary components, see [41]. It is interesting to observe that essentially the same proof as above gives a similar result for the non-compact static triples considered in that paper.

We have now all ingredients to prove Theorem B.

Theorem 16. *Let (M^3, g, V) be a compact oriented static triple with positive scalar curvature. Then*

- i) The universal cover of M is compact.*
- ii) If ∂M contains an unstable component, then ∂M contains exactly one unstable component. In this case, M is simply connected.*
- iii) Each connected component of ∂M is diffeomorphic to a sphere.*

Proof. As usual, we may assume $R = 6$.

i) Since M is compact, $|\nabla V| + V^2$ attains its maximum, necessarily at a boundary component $\partial_1 M$ (Proposition 5). By Proposition 8, $\partial_1 M$ has Gaussian curvature $K \geq 1$ and, therefore, is diffeomorphic to a two-sphere. This implies that the boundary of the universal covering \tilde{M} has also a component $\partial_1 \tilde{M}$ diffeomorphic to a two-sphere (in fact, every component covering $\partial_1 M$ satisfies this property). Now, any point in the regular set of the associated singular Einstein manifold (\mathcal{N}^4, h) can be joined to the compact set $\partial_1 \tilde{M}$ by a geodesic minimizing the distance of p to $\partial_1 M$ and touching this hypersurface orthogonally. Estimating the distance to the first focal point as in the classical proof of Bonnet–Myers theorem, we conclude by the generalization of the Hopf–Rinow theorem to the singular space $(\tilde{\mathcal{N}}^4, h)$ ([51], section 2.4) that $\tilde{\mathcal{N}}$ must be compact, as any of its points must be at distance at most $\pi/2$ from the compact set $\partial_1 \tilde{M}$.

ii) By item *i)*, without loss of generality we can assume M is simply connected and show that ∂M has exactly one unstable component. (M^3, g) cannot be the standard cylinder because it contains by hypothesis at least one unstable boundary component that we denote by $\partial_u M$. We remark also that all embedded closed surfaces in M are two-sided (as M is simply-connected) and that we can assume $[\partial_u M] \neq 0$ in $H_2(M; \mathbb{Z})$ (otherwise, $\partial M = \partial_u M$). Since the standard cylinder is ruled out, arguing as in Proposition 15 using the results of the Geometric Measure Theory about the existence and regularity of area-minimizing currents in an integral homology class (see [58]) and Proposition 14, item *i)*, one shows that the minimization of area in the homology class of $\partial_u M$ produces a smooth embedded oriented stable minimal surface (possibly disconnected and with integer multiplicities) homologous to $\partial_u M$ that is necessarily contained inside ∂M . No connected component of such surface can be contained in an unstable component of ∂M . Reasoning on the homology relations between the boundary components (as in [41], Lemma 3.3), the existence of an unstable component other than $\partial_u M$ is seen to lead to a contradiction.

iii) Each stable component of ∂M is diffeomorphic to a sphere because M is oriented and has positive scalar curvature [55]. On the other hand, if there is an unstable component of ∂M , M must be simply connected by *ii)*. The conclusion follows from Proposition 15. q.e.d.

To conclude this section, we describe the two static triples that can be obtained from the standard cylinder. Besides the (non-orientable)

static quotient of the Schwarzschild–de Sitter spaces, they are the only static quotients of the locally conformally flat examples described in Theorem 1.

Given

$$(M^3, g, V) = \left(\left[0, \frac{\pi}{\sqrt{3}} \right] \times S^2, g_{prod} = dt^2 + \frac{1}{3}g_{can}, \frac{1}{\sqrt{3}} \sin(\sqrt{3}t) \right),$$

the maps

$$A_1 : (t, x) \in M \mapsto (t, -x) \in M, \quad A_2 : (t, x) \in M \mapsto \left(\frac{\pi}{\sqrt{3}} - t, -x \right) \in M$$

are isometric involutions without fixed points such that $V \circ A_i = V$, $i = 1, 2$. Hence, the quotients $(M^3, g, V)/A_i$, $i = 1, 2$, are static triples.

A_1 inverts orientation and the first quotient M/A_1 is diffeomorphic to $[0, \pi/\sqrt{3}] \times \mathbb{RP}^2$. On the other hand, M/A_2 is oriented, contains an embedded copy of \mathbb{RP}^2 and has a *connected* boundary diffeomorphic to S^2 (in fact, M/A_2 is diffeomorphic to \mathbb{RP}^3 minus a ball). The associated Einstein manifolds are both isometric to the standard product

$$(S^2 \times \mathbb{RP}^2, \frac{1}{3}(g_{can} + g_{can})),$$

but the corresponding S^1 actions are generated by rotations around a fixed axis of different factors of $S^2 \times \mathbb{RP}^2$: S^2 in the case of $(M, g, V)/A_1$, \mathbb{RP}^2 in the case of $(M, g, V)/A_2$. We stress the subtle fact that, despite M/A_2 being orientable, the associated Einstein manifold to M/A_2 is non-orientable.

8. Solution to a Yamabe type problem

Given (M^3, g, V) a compact static triple with scalar curvature 6, we consider (\mathcal{N}^4, h) the associated singular Einstein manifold (as described in Section 6). As before, we denote by Ω the open dense subset of \mathcal{N} where the Riemannian metric h is defined, i.e., $\Omega = \mathcal{N} \setminus \partial M$.

The metric h is Einstein with scalar curvature 12. Let $|W^+|_h$ be the norm of self-dual part of the Weyl curvature tensor of h . As remarked before, $|W^+|_h = |\mathring{Ric}_g|_g$.

We consider the equation

$$(29) \quad -\Delta_h u + \frac{1}{6}(12 - 2\sqrt{6}|W^+|_h)u = \frac{1}{6}\lambda u^3,$$

where λ is some constant. As showed in [31], the right hand side defines a second order differential operator that has some conformal covariance properties. To explain the geometric relevance of (29), we observe that a positive function $u \in C^2(\Omega)$ that solves (29) gives rise to a conformal metric $\tilde{h} = u^2 h$ on Ω such that

$$\tilde{R} - 2\sqrt{6}|\tilde{W}^+|_{\tilde{h}} = \lambda$$

is constant (see [31]).

The aim of this section is to prove that, except in the case where $W^+ = 0$, there exists for some *non-positive* λ a mildly regular solution to (29) on (\mathcal{N}^4, h) that is invariant under the S^1 action on \mathcal{N} .

Theorem 17. *Let (\mathcal{N}^4, h) be the associated singular Einstein manifold to a compact static triple with positive scalar curvature (M^3, g, V) that is not equivalent to the standard hemisphere.*

There exists a constant $\lambda \leq 0$ and a positive function $\phi \in W^{1,2}(\mathcal{N}) \cap L^\infty(\mathcal{N})$, invariant under the S^1 action on \mathcal{N} , that is a weak solution to (29). Moreover, $\phi \in C^{0,\mu}(\mathcal{N}) \cap C^{2,\alpha}(\Omega)$ for some $\mu, \alpha \in (0, 1)$ and ϕ is a strong solution to (29) on Ω .

We remark that the Sobolev space $W^{1,2}(\mathcal{N})$ is the completion of $C_0^\infty(\Omega)$ with respect to the usual Sobolev norm.

In order to deal with this modified Yamabe problem in the singular space (\mathcal{N}^4, h) , we apply the machinery developed by Akutagawa, Carron and Mazzeo in [1] to solve the Yamabe problem in singular spaces. In fact, once we reformulate the problem considering the required S^1 -invariance of the solution, the proof of Theorem 17 is essentially the same proof of Theorems 1.12 and 1.15 in [1], combined with the regularity results of [2].

As in the study of the classical Yamabe problem, one uses the variational approach. Consider the functional

$$\mathcal{W}(\phi) = \frac{\int_{\mathcal{N}} |\nabla^h \phi|_h^2 + \frac{1}{6}(12 - 2\sqrt{6}|W^+|)\phi^2 d\mu_h}{(\int_{\mathcal{N}} \phi^4 d\mu_h)^{1/2}}$$

defined for every function $\phi \in W^{1,2}(\mathcal{N})$, $\phi \neq 0$. This functional is well-defined because the usual Sobolev embedding $W^{1,2}(\mathcal{N}) \subset L^4(\mathcal{N})$ holds in (\mathcal{N}^4, h) (see [1], section 2.2). It is straightforward to verify that \mathcal{W} is bounded from below.

Let A denote the set consisting of functions on $W^{1,2}(\mathcal{N})$ that are invariant under the S^1 action on (\mathcal{N}^4, h) . If $\phi \in A$ is not identically zero, we have

$$\mathcal{W}(\phi) = (2\pi)^{1/2} \frac{\int_M (|\nabla^g \phi|_g^2 + \frac{1}{6}(12 - 2\sqrt{6}|Ric_g^\circ|)\phi^2) V d\mu_g}{(\int_M \phi^4 V d\mu_g)^{1/2}}.$$

Now we define

$$\mathcal{W}^M(\mathcal{N}, [h]) := \inf\{\mathcal{W}(\phi); \phi \in A, \phi \geq 0, \phi \neq 0\}.$$

A non-negative function $\phi \in A$ attaining the infimum $\mathcal{W}^M(\mathcal{N}, [h])$ has to be a weak solution of the following equation on (M^3, g) :

$$(30) \quad -div_g(V\nabla^g \phi) + \frac{1}{6}(12 - 2\sqrt{6}|Ric_g^\circ|)V\phi = \frac{1}{6}\mathcal{W}^M(\mathcal{N}, [h])V\phi^3.$$

This is precisely equation (29) for the S^1 -invariant function $\phi \in A$, where $\lambda = \mathcal{W}^M(\mathcal{N}, [h])$. It is a quasi-linear degenerated elliptic equation

(see Appendix C for more information about this type of degenerated equation).

We first argue that the Yamabe type invariant $\mathcal{W}^M(\mathcal{N}, [h])$ is *non-positive*. The next proposition contains the key inequality, which should be compared with the refined Kato inequality for Codazzi tensors [15].

Proposition 18. *Let (M^3, g, V) be a static triple. At a point $p \in M$ where \mathring{Ric} does not vanishes,*

$$(31) \quad |\nabla|\mathring{Ric}||^2 \leq \frac{3}{5} \left(|\nabla\mathring{Ric}|^2 + \frac{|C|^2}{2} \right).$$

The proof is postponed to the Appendix A, where it is shown more generally that a similar inequality holds true for every symmetric two-tensor that has constant trace and zero divergence.

Following [31], Proposition 3.4, the idea to prove that $\mathcal{W}^M(N, [h])$ is non-positive is to manipulate the Bochner type formula (16) to get rid of the gradient terms using the refined inequality (31). To a systematic recent use of this idea, we refer the reader to the beautiful paper [13].

Proposition 19. *Let (M^3, g, V) be a compact oriented static triple with scalar curvature $R = 6$. If $\mathring{Ric} \neq 0$ on M , then*

$$\inf_{\{\phi \in C^\infty(M); \phi > 0\}} \int_M \left(6|\nabla\phi|^2 + (12 - 2\sqrt{6}|\mathring{Ric}|)\phi^2 \right) V d\mu_g \leq 0.$$

Proof. For each $\epsilon > 0$, we consider a suitable power of the positive smooth function $\phi_\epsilon = (|\mathring{Ric}|^2 + \epsilon)^{1/2}$ and take the limit as ϵ goes to zero.

By hypothesis, ϕ_ϵ is not constant. Given $p > 0$,

$$(32) \quad \begin{aligned} \operatorname{div}(V\nabla(\phi_\epsilon^p)) &= p \operatorname{div}(\phi_\epsilon^{p-1}(V\nabla\phi_\epsilon)) \\ &= p\phi_\epsilon^{p-1} \operatorname{div}(V\nabla\phi_\epsilon) + p(p-1)V\phi_\epsilon^{p-2}|\nabla\phi_\epsilon|^2 \\ &= p\phi_\epsilon^{p-2}(\phi_\epsilon \operatorname{div}(V\nabla\phi_\epsilon) + (p-1)V|\nabla\phi_\epsilon|^2) \\ &= p\phi_\epsilon^{p-2} \left(\operatorname{div} \left(\frac{V\nabla(\phi_\epsilon^2)}{2} \right) + (p-2)V|\nabla\phi_\epsilon|^2 \right). \end{aligned}$$

Since $\nabla(|\mathring{Ric}|^2) = \nabla(\phi_\epsilon^2) = 2\phi_\epsilon \nabla\phi_\epsilon$, at a point where $|\mathring{Ric}| \neq 0$ we have

$$|\nabla\phi_\epsilon|^2 = \frac{1}{4\phi_\epsilon^2} |\nabla|\mathring{Ric}|^2|^2 = \frac{|\mathring{Ric}|^2}{|\mathring{Ric}|^2 + \epsilon} |\nabla|\mathring{Ric}||^2 \leq |\nabla|\mathring{Ric}||^2.$$

By Proposition 18, we conclude that

$$(33) \quad |\nabla\phi_\epsilon|^2 \leq \frac{3}{5} \left(|\nabla\mathring{Ric}|^2 + \frac{|C|^2}{2} \right).$$

Inequality (33) holds even at a point where $\mathring{Ric} = 0$, since at such point ϕ_ϵ attains a minimum. Combining (32), the Bochner type formula

(16) and inequality (33), we obtain

$$\begin{aligned} & \operatorname{div}(V\nabla(\phi_\epsilon^p)) \\ & \geq p\phi_\epsilon^{p-2}V\left(6|\mathring{Ric}|^2 + 18\det(\mathring{Ric}) + \left(\frac{5}{3} + p - 2\right)|\nabla\phi_\epsilon|^2\right). \end{aligned}$$

Choosing $p = 1/3$, the last term vanishes. Recalling inequality (25), we finally obtain

$$(34) \quad \operatorname{div}(V\nabla(\phi_\epsilon^{1/3})) \geq \frac{1}{3}V(6 - \sqrt{6}|\mathring{Ric}|)|\mathring{Ric}|^2\phi_\epsilon^{-5/3}.$$

Multiplying (34) by the $\phi_\epsilon^{1/3} \in C^\infty(M)$ and integrating by parts, we have

$$\int_M (3|\nabla(\phi_\epsilon^{1/3})|^2 + (6 - \sqrt{6}|\mathring{Ric}|)|\mathring{Ric}|^2\phi_\epsilon^{-4/3})Vd\mu_g \leq 0.$$

Hence, for all $\epsilon > 0$,

$$\begin{aligned} & \int_M \left(3|\nabla(\phi_\epsilon^{1/3})|^2 + (6 - \sqrt{6}|\mathring{Ric}|)(\phi_\epsilon^{1/3})^2\right)Vd\mu_g \\ & \leq \int_M (6 - \sqrt{6}|\mathring{Ric}|)\phi_\epsilon^{-4/3}(\phi_\epsilon^2 - |\mathring{Ric}|^2)Vd\mu_g \leq D\epsilon^{1/3}, \end{aligned}$$

for some constant $D > 0$. The proposition follows. q.e.d.

Proposition 19 implies that

$$\mathcal{W}^M(\mathcal{N}, [h]) \leq 0,$$

when the compact static triple (M^3, g, V) is not equivalent to the standard hemisphere.

Now, following [1], for each open set $\mathcal{U} \subset \mathcal{N}$ we set

$$\mathcal{S}(\mathcal{U}) = \inf\left\{\int_{\mathcal{N}} |\nabla^h\phi|_h^2 d\mu_h; \phi \in W_0^{1,2}(\mathcal{U} \cap \Omega), \|\phi\|_4^2 = 1\right\},$$

$$\mathcal{W}(\mathcal{U}) = \inf\{\mathcal{W}(\phi); \phi \in W_0^{1,2}(\mathcal{U} \cap \Omega), \|\phi\|_4^2 = 1\},$$

and define the local Sobolev and the local (modified) Yamabe constant by

$$\mathcal{S}_\ell(\mathcal{N}, h) = \inf_{p \in \mathcal{N}} \lim_{r \rightarrow 0} \mathcal{S}(B(p, r)) \quad \text{and} \quad \mathcal{W}_\ell(\mathcal{N}, [h]) = \inf_{p \in \mathcal{N}} \lim_{r \rightarrow 0} \mathcal{W}(B(p, r)).$$

Accordingly to [1],

$$\mathcal{S}_\ell(\mathcal{N}, h) = \mathcal{W}_\ell(\mathcal{N}, [h]) > 0,$$

see Lemma 1.3, Proposition 1.4 and Theorem 2.2 in [1]. The explicit value of the local Sobolev constant of an edge space like (\mathcal{N}^4, h) was computed in [48]. At this point, it may be useful to remark that $R_h - 2\sqrt{6}|W^+|_h$ is a bounded function on \mathcal{N} (in fact, it is a Lipschitz function) so that all conditions *iv*) *a*), *b*), *c*) described in [1] for $Scal_g$ (their

notation) are satisfied by $R_h - 2\sqrt{6}|W^+|_h$. It is straightforward to verify that their arguments apply for such modified potential instead of $Scal_g$.

With this information at hands, we now sketch the proof of Theorem 17 (see [1], proof of Theorem 1.12). The idea to obtain a solution to (30) is first to produce minimizers $\phi_p \in A$ of the modified functional

$$\mu_p(\phi) = \frac{\int_{\mathcal{N}} |\nabla^h \phi|_h^2 + \frac{1}{6}(12 - 2\sqrt{6}|W^+|)\phi^2 d\mu_h}{(\int_{\mathcal{N}} \phi^p d\mu_h)^{2/p}},$$

on the space of non-negative functions $\phi \in A$ such that $\|\phi\|_{L^{2p/(p-2)}(\mathcal{N})} = 1$ by the direct method of the calculus of variations and then, using a priori estimates, obtain a sequence ϕ_{p_i} with $p_i \rightarrow 4$ converging to a function $\phi \in A$ such that $\mu_4(\phi) = \mathcal{W}^M(\mathcal{N}, [h])$.

In [1] it was shown that this approach works if the local Yamabe invariant is *strictly greater* than $\mathcal{W}^M(\mathcal{N}, h)$. Since we proved $\mathcal{W}^M(\mathcal{N}, h) \leq 0$, the proof of Theorem 1.12 in [1] applies.

The conclusion is that there exists a weak solution $\phi \in W^{1,2}(\mathcal{N}) \cap L^\infty(\mathcal{N})$ to (30) that is non-negative and invariant under the S^1 -action. The positiveness of ϕ follows by [1], Proposition 1.15. Concerning the regularity of the solution, the $C^{2,\alpha}$ regularity on Ω follows by standard results for quasi-linear elliptic equations [29] and the Hölder continuity on \mathcal{N} follows by [2].

We conclude this section with a remark concerning the Yamabe problem on (\mathcal{N}^4, h) . We claim that the Yamabe constant of (\mathcal{N}^4, h) ,

$$\mathcal{Y}(\mathcal{N}, [h]) = \inf_{\{\psi \in W^{1,2}(\mathcal{N}); \psi \neq 0\}} \frac{\int_{\mathcal{N}} |\nabla^h \psi|^2 + 2\psi^2 d\mu_h}{(\int_{\mathcal{N}} \psi^4 d\mu_h)^{1/2}}$$

is positive. In fact, suppose by contradiction that $\mathcal{Y}(\mathcal{N}, [h]) \leq 0$. Since the local Sobolev constant of (\mathcal{N}^4, h) is positive, Theorems 1.12 and 1.15 of [1] implies the existence of a positive function $\psi_0 \in L^\infty(\mathcal{N}) \cap W^{1,2}(\mathcal{N})$ that satisfies $-\Delta_h \psi_0 + 2\psi_0 = \mathcal{Y}(\mathcal{N}, [h])\psi_0^3$ weakly on \mathcal{N} . Thus, $\psi_0 > 0$ satisfies an equation of the form $-\Delta_h \psi_0 + G\psi_0 = 0$ for some positive function G and we have $\int_{\mathcal{N}} |\nabla^h \psi_0|_h^2 d\mu_h = -\int_{\mathcal{N}} G\psi_0^2 d\mu_h < 0$, a contradiction.

In particular, the solution ϕ obtained in Theorem 17 is such that the integral of the scalar curvature of the metric $\hat{h} = \phi^2 h$ is positive, since by the definition of the Yamabe constant, $\int_{\mathcal{N}} \hat{R} d\mu_{\hat{h}} \geq \mathcal{Y}(\mathcal{N}, [h])\|\phi\|_{L^4(\mathcal{N})}^2 > 0$. This remark will be used in the next section (see Proposition 21).

9. Proof of Theorem C

We start with two propositions that are the key results needed for the proof of Theorem C. The first one is a version of the Gauss–Bonnet–Chern formula.

Lemma 20. *Let (\mathcal{N}^4, h) be an oriented associated singular Einstein manifold to a compact static triple (M^3, g, V) with positive scalar curvature.*

Given $u \in C^4(M) \cap C^\infty(\text{int}(M))$, consider the Riemannian metric $\tilde{h} = u^2 h = u^2(V^2 d\theta^2 + g)$ on Ω . Then the singular space $(\mathcal{N}^4, \tilde{h})$ is such that

$$(35) \quad \frac{1}{48} \int_{\mathcal{N}} \tilde{R}^2 d\mu_{\tilde{h}} + \int_{\mathcal{N}} |\tilde{W}^+|_{\tilde{h}}^2 d\mu_{\tilde{h}} = 8\pi^2 \sum_{i=1}^r k_i + \frac{1}{4} \int_{\mathcal{N}} |\overset{\circ}{\text{Ric}}|_{\tilde{h}}^2 d\mu_{\tilde{h}},$$

where k_i denotes the constant value of $|\nabla V|$ on a connected component of $\partial M = \partial_1 M \cup \dots \cup \partial_r M$ for all $i = 1, \dots, r$.

Proof. First, we study the behavior of the metric $\tilde{h} = u^2(V^2 d\theta^2 + g)$ near its singular set ∂M . In order to do this, we introduce Fermi coordinates on M with respect to the metric $\tilde{g} = u^2 g$. Since by hypothesis this is a C^4 metric on M , the normal exponential map of ∂M is only of class C^3 .

Let s and \tilde{s} denote the distance function to the boundary ∂M in (M^3, g) and (M^3, \tilde{g}) , respectively. Let $x = (x^2, x^3)$ be local coordinates in ∂M . We have $\tilde{s}(s, x) = u(0, x)s + O(s^2)$. Since $V(0, x) = 0$ and $\partial_s V(0, x) = k_i$ at a point $(0, x) = (0, x^2, x^3)$ belonging to the connected component $\partial_i M$ (Proposition 7), we also have

$$(uV)(\tilde{s}, x^2, x^3) = (u(0, x) + O(\tilde{s}))(\partial_{\tilde{s}} V(0, x^2, x^3)\tilde{s} + O(\tilde{s}^2)) = k_i \tilde{s} + O(\tilde{s}^2).$$

Using the local coordinates $(x^0 = \theta, x^1 = \tilde{s}, x^2, x^3)$ near $\partial_i M$, we conclude that $\tilde{h} = (uV)^2 d\theta^2 + \tilde{g}$ has the expansion

$$\begin{aligned} \tilde{h}_{ij}(\tilde{s}, x^2, x^3) &= d\tilde{s}^2 + k_i^2 \tilde{s}^2 d\theta^2 + \sum_{i,j=2}^3 \tilde{g}_{ij}(0, x^2, x^3) dx^i dx^j + \\ &+ 2 \sum_{i,j=2}^3 \tilde{A}_{ij}(0, x^2, x^3) \tilde{s} dx^i dx^j + \sum_{i,j=2}^3 \tilde{G}_{ij}(0, x^2, x^3) \tilde{s}^2 dx^i dx^j + E, \end{aligned}$$

where $\tilde{A}(0, x^2, x^3)$ is the second fundamental form of ∂M in $(M^3, \tilde{g} = u^2 g)$ at $x = (0, x^2, x^3)$. Moreover, the remainder E satisfies $|E| \leq C\tilde{s}^3$ for some constant $C > 0$ and $E(\partial_{\tilde{s}}, \partial_{x^A}) = 0$ for all $A = 0, 1, 2, 3$.

Thus, \tilde{h} is singular in the sense of [43]. Accordingly to [43] (compare with [8]), a Gauss–Bonnet–Chern formula holds for the singular space $(\mathcal{N}^4, \tilde{h})$ and is given by

$$\begin{aligned} \frac{1}{48} \int_{\mathcal{N}} \tilde{R}^2 d\mu_{\tilde{h}} + \frac{1}{8} \int_{\mathcal{N}} |\tilde{W}^+|_{\tilde{h}}^2 d\mu_{\tilde{h}} - \frac{1}{4} \int_{\mathcal{N}} |\overset{\circ}{\text{Ric}}|_{\tilde{h}}^2 d\mu_{\tilde{h}} &= \\ &= 4\pi^2(\chi(\mathcal{N}) - \chi(\partial M)) + \sum_{i=1}^r k_i \chi(\partial_i M). \end{aligned}$$

In the above formula $|\tilde{W}|_{\tilde{h}}$, denotes the norm of the Weyl tensor \tilde{W} . Considering the self-dual/anti-self-dual decomposition $\tilde{W}^+ + \tilde{W}^-$ of \tilde{W} viewed as a symmetric endomorphism of the space of two-forms, we have $|\tilde{W}|_{\tilde{h}}^2 = 4(|\tilde{W}^+|_{\tilde{h}}^2 + |\tilde{W}^-|_{\tilde{h}}^2)$.

As observed in Section 6, \mathcal{N}^4 and the fixed point set of the S^1 action on \mathcal{N} , i.e., ∂M , have the same Euler characteristic. Moreover, since the map $\theta \mapsto -\theta$ is an orientation inverting isometry of $(\mathcal{N}^4, \tilde{h})$, $|\tilde{W}|_{\tilde{h}}^2 = 8|\tilde{W}^+|_{\tilde{h}}^2$. By Theorem B, *iii*), each connected component of ∂M is a two-sphere. Formula (35) follows. q.e.d.

REMARK 9. The associated Einstein manifold $(\mathcal{N}^4, h = V^2 d\theta^2 + g)$ itself is such that

$$\frac{1}{48} \int_{\mathcal{N}} R_{\tilde{h}}^2 d\mu_h = 2\pi \sum_{i=1}^r k_i |\partial_i M|, \quad \int_{\mathcal{N}} |W^+|_{\tilde{h}}^2 d\mu_h = 2\pi \int_M |Ric_g|_g^2 V d\mu_g.$$

Thus, formula (35) for (\mathcal{N}^4, h) is precisely the fundamental formula (9) for (M^3, g, V) (see Proposition 6). This curious observation can be used to give a different proof of the above lemma using only the conformal invariance of the Weyl term in (35) and the conformal transformations laws of the Q -curvature (for the definition and properties of the Q -curvature, see, for example, [32]).

The second key proposition is a consequence of the existence result Theorem 17 (compare with [31], Proposition 3.5).

Proposition 21. *Let (\mathcal{N}^4, h) be the associated singular Einstein manifold to a compact static triple (M^3, g, V) with positive scalar curvature. Assume that (M^3, g, V) is not equivalent to the standard hemisphere.*

Then

$$\inf_{\substack{\tilde{h}=u^2h, u>0 \\ u \in C^\infty(M)}} \int_{\mathcal{N}} \tilde{R}^2 d\mu_{\tilde{h}} \leq 24 \int_{\mathcal{N}} |W^+|_{\tilde{h}}^2 d\mu_h.$$

Moreover, assume the equality holds. Let ϕ be the solution to the modified Yamabe problem given in Theorem 17. Then $\phi \in C^\infty(M)$ and the metric $\hat{h} = \phi^2 h$ on Ω is such that $\tilde{R} = 2\sqrt{6}|\tilde{W}^+|_{\tilde{h}}$ is a positive constant.

Proof. By Theorem 17, there exists a Riemannian metric $\hat{h} = \phi^2 h$ on $\Omega = \mathcal{N} \setminus \partial M$, where $\phi \in C^0(M) \cap C^{2,\alpha}(\text{int}(M))$ is positive, such that

$$\hat{R} - 2\sqrt{6}|\hat{W}^+|_{\hat{h}}$$

is a constant $\lambda \leq 0$. Multiplying by \hat{R} and integrating on \mathcal{N} , we obtain

$$\int_{\mathcal{N}} \hat{R}^2 d\mu_{\hat{h}} = 2\sqrt{6} \int_{\mathcal{N}} \hat{R} |\hat{W}^+|_{\hat{h}} d\mu_{\hat{h}} + \lambda \int_{\mathcal{N}} \hat{R} d\mu_{\hat{h}}$$

$$\leq 2\sqrt{6}\left(\int_{\mathcal{N}} \hat{R}^2 d\mu_{\hat{h}}\right)^{1/2} \left(\int_{\mathcal{N}} |\hat{W}^+|_{\hat{h}}^2 d\mu_{\hat{h}}\right)^{1/2} + \lambda \int_{\mathcal{N}} \hat{R} d\mu_{\hat{h}}.$$

As remarked in the end of the previous section, $\int_{\mathcal{N}} \hat{R} d\mu_{\hat{h}} > 0$. Therefore,

$$\int_{\mathcal{N}} \hat{R}^2 d\mu_{\hat{h}} \leq 24 \int_{\mathcal{N}} |\hat{W}^+|_{\hat{h}}^2 d\mu_{\hat{h}} = \int_{\mathcal{N}} |W^+|_h^2 d\mu_h,$$

and equality holds if and only if $\lambda = 0$. The last equality follows from the conformal invariance of $|W^+|_h^2 d\mu_h$ in dimension four.

Since an approximation argument implies that

$$\inf_{\substack{\hat{h}=u^2h, u>0 \\ u \in C^\infty(M)}} \int_{\mathcal{N}} \tilde{R}^2 d\mu_{\tilde{h}} = \inf_{\substack{\tilde{h}=u^2h, u>0 \\ u \in C^0(M) \cap C^{2,\alpha}(\text{int}(M))}} \int_{\mathcal{N}} \tilde{R}^2 d\mu_{\tilde{h}},$$

we conclude that

$$\inf_{\substack{\tilde{h}=u^2h, u>0 \\ u \in C^\infty(M)}} \int_{\mathcal{N}} \tilde{R}^2 d\mu_{\tilde{h}} \leq \int_{\mathcal{N}} \hat{R}^2 d\mu_{\hat{h}} \leq \int_{\mathcal{N}} |W^+|_h^2 d\mu_h.$$

If equality holds, then $\hat{R} = 2\sqrt{6}|\hat{W}^+|_{\hat{h}}$ and the infimum of the functional $\hat{h} \mapsto \int_{\mathcal{N}} \hat{R}^2 d\mu_{\hat{h}}$ on the set of metrics $\{\hat{h} = \psi^2 h; \psi \in C^0(M) \cap C^{2,\alpha}(\text{int}(M)), \psi > 0\}$ is attained by ϕ . Considering variations $\hat{h}_t = (\phi + t\phi^{-1}v)^2 h$, $v \in C_0^\infty(\text{int}(M))$, one can show that \hat{R} must be constant (see [10], 4.67). Therefore, $|\hat{W}^+|_{\hat{h}}$ is constant and ϕ satisfies the equation

$$-6\text{div}_g(V\phi) + 12V\phi = \hat{R}V\phi^3.$$

In Appendix C, we argue that ϕ is smooth up to the boundary. The proposition follows. q.e.d.

We have now all ingredients we need to prove Theorem C.

Theorem 22. *Let (M^3, g, V) be a compact simply connected static triple with scalar curvature 6. One of the following alternatives holds:*

- i) (M^3, g, V) is equivalent to the standard hemisphere; or
- ii) (M^3, g, V) is equivalent to the standard cylinder; or
- iii) Denoting by $\partial M = \partial_1 M \cup \dots \cup \partial_r M$ the connected components of the boundary and by k_i the value of $|\nabla V|$ on $\partial_i M$, the following inequality holds:

$$\sum_{i=1}^r k_i |\partial_i M| < \frac{4\pi}{3} \sum_{i=1}^r k_i.$$

Proof. Let (\mathcal{N}^4, h) be the associated singular Einstein manifold to (M^3, g, V) . Assume (M^3, g, V) is not equivalent to the standard hemisphere. As a consequence of Proposition 21, there exists a positive

function $u \in C^\infty(M)$ such that $(\mathcal{N}^4, \tilde{h} = u^2 h)$ satisfies

$$\frac{1}{24} \int_{\mathcal{N}} \tilde{R}^2 d\mu_{\tilde{h}} \leq \int_{\mathcal{N}} |\tilde{W}^+|_{\tilde{h}}^2 d\mu_{\tilde{h}} = \int_{\mathcal{N}} |W^+|_h^2 d\mu_h.$$

Moreover, we can assume equality holds if and only if $\tilde{R} = 2\sqrt{6}|\tilde{W}^+|_{\tilde{h}}$ is a positive constant.

Combining the above inequality with the Gauss–Bonnet–Chern formula (35) for $(\mathcal{N}^4, \tilde{h})$, we conclude that

$$(36) \quad \frac{3}{2} \int_{\mathcal{N}} |W^+|_h^2 d\mu_h \geq 8\pi^2 \sum_{i=1}^r k_i + \frac{1}{4} \int_M |\overset{\circ}{Ric}|_{\tilde{h}}^2 d\mu_{\tilde{h}} \geq 8\pi^2 \sum_{i=1}^r k_i.$$

Now, $(\mathcal{N}^4, h = V^2 d\theta^2 + g)$ satisfies

$$\int_{\mathcal{N}} |W^+|_h^2 d\mu_h = 2\pi \int_M |\overset{\circ}{Ric}_g|_g^2 V d\mu_g.$$

By the fundamental formula (9) for (M^3, g, V) (see Proposition 6 and Theorem B, *iii*), we also have

$$\sum_{i=1}^r k_i |\partial_i M| + \int_M |\overset{\circ}{Ric}_g|_g^2 V d\mu_g = 4\pi \sum_{r=1}^r k_i.$$

Therefore,

$$\sum_{i=1}^r k_i |\partial_i M| \leq \frac{4\pi}{3} \sum_{i=1}^r k_i.$$

If equality holds above, the conformally related metric $\tilde{h} = u^2 h$ must be Einstein (by (36)) and $\tilde{R} = 2\sqrt{6}|\tilde{W}^+|_{\tilde{h}}$ must be a positive constant. As shown in [31], using the formulas relating the divergence of the Weyl tensors of conformally related metrics one can then conclude that u must be constant. Thus, the metric $h = V^2 d\theta^2 + g$ also satisfies $R_h = 2\sqrt{6}|W^+|_h$, that is, $|\overset{\circ}{Ric}_g|_g = \sqrt{6}$. And then Theorem A implies that (M^3, g, V) is the standard cylinder. q.e.d.

Appendix A. An inequality for transverse trace-free tensors

We prove a generalization of the refined Kato inequality for Codazzi tensors on three-manifolds (see [15]). The proof is similar to [61], Lemma 5.1. We hope this inequality will be useful in other contexts.

Proposition 23. *Let (M^3, g) be a Riemannian three-manifold and T a symmetric 2-tensor that has constant trace and zero divergence. Denote by C the 3-tensor defined by*

$$C(X, Y, Z) = \nabla T(X, Y, Z) - \nabla T(X, Z, Y) \quad \text{for all } X, Y, Z \in \mathcal{X}(M).$$

Then, at every $p \in M$ such that $|T|(p) \neq 0$,

$$|\nabla|T||^2 \leq \frac{3}{5} \left(|\nabla T|^2 + \frac{|C|^2}{2} \right).$$

Proof. Let e_1, e_2, e_3 be a local orthonormal basis that diagonalizes T at a point $p \in M$ where $|T| \neq 0$. Then $|\nabla|T|^2| = 4|T|^2|\nabla|T|^2|$ implies

$$|\nabla|T|^2|^2 = \frac{1}{|T|^2} \sum_{k=1}^3 \left(\sum_{i,j=1}^3 T_{ij} T_{ij;k} \right)^2 = \frac{1}{|T|^2} \sum_{k=1}^3 \left(\sum_{i=1}^3 T_{ii} T_{ii;k} \right)^2.$$

By Cauchy–Schwartz inequality,

$$(37) \quad |\nabla|T|^2|^2 \leq \frac{1}{|T|^2} \sum_{k=1}^3 \left(\sum_{i=1}^3 T_{ii}^2 \right) \left(\sum_{i=1}^3 T_{ii;k}^2 \right) = \sum_{i,k=1}^3 T_{ii;k}^2.$$

Let $A, B, C \in \{1, 2, 3\}$ denote different indexes. Then

$$(38) \quad \begin{aligned} C_{BBA} + C_{CCA} &= \sum_{i=1}^3 C_{iiA} = \sum_{i=1}^3 T_{ii;A} - \sum_{i=1}^3 T_{iA;i} \\ &= d(\text{tr}(T))_A - \text{div}(T)_A = 0, \end{aligned}$$

by hypothesis. This implies

$$\begin{aligned} T_{AA;A} &= -T_{BB;A} - T_{CC;A} = -T_{BA;B} - T_{CA;C}, \\ T_{BB;A} &= T_{BA;B} + C_{BBA}, \\ T_{CC;A} &= T_{CA;C} + C_{CCA} = T_{CA;C} - C_{BBA}. \end{aligned}$$

Thus, rewriting inequality (37) as

$$|\nabla|T|^2|^2 \leq \sum_{i=1}^3 (T_{11;i}^2 + T_{22;i}^2 + T_{33;i}^2),$$

each of the above summand is of the form $(a+b)^2 + (a-c)^2 + (b+c)^2$ for some $a, b, c \in \mathbb{R}$. Because of the algebraic identity

$$(a+b)^2 + (a-c)^2 + (b+c)^2 + (a-b+c)^2 = 3(a^2 + b^2 + c^2),$$

we then conclude that

$$(39) \quad \begin{aligned} |\nabla|T|^2|^2 &\leq 3(T_{21;2}^2 + T_{31;3}^2 + C_{221}^2) + 3(T_{12;1}^2 + C_{112}^2 + T_{32;3}^2) \\ &\quad + 3(C_{223}^2 + T_{13;1}^2 + T_{23;2}^2) \\ &= 3 \sum_{i \neq k} T_{ik;i}^2 + \frac{3}{2} ((C_{221}^2 + C_{331}^2) + (C_{112}^2 + C_{332}^2) + (C_{223}^2 + C_{113}^2)) \\ &= \frac{3}{2} \left(\sum_{i \neq k} T_{ik;i}^2 + \sum_{i \neq k} T_{ki;i}^2 + \sum_{i \neq k} C_{ik}^2 \right), \\ &= \frac{3}{2} \left(\sum_{i \neq k} T_{ik;i}^2 + \sum_{i \neq k} T_{ki;i}^2 + \frac{1}{2} \left(\sum_{i \neq k} C_{ik}^2 + \sum_{i \neq k} C_{ki}^2 \right) \right), \end{aligned}$$

where in the last three lines we used (38), the symmetry of T and the antisymmetry of C in the last two variables. The proposition follows by combining inequalities (37) and (39). q.e.d.

Appendix B. Solutions to an equation of type $HessV = 0$

Let (M^n, g) be a Riemannian manifold admitting a non-constant solution to the equation $HessV = (\Delta V/n)g$. These manifolds have been classified by Kühnel [38]. In this section, we consider a particular case of this equation on compact surfaces (Σ^2, g) ,

$$HessV = \frac{1}{2}(K - 3\epsilon)Vg,$$

where K is the Gaussian curvature of (Σ, g) and $\epsilon \in \{-1, 0, 1\}$. It has appeared in the proof of Proposition 27 in Section 8.

For completeness, we prove here the classification of positive solutions of this equation. The argument follows [24], Lemma 3.

Proposition 24. *Let (Σ^2, g) be a closed surface with Gaussian curvature K and $\epsilon \in \{-1, 0, 1\}$. If there exists a positive solution $V \in C^\infty(M)$ of the equation*

$$(40) \quad Hess_\Sigma V = \frac{1}{2}(K - 3\epsilon)Vg,$$

then V is constant and Σ has constant Gaussian curvature $K = 3\epsilon$.

Proof. Taking the divergence of (40), we obtain the following integrability condition:

$$(41) \quad d((K - \epsilon)V^3) = 0.$$

Therefore, $K = cV^{-3} + \epsilon$ for some constant c . Integrating the trace of (40), we have that c is positive, zero or negative if ϵ equals 1, 0 or -1 , respectively. We analyze each case separately.

If $\epsilon = 0$, then $c = 0$ and, therefore, $K \equiv 0$ and $HessV = 0$. The theorem follows.

If $\epsilon = -1$, then $c < 0$ and, therefore, $K < -1$. Using this and the integrability condition, we conclude that a point $p \in \Sigma$ is a critical point of K if and only if it is a critical point of V . Moreover,

$$\Delta K = c \left(\frac{12}{V^5} |\nabla V|^2 - \frac{3}{V^4} \Delta V \right) = c \left(\frac{12}{V^5} |\nabla V|^2 - \frac{3}{V^3} (K + 3) \right).$$

Hence, at a point $p_0 \in \Sigma$ of absolute minimum of K ,

$$0 \leq \Delta K(p_0) = -\frac{3c}{V^3(p_0)}(K(p_0) + 3) \Rightarrow K(p_0) \geq -3,$$

since $c < 0$ and V is a positive function. Similarly, at a point $p_1 \in \Sigma$ of absolute maximum of K , $K(p_1) \leq -3$. Therefore, $K \equiv -3$ and $HessV = 0$. The result follows.

If $\epsilon = 1$, then $c > 0$ and, therefore, $K > 1$. We can assume Σ is homeomorphic to the sphere S^2 .

Suppose, by contradiction, that V is not constant. By [38], the existence of a non-constant solution V to the equation

$$\mathring{H}essV = 0,$$

in a closed surface implies that V has precisely two non-degenerated critical points (the point of minimum p_0 and the point of maximum p_1), that V depends only on the distance to a critical point, and that the metric on $\Sigma \setminus \{p_0, p_1\}$ can be written in the warped-product form

$$g = du^2 + \left(\frac{V'(u)}{V''(0)} \right)^2 d\theta^2,$$

where $u \in (0, u_*)$ is the distance to p_0 and θ is a 2π -periodic variable (see [38]). V is an increasing function of u defined on the interval $[0, u_*]$.

For the above metric, $Hess_g V = V''g$, so equation (40) is equivalent to

$$(42) \quad 2V'' = (K - 3)V = \frac{c}{V^2} - 2V.$$

This second order differential equation can be expressed in the Hamiltonian form using the energy

$$H(x, y) = y^2 + \frac{c}{x} + x^2,$$

for $x = V > 0$, $y = V'$. Sketching the level sets, one can verify that the constant solution of (42) corresponds to the point $(x, y) = ((c/2)^{1/3}, 0)$ and that the solution V of (42) originated in our problem is such that $V(0) < (c/2)^{1/3} < V(u_*)$.

Now we use the Gauss–Bonnet theorem and the explicit expression of the metric to compute:

$$4\pi = \int_{\Sigma} K d\Sigma = \int_0^{2\pi} \int_0^{u_*} \left[1 + \frac{c}{V^3} \right] \frac{V'}{V''(0)} dud\theta = \frac{2\pi}{V''(0)} \left[V - \frac{c}{2V^2} \right]_0^{u_*}.$$

By (42), we conclude that $c = 2V^3(u_*)$. And this is a contradiction. q.e.d.

Proposition 25. *Let (Σ^2, g) be a compact surface with non-empty boundary. Let K denote its Gaussian curvature and let $\epsilon \in \{-1, 0, 1\}$. Assume that there exists a non-trivial non-negative solution $V \in C^\infty(M)$ to the problem*

$$Hess_{\Sigma} V = \frac{1}{2}(K - 3\epsilon)Vg \quad \text{in } \Sigma, \quad V = 0 \quad \text{in } \partial\Sigma.$$

Then $\epsilon = 1$, Σ is isometric to the round hemisphere (S_+^2, g_{can}) with constant Gaussian curvature $K = 1$ and V is a static potential on (Σ^2, g) .

Proof. The integrability condition (41) and the hypothesis that V vanishes on $\partial\Sigma$ implies that $K = \epsilon$ on Σ . Taking the trace of equation (40), we obtain

$$(43) \quad \Delta V + 2\epsilon V = 0.$$

Hence, multiplying by V and integrating by parts on Σ using that V vanishes on $\partial\Sigma$ one concludes that ϵ must be positive. Therefore, $V \neq 0$ satisfies the static equations (40) and (43), since (Σ^2, g) has constant Gaussian curvature $K = \epsilon = 1$. q.e.d.

Appendix C. Regularity of solutions to a degenerated elliptic problem

Given (M^3, g, V) a compact orientable static triple, let (\mathcal{N}^4, h) be its associated singular Einstein manifold. We want to study the regularity properties of a positive solution $\phi \in C^{0,\mu}(\mathcal{N}) \cap W^{1,2}(\mathcal{N}) \cap C^{2,\alpha}(\text{int}(M))$ to the Yamabe equation

$$(44) \quad -\Delta_h \phi + 2\phi = 2\phi^3,$$

on the singular space (\mathcal{N}^4, h) . Equation (44) means that the conformally related metric $\hat{h} = \phi^2 h$ on $\mathcal{N} \setminus \partial M$ has constant scalar curvature 12.

The first observation is that $\phi \in C^\infty(\text{int}(M))$ by standard elliptic regularity [29]. The problem is to understand the behavior of ϕ near ∂M . In order to analyze this point, we first deduce some formulas in an appropriated coordinate system near ∂M .

Let $(x^0 = s, x^1, x^2)$ be Fermi coordinates on a small tubular neighborhood $\mathcal{U} = \{p \in M; s(p) = d_{\partial M}(p) < \epsilon\}$ of ∂M , where (x^1, x^2) are local coordinates on ∂M . We write $M_s = \{p \in M; d_{\partial M}(p) = s\}$ and use the subscript s to denote the geometric quantities related to M_s in (M, g) . Using the fact that ∂M is totally geodesic in (M, g) and the expansion of V near ∂M (Proposition 7), we conclude that the Laplacian of the metric $h = V^2 d\theta^2 + g$ takes the form:

$$\begin{aligned} \Delta_h \phi &= \Delta_g \phi + \frac{1}{V} g(\nabla^g V, \nabla^g \phi) \\ &= (\partial_s^2 \phi + \Delta_s \phi + H_s \partial_s \phi) + \frac{\partial_s V}{V} \partial_s \phi + \frac{1}{V} g^{ij} \partial_i V \partial_j \phi \\ &= \partial_s^2 \phi + \frac{1}{s} \partial_s \phi + \Delta_0 \phi + Q(\phi), \end{aligned}$$

where, for indexes i, j running through $\{1, 2\}$ and indexes A through $\{0, 1, 2\}$,

$$\begin{aligned} \Delta_0 \phi &= g^{ij}(0, x)(\partial_i \partial_j \phi - \Gamma_{ij}^k(0, x) \partial_k \phi), \quad \text{and} \\ Q(\phi) &= Q^{ij} s^2 \partial_i \partial_j \phi + Q^A s \partial_A \phi \quad \text{for } Q^{i,j}, Q^A \in C^\infty([0, \epsilon) \times \partial M). \end{aligned}$$

Hence,

$$(45) \quad \partial_s \phi^2 + \frac{1}{s} \partial_s \phi + \Delta_0 \phi + Q(\phi) = 2\phi - 2\phi^3.$$

Now, considering the metric $\tilde{h} = s^{-2}h$ for some smooth positive function s on $M \setminus \partial M$ that coincides with the distance function to ∂M in \mathcal{U} , the function $\eta = s\phi$ satisfies

$$(46) \quad -\Delta_{\tilde{h}} \eta + \tilde{R} \eta = 2\eta^3.$$

In the above equation, \tilde{R} denotes the scalar curvature of (\mathcal{N}, \tilde{h}) . Using the formulas relating the geometric quantities associated to conformal metrics (see [10]), one observes that $\Delta_{\tilde{h}} \eta = s^2 \Delta_h \eta - 2ds(\nabla^h \eta)$ and that $\tilde{R} \in C^\infty(M)$ satisfies $\tilde{R} = -12 + O(s^2)$ as s goes to zero.

Since η is invariant under the S^1 action on (\mathcal{N}^4, h) , this equation reduces to a uniformly degenerated elliptic equation (in the sense of [30]) on the conformally compact manifold $(\bar{M} = M \setminus \partial M, \bar{g} = s^{-2}g)$. In fact, by the above formulas, we have

$$s^2 \partial_s \eta - s \partial_s \eta + s^2 \Delta_s \eta + \tilde{Q}^A s \partial_A \eta = \frac{1}{6} \hat{R} \eta - 2\eta^3,$$

where $\Delta_s \eta = g^{ij}(\partial_i \partial_j \phi - \Gamma_{ij}^k \partial_k \phi)$ for indexes i, j running through the set $\{1, 2\}$ and $\tilde{Q}^A \in C^\infty([0, \epsilon) \times \partial M)$ satisfies $\tilde{Q}^A = O(s^2)$ as s goes to zero for all $A \in \{0, 1, 2\}$. Observing that $\eta = s\phi$, $\phi \in C^{0,\mu}(M) \cap C^\infty(\text{int}(M))$, Schauder type estimates for η implies that $\phi = s^{-1}\eta$ is such that $s^{|\alpha|} \partial_x^\alpha \phi$ is uniformly bounded for every multi-index α (see [30], Proposition 3.4).

Accordingly to [1], one can go further and show that *i)* ϕ is, in fact, conormal (i.e., the derivatives $s^i \partial_s^i \partial_{x^1}^j \partial_{x^2}^k \phi$ are bounded on M for all non-negative integers i, j, k); and *ii)* has an expansion in s as a sum involving terms like $a_{pq} s^p \log^q(s)$ where $a_{pq} \in C^\infty(\partial M)$. The reader will find general theories about this kind of degenerated elliptic equation developed in [45] (the ‘‘edge-calculus’’) and [6] (we also refer the reader to [7], where some theory is developed in a very concrete situation).

Using the information that ϕ has an expansion in s near ∂M , we compute it considering that ϕ satisfies

$$(47) \quad \partial_s^2 \phi + \frac{1}{s} \partial_s \phi = F,$$

for $F = F(\phi) = 2\phi - 2\phi^3 - \Delta_0 \phi - Q(\phi)$ smooth on $(0, \epsilon) \times \partial M$ and such that $s^k \partial_s^{(k)} F$ is bounded on M for all positive integer k .

For every $x \in \partial M$ and every $s \in (0, \epsilon)$, there exists $\zeta_s \in (0, \epsilon)$ such that $\phi(s, x) = \phi(0, x) + \partial_s \phi(\zeta_s, x)s$, by Taylor’s Theorem. Since ϕ is smooth on the interior of M , from this formula for $x \mapsto \phi(0, x)$ it follows that ϕ restricted to ∂M belongs to $C^\infty(\partial M)$.

We will compute the expansion of ϕ using the classical method of separation of variables. We are grateful to Andr e Neves for suggesting this approach.

Let $\lambda_0 = 0 < \lambda_1 \leq \lambda_2 \leq \dots$ denote the eigenvalues of the Laplacian Δ_0 on $(\partial M, g_{\partial M})$ and let $\{\xi_i\}$ be a L^2 -orthonormal basis of eigenfunctions,

$$\Delta_0 \xi_i + \lambda_i \xi_i = 0.$$

For each $s \in (0, \epsilon)$, we can write

$$\phi(s, x) = \sum_{i=0}^{+\infty} \alpha_i(s) \xi_i(x) \quad \text{where} \quad \alpha_i(s) = \int_{\partial M} \phi(s, x) \xi_i(x) d\mu_0(x).$$

The convergence of the above series is to be understood in the space $L^2(\partial M, d\mu_b)$. Since ϕ is smooth in the interior of M and continuous up to ∂M , we have that $\alpha_i \in C^0([0, \epsilon]) \cap C^\infty((0, \epsilon))$ and, in fact, for all $s \in (0, \epsilon)$,

$$\partial_s^{(k)} \phi(s, x) = \sum_{i=0}^{+\infty} \alpha_i^{(k)}(s) \xi_i(x) \quad \text{where} \quad \alpha_i^{(k)}(s) = \frac{d^k \alpha_i}{ds^k}(s).$$

Integrating (47) against each ξ_i for each fixed $s \in (0, \epsilon)$ we conclude that each $\alpha_i \in C^0([0, \epsilon]) \cap C^\infty((0, \epsilon))$ satisfies an equation of the form

$$(48) \quad \alpha'' + \frac{1}{s} \alpha' - \lambda \alpha = F,$$

for some constant λ and a function $F \in C^0([0, s]) \cap C^\infty((0, s))$ depending on ϕ . In fact, α_i satisfies (48) with $\lambda = \lambda_i + 2$ and

$$F(s) = F_i(\phi)(s) = \int_{\partial M} (-2\phi^3(s, x) - Q(\phi)(s, x)) \xi_i(x) d\mu_0(x).$$

Lemma 26. *Let k be a non-negative integer. If $\alpha \in C^0([0, \epsilon]) \cap C^\infty((0, \epsilon))$ is a solution to (48) for some $F \in C^k([0, \epsilon]) \cap C^\infty((0, \epsilon))$, then $\alpha \in C^{k+2}([0, \epsilon])$.*

Proof. We first explain the argument for $k = 0$. From (48), we have

$$(s\alpha')' = \lambda s\alpha + sF.$$

Integrating from $\delta > 0$ to $s < \epsilon$ we obtain

$$s\alpha'(s) - \delta\alpha'(\delta) = \lambda \int_{\delta}^s t\alpha(t) dt + \int_{\delta}^s tF(t) dt.$$

Since α is bounded, $\lim_{\delta \rightarrow 0} \delta\alpha'(\delta) = 0$ (otherwise, $\alpha(s)$ would grow like $|\log(s)|$ as s goes to zero). Since α and F are bounded, we can take the limit as δ goes to zero and conclude that, for all $s \in (0, \delta)$,

$$(49) \quad \alpha'(s) = \frac{\lambda}{s} \int_0^s t\alpha(t) dt + \frac{1}{s} \int_0^s tF(t) dt.$$

In particular, $\lim_{s \rightarrow 0} \alpha'(s) = 0$ because α and tF are continuous. Moreover, for every $0 < \delta < s < \epsilon$, we have

$$\frac{\alpha(s) - \alpha(\delta)}{s - \delta} = \frac{\lambda}{s - \delta} \int_{\delta}^s \frac{1}{u} \int_0^u t\alpha(t) dt du + \frac{1}{s - \delta} \int_{\delta}^s \frac{1}{u} \int_0^u tF(t) dt du,$$

so that taking the limit as δ goes to zero we obtain

$$\frac{\alpha(s) - \alpha(0)}{s} = \frac{\lambda}{s} \int_0^s \frac{1}{u} \int_0^u t\alpha(t) dt du + \frac{1}{s} \int_0^s \frac{1}{u} \int_0^u tF(t) dt du = O(s).$$

It follows that α is differentiable at the origin and satisfies $\alpha'(0) = 0$. Thus, $\alpha \in C^1([0, \epsilon))$.

For every $s \in (0, s)$, integration by parts now gives

$$\int_0^s t\alpha(t) dt = \frac{s^2}{2}\alpha(s) - \frac{1}{2} \int_0^s t^2\alpha'(t) dt.$$

We claim that a similar formula also holds for the integral of $tF(t)$. In fact, observe that $tF'(t)$ is bounded (as before, we have $\lim_{s \rightarrow 0} sF'(s) = 0$). For every $0 < \delta < s < \epsilon$, integration by parts gives

$$\int_\delta^s tF(t) dt = \left(\frac{s^2}{2}F(s) - \frac{\delta^2}{2}F(\delta) \right) - \frac{1}{2} \int_\delta^s t^2F'(t) dt.$$

Using that F is continuous and $tF'(t)$ is bounded, we can make δ go to zero to obtain

$$\int_0^s tF(t) dt = \frac{s^2}{2}F(s) - \frac{1}{2} \int_0^s t^2F'(t) dt,$$

as claimed.

Differentiating (49) at $s \in (0, \epsilon)$, we obtain

$$\begin{aligned} \alpha''(s) &= -\frac{\lambda}{s^2} \int_0^s t\alpha(t) dt + \lambda\alpha(s) - \frac{1}{s^2} \int_0^s tF(t) dt + F(s) \\ &= \frac{\lambda}{2}\alpha(s) + \frac{1}{2}F(s) + \frac{1}{2s^2} \int_0^s t^2(\lambda\alpha'(t) + F'(t)) dt. \end{aligned}$$

Observe that, as s goes to zero,

$$\left| \frac{1}{2s^2} \int_0^s (t^2(\lambda\alpha'(t) + F'(t))) dt \right| \leq C \left(\sup_{t \in (0, s)} |t\alpha'(t)| + \sup_{t \in (0, s)} |tF'(t)| \right) \rightarrow 0.$$

Therefore, by continuity of α and F , $\lim_{s \rightarrow 0} \alpha''(s) = (\lambda\alpha(0) + F(0))/2$ and, similarly, we obtain

$$\frac{\alpha'(s) - \alpha'(0)}{s} = \frac{\alpha'(s)}{s} = \frac{1}{2s} \int_0^s (\lambda\alpha(t) + F(t)) dt + o(1),$$

so that α is differentiable at $s = 0$ with $\alpha''(0) = (\lambda\alpha(0) + F(0))/2$. Thus, $\alpha \in C^2([0, \epsilon))$ as we wanted to prove.

For the general case, one can combine successive derivatives of (49) and integration by parts to prove by induction that for every $i = 0, 1, \dots, k+2$ and every $s \in (0, \epsilon)$,

$$\alpha^{(i)}(s) = \lambda \sum_{j=0}^{i-2} a_j \alpha^{(j)}(s) + \sum_{j=0}^{i-2} b_j F^{(j)}(s) + \frac{c_i}{t^i} \int_0^s t^i (\alpha^{(i-1)}(t) + F^{(i-1)}(t)) dt,$$

for some constants a_j, b_j and c_i . Moreover, the last term behaves as $o(1)$ as s goes to zero, $\alpha^{(i)}$ is differentiable at the origin and

$$\alpha^{(i)}(0) = \lim_{s \rightarrow 0} \alpha^{(i)}(s) = \lambda a_i \sum_{j=0}^{i-2} \alpha^{(j)}(0) + b_i \sum_{j=0}^{i-2} F^{(j)}(0),$$

for every $i = 0, 1, \dots, k+1$. The result follows. q.e.d.

To prove the claimed regularity of the solution $\phi \in C^0(M) \cap C^\infty(\text{int}(M))$ (see the end of Proposition 21), we would have to observe that for every non-negative integer k , $\alpha_i \in C^k([0, \epsilon))$ for all non-negative i implies that $F = F(\phi) \in C^k(M) \cap C^\infty(\text{int}M)$ and apply induction using the above lemma.

References

- [1] K. Akutagawa, G. Carron and R. Mazzeo, *The Yamabe problem on stratified spaces*, *Geom. Funct. Anal.* **24** (2014), no. 4, 1039–1079, MR 3248479, Zbl 1310.53034.
- [2] K. Akutagawa, G. Carron and R. Mazzeo, *Hölder regularity of solutions for Schrödinger operators on stratified spaces*, *J. Funct. Anal.* **269** (2015), no. 3, 815–840, MR3350731, Zbl 1316.35060.
- [3] B. Ammann, R. Lauter and V. Nistor, *On the geometry of Riemannian manifolds with a Lie structure at infinity*, *Int. J. Math. Math. Sci.* (2004), nos 1–4, 161–193, MR 2038804, Zbl 1297.01042.
- [4] M. Anderson, *On the structure of solutions to the static vacuum Einstein equations*, *Ann. Henri Poincaré* **1** (2000), no. 6, 995–1042, MR 1809792, Zbl 0981.83013.
- [5] M. Anderson, P. Chruściel and E. Delay, *Non-trivial, static, geodesically complete, vacuum space-times with a negative cosmological constant*, *J. High Energy Phys.* (2002), no. 10, 063, 27 pp., MR 1951922, Zbl 1075.83031.
- [6] L. Andersson and P. Chruściel, *Solutions of the constraint equations in general relativity satisfying “hyperboloidal boundary conditions”*, *Dissertationes Math. (Rozprawy Mat.)* **355** (1996), 100 pp, MR 1405962, Zbl 0873.35101.
- [7] L. Andersson, P. Chruściel and H. Friedrich, *On the regularity of solutions to the Yamabe equation and the existence of smooth hyperboloidal initial data for Einstein’s field equations*, *Comm. Math. Phys.* **149** (1992), no. 3, 587–612, MR 1186044, Zbl 0764.53027.
- [8] M. Atiyah and C. LeBrun, *Curvature, cones and characteristic numbers*, *Math. Proc. Cambridge Philos. Soc.* **155** (2013), no. 1, 13–37, MR 3065256, Zbl 1273.53041.
- [9] T. Aubin, *Some nonlinear problems in Riemannian geometry*, *Springer Monographs in Mathematics*, Springer-Verlag, Berlin, 1998, MR 1636569, Zbl 0896.53003.
- [10] A. Besse, *Einstein manifolds*, Springer-Verlag, Berlin, 1987, MR 2371700, Zbl 1147.53001.
- [11] W. Boucher and G. Gibbons, *Cosmic baldness*, in *The very early universe*, (Cambridge, 1982), 273–278, Cambridge Univ. Press, Cambridge, 1983, MR 0746321 (arXiv:1109.3535).

- [12] W. Boucher, G. Gibbons and G. Horowitz, *Uniqueness theorem for anti-de Sitter spacetime*, Phys. Rev. D (3) **30** (1984), no. 12, 2447–2451, MR 0771446.
- [13] V. Bour and G. Carron, *Optimal Integral Pinching Results*, Ann. Sci. Éc. Norm. Supér. (4) **48** (2015), no. 1, 41–70, MR 3335838, Zbl 1317.58021.
- [14] J.-P. Bourguignon, *Une stratification de l'espace des structures riemanniennes*, Compositio Math. **30** (1975), 1–41, MR 0418147, Zbl 0301.58015.
- [15] J.-P. Bourguignon, *The “magic” of Weitzenböck formulas*, Variational Methods Paris (1998), Progress in Nonlinear Differential Equations and Applications IV, Birkhäuser, 1990, MR 1205158, Zbl 0774.35003.
- [16] S. Brendle, *Constant mean curvature surfaces in warped product manifolds*, Publ. Math. Inst. Hautes Études Sci. **117** (2013), 247–269, MR 3090261, Zbl 1273.53052.
- [17] S. Brendle, F. Marques and A. Neves, *Deformations of the hemisphere that increase scalar curvature*, Invent. Math. **185** (2011), no. 1, 175–197, MR 2810799, Zbl 1227.53048.
- [18] G. Bunting and A. Masood-ul-Alam, *Nonexistence of multiple black holes in asymptotically Euclidean static vacuum space-time*, Gen. Relativ. Gravitation **19** (1987), 147–154, MR 0876598, Zbl 0615.53055.
- [19] X. Chen, C. LeBrun and B. Weber, *On conformally Kähler, Einstein manifolds*, J. Amer. Math. Soc. **21** (2008), no. 4, 1137–1168, MR 2425183, Zbl 1208.53072.
- [20] P. Chruściel and W. Simon, *Towards the classification of static vacuum spacetimes with negative cosmological constant*, J. Math. Phys. **42** (2001), no. 4, 1779–1817, MR 1820431, Zbl 1009.83009.
- [21] J. Corvino, *Scalar curvature deformation and a gluing construction for the Einstein constraint equations*, Comm. Math. Phys. **214** (2000), no. 1, 137–189, MR 1794269, Zbl 1031.53064.
- [22] M. Eichmair, *The size of isoperimetric surfaces in 3-manifolds and a rigidity result for the upper hemisphere*, Proc. Amer. Math. Soc. **137** (2009), no. 8, 2733–2740, MR 2497486, Zbl 1031.53064.
- [23] A. Fischer and J. Marsden, *Linearization stability of nonlinear partial differential equations*, Proc. Symp. Pure Math. **27** (1975), 219–262, MR 0383456, Zbl 0314.53031.
- [24] G. Galloway, *On the topology of black holes*, Comm. Math. Phys. **151** (1993), no. 1, 53–66, MR 1201655, Zbl 0314.53031.
- [25] G. Galloway and P. Miao, *Variational and rigidity properties of static potentials*, arXiv:1412.1062.
- [26] G. Galloway, S. Surya and E. Woolgar, *On the geometry and mass of static, asymptotically AdS spacetimes, and the uniqueness of the AdS soliton*, Comm. Math. Phys., **241** (2003), no. 1, 1–25, MR 2013749, Zbl 1149.83309.
- [27] G. Galloway and E. Wolgar, *On static Poincaré–Einstein metrics*, J. High Energy Phys. (2015), no. 6, 051, front matter+17 pp.
- [28] G. Gibbons and S. Hawking, *Classification of gravitational instanton symmetries*, Comm. Math. Phys. **66** (1979), no. 3, 291–310, MR 0535152.
- [29] D. Gilbarg and N. Trudinger, *Elliptic partial differential equations of second order*, Classics in Mathematics. Springer-Verlag, Berlin, 2001, MR 1814364, Zbl 1042.35002.

- [30] C. Graham and J. Lee, *Einstein metrics with prescribed conformal infinity on the ball*, Adv. Math. **87** (1991), no. 2, 186–225, MR 1112625, Zbl 0765.53034.
- [31] M. Gursky, *Four-manifolds with $\delta W_+ = 0$ and Einstein constants of the sphere*, Math. Ann. **318** (2000), no. 3, 417–431, MR 1800764, Zbl 1034.53032.
- [32] M. Gursky, *The principal eigenvalue of a conformally invariant differential operator, with an application to semilinear elliptic PDE*, Comm. Math. Phys. **207** (1999), no. 1, 131–143, MR 1724863, Zbl 0988.58013.
- [33] F. Hang and X. Wang, *Rigidity theorems for compact manifolds with boundary and positive Ricci curvature*, J. Geom. Anal. **19** (2009), 628–642, MR 2496569, Zbl 1175.53056.
- [34] O. Hijazi, S. Montiel and S. Raulot, *Uniqueness of the de Sitter spacetime among static vacua with positive cosmological constant*, Ann. Glob. Anal. Geom. **47** (2014), no. 2, 167–178, MR 3313139, Zbl 1318.53054.
- [35] W. Israel, *Event horizons in static vacuum space-times*, Phys. Rev. **164** (1967), 1776–1779.
- [36] O. Kobayashi, *A differential equation arising from scalar curvature function*, J. Math. Soc. Japan **34** (1982), no. 4, 665–675, MR 0669275, Zbl 0486.53034.
- [37] O. Kobayashi, *Fixed points of isometries*, Nagoya Math. J. **13** (1958), 63–68, MR 0103508, Zbl 0084.18202.
- [38] W. Kühnel, *Conformal transformations between Einstein spaces*, in *Conformal geometry* (Bonn, 1985/1986), 105–146, Aspects Math., E12, Vieweg, Braunschweig, 1988, MR 0979791, Zbl 0667.53039.
- [39] J. Lafontaine, *Sur la géométrie d’une généralisation de l’équation d’Obata*, J. Math. Pures Appliquées **62** (1983), 63–72, MR 0700048, Zbl 0513.53046.
- [40] J. Lafontaine and J.-L. Rozoy, *Courbure scalaire et trous noirs*, Séminaire de Théorie Spectrale et Géométrie, Vol. 18, Année 1999–2000, 69–76, MR 1812213, Zbl 0990.53045.
- [41] D. Lee and A. Neves, *The Penrose inequality for asymptotically locally hyperbolic spaces with nonpositive mass*, Comm. Math. Phys. **339** (2015), no. 2, 327–352, MR 3370607, Zbl 1322.53038.
- [42] L. Lindblom, *Some properties of static general relativistic stellar models*, J. Math. Phys. **21**, no. 6 (1980), 1455–1459.
- [43] Z.-D. Liu and Z. Shen, *Riemannian geometry of conical singular sets*, Ann. Global Anal. Geom. **16** (1998), no. 1, 29–62, MR 1616602, Zbl 0967.53026.
- [44] F. Marques and A. Neves, *Rigidity of min-max minimal spheres in three-manifolds*, Duke Math. J. **161** (2012), no. 14, 2725–2752, MR 2993139, Zbl 1260.53079.
- [45] R. Mazzeo, *Elliptic theory of differential edge operators. I.*, Comm. Partial Differential Equations **16** (1991), no. 10, 1615–1664, MR 1133743, Zbl 0745.58045.
- [46] W. Meeks and S.-T. Yau, *The existence of embedded minimal surfaces and the problem of uniqueness*, Math. Z. **179** (1982), no. 2, 151–168, MR 0645492, Zbl 0479.49026.
- [47] P. Miao and L.-F. Tam, *Static potentials on asymptotically flat manifolds*, Ann. Henri Poincaré **16** (2015), no. 10, 2239–2264, MR 3385979, Zbl 1326.83052.
- [48] I. Mondello, *The local Yamabe constant of Einstein stratified spaces*, arXiv:1411.7996.

- [49] H. Müller zum Hagen, D. Robinson and H. Seifert, *Black holes in static vacuum space-times*, Gen. Relativ. Gravitation **4** (1973), 53–78, MR 0398432.
- [50] M. Obata, *Certain conditions for a Riemannian manifold to be isometric with a sphere*, J. Math. Soc. Japan **14** (1962), 333–340, MR 0142086, Zbl 0115.39302.
- [51] M. Pflaum, *Analytic and geometric study of stratified spaces*, Lecture Notes in Mathematics, 1768. Springer-Verlag, Berlin, 2001, MR 1869601, Zbl 0988.58003.
- [52] J. Qing, *On the rigidity for conformally compact Einstein manifolds*, Int. Math. Res. Not. **21** (2003), 1141–1153, MR 1962123, Zbl 1042.53031.
- [53] J. Qing and W. Yuan, *A note on static spaces and related problems*, J. Geom. Phys. **74** (2013), 18–27, MR 3118569, Zbl 1287.83016.
- [54] D. Robinson, *A simple proof of the generalization of Israel’s Theorem*, Gen. Relativ. Gravitation **8** (1977), no. 8, 695–698, Zbl 0429.53043.
- [55] R. Schoen and S.T. Yau, *Existence of incompressible minimal surfaces and the topology of three dimensional manifolds with non-negative scalar curvature*, Ann. of Math (2) **110** (1979), no. 1, 127–142, MR 0541332, Zbl 0431.53051.
- [56] H. Seshadri, *On Einstein four-manifolds with S^1 -actions*, Math. Z. **247** (2004), no. 3, 487–503, MR 2114424, Zbl 1073.53063.
- [57] Y. Shen, *A note on Fischer–Marsden’s conjecture*, Proc. Amer. Math. Soc. **125** (1997), no. 3, 901–905, MR 1353399, Zbl 0867.53035.
- [58] L. Simon, *Lectures on geometric measure theory*, Proceedings of the Centre for Mathematical Analysis, Australian National University, Canberra, 1983, MR 0756417, Zbl 0546.49019.
- [59] R. Wald, *General relativity*, University of Chicago Press, Chicago, 1984, Zbl 0549.53001, MR 1243449, Zbl 0549.53001.
- [60] X. Wang, *On the uniqueness of the AdS spacetime*, Acta Math. Sin. (Engl. Ser.) **21** (2005), no. 4, 917–922, MR 2156971, Zbl 1082.53040.
- [61] D.-G. Yang, *Rigidity of Einstein 4-manifolds with positive curvature*, Invent. Math. **142** (2000), no. 2, 435–450, MR 1794068, Zbl 0981.53025.

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