# THE SPLITTING THEOREM FOR SPACE-TIMES WITH STRONG ENERGY CONDITION 

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## 1. Introduction

Our aim is the proof of the following theorem:
Theorem. Let $(M, g)$ be a connected, time oriented, globally hyperbolic Lorentzian manifold which is timelike geodesically complete ("tgc") and satisfies $\operatorname{Ric}(v, v) \geqslant 0$ for every timelike tangent vector $v$, where Ric denotes the Ricci tensor of $g$. Let $\gamma: \mathbb{R} \rightarrow M$ be a line, i.e. a timelike geodesic which realizes the distance between any two of its points. Then $(M, g)$ is isometric to $(\mathbb{R} \times H$, $\left.-d t^{2} \otimes h\right)$ where $(H, h)$ is a complete Riemannian manifold, and the factor $\left(\mathbb{R},-d t^{2}\right)$ is represented by $\gamma$.

This is the Lorentzian version of the Cheeger-Gromoll Splitting Theorem for Riemannian manifolds of nonnegative Ricci curvature [5] which solves a problem raised by S. T. Yau [12]. Our result extends earlier work of Galloway [8] and Beem et al. [2], [3]. Galloway [8] has proved the theorem under the additional assumption that $M$ admits a smooth function whose level sets are compact spacelike Cauchy hypersurfaces. Beem, Ehrlich, Markvorsen, and Galloway [2], [3] proved a Toponogov type splitting theorem [10] for Lorentzian manifolds, i.e. they assumed $g(R(w, v) v, w) \geqslant 0$ for any timelike vector $v$ and any $w \perp v$. The weaker Ricci curvature assumption of our theorem, called the strong energy condition (cf. [9]), is of particular interest in General Relativity. However, we need to assume the tgc property which can be concluded from the curvature assumption in the case of Beem et al.

The proofs of Cheeger and Gromoll [5] and H. Wu [11] for the Riemannian case apply the theory of elliptic operators to the Laplace operator $\Delta$ on the manifold. This fails in the Lorentzian case. In [6], a new proof was given which is based on the following idea: The Busemann functions $b^{+}$and $b^{-}$of the given line $\gamma$ satisfy $b^{+}+b^{-} \leqslant 0$ with equality along $\gamma$, by the triangle inequality. On the other hand, by Ric $\geqslant 0$, we have $\Delta b^{ \pm} \geqslant 0$ in the sense of support
functions. Thus we may apply the Hopf-Calabi maximum principle [4] to $b^{+}+b^{-}$to show that $b^{+}+b^{-}=0$ which is the core of the proof. Unfortunately, the maximum principle also fails for the Laplacian of a Lorentzian manifold. Therefore, instead of $b^{+}+b^{-}$we consider $b^{-} \mid\left\{b^{+}=a\right\}$ for a suitable value $a$. The level set $\left\{b^{+}=a\right\}$ can be approximated by smooth spacelike hypersurfaces where we can use the Riemannian Laplacian.

Another difficulty in the Lorentzian case is that the Busemann function might be uncontinuous. In [3], continuity has been shown by using a triangle comparison theorem which is not valid under the weaker curvature assumption. Instead, we show the continuity of the Busemann function near the line $\gamma$ (§3), and we get the splitting in a small neighborhood of $\gamma$ (§6). An "open and closed subset" argument then finishes the proof (§7).

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## 2. Rays and co-rays

2.1. Let $(M, g)$ be a connected, time oriented, globally hyperbolic, timelike geodesically complete ("tgc") Lorentzian manifold of dimension $n+1, n \geqslant 1$. We agree that $g$ is of type $(-,+, \cdots,+)$. Let $d$ be the Lorentzian distance function (see [9]) and $\ll$ the timelike future relation on $M$ : For $p, q \in M$ we say $p \ll q$ or $p \in I^{-}(q)$ or $q \in I^{+}(p)$ if there is a future oriented timelike curve from $p$ to $q$. The Lorentzian distance is continuous and satisfies the reverse triangle inequality: If $p, q, r \in M$ with $p \ll q \ll r$, then

$$
\begin{equation*}
d(p, q)+d(q, r) \leqslant d(p, r) \tag{T}
\end{equation*}
$$

A ( future oriented) ray is, by definition, a timelike, future oriented unit speed geodesic $\gamma:[a, \infty) \rightarrow M$ (for some $a \in \mathbb{R}$ ) which is maximal, i.e. $d(\gamma(t), \gamma(s))$ $=s-t$ for all $s, t$ with $a \leqslant t \leqslant s$. Everything will work as well for past oriented rays; we just have to reverse the time orientation of $M$. For a given ray $\gamma$ let us define functions $b_{s}: M \rightarrow \mathbb{R}$ for any $s \geqslant a$ by

$$
b_{s}(x)=s-d(x, \gamma(s)) .
$$

Then (T) implies for any $x \ll \gamma(s)$ and $r>s$

$$
\begin{equation*}
b_{r}(x) \leqslant b_{s}(s) \tag{1}
\end{equation*}
$$

On the other hand, if $\gamma(a) \ll x \ll \gamma(s)$ for some $s>a$, then

$$
\begin{equation*}
b_{s}(x) \geqslant d(\gamma(a), x)+a . \tag{2}
\end{equation*}
$$

Let $I(\gamma)=\{x \in M ; \gamma(a) \ll x \ll \gamma(s)$ for some $s>a\}$. This is an open neighborhood of $\gamma((a, \infty))$. By (1) and (2), $b_{s}(x)$ is a monotonously decreasing and bounded function of $s$ for any $x \in I(\gamma)$, so $b(x):=\lim _{s \rightarrow \infty} b_{s}(x)$ exists and defines a function $b: I(\gamma) \rightarrow \mathbb{R}$, called the Busemann function of $\gamma$. In general, the convergence need not be locally uniform and $b$ is ıot necessarily continuous.
2.2. However, the triangle inequality ( T ) gives a lower bound for the increase of $b$. For $p, q \in I(\gamma)$ with $p \ll q$, we have

$$
\begin{equation*}
b(q) \geqslant b(p)+d(p, q) \tag{3}
\end{equation*}
$$

Suppose we find two points $p_{0} \ll q_{0}$ where we have equality in (3). Then the function $f: I(\gamma) \cap I^{-}\left(q_{0}\right) \rightarrow \mathbb{R}$ with

$$
f(p)=b\left(q_{0}\right)-d\left(p, q_{0}\right)
$$

is an upper support function of $b$ at $p_{0}$, that means that $f$ is defined and continuous near $p_{0}$, and we have $f \geqslant b$ with equality at $p_{0}$. Likewise, $e$ : $I(\gamma) \cap I^{+}\left(p_{0}\right) \rightarrow \mathbb{R}$ with

$$
e(q)=b\left(p_{0}\right)+d\left(p_{0}, q\right)
$$

is a lower support function of $b$ at $q_{0}$.
If $\beta$ is a maximal geodesic connecting $p_{0}$ and $q_{0}$, then equality in (3) also holds for any two points on $\beta$ between $p_{0}$ and $q_{0}$. A ray $\beta:[0, \infty) \rightarrow I(\gamma)$ is called a co-ray of $\gamma$ if equality in (3) holds along $\beta$, i.e. for all $s>t \geqslant 0$,

$$
\begin{equation*}
b(\beta(s))-b(\beta(t))=s-t . \tag{4}
\end{equation*}
$$

(In fact it is sufficient to check (4) for $t=0$.) Note that the definition of a co-ray is slightly different from that in [2], [3]. Now for any $s>0$, the function $b_{\beta, s}: I(\gamma) \cap I^{-}(\beta(s)) \rightarrow \mathbb{R}$,

$$
b_{\beta, s}(x)=b(\beta(s))-d(x, \beta(s))=b(\beta(0))+s-d(x, \beta(s)),
$$

is an upper support function of $b$ at $\beta(0)$ which is smooth near $\beta(0)$ since $\beta(0)$ is not in the cut locus of $\beta(s)$ (cf. [1, p. 105]). In particular, for the Busemann functions $b_{\beta}$ of $\beta$ which is defined on $I(\beta) \subset I(\gamma)$, we have

$$
\begin{equation*}
b_{\beta} \geqslant b-b(\beta(0)) . \tag{5}
\end{equation*}
$$

Moreover, at any point $\beta(t)$ for $t>0$ there are smooth upper and lower support functions $f$ and $e$, namely $f(x)=b(\beta(s))-d(x, \beta(s))$ for any fixed $s>t$ and $e(x)=b(\beta(0))+d(\beta(0), x)$. Hence $b$ is differentiable at $\beta(t)$ with gradient $\nabla b(\beta(t))=-\beta^{\prime}(t)$, and, therefore, $\beta$ is the only co-ray passing through $\beta(t)$.
2.3. Does there exist a co-ray $\beta$ with $\beta(0)=q$ for any $q \in I(\gamma)$ ? For sufficiently large $s$ consider a maximal unit speed geodesic segment $\beta_{s}$ connecting $q$ to $\gamma(s)$ (which exists by global hyperbolicity), say $\beta_{s}(0)=q$ and
$\beta_{s}(r)=\gamma(s)$. Then we have $b_{s}(q)=s-r$, and in particular, $r=r(s) \rightarrow \infty$ as $s \rightarrow \infty$. Moreover, for any $0<u<r$ and $x \ll \beta_{s}(u)$,

$$
u-d\left(x, \beta_{s}(u)\right) \geqslant r-d\left(x, \beta_{s}(r)\right)=r-s+b_{s}(x)
$$

by (T). So we get

$$
\begin{equation*}
b_{s}(q)+u-d\left(x, \beta_{s}(u)\right) \geqslant b_{s}(x) \tag{6}
\end{equation*}
$$

Now suppose that the subset $\left\{\beta_{s}^{\prime}(0) ; s>0\right\}$ of $T_{q} M$ is bounded. This is true near $\gamma$ as we will see below. Then there is a sequence $s(i) \rightarrow \infty$ such that the vectors $\beta_{s(i)}^{\prime}(0)$ converge to a future oriented unit vector in $T_{q} M$. Hence the geodesics $\beta_{s(i)}$ converge to a geodesic $\beta$ which is a ray since $r(s(i)) \rightarrow \infty$ and $M$ is $\operatorname{tgc}$. If $\beta_{s}^{\prime}(u)$ is close enough to $\beta^{\prime}(u)$, then $\beta_{s}$ enters $I^{+}(\beta(u))$ for any $u \geqslant 0$. Thus $\beta$ lies in $I(\gamma)$. So we get from (6) for $x=\beta(u), u>0$, that

$$
b(\beta(0))+u \geqslant b(\beta(u))
$$

Together with (3) we get equality which shows that $\beta$ is a co-ray of $\gamma$.
2.4. Let $\beta_{k}$ be a sequence of co-rays with $\beta_{k}(0)=p_{k}, \beta_{k}^{\prime}(0)=v_{k}$. Assume $p_{k} \rightarrow p$ and $v_{k} \rightarrow v$. Then the rays $\beta_{k}$ converge to a ray $\beta$.

Claim. If $b\left(p_{k}\right) \rightarrow b(p)$, then $\beta$ is a co-ray. In fact, let $b_{k, s}$ be the upper support functions of $b$ corresponding to $\beta_{k}$ as defined in 2.2. Then $b_{k, s} \rightarrow b_{\beta, s}$. Thus $b \leqslant b_{\beta, s}$. In particular, for $0<t<s$,

$$
b(\beta(t)) \leqslant b_{\beta, s}(\beta(t))=b(p)+t .
$$

Comparing with (3) we get equality, and therefore, $\beta$ is a co-ray.

## 3. The Busemann function near its ray

Let $\gamma:[a, \infty) \rightarrow M$ be a ray with $a<0$. We want to study the functions $b_{s}$ and $b$ near $\gamma(0)$. Consider a small closed tubular neighborhood $U=U(R, T)$ $\subset I(\gamma)$ of radius $R$ around a segment $\gamma \mid[-T, T]$ with $-T>a$. Using Fermi coordinates, we may identify $U$ with $I \times K$ where $I=[-T, T]$ and $K=K(R)$ is the closed ball of radius $R$ around 0 in $\mathbb{R}^{n}$ : We choose a parallel orthonormal basis $\gamma^{\prime}=e_{0}, e_{1}, \cdots, e_{n}$ along $\gamma \mid[-T, T]$ and consider the map $\Phi: I \times K$ $\rightarrow U$,

$$
\Phi\left(x_{0}, x_{1}, \cdots, x_{n}\right)=\exp \left(\sum_{i=1}^{n} x_{i} e_{i}\left(x_{0}\right)\right) .
$$

If $R$ is small enough, $\Phi$ is a diffeomorphism and $\Phi^{-1}=x=\left(x_{0}, \cdots, x_{n}\right)$ is a coordinate chart on $U$. Moreover, for small $R$ the metric $g$ is close to the flat Minkowski metric $g_{0}=\|d \mathbf{x}\|^{2}-d x_{0}^{2}$ where $\mathbf{x}=\left(x_{1}, \cdots, x_{n}\right)$, and the corresponding connections $D$ and $D^{0}$ are close to each other.

Lemma 3.1. If $U$ is small enough, there is a constant $C>0$ such that $\left|b_{s}-x_{0}\right| \leqslant C\|\mathbf{x}\|^{2}$ on $U$ for all $s>2 T$.

Proof. Choose $R_{1}, T_{1}$ as above and put $U_{1}=U\left(R_{1}, T_{1}\right)$. Let $T=T_{1} / 2$. If we choose $R<R_{1}$ small enough, then the functions

$$
f(x)=2 T-d(x, \gamma(2 T)), \quad e(x)=-2 T+d(\gamma(-2 T), x)
$$

are defined and smooth on $U=U(R, T)$ with $f(\gamma(t))=e(\gamma(t))=t$ for $-T \leqslant t \leqslant T$, and $D f_{\gamma(t)}$ and $D e_{\gamma(t)}$ vanish on the normal space of $\gamma$. Hence there is a constant $C>0$ such that

$$
\left|f-x_{0}\right| \leqslant C\|\mathbf{x}\|^{2}, \quad\left|e-x_{0}\right| \leqslant C\|\mathbf{x}\|^{2} .
$$

By (T), we have $f \geqslant b_{s} \geqslant e$ for any $s>2 T$ (compare 2.2), so we get the result.
On $U$ we also have the euclidean metric $g_{1}=\|d x\|^{2}=\|d \mathbf{x}\|^{2}+d x_{0}^{2}$ which has the same connection $D^{0}$ as $g_{0}$. Let $\left\|\|_{1}\right.$ denote the norm with respect to $g_{1}$.

Lemma 3.2. If $U$ is small enough, there is a constant $M>0$ with the following property: If $q \in U$ and $s>2 T$ then every maximal unit speed geodesic segment $\alpha$ from $q$ to $\gamma(s)$ satisfies $\left\|\alpha^{\prime}(0)\right\|_{1} \leqslant M$.

Proof. Start with $U_{0}=U\left(R_{0}, T_{0}\right)$ which satisfies the assumption of Lemma 3.1 and $R_{0}<2 T_{0}$. Choose $R_{1} \leqslant R_{0}$ so that $\mu:=C \cdot R_{1}^{2} \leqslant R_{1} / 10$. Put $U_{1}=$ $U\left(R_{1}, T_{0}\right)$. Let $R=R_{1} / 2, T=T_{0} / 2$, and $U=U(R, T)$. For any $q \in U$ let $\alpha$ be a maximal unit speed geodesic segment from $q$ to $\gamma(s), s>2 T$. While $\alpha$ is in $U_{1}$, we have $\alpha(\tau)=\left(\alpha_{0}(\tau), \alpha(\tau)\right)$. Since $\alpha$ is timelike future oriented and $g$ is close to $g_{0}$, we may assume that

$$
2 \alpha_{0}^{\prime}(\tau)>\left\|\alpha^{\prime}(\tau)\right\|
$$

In particular, the function $t(\tau)=\alpha_{0}(\tau)-\alpha_{0}(0)$ has positive derivative. Choose $t_{0} \in\left(4 \mu, R_{1} / 2\right)$ and let $\tau_{0}$ be such that $t\left(\tau_{0}\right)=t_{0}$. Then $\alpha\left(\left[0, \tau_{0}\right]\right) \subset U_{1}$.

Since $\alpha$ is a maximal geodesic through $\gamma(s)$, we have

$$
\tau_{0}=b_{s}\left(\alpha\left(\tau_{0}\right)\right)-b_{s}(q)
$$

On the other hand, by 3.1,

$$
\left|b_{s}\left(\alpha\left(\tau_{0}\right)\right)-b_{s}(q)-t_{0}\right| \leqslant 2 \mu
$$

hence in particular

$$
\tau_{0} \geqslant t_{0}-2 \mu>t_{0} / 2
$$

Now by the mean value theorem, there exist $\tau_{1} \in\left(0, \tau_{0}\right)$ such that $\alpha_{0}^{\prime}\left(\tau_{1}\right)<2$, thus $\left\|\alpha^{\prime}\left(\tau_{1}\right)\right\|<4$ and $\left\|\alpha^{\prime}\left(\tau_{1}\right)\right\|_{1}<5$.

Suppose that $\left\|D-D^{0}\right\|_{1} \leqslant A$ on $U_{1}$. Then the function $f(\tau)=\left\|\alpha^{\prime}(\tau)\right\|_{1}$ satisfies $\left|f^{\prime}\right| \leqslant A \cdot f^{2}$ on $\left[0, \tau_{1}\right]$, whence $\left|(1 / f)^{\prime}\right| \leqslant A$ and

$$
1 / f(0) \geqslant 1 / f\left(\tau_{1}\right)-A \cdot \tau_{1} \geqslant 1 / 5-A \cdot \tau_{0} .
$$

Note that $\tau_{0}<t_{0}+2 \mu<T_{0}+2 \mu$ is uniformly bounded. So we get $\left\|\alpha^{\prime}(0)\right\|_{1}$ $<M:=1 / 10$ if $R_{1}$ was chosen so small that $A \cdot\left(T_{0}+2 \mu\right) \leqslant 1 / 10$.

Lemma 3.3. If $U$ is small enough, the Busemann function is Lipschitz continuous on $U$ with respect to the metric $g_{1}$. Moreover, for any $p \in U$ there is a co-ray $\beta$ starting at $p$ and therefore an upper support function $b_{p, s}:=b_{\beta, s}$ of $b$ at $p$.

Proof. We will show that the functions $b_{s}=s-d_{s}$ with $d_{s}(x)=d(x, \gamma(s))$ are equi-Lipschitz continuous. The function $d_{s}$ on $I^{-}(\gamma(s))$ is continuous but not smooth everywhere. However, for any $p \in I^{-}(\gamma(s))$, we get a smooth lower support function $e$ of $d_{s}$ at $p$ as follows: Let $\alpha$ be a maximal unit speed geodesic segment from $p$ to $\gamma(s)$ and choose $q$ on $\alpha$ between $p$ and $\gamma(s)$. Put $e(x)=d(x, q)+d(q, \gamma(s))$. Then $e$ is smooth near $p$ because there is no cut point of $q$ near $p$ and, by (T), $e$ is a lower support function of $d_{s}$ at $p$. Thus we only have to show that $\left\|D e_{p}\right\|_{1} \leqslant L$ for all $p \in U$ with a constant $L$ not depending on $s$ (see the Appendix). Now for any $v \in T_{p} M$,

$$
\left|D e_{p}(v)\right|=\left|g\left(\alpha^{\prime}(0), v\right)\right| \leqslant G \cdot M \cdot\|v\|_{1},
$$

where $M$ is as in 3.2 and $G$ is an upper bound of $\|g\|_{1}$ on $U$. If we put $L=G \cdot M$, then $b_{s}$ and hence $b$ are Lipschitz continuous with Lipschitz constant $L$. The existence of co-rays is clear from 3.2 and 2.3.

Lemma 3.4. Let $\left(p_{k}\right)$ be a sequence in $U$ with $p_{k} \rightarrow p:=\gamma(0)$. Let $\beta_{k}$ be a co-ray starting at $p_{k}$ with initial vector $v_{k}$. Then $v_{k} \rightarrow \gamma^{\prime}(0)$.

Proof. Since $\left\|v_{k}\right\|_{1}$ is bounded, we may assume $v_{k} \rightarrow v$ for some $v \in T_{p} M$. Thus, by 2.4 , the co-rays $\beta_{k}$ converge to a co-ray $\beta$ starting at $p$. But by 2.3, since $a<0, \gamma$ is the only co-ray passing through $p$. Thus $\beta=\gamma \mid[0, \infty)$ which finishes the proof.

## 4. Lines

A complete geodesic $\gamma: \mathbb{R} \rightarrow M$ is called a line if $\gamma \dot{\mathrm{i}}[a, \infty)$ is a ray for any $a \in \mathbb{R}$. In other words, we have $d(\gamma(t), \gamma(s))=s-t$ for any $t<s$. Let $\gamma^{-}$: $\mathbb{R} \rightarrow M, \gamma^{-}(t)=\gamma(-t)$. Then $\gamma^{-} \mid[a, \infty)$ is a past oriented ray for any $a$. Besides $b_{s}^{+}:=b_{s}$, we have the analogue functions for $\gamma^{-}$,

$$
b_{s}^{-}(x)=s-d(\gamma(-s), x)
$$

and the two Busemann functions $b^{ \pm}=\lim _{s \rightarrow \infty} b_{s}^{ \pm}$, defined on

$$
I(\gamma)=\{x \in M ; \gamma(t) \ll x \ll \gamma(s) \text { for some } s, t \in \mathbb{R}\} .
$$

The triangle inequality ( T ) implies that $-b_{r}^{-} \leqslant b_{s}^{+}$with equality on $\gamma \mid[-r, s]$ for any $r, s>0$, and in particular we get $b^{+}+b^{-} \geqslant 0$ with equality along $\gamma$.

Lemma 4.1. Let $q \in I(\gamma)$ be such that $\left(b^{+}+b^{-}\right)(q)=0$. Assume further that there are co-rays $\beta^{+}$of $\gamma$ and $\beta^{-}$of $\gamma^{-}$starting at $q$. Then $\beta^{+}$and $\beta^{-}$fit together to a line $\beta$ whose Busemann functions $b_{\beta}^{ \pm}$satisfy $b_{\beta}^{+}+b_{\beta}^{-} \geqslant b^{+}+b^{-}$ on $I(\beta) \subset I(\gamma)$ with equality along $\beta$.

Proof. Let $b^{+}(q)=-b^{-}(q)=: a$. We have $b^{-}\left(\beta^{-}(t)\right)=-a+t$ and by (3) (§2) $b^{+}\left(\beta^{-}(t)\right) \leqslant a-t$ for all $t \geqslant 0$. Hence $\left(b^{+}+b^{-}\right)\left(\beta^{-}(t)\right) \leqslant 0$ and so we have equality and in particular, $b^{+}\left(\beta^{-}(t)\right)=a-t$. Let us define the broken geodesic $\beta: \mathbb{R} \rightarrow M$ by $\beta(t)=\beta^{+}(t)$ for $t \geqslant 0$ and $\beta(t)=\beta^{-}(-t)$ for $t \leqslant 0$. Then $b^{+}(\beta(t))=a+t$ and, similarly, $b^{-}(\beta(t))=-(a+t)$ for all $t \in \mathbb{R}$, and it follows from (3) (§2) that $\beta$ is an unbroken maximal geodesic, hence a line. By (5) (§2) we have $b_{\beta}^{ \pm} \geqslant \pm a+b^{ \pm}$which finishes the proof.

Lemma 4.2. Let $\gamma: \mathbb{R} \rightarrow M$ be a line. Then there is a neighborhood $W$ of $\gamma(\mathbb{R})$ such that the Busemann functions $b^{+}$and $b^{-}$are continuous on $W$, and for any $q \in W$ there exist co-rays of $\gamma$ and $\gamma^{-}$starting at $q$.

Proof. This follows from Lemma 3.3.

## 5. The Ricci curvature condition

Let $\left(M^{n+1}, g\right)$ be a time oriented Lorentzian manifold, $M^{\prime} \subset M$ an open subset, and $f: M^{\prime} \rightarrow \mathbb{R}$ a smooth function with $g(\nabla f, \nabla f)=-1$. Then the gradient lines of $f$ are timelike unit speed geodesics and the level sets of $f$ are spacelike hypersurfaces, and $f$ can be viewed as a signed distance function of each of its level hypersurfaces, up to a constant. Moreover, if $\gamma$ is a gradient line and $U(t)=\left.D \nabla f\right|_{\gamma(t)}$ is the Hessian tensor along $\gamma$, then this tensor field along $\gamma$ satisfies the Riccati equation

$$
\begin{equation*}
U^{\prime}+U^{2}+R=0 \tag{1}
\end{equation*}
$$

with $R(t) v:=R\left(v, \gamma^{\prime}(t)\right) \gamma^{\prime}(t)$. Taking the trace and putting $u(t)=\operatorname{trace} U(t)$ $=\Delta f\left(\gamma(t)\right.$ ), we get by the Schwarz inequality (note that $U\left(\left(\gamma^{\prime}\right)^{\perp}\right) \subset\left(\gamma^{\prime}\right)^{\perp}$ and $U\left(\gamma^{\prime}\right)=0$ )

$$
\begin{equation*}
u^{\prime}+u^{2} / n+\operatorname{Ric}\left(\gamma^{\prime}, \gamma^{\prime}\right) \leqslant 0 \tag{2}
\end{equation*}
$$

with equality if and only if $U \mid\left(\gamma^{\prime}\right)^{\perp}$ is a multiple of the identity (cf. [6], [7]).
From now on suppose $\operatorname{Ric}(v, v) \geqslant 0$ for any timelike tangent vector $v$. Then, by (2), the function $\phi=1 / u$ satisfies $\phi^{\prime} \geqslant 1 / n$. Let $f=d_{q}:=d(\cdot, q)$, defined on $I^{-}(q)$ outside the cut locus of $q$, for some point $q \in M$. Then $f$ is as described above and any gradient line $\gamma$ is future oriented with future end point $q$. Let $q=\gamma(s)$. Then $u(t) \rightarrow \infty$ as $t \rightarrow s$ and $t<s$, whence $\phi(s)=0$ and so $-\phi(t) \geqslant(s-t) / n$. So we get

$$
\Delta d_{q} \geqslant-n / d_{q} .
$$

Therefore, for any ray $\gamma$, the functions $b_{s}=s-d_{\gamma(s)}$ satisfy

$$
\Delta b_{s} \leqslant n / d_{\gamma(s)}
$$

In particular, we get from 2.2:
Lemma 5.1. Let $\gamma$ be a ray and $\beta$ a co-ray with $\beta(0)=q$. Then for any $\varepsilon>0$ there exists $r>0$ such that $\Delta b_{\beta, s} \leqslant \varepsilon$ near $q$ for any $s \geqslant r$.

Lemma 5.2. Let $\gamma: \mathbb{R} \rightarrow M$ be a line. Then for any $t \in \mathbb{R}$ we have $\left.D D\left(b_{s}^{ \pm}\right)\right|_{\gamma(t)} \rightarrow 0$ as $s \rightarrow \infty$, and $R\left(, \gamma^{\prime}(t)\right) \gamma^{\prime}(t)=0$ for any $t \in \mathbb{R}$.

Proof. Let $U_{s}(t)=\left.D \nabla\left(b_{s}^{+}\right)\right|_{\gamma(t)}$ for $t<s$. For fixed $v \in T_{\gamma(t)} M$, the function $s \rightarrow D D\left(b_{s}^{+}\right)(v, v)=g\left(U_{s}(t) v, v\right)$ is monotonously decreasing and bounded from below by $-D D\left(b_{r}^{-}\right)(v, v)$ for arbitrary $r>|t|$ (see (1), (7) in §2). Thus $U(t):=\lim _{s \rightarrow \infty} U_{s}(t)$ exists for any $t \in \mathbb{R}$ and solves the Riccati equation (1). Thus $u:=\operatorname{trace} U \equiv 0$ since otherwise $\phi=1 / u$ would satisfy $\phi^{\prime} \geqslant 1 / n$ and therefore $\phi$ would have a zero which is impossible. Now $U \equiv 0$ follows from the equality discussion of (2). So we get the result for $b_{s}^{+}$and similarly for $b_{s}^{-}$. In particular we get from $U=0$ that $R\left(, \gamma^{\prime}\right) \gamma^{\prime}=0$.

## 6. Local splitting

Let $\gamma: \mathbb{R} \rightarrow M$ be a line and $W$ a neighborhood of $\gamma(\mathbb{R})$ as in Lemma 4.2.
Proposition 6.1. There is a neighborhood $W_{0} \subset W$ of $\gamma(\mathbb{R})$ such that $b^{+}+$ $b^{-} \equiv 0$ on $W_{0}$.

Proof. Assume the contrary. Then the open set

$$
P=\left\{b^{+}+b^{-}>0\right\} \cap W
$$

is nonempty (recall that $b^{+}+b^{-} \geqslant 0$ ) and has a boundary point $p$ on $\gamma(\mathbb{R})$. We may shift the parameter so that $p=\gamma(0)$.

We may assume without restriction of generality that $P$ contains an open coordinate ball $B$ with $p \in \partial B$. Namely, otherwise let $V \subset W$ be a coordinate neighborhood around $p$, and choose $q \in P \cap V$ near $p$. Since $P \cap V$ is open, there exists an open coordinate ball $B^{\prime} \subset P \cap V$ centered at $q$. By enlarging the radius of $B^{\prime}$ we will meet $\partial P \cap V$ at last. Let $B^{\prime \prime}$ be the smallest of these balls such that $\partial B^{\prime \prime}$ intersects $\partial P$, say at $p_{1}$. Then we still have $B^{\prime \prime} \subset P$, but at $p_{1} \in \partial B^{\prime \prime}$ we have $\left(b^{+}+b^{-}\right)\left(p_{1}\right)=0$. Thus by 4.1 there is a line $\gamma_{1}$ passing through $p_{1}$ whose Busemann functions $b_{1}^{ \pm}$satisfy $b_{1}^{+}+b_{1}^{-} \geqslant b^{+}+b^{-}$. Let $W_{1}$ be a neighborhood of $\gamma_{1}(\mathbb{R})$ as in 4.2. So the set $P_{1}:=\left\{b_{1}^{+}+b_{1}^{-}>0\right\} \cap W \cap$ $W_{1}$ contains $P \cap W_{1}$ and hence $B^{\prime \prime} \cap W_{1}$. Since $W_{1}$ is a neighborhood of $p_{1}$, there is a smaller coordinate ball $B \subset B^{\prime \prime} \cap W_{1}$ with $p_{1} \in \partial B$. Now we replace $\gamma$ and $P$ with $\gamma_{1}$ and $P_{1}$ and our assumption is satisfied.

Let $U=U(R, T) \subset W$ be a closed tubular neighborhood of radius $R$ around $\gamma \mid[-T, T]$ as in Lemma 3.3, where $R$ and $T$ are chosen small enough. For suitable constants $\alpha, \mu>0$ let

$$
\begin{aligned}
& A=\left\{x \in U ;\|\mathbf{x}\|^{2}>\alpha \cdot x_{0}^{2}\right\} \subset U \\
& V=\left\{v \in T U ;\|\mathbf{v}\|^{2}<\mu \cdot v_{0}^{2}\right\} \subset T U .
\end{aligned}
$$

In the subsequent Lemma 6.2, we will construct a smooth function $h$ on $U$ with the following properties:

$$
\begin{gather*}
h>0 \quad \text { on } A \backslash B,  \tag{1}\\
h(p)=0,  \tag{2}\\
\Delta h<0, \tag{3}
\end{gather*}
$$

$$
\begin{equation*}
\operatorname{DDh}(v, v)<0 \text { for every } v \in V \tag{4}
\end{equation*}
$$

If $R$ is small enough, we get from Lemma 3.1:

$$
\begin{gathered}
\left\{b^{ \pm}=0\right\} \cap U \subset A, \\
\left\{b^{+} \leqslant 0\right\} \cap U \subset\left\{x_{0} \leqslant 0\right\} \cup A, \\
b^{-}>0 \quad \text { on }\left\{x_{0} \leqslant 0\right\} \backslash A .
\end{gathered}
$$

For any $r \in(0, R]$ let $U_{r}=U(r, T)$. If we choose $r$ sufficiently small, then, by Lemma 3.4, every coray $\beta^{-}$of $\gamma^{-}$with $\beta^{-}(0) \in U_{r}$ is very close to $\gamma^{-}$. Thus for a support function $b_{q, s}^{-}$with $q \in U_{r}$ corresponding to a coray $\beta^{-}$, we get that $v^{-}=-\nabla\left(b_{q, s}^{-}\right)(q)$ is close to the coordinate vector $\left.\left(\partial / \partial x_{0}\right)\right|_{q}$. Moreover, for any smooth vector field $W$, we have $D D\left(b_{q, s}^{-}\right)\left(W_{q}, W_{q}\right) \leqslant C$ with a constant $C$ independent of $q$ and $s$ (recall that $D D\left(b_{q, s}^{-}\right)(q)$ is monotonously decreasing with $s$; compare the proof of Lemma 5.2).

For small enough $r$ let $S_{r}=\left(\partial U_{r}\right) \cap\left\{b^{+} \leqslant 0\right\}$. Then $b^{-} \geqslant 0$ on $S_{r}$ and $S_{r} \subset\left\{b^{-}>0\right\} \cup\{h>0\}$. Thus there exists $\varepsilon_{r}>0$ such that $f:=b^{-}+\varepsilon \cdot h>$ 0 on $S_{r}$ for any $\varepsilon \in\left(0, \varepsilon_{r}\right]$. Let $U_{r}^{-}=U_{r} \cap\left\{b^{+} \leqslant 0\right\}$. Then

$$
\partial U_{r}^{-} \subset S_{r} \cup\left(\left\{b^{+}=0\right\} \cap \operatorname{Int}(U)\right)
$$

We have $f(p)=0$ by (2), and $p \in U_{r}^{-}$. Thus $f$ takes a minimum $m \leqslant 0$ on $U_{r}^{-}$, say at $q$. Since $f \mid S_{r}>0$, we have

$$
q \in \operatorname{Int}\left(U_{r}^{-}\right) \cup\left(\left\{b^{+}=0\right\} \cap \operatorname{Int}\left(U_{r}\right)\right) .
$$

By Lemma 3.3, there is a co-ray $\beta^{-}$of $\gamma^{-}$starting at $q$, and $b^{-}$has smooth upper support functions $b_{q, s}^{-}$at $q$. Thus $f_{s}:=b_{q, s}^{-}+\varepsilon \cdot h$ is a smooth upper support function of $f$ at $q$, and therefore, also, $f_{s}$ takes its minimum on $U_{r}^{-}$ at $q$.

By (3) and Lemma 5.1 we have $\Delta f_{s}(q)<0$ if $s$ is large enough, depending on $\varepsilon$. Moreover, we claim that $D D f_{s}(v, v)<0$ for $v=\nabla f_{s}(q)$ for arbitrary $s \geqslant 1$, provided that $\varepsilon$ is small enough. In fact, we have $v=-v^{-}+\varepsilon \cdot w$ where $w=\nabla h(q)$ and $v^{-}=-\nabla\left(b_{q, s}^{-}\right)(q)=\left(\beta^{-}\right)^{\prime}(0)$. Since $r$ is small, $v^{-}$and hence $-v$ are close enough to the coordinate vector $\partial / \partial x_{0}$ to satisfy the assumption of (4) (see Lemma 3.4). In fact, we may assume that $D D h(v, v)<-\delta$ for some $\delta>0$ depending only on $r$. Further, $v^{-}$is in the kernel of the Hessian $D D\left(b_{q, s}^{-}\right)$at $q$ for any $s$, so

$$
D D\left(b_{q, s}^{-}\right)(v, v)=\varepsilon^{2} D D\left(b_{q, s}^{-}\right)(w, w) \leqslant C \cdot \varepsilon^{2} .
$$

Thus

$$
\begin{aligned}
D D f_{s}(v, v) & =\varepsilon^{2} \cdot D D\left(b_{q, s}^{-}\right)(w, w)+\varepsilon \cdot \operatorname{Dh}(v, v) \\
& \leqslant \varepsilon^{2} C-\varepsilon \delta<0
\end{aligned}
$$

if $\varepsilon$ is small enough.
Let $a:=b^{+}(q) \leqslant 0$. Put $g_{s}:=b_{q, s}^{+}$and let $H_{s}$ be the level set $\left\{g_{s}=a\right\}$ passing through $q$. This is a smooth spacelike hypersurface near $q$. We have $g_{s} \geqslant b^{+}$with equality at $q$. Therefore $H_{s} \subset U_{r}^{-}$and $f_{s} \mid H_{s}$ takes a minimum at $q$. Hence the gradients of $f_{s}$ and $g_{s}$ at $q$ are linearly dependent. Since the co-rays $\beta^{ \pm}$starting at $q$ are integral curves of $-\nabla\left(b_{q, s}^{ \pm}\right)$and $\beta^{+}\left(\beta^{-}\right)$is future (past) oriented, we get $\nabla f_{s}(q)=-\lambda \cdot \nabla g_{s}(q)$ for some positive $\lambda$ close to 1 . On the other hand, along $H_{s}$ we have $\Delta g_{s}=-g\left(\eta, \nabla g_{s}\right)$ where $\eta$ denotes the mean curvature normal field on $H_{s}$ (recall that $g_{s}$ is the distance function of $H_{s}$, up to a constant) and

$$
\Delta f_{s}=\Delta_{s} f_{s}-g\left(\eta, \nabla f_{s}\right)-D D f_{s}(N, N)
$$

where $N$ is the unit normal field on $H_{s}$ and $\Delta_{s}$ is the Laplacian of the induced (Riemannian) metric on $H_{s}$. So

$$
\Delta_{s} f_{s}(q)=\Delta f_{s}(q)+\lambda \cdot \Delta g_{s}(q)+D D f_{s}\left(N_{q}, N_{q}\right)<0
$$

for sufficiently large $s$, since $N_{q}$ and $v=\nabla f_{s}(q)$ are linearly dependent. But this is a contradiction to the minimality of $f_{s} \mid H_{s}$ at $q$. This finishes the proof of Proposition 6.1.

Lemma 6.2. There exists a function $h$ on $U$ with the properties (1)-(4) as above.

Proof. As before, let $x=\left(x_{0}, \cdots, x_{n}\right): U \rightarrow \mathbb{R}^{n+1}$ be the Fermi coordinate system on $U$, and put $\mathbf{x}=\left(x_{1}, \cdots, x_{n}\right)$ and $\mathbf{x}^{*}=\left(x_{1}, \cdots, x_{n-1}\right)$. We may assume that $B \cap U=\{\psi<0\}$ with

$$
\psi=x_{0}^{2}+\left\|\mathbf{x}^{*}\right\|^{2}+x_{n}^{2}-2 S \cdot x_{n}
$$

for some radius $S>0$. Assume $R<S / 12 \ll 1$ and $T<2 R$. Choose $\mu \in$ $(R / 2, S / 12)$ and $\alpha>1+4 \mu$. Put

$$
\phi=\left\|\mathbf{x}^{*}\right\|^{2} /(2 \mu)-x_{n}-x_{0}^{2} .
$$

Claim. If $x \in A$ and $\phi(x) \leqslant 0$, then $\psi(x)<0$.
In fact, from
(a)

$$
\left|\mathbf{x}^{*}\right|^{2}>\alpha \cdot x_{0}^{2}-x_{n}^{2}
$$

$$
\begin{equation*}
\left|\mathbf{x}^{*}\right|^{2} \leqslant 2 \mu\left(x_{n}+x_{0}^{2}\right) \tag{b}
\end{equation*}
$$

we get $x_{n}^{2}+2 \mu \cdot x_{n}>(\alpha-2 \mu) x_{0}^{2}>(1+2 \mu) x_{0}^{2}$. Thus either $x_{n}<-2 \mu<-R$ (which is impossible) or

$$
x_{n}>\mu \cdot\left\{\left[1+x_{0}^{2}(1+2 \mu) / \mu^{2}\right]^{1 / 2}-1\right\}
$$

Since $x_{0}^{2}<T^{2}<4 R^{2}<16 \mu^{2}$, we may assume

$$
x_{0}^{2}(1+2 \mu) / \mu^{2}<24
$$

and since the slope of the square root function is bigger than $1 / 10$ on the interval ( 1,25 ), the inequality above gives

$$
x_{n}>x_{0}^{2}(1+2 \mu) /(10 \mu)
$$

Now using (b) we have

$$
\begin{aligned}
\psi(x) & \leqslant x_{0}^{2}(1+2 \mu)+2 \mu \cdot x_{n}+x_{n}^{2}-2 S \cdot x_{n} \\
& <\left(12 \mu+x_{n}-2 S\right) \cdot x_{n}<\left(x_{n}-S\right) \cdot x_{n}<0
\end{aligned}
$$

since $x_{n} \leqslant R<S$. This proves the claim. So we get $\phi>0$ on $A \backslash B$.
Now we put $h=1-e^{-\sigma \cdot \phi}$ for some sufficiently big constant $\sigma$. Then (1) and (2) are true. Since $\phi$ satisfies $g_{0}(D \phi, D \phi)>0$ and thus $g(D \phi, D \phi)>0$ if $R$ is small enough, we have

$$
\Delta h=e^{-\sigma \phi}\left(-\sigma^{2} \cdot g(\nabla \phi, \nabla \phi)+\sigma \cdot \Delta \phi\right)<0
$$

if $\sigma$ is big enough, so (3) holds. Moreover,

$$
D D h(v, v)=e^{-\sigma \phi}\left(-\sigma^{2} \cdot g(\nabla \phi, v)^{2}+\sigma \cdot D D \phi(v, v)\right)
$$

Since the Hessian of $\phi$ with respect to $g_{0}$, namely

$$
D^{0} D \phi(v, v)=\left\|\mathbf{v}^{*}\right\|^{2} / \mu-2 v_{0}^{2}
$$

is negative if $\left\|\mathbf{v}^{*}\right\|^{2} \leqslant\|\mathbf{v}\|^{2}<2 \mu \cdot v_{0}^{2}$, we get $D D \phi(v, v)<0$ if $\left\|\mathbf{v}^{*}\right\|^{2}<\mu \cdot v_{0}^{2}$ provided that $R$ is small enough. This shows (4).

Proposition 6.3. There is a neighborhood $W_{1}$ of $\gamma(R)$ which splits. More precisely, there exists an n-dimensional Riemannian manifold $H$ and an isometry $j: \mathbb{R} \times H \rightarrow W_{2}$ such that $j \mid \mathbb{R} \times\{p\}=\gamma$ for some $p \in H$.

Proof. The argument is similar to the Riemannian case (cf. [6, Chapter 4], and also [2], [3]). Let $W_{0}$ be as in Proposition 6.1. By 4.1 and 5.2, at any $q \in W_{0}$ there are smooth upper support functions $b_{q, s}^{ \pm}$of $b^{ \pm}$with $\left.D D\left(b_{q, s}^{ \pm}\right)\right|_{q}$ arbitrarily close to zero. Thus for any geodesic $c: I \rightarrow W_{0}$, the functions $b^{ \pm} \circ c$ have upper support functions with second derivatives arbitrarily close to zero, at any $t \in I$. The same is true for $b^{ \pm} \circ c-g$ where $g: I \rightarrow \mathbb{R}$ is the affine function with $b^{ \pm} \circ c=g$ on $\partial I$. Thus $b^{ \pm} \circ c \geqslant g$ by the one-dimensional maximum principle. So $b^{ \pm} o c$ is concave, and therefore, $b^{+}=-b^{-}$is concave and convex. Hence on $W_{0}$, the level sets of $b^{+}$are totally geodesic and $b^{+}$is smooth with parallel gradient field $\nabla b^{+}=: V$. For any $q \in W_{0}$, let $\gamma_{q}$ be the line passing through $q$ (see 4.1). Put

$$
H=\left\{b^{+}=0\right\} \cap W_{0} .
$$

Let $j: \mathbb{R} \times H \rightarrow M, j(t, q)=\gamma_{q}(t)=\exp _{q}\left(t \cdot V_{q}\right)$. Since $\gamma_{q}$ is a line, $j$ is defined on all of $\mathbb{R} \times H$; moreover we have $R\left(, \gamma_{q}^{\prime}\right) \gamma_{q}^{\prime}=0$ by 5.2. So for any $v \in T_{q} H$, the Jacobi field $J(t)=D j_{(t, q)}(v)$ along $\gamma_{q}$ with initial values $J(0)=v$, $J^{\prime}(0)=D_{v}\left(\nabla b^{+}\right)=0$ is parallel. Therefore, $j$ is a local isometry. Since the lines $\gamma_{q}$ are co-rays of $\gamma$, they do not intersect each other, so $j$ is one-to-one. Thus putting $W_{1}=j(\mathbb{R} \times H)$, we finish the proof.

## 7. Global splitting

By a flat strip, we mean a totally geodesic isometric immersion $f$ of $(\mathbb{R} \times I$, $-d t^{2}+d s^{2}$ ) into $M$ for some real interval $I$, such that $f \mid \mathbb{R} \times\{s\}$ is a line for any $s \in I$.

Proposition 7.1. Let $\gamma$ be a line and $c:[0,1] \rightarrow M$ any geodesic with $c(0)=\gamma(0)$. Then there is a flat strip containing $\gamma$ and $c$.

Proof. We may assume that $c$ is not contained in $\gamma$. Let $F$ be the set of all parameters $u \in[0,1]$ such that there is a flat strip containing $\gamma$ and $c \mid[0, u]$. Obviously $0 \in F$. Let $v=\sup F$. It follows from the local splitting (Proposition 6.3) that $v>0$. There is a flat strip $f: \mathbb{R} \times[0, a) \rightarrow M$ containing $\gamma$ and $c \mid[0, v)$ for some $a \in(0, \infty]$. More precisely, $\gamma(t)=f(t, 0)$ and $c(u)=$ $f(k \cdot u, m \cdot u)$ for $0 \leqslant u<v$, where $k, m \in \mathbb{R}$, and $a=m \cdot v<\infty$. Since $c$ is not contained in $\gamma$, we have $m \neq 0$.

Claim. $\quad f$ can be extended to $\mathbb{R} \times[0, a]$.
In fact, put $\gamma_{u}(t)=f(k \cdot u+t, m \cdot u)$ for $0 \leqslant u<v$. This is a line. Let $X$ denote the parallel vector field along $c$ with $X(0)=\gamma^{\prime}(0)$. Then $X(u)=\gamma_{u}^{\prime}(0)$
for $0<u<v$. Let $\gamma_{v}$ be the geodesic with $\gamma_{v}^{\prime}(0)=X(v)$. Then $\gamma_{v}(t)=$ $\lim _{u \rightarrow v} \gamma_{u}(t)$, hence $\gamma_{v}$ is a line. Now for $0 \leqslant s \leqslant a$ and $t \in \mathbb{R}$ we may put

$$
f(t, s)=\gamma_{s / m}(t-k \cdot s / m)
$$

this is a smooth extension of the previous map $f$ and hence a flat strip containing $\gamma$ and $c \mid[0, v]$.

Now the local splitting 6.3 implies again that the flat strip can be extended beyond $\gamma_{v}$. Ths is a contradiction to the choice of $v$ unless we have $v=1$ which finishes the proof.

Two lines $\gamma_{1}$ and $\gamma_{2}$ are called strongly parallel if they bound a flat strip, i.e. there is a flat strip $f: \mathbb{R} \times\left[a_{1}, a_{2}\right] \rightarrow M$ with $\gamma_{i}=f \mid\left(\mathbb{R} \times\left\{a_{i}\right\}\right)$ for $i=1,2$. They are called parallel if there exist lines $\gamma_{1}=\beta_{0}, \beta_{1}, \cdots, \beta_{k}=\gamma_{2}$ such that $\beta_{j-1}$ and $\beta_{j}$ are strongly parallel for $j=1, \cdots, k$.

Lemma 7.2. If $\gamma_{1}, \gamma_{2}$ are parallel lines, then $I\left(\gamma_{1}\right)=I\left(\gamma_{2}\right)$ and the Busemann functions $b_{1}^{ \pm}$and $b_{2}^{ \pm}$of $\gamma_{1}$ and $\gamma_{2}$ agree.

Proof. We may assume that $\gamma_{1}$ and $\gamma_{2}$ are strongly parallel and bound the flat strip $f: \mathbb{R} \times[0, a] \rightarrow M$. Then $\gamma_{1} \subset I\left(\gamma_{2}\right)$ and $\gamma_{2} \subset I\left(\gamma_{1}\right)$, hence $I\left(\gamma_{1}\right)=$ $I\left(\gamma_{2}\right)$. Further, for $t>a$ and $0 \leqslant s \leqslant a$ we have $d\left(f(0, s), \gamma_{i}(t)\right)^{2} \geqslant t^{2}-a^{2}$, thus $b_{i}^{+}(f(0, s)) \leqslant 0$. Likewise, $b_{i}^{-}(f(0, s)) \leqslant 0$, and therefore $b_{i}^{ \pm}(f(0, s))=0$, due to $b_{i}^{+}+b_{i}^{-} \geqslant 0$. Moreover, $\gamma_{1}^{ \pm}$is a co-ray of $\gamma_{2}^{ \pm}$and vice versa. So, by (5) in §2, we have $b_{1}^{ \pm} \geqslant b_{2}^{ \pm} \geqslant b_{1}^{ \pm}$and we get the result.

Now consider a fixed line $\gamma$. Let $P_{\gamma} \subset M$ be the set of points which lie on a parallel line. It follows from 7.2 and 2.2 that $b^{+}$is differentiable at any point $q \in P_{\gamma}$ and that there is exactly one parallel line $\gamma_{q}$ passing through $q$.

Lemma 7.3. $\quad P_{\gamma}$ is a connected component of $M$.
Proof. $\quad P_{\gamma}$ is open by the local splitting 6.3. We show that $\partial P_{\gamma}$ is empty. Suppose that there exists a point $q \in \partial P_{\gamma}$. Let $B_{q}$ be a geodesically convex open coordinate ball around $q$. Then $B_{q} \cap P_{\gamma}$ is open and nonempty. Choose $p \in B_{q} \cap P_{\gamma}$ close to $q$. There is a geodesically convex ball $B_{p} \subset B_{q} \cap P_{\gamma}$ such that $\partial B_{p}$ hits $\left(\partial P_{\gamma}\right) \cap B_{q}$, say at $q^{\prime}$. Let $c:[0,1] \rightarrow B_{q}$ be a geodesic segment with $c(0)=p, c(1)=q^{\prime}$. There is a line $\gamma_{1}$ passing through $p$ which is parallel to $\gamma$. By 7.1, there is a line $\gamma_{2}$ through $q^{\prime}$ which is strongly parallel to $\gamma_{1}$. Thus $q^{\prime} \in P_{\gamma}$. But since $P_{\gamma}$ is open, it does not intersect its boundary, a contradiction.

Now we can prove the Theorem. Since $M$ is connected, through every point $q \in M$ there is exactly one line $\gamma_{q}$ which is parallel to the given line $\gamma$, by 7.3. Let $V_{q}$ be the tangent vector of $\gamma_{q}$ at $q$. By the local splitting 6.3, this defines a parallel timelike vector field $V$ on $M$. Thus $V^{\perp}$ is a parallel distribution; in particular, it is integrable. Let $H$ be the maximal integral leave through
$p=\gamma(0)$. Then the map

$$
j: \mathbb{R} \times H \rightarrow M, \quad j(t, q)=\gamma_{q}(t)
$$

is the desired isometry. So in particular, $\mathbb{R} \times H$ with the product metric is globally hyperbolic which implies that $H$ is a complete Riemannian manifold (cf. [1, p. 65]).

## Appendix

Proposition. Let $U$ be an open convex domain in $\mathbb{R}^{n}$ and $f: U \rightarrow \mathbb{R}$ a continuous function. Assume that for any $p \in U$ there is a smooth lower support function $f_{p}$ at $p$, i.e. $f_{p}$ is defined in a neighborhood of $p$ with $f_{p} \leqslant f$ and $f_{p}(p)=f(p)$, and that $\left\|D\left(f_{p}\right)_{p}\right\| \leqslant L$. Then $f$ is Lipschitz with Lipschitz constant L, i.e. for all $x, y \in U$ we have

$$
|f(x)-f(y)| \leqslant L \cdot\|x-y\| .
$$

Proof. Case 1: $n=1$. Then $U$ is an open interval $I$. Assume that there are $x, y \in I$ such that $|f(x)-f(y)|>L \cdot|x-y|$. We may assume $x<y$ and $f(x)<f(y)$ (otherwise replace $f$ with $-b$ ). Let $l: \mathbb{R} \rightarrow \mathbb{R}$ be an affine function with $l(x)>f(x), l(y)<f(y)$, and slope $l^{\prime}=L_{0}>L$. Let $p=\sup \{t \in[x, y]$; $l(t)>f(t)\}$. Then $l(p)=f(p)=f_{p}(p)$ and $f_{p}(t) \leqslant f(t)<l(t)$ for $t<p$. But this implies $f_{p}^{\prime}(p) \geqslant L_{0}>L$ which is a contradiction.

Case 2: $n>1$. For given $x, y \in U$ let $c$ be the line $c(t)=x+t(y-x)$, restricted to the interval $I=c^{-1}(U)$. Now apply Case 1 to the function $f \circ c$ : $I \rightarrow \mathbb{R}$.

Added in proof. Very recently, G. J. Galloway succeeded in proving the splitting theorem without assuming the timelike geodesical completeness, but using the existence and regularity of certain maximal spacelike hypersurfaces obtained by R. Bartnik; see G. J. Galloway, The Lorentzian splitting theorem without completeness assumption, Preprint, University of Miami, Coral Gables, FL, 1987.

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