# HOLOMORPHIC MAPPINGS OF COMPLEX MANIFOLDS 

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## 1. Introduction

Schwarz's lemma, as formulated by Pick, can be stated as follows: Every holomorphic map $f$ of the open unit disk $D$ into itself is distance-decreasing with respect to the Poincáre-Bergman metric $d s^{2}$, i.e. $f *\left(d s^{2}\right) \leq d s^{2}$, where the equality holding at one point of $D$, implies that $f$ is an isometry. Bochner and Martin proved in their book [2] the following generalization of Schwarz's lemma to higher dimensions. Let $D_{n}$ be the $n$-dimensional open unit ball. If $f$ is a holomorphic map of $D_{m}$ into $D_{n}$ such that $f(0)=0$, then $f(z) \leq z$ for all $z \in D_{m}$. In other words, every holomorphic map of $D_{m}$ into $D_{n}$ is distance-decreasing with respect to the Bergman metric $d s_{D_{m}}^{2}$ and $d s_{D_{n}}^{2}$ of $D_{m}$ and $D_{n}$ respectively. Koranyi proved [9] that if $M$ is a hermitian symmetric space of non-compact type with the Bergman metric $d s^{2}$, and $f$ is a holomorphic map of $M$ into itself, then $f^{*}\left(d s^{2}\right) \leq k d s^{2}$, where $k$ denotes the rank of $M$. This is another generalized Schwarz's lemma. Ahlfors was the first to generalize Schwarz's lemma by essentially considering the curvature; his result can be stated as the following: Let $M$ be a Riemann surface with hermitian metric $d s_{M}^{2}$ whose Gaussian curvature is bounded above by a negative constant $-B$, and $D$ the unit disk in $C$ with an invariant metric $d s_{D}^{2}$ whose Gaussian curvature is a negative constant $-A$, then every holomorphic map $f: D \rightarrow M$ satisfies $f^{*}\left(d s^{2}\right) \leq \frac{A}{B} d s_{D}^{2}$. Kobayashi generalized this result to higher dimensional case in his recent paper [6].

Recently Chern [5] has shown that a holomorphic map $f$ of $D_{n}$ into a $n$ dimensional hermitian Einstein manifold $N$ with scalar curvature less than or equal to $-2 n(n+1)$ is volume-decreasing. Kobayashi [8] generalized the result of Chern to the case of a holomorphic mapping $f$ from a more generalized domain $M$ (of dimension $n$ ) into a more general image manifold $N$ (of dimension $n$ ). This paper is devoted to a generalization of Schwarz's lemma as well as one of Chern's results in his paper [5] concerning a Laplacian formula for the ratio function of top volume elements between hermitian manifolds.

Let $M$ and $N$ be two hermitian manifolds with dimensions $m$ and $n$ respectively. The simplest metrical invariant is the ratio of distance and volume elements. In this case, in which we have different dimensions, the ratio of volume elements of $M$ and $N$ is called the ratio of the intermediate volume elements. This ratio has an important effect on the holomorphic mappings between $M$ and $N$, and the study of these ratios will be the main topic of this paper.

In this paper we assume some basic knowledge in hermitian geometry, such as complex and hermitian structures on a complex manifold and the notions of connection and curvature of a holomorphic vector bundle, which can be found in [3], [4], or the author's dissertation [7]. In $\S 2$ we discuss the Ricci and scalar curvatures in detail, and introduce the Laplacian on an hermitian manifold. $\S \S 3$ and 4 are devoted to defining some general elementary symmetric functions, such as the ratio function of distances, the ratio function of intermediate volume elements and, in the case of manifolds of the same dimension, the ratio function of top value elements, and to deriving Laplacian formulas for those functions. These formulas are significant because from them we find that the Laplacian of those functions will involve connection form and curvatures only; hence they will give us the desired geometrical result.

From the result of $\S 4$ and by using the maximum principle, we will discuss the properties of holomorphic mappings from a compact hermitian manifold into an hermitian manifold with certain curvature restrictions. The main result in the last section will be stated as follows:

Let $F: D_{m} \rightarrow N$ be an holomorphic mapping, where $D_{m}$ is the unit $m$-ball in $C^{m}$ with the standard Kaehler metric, and $N$ is an $n$-dimensional hermitian manifold with negative constant holomorphic sectional curvature ( $=-2 m(m+1)$ ). Then $f$ is distance-decreasing.

## 2. Foundations of hermitian geometry

From now on we will fix the notations throughout this paper as follows:

$$
\begin{gathered}
1 \leq i, j, k, l, i_{1}, \cdots, i_{\mu} \leq m \\
1 \leq \alpha, \beta, \gamma, \eta, \alpha_{1}, \cdots, \alpha_{\mu} \leq n \\
m+1 \leq \sigma, \rho, \tau \leq n \\
1 \leq r, s, t \leq \mu
\end{gathered}
$$

In this section we will discuss curvature tensors in detail and derive a Laplacian formula for a real-valued $C^{\infty}$-function on an hermitian manifold.

First we consider curvature tensors. Let $M$ be a hermitian manifold of dimension $m$, that is, let $M$ be a complex manifold with a hermitian structure
on its tangent bundle $T$. The hermitian metric defines an hermitian inner product in $T$ and a type $(1,0)$ connection which is invariant under parallel translation. In the following we will discuss the tangent and cotangent bundles. We know that a unitary frame is an ordered set of $m$ tangent vectors $\left\{s_{i}\right\}$ with the same origin such that $\left\langle s_{i}, s_{j}\right\rangle=\delta_{i j}$. The dual basis to a unitary frame $\left\{s_{i}\right\}$ in the cotangent bundle is a unitary coframe which consists of $m$ complex-valued linear differential forms $\theta^{i}$ of type $(1,0)$ such that

$$
\begin{equation*}
d s_{t g}^{2}=\sum_{i} \theta^{i} \bar{\theta}^{i} \tag{2.1}
\end{equation*}
$$

Let $B$ be the bundle of all unitary frames of $M$. Then the forms $\theta^{i}$ are forms in $B$. It is known [6] that there exist connection forms $\theta_{j}^{i}$ in $B$ such that

$$
\begin{equation*}
d \theta^{i}=\sum_{j} \theta^{j} \wedge \theta_{j}^{i}+\Theta^{i} \tag{2.2}
\end{equation*}
$$

with

$$
\begin{gather*}
\theta_{j}^{i}+\bar{\theta}_{i}^{j}=0  \tag{2.3}\\
\Theta^{i}=\frac{1}{2} \sum_{j, k} T_{j k}^{i} \theta^{j} \wedge \theta^{k} \tag{2.4}
\end{gather*}
$$

The $\theta^{i}$ 's are called torsion forms, and $T_{j k}^{i}$ 's torsion tensors.
By taking exterior derivative of (2.2) and using (2.2) again, we find ([3] or [4])

$$
\begin{equation*}
d \Theta^{i}=\sum_{j} \theta^{j} \wedge \Theta_{j}^{i}-\sum_{j} \Theta^{j} \wedge \theta_{j}^{i} \tag{2.5}
\end{equation*}
$$

where

$$
\begin{equation*}
\Theta_{j}^{i}=d \theta_{j}^{i}+\sum_{k} \theta_{k}^{i} \wedge \theta_{j}^{k} \tag{2.6}
\end{equation*}
$$

are the $(1,1)$ curvature forms satisfying

$$
\begin{equation*}
\Theta_{j}^{i}+\Theta_{i}^{j}=0 \tag{2.7}
\end{equation*}
$$

Those (1,1) forms $\Theta_{j}^{i}$ can be written as

$$
\begin{equation*}
\Theta_{j}^{i}=\frac{1}{2} \sum_{k, l} R_{j k l}^{i} \theta^{k} \wedge \bar{\theta}^{l} \tag{2.8}
\end{equation*}
$$

In terms of the coefficients $R_{j k l}^{i}$, the symmetry properties (2.7) can be written as

$$
\begin{equation*}
R_{j k l}^{i}=\bar{R}_{i k l}^{j} . \tag{2.9}
\end{equation*}
$$

The quantities $R_{j k l}^{i}$ are functions in $B$ and constitute the curvature tensors of the hermitian metric.

We define the Ricci tensors $R_{k l}$ as

$$
\begin{equation*}
R_{k l}=\sum_{i} R_{i k l}^{i}=\bar{R}_{l k} \tag{2.10}
\end{equation*}
$$

and the scalar curvature as

$$
\begin{equation*}
R=\sum_{k} R_{k k} \tag{2.11}
\end{equation*}
$$

where $R$ is a real-valued function in $M$.
Now we derive a Laplacian formula for a real-valued $C^{\infty}$ function on a hermitian manifold. Let $M$ and $N$ be two arbitrary hermitian manifolds with dimensions $m$ and $n$ respectively. In the previous paragraphs we have already defined

$$
\theta^{i}, \theta_{j}^{i}, \Theta^{i}, \Theta_{j}^{i}, R_{j k l}^{i}, R_{k l}, R
$$

in $M$. Here we list the corresponding ones for $N$ :

$$
\omega^{\alpha}, \omega_{\beta}^{\alpha}, \Omega^{\alpha}, \Omega_{\beta}^{\sim}, S_{\beta_{r}}^{\alpha}, S_{r \eta}, S .
$$

Naturally all the general formulas discussed above remain valid by inserting these quantities.

In the following we shall make no distinction between $\theta^{i}$ and $\theta_{i}, \theta_{j}^{i}$ and $\theta_{i j}$ etc., because we will select unitary fields and unitary coframe fields.

Let $u$ be any real-valued $C^{\infty}$ function on $M$. By choosing a local unitary frame field and restricting ourselves to a neighborhood of $M$, we let

$$
\begin{equation*}
d u=\sum_{i}\left(u_{i} \theta_{i}+\bar{u}_{i} \bar{\theta}_{i}\right) \tag{2.12}
\end{equation*}
$$

Taking the exterior derivative of (2.12) we get

$$
\begin{gather*}
\sum_{i}\left(d u_{i}-\sum_{k} u_{k} \theta_{k i}\right) \wedge \theta_{i}+\sum_{i}\left(d \bar{u}_{i}-\sum_{k} \bar{u}_{k} \bar{\theta}_{k i}\right) \wedge \bar{\theta}_{i}  \tag{2.13}\\
+\sum_{i} u_{i} \theta_{i}+\sum_{i} \bar{u}_{i} \bar{\theta}_{i}=0 .
\end{gather*}
$$

Then let

$$
\begin{equation*}
d u_{i}-\sum_{k} u_{k} \theta_{k i}=\sum_{k}\left(u_{i k}^{\prime} \theta_{k}+u_{i k} \bar{\theta}_{k}\right) \tag{2.14}
\end{equation*}
$$

By considering the type of forms in (2.13), we have

$$
\begin{equation*}
\sum_{i, k} u_{i k}^{\prime} \theta_{k} \wedge \theta_{i}+\sum_{i} u_{i} \theta_{i}=0 \tag{2.15}
\end{equation*}
$$

or

$$
\begin{equation*}
\sum_{i} d\left(u_{i} \theta_{i}\right)=\sum_{i, k} u_{i k} \bar{\theta}_{k} \wedge \theta_{i} \tag{2.16}
\end{equation*}
$$

From (2.12) we get

$$
\begin{equation*}
d^{\prime} u=\sum_{i} u_{i} \theta_{i} \tag{2.17}
\end{equation*}
$$

Using (2.16) we have

$$
\begin{equation*}
d^{\prime} d^{\prime \prime} u=-d d^{\prime} u=\sum_{i, j} u_{i j} \theta_{i} \wedge \bar{\theta}_{j} \tag{2.18}
\end{equation*}
$$

We define the Laplacian of $u$ to be

$$
\begin{equation*}
\Delta u=4 \sum_{i} u_{i i} . \tag{2.19}
\end{equation*}
$$

If $u>0$, we find

$$
\begin{equation*}
\Delta \log u=\frac{1}{u} \Delta u-\frac{4}{u^{2}} \sum_{i} u_{i} \bar{u}_{k} \tag{2.20}
\end{equation*}
$$

Remark. Here we should emphasize that $u_{i k}$ are the coefficients of $\bar{\theta}_{k}$ in the equation of $d u_{i}$, and that $u_{i}$ are the coefficients of $\theta_{i}$ in the equation of $d u$. This remark will be very useful later.

## 3. General elementary symmetric functions

We shall now derive a formula for the Laplacian of general elementary symmetric functions of a mapping $f: M \rightarrow N$ between two hermitian manifolds. By the latter we mean a symmetric function of the eigenvalues of the linear transformation $B=A^{t} \bar{A}$ where

$$
\begin{equation*}
A=f^{*}: T_{N}^{*} \rightarrow T_{N}^{*} \tag{3.1}
\end{equation*}
$$

in the linear map on the cotangent spaces induced by $f$.
First of all, we will discuss the algebraic situation.
Suppose $V, W$ be two complex vector spaces of dimensions $m$ and $n$ respectively. Let $\left\{e_{1}, \cdots, e_{m}\right\}$ be a basis for $V$, and $\left\{f_{1}, \cdots, f_{n}\right\}$ a basis for $W$. If $\mu \leq \min \{m, n\}$, then $\left\{\left(e_{i_{1}} \wedge e_{i_{2}} \wedge \cdots \wedge e_{i_{\mu}}\right) \mid 1 \leq i_{1}<i_{2}<\cdots<i_{\mu} \leq m\right\}$ forms a basis for $\Lambda^{\mu} V$, and $\left\{\left(f_{\alpha_{1}} \wedge f_{\alpha_{2}} \wedge \cdots \wedge f_{\alpha_{\mu}}\right) \mid 1 \leq \alpha_{1}<\alpha_{2}<\cdots<\alpha_{\mu} \leq n\right\}$ forms a basis for $\Lambda^{\mu} W$, where $\Lambda^{\mu} V, \Lambda^{\mu} W$ are exterior powers of vector spaces $V$ and $W$ respectively.

Let $A: W \rightarrow V$ be a linear transformation such that

$$
\begin{align*}
& A\left(f_{\alpha_{\tau}}\right)=\sum_{i} a_{\alpha_{\tau} i} e_{i}  \tag{3.2}\\
& { }^{t} \bar{A}\left(e_{i_{s}}\right)=\sum_{\alpha} a_{\alpha i_{s}} f_{\alpha} \tag{3.3}
\end{align*}
$$

We repeat that $1 \leq r, s, t \leq \mu$. Define

$$
\begin{equation*}
\Lambda^{\mu} A\left(f_{\alpha_{1}} \wedge \cdots \wedge f_{\alpha_{\mu}}\right) \equiv A\left(f_{\alpha_{1}}\right) \wedge \cdots \wedge A\left(f_{\alpha_{\mu}}\right) \tag{3.4}
\end{equation*}
$$

Then

$$
\begin{equation*}
\Lambda^{\mu} A\left(f_{\alpha_{1}} \wedge \cdots \wedge f_{\alpha_{\mu}}\right)=\sum_{i_{1} \wedge \cdots \wedge i_{\mu}} \operatorname{det}\left(a_{\alpha r i_{s}}\right) e_{i_{1}} \wedge \cdots \wedge e_{i_{\mu}} \tag{3.5}
\end{equation*}
$$

For the purpose of simplification, we adopt the following notations:
Let $I$ denote $\left(i_{1}, \cdots, i_{\mu}\right)$ with $1 \leq i_{1}<i_{2}<\cdots<i_{\mu} \leq m, J$ denote ( $\alpha_{1}, \cdots, \alpha_{\mu}$ ) with $1 \leq \alpha_{1}<\alpha_{2}<\cdots<\alpha_{\mu} \leq n$, and $D_{I}^{J} \operatorname{denote} \operatorname{det}\left(a_{\alpha \tau i_{s}}\right)$ where $\alpha_{r}$ 's are components of $J$ and $i_{s}$ 's are components of $I$. Then

$$
\begin{equation*}
\Lambda^{\mu} A\left(f_{\alpha_{1}} \wedge \cdots \wedge f_{a_{\mu}}\right)=\sum_{I} D_{I}^{J}\left(e_{i_{1}} \wedge \cdots \wedge e_{i_{\mu}}\right) \tag{3.6}
\end{equation*}
$$

where the sum is over all possible I's.
Define

$$
\begin{equation*}
B=A^{t} \bar{A} \tag{3.7}
\end{equation*}
$$

Then

$$
\begin{equation*}
B\left(e_{i_{r}}\right)=\sum_{i_{t}} \sum_{\alpha_{s}} a_{\alpha_{s} j_{l}} \bar{a}_{\alpha_{s} i_{r}} e_{j_{\iota}} . \tag{3.8}
\end{equation*}
$$

The elementary symmetric function of $f$ is defined to be the elementary symmetric function of $B$, in other words, is defined to be the trace of $\Lambda^{\mu} B$ which will be denoted as $u$. Before going further, we shall state the following lemma in linear algebra; its proof is in [7].

Lemma. If $P$ is any matrix with n-columns, and $Q$ any matrix with n-rows, then any t-rowed determinant of the matrix $P Q$ is equal to a sum of terms, each of which is the product of a t-row determinant of $P$ and a $t$-columned determinant of $B$.

By using this lemma, we can write

$$
\begin{equation*}
u=\operatorname{tr} \Lambda^{\mu} B=\sum_{I} \sum_{J}\left|D_{I}^{J}\right|^{2} \tag{3.9}
\end{equation*}
$$

where

$$
\begin{equation*}
D_{I}^{J}=\operatorname{det}\left(a_{\alpha r i_{s}}\right) \tag{3.10}
\end{equation*}
$$

We shall now discuss some examples of the ratios of distance, volume elements and intermediate volume elements.

Let $\left\{\theta_{i}\right\},\left\{\omega_{\alpha}\right\}$ be the unitary coframe fields of two hermitian manifolds $M$ and $N$ respectively, and $f: M \rightarrow N$ be a holomorphic mapping such that

$$
\begin{equation*}
f^{*} \omega_{\alpha}=\sum_{i} a_{\alpha i} \theta_{i} . \tag{3.11}
\end{equation*}
$$

Then

$$
\begin{equation*}
f^{*}\left(\omega_{\alpha} \wedge \bar{\omega}_{\alpha}\right)=\sum_{i, j} a_{a i} \bar{a}_{\alpha j} \theta_{i} \wedge \bar{\theta}_{j} \tag{3.12}
\end{equation*}
$$

From now on we shall omit $f^{*}$ without fear of confusion. It is clear that the $m \times m$ matrix

$$
\begin{equation*}
\left(b_{i j}\right)=\left(\sum_{\alpha} a_{\alpha i} \bar{a}_{\alpha j}\right) \tag{3.13}
\end{equation*}
$$

is self-adjoint and positive semi-definite. Let $\lambda_{1}, \cdots, \lambda_{m}$ be its eigenvalues. Then

$$
\begin{gather*}
\lambda_{i} \geq 0, \quad \text { for } 1 \leq i \leq m,  \tag{3.14}\\
d s_{N}^{2}=\sum_{\alpha} \omega_{\alpha} \bar{\omega}_{\alpha}=\lambda_{1} \theta_{1} \bar{\theta}_{1}+\cdots+\lambda_{m} \theta_{m} \bar{\theta}_{m}  \tag{3.15}\\
\leq\left(\lambda_{1}+\cdots+\lambda_{m}\right)\left(\theta_{1} \bar{\theta}_{1}+\cdots+\theta_{m} \bar{\theta}_{m}\right) .
\end{gather*}
$$

Let

$$
\begin{equation*}
u=\sum_{i} \lambda_{i}=\sum_{\alpha, i} a_{\alpha i} \bar{a}_{\alpha i} \tag{3.16}
\end{equation*}
$$

Then

$$
\begin{equation*}
d s_{N}^{2}<u d s_{m}^{2} \tag{3.17}
\end{equation*}
$$

where $u$ is a well defined function on $M$, and gives, by looking at (3.17), a good control of the ratio between the distances of $M$ and $N$.

By raising (3.12) to the $m^{\text {th }}$ power, let

$$
\begin{equation*}
v=\frac{f^{*}\left(\sum_{\alpha} \omega_{\alpha} \wedge \bar{\omega}_{\alpha}\right)^{m}}{\left(\sum_{\alpha} \theta_{i} \wedge \bar{\theta}_{i}\right)^{m}}=\sum_{J}\left|D^{J}\right|^{2}, \quad \text { for } m \leq n \tag{3.18}
\end{equation*}
$$

where

$$
\begin{equation*}
D^{J}=\operatorname{det}\left(a_{\alpha_{i} j}\right) \tag{3.19}
\end{equation*}
$$

Then $v$ is geometrically the ratio of intermediate volume elements.
For the special case in which $m=n, v$ is the ratio of volume elements and

$$
\begin{equation*}
v=D \bar{D} \tag{3.20}
\end{equation*}
$$

where

$$
\begin{equation*}
D=\operatorname{det}\left(a_{\alpha \beta}\right) \tag{3.21}
\end{equation*}
$$

## 4. Holomorphic mappings of hermitian manifolds

As before, let $f: M \rightarrow N$ be a holomorphic map such that

$$
f^{*} \omega_{\alpha}=\sum_{i} a_{a i} \theta_{i},
$$

or in short,

$$
\begin{equation*}
\omega_{\alpha}=\sum_{i} a_{\alpha i} \theta_{i} \tag{4.1}
\end{equation*}
$$

By taking exterior derivatives

$$
\begin{equation*}
d \omega_{\alpha}=\sum_{i}\left(d a_{\alpha i} \wedge \theta_{i}+a_{\alpha i} d \theta_{i}\right) \tag{4.2}
\end{equation*}
$$

and making use of (2.2) and its corresponding formulas for $d \theta_{i}$, we get

$$
\begin{equation*}
\sum_{i}\left(d a_{\alpha i}-\sum_{j} a_{\alpha j} \theta_{j i}+\sum_{\beta} a_{\beta i} \omega_{\alpha \beta}\right) \wedge \theta_{i}+\sum_{i} a_{\alpha i} \theta_{i}-\Omega_{\alpha}=0 . \tag{4.3}
\end{equation*}
$$

Since the torsion terms are of type $(2,0)$, it allows us to put

$$
\begin{equation*}
d a_{\alpha i}-\sum_{j} a_{\alpha j} \theta_{j i}+\sum_{\beta} a_{\beta i} \omega_{\alpha \beta}=\sum_{k} a_{\alpha i k} \theta_{k}, \tag{4.4}
\end{equation*}
$$

where $a_{\text {aik }}$ satisfies the relations

$$
\begin{equation*}
\sum_{i, k} a_{\alpha i k} \theta_{k} \wedge \theta_{i}+\sum_{i} a_{\alpha i} \theta_{i}-\Omega_{\alpha}=0 \tag{4.5}
\end{equation*}
$$

Similarly, we take the exterior derivative of (4.4) and obtain

$$
\begin{gather*}
-\sum_{j} d a_{\alpha j} \wedge \theta_{j i}-\sum_{j} a_{\alpha j} d \theta_{j i}+\sum_{\beta} d a_{\beta i} \wedge \omega_{\alpha \beta}+\sum_{\beta} a_{\beta i} d \omega_{\alpha \beta}  \tag{4.6}\\
=\sum_{k} d a_{\alpha i k} \wedge \theta_{k}+\sum_{k} a_{\alpha i k} d \theta_{k} .
\end{gather*}
$$

Using (4.4), (2.2), (2.5) and simplifying we get

$$
\begin{align*}
\sum_{k}\left(d a_{\alpha i k}\right. & \left.-\sum_{j} a_{\alpha i j} \theta_{j k}-\sum_{j} a_{\alpha j k} \theta_{j i}+\sum a_{\beta i k} \omega_{\alpha \beta}\right) \wedge \theta_{k}  \tag{4.7}\\
& =-\sum_{k} a_{\alpha i k} \Theta_{k}+\sum_{j} a_{\alpha j} \Theta_{j i}-\sum_{\beta} a_{\beta i} \Omega_{\alpha \beta}
\end{align*}
$$

Suppose that

$$
\begin{equation*}
-\sum_{k} a_{a i k} \Theta_{k}=\sum_{k, l} a_{a i k l} \theta_{k} \wedge \theta_{l} \tag{4.8}
\end{equation*}
$$

$$
\begin{equation*}
\sum_{j} a_{a j} \Theta_{j i}-\sum_{\beta} a_{A i} \Omega_{\alpha \beta}=\sum_{k, l} b_{a i k l} \theta_{k} \wedge \bar{\theta}_{l} \tag{4.9}
\end{equation*}
$$

where

$$
\begin{equation*}
b_{a i k l}=\frac{1}{2}\left(\sum_{j} a_{a j} R_{j i k l}-\sum_{\beta, \gamma, \eta} a_{\beta i} a_{r k} \bar{a}_{\eta k} S_{a \beta_{r \eta}}\right) . \tag{4.10}
\end{equation*}
$$

Then

$$
\begin{align*}
d a_{\alpha i k} & -\sum_{j} a_{\alpha i j} \theta_{j k}-\sum_{j} a_{\alpha j k} \theta_{j i}+\sum_{\beta} a_{\beta i k} \omega_{\alpha \beta}  \tag{4.11}\\
& =\sum_{i} a_{\alpha i k l} \theta_{l}+\sum_{l} b_{\alpha i k l} \bar{\theta}_{l}
\end{align*}
$$

which will be useful in proving

Theorem 4.1. Let $f: M \rightarrow N$ be a holomorphic mapping as in (4.1), and $u$ the function on $M$ defined in (3.16). Then

$$
\begin{align*}
\frac{1}{2} \Delta u= & 2 \sum_{\alpha, i, k}\left|a_{\alpha i k}\right|^{2}+\sum_{\alpha, i, j} a_{\alpha i} \bar{a}_{\alpha j} R_{j i}  \tag{4.12}\\
& -\sum_{i, k} \sum_{\alpha, \beta, \gamma, \eta} \bar{a}_{\alpha i} \bar{a}_{\beta i} a_{r k} \bar{a}_{\eta k} S_{\alpha \beta \gamma \eta}
\end{align*}
$$

or

$$
\begin{equation*}
\frac{1}{2} \Delta \log u=\frac{2}{u}\left(\sum_{\alpha, i, j} \bar{a}_{\alpha i} a_{\alpha j} R_{j i}-\sum_{i, k} \sum_{\alpha, \beta, \gamma, \eta} \bar{a}_{\alpha i} a_{\beta i} a_{r k} \bar{a}_{\eta k} S_{\alpha \beta \gamma \eta}\right) \tag{4.13}
\end{equation*}
$$

Proof. By taking exterior derivative of (3.16) and using (4.4),

$$
\begin{align*}
d u= & \sum_{\alpha, i} \bar{a}_{\alpha i}\left(\sum_{j} a_{\alpha i j} \theta_{j}+\sum_{j} a_{\alpha j} \theta_{j i}-\sum_{\beta} a_{\beta i} \omega_{\alpha \beta}\right)  \tag{4.14}\\
& +\sum_{\alpha, i} a_{\alpha i}\left(\sum_{j} \bar{a}_{\alpha i j} \bar{\theta}_{j}+\sum_{j} \bar{a}_{\alpha j} \bar{\theta}_{j i}-\sum_{\beta} \bar{a}_{\beta i} \bar{\omega}_{\alpha \beta}\right)
\end{align*}
$$

Since we select unitary coframes, we have

$$
\begin{align*}
& \theta_{i j}+\bar{\theta}_{j i}=0 \\
& \omega_{\alpha \beta}+\bar{\omega}_{\alpha \beta}=0 \tag{4.15}
\end{align*}
$$

By interchanging appropriate indices, $i$ into $j$ and $\alpha$ into $\beta$, and using (4.15) we get

$$
\begin{equation*}
d u=\sum_{\alpha, i, j} \bar{a}_{\alpha i} a_{\alpha i j} \theta_{j}+\sum_{\alpha, i, j} a_{\alpha i} \bar{a}_{\alpha i j} \bar{\theta}_{j} \tag{4.16}
\end{equation*}
$$

Comparing (4.16) with (2.20) we have

$$
\begin{equation*}
u_{j}=\sum_{\alpha, i} \bar{a}_{\alpha i} a_{\alpha i j} \tag{4.17}
\end{equation*}
$$

Remark. We should realize that (4.17) is obvious even without using (4.15), since $u$ and $u_{j}$ are functions on $M$, and $\theta_{i j}, \omega_{\alpha \beta}$, etc. are pull-back forms of the principal bundle, which should cancel out each other automatically. Then by considering the remark at the end of $\S 2$ and comparing (4.16) with (2.2), (4.17) can be written out in one step.

Now making use of the remark at the end of $\S 2$, and (4.10), (4.11), and applying the above technique to

$$
\begin{equation*}
d u_{j}=\sum_{\alpha, i} \bar{a}_{\alpha i} d a_{\alpha i j}+\sum_{\alpha, i} a_{\alpha i j} d \bar{a}_{\alpha i} \tag{4.18}
\end{equation*}
$$

we get

$$
\begin{equation*}
u_{j k}=\sum_{\alpha, i} a_{\alpha i j} \bar{a}_{\alpha i k}+\sum_{\alpha, i} a_{\alpha i} b_{\alpha i j k} \tag{4.19}
\end{equation*}
$$

It follows that

$$
\begin{align*}
u_{k k}= & \sum_{\alpha, i}\left|a_{\alpha i k}\right|^{2}+\sum_{\alpha, i} a_{\alpha i} b_{\alpha i k k} \\
= & \sum_{\alpha, i}\left|a_{\alpha i k}\right|^{2}+\frac{1}{2} \sum_{\alpha, i, j} a_{\alpha i} a_{\alpha j} R_{j i k k}  \tag{4.20}\\
& -\frac{1}{2} \sum_{i, k, r, \eta} \bar{a}_{\alpha i} a_{\beta i} a_{r k} \bar{a}_{\eta k} S_{\alpha \beta r \eta} .
\end{align*}
$$

From (2.27) and (2.28) we have (4.12) and (4.13), since

$$
\begin{gather*}
\sum_{k} R_{j i k k}=R_{j i}  \tag{4.21}\\
u \sum_{\alpha, i, k}\left|a_{\alpha i k}\right|^{2}-\sum_{k} u_{k} \bar{u}_{k}=0
\end{gather*}
$$

In general, we can prove
Theorem 4.2. Let $f: M \rightarrow N$ be a holomorphic mapping as in (4.1), and $u$ the general elementary symmetric function of $f$ as in (3.10). Then

$$
\begin{equation*}
\frac{1}{2} \Delta u=E+R(\xi, \xi)-S(\xi, \xi), \tag{4.22}
\end{equation*}
$$

where $E$ is a non-negative quantity, $R(\xi, \xi)$ and $S(\xi, \xi)$ are sums of Ricci tensors of $M$ and curvature tensors of $N$ respectively as given by (4.32) and (4.33).

Proof. In this case $u=\sum_{I} \sum_{J}\left|D_{I}^{J}\right|^{2}$. Define

$$
\begin{gather*}
u_{I}^{J}=D_{I}^{J} \bar{D}_{I}^{J},  \tag{4.23}\\
D_{I, k}^{J}=\sum_{s=1}^{\mu} \sum_{B} \varepsilon_{\alpha_{2} \cdots \alpha_{\mu}}^{\beta_{1} \ldots \beta_{\mu}} a_{\beta_{1} i_{1}} \cdots\left(a_{\beta_{s} i_{s} k}\right) \cdots a_{\beta_{\mu} i_{\mu}},  \tag{4.24}\\
D_{I, k l}^{J}=\sum_{s=1}^{\mu} \sum_{B} \varepsilon_{\alpha_{1} \cdots \alpha_{\mu}}^{\beta_{1} \ldots \beta_{\mu} \mu} a_{\beta_{1} i_{1}} \cdots\left(b_{\beta_{s} i_{s} k l}\right) \cdots a_{\beta_{\mu} i_{\mu}}, \tag{4.25}
\end{gather*}
$$

where $a_{\beta_{s} i_{s} k}$ and $b_{\beta_{s} i_{s} k l}$ are defined in (7.4) and (7.10). Further define

$$
\begin{align*}
& {\left[D_{I}^{J}\right]_{j}^{i_{s}}=\sum_{B} \varepsilon_{\alpha_{1} \cdots \alpha_{\mu}}^{\beta_{1} \cdots \beta_{\mu} a_{\beta_{1} i_{1}} \cdots\left(a_{\beta_{s}}\right) \cdots a_{\beta_{s} i_{s}}},}  \tag{4.26}\\
& {\left[D_{I}^{J}\right]_{\alpha}^{\beta_{s}}=\sum_{B} \varepsilon_{\alpha_{1} \cdots \alpha_{\mu}}^{\beta_{1} \cdots \beta_{\mu} a_{\beta_{1} i_{1}} \cdots\left(a_{\alpha i_{s}}\right) \cdots a_{\beta_{s} i_{s}} .}} \tag{4.27}
\end{align*}
$$

Using the same technique as in Theorem 4.1, it is easy to see

$$
\begin{align*}
\frac{1}{2} \Delta u_{I}^{J}= & 2 \sum_{k=1}^{m}\left|D_{I, k}^{J}\right|^{2}+\bar{D}_{I}^{J} \sum_{s=1}^{\mu} \sum_{j}\left[D_{I}^{J}\right]_{j}^{i_{s}} R_{j i_{s}}  \tag{4.28}\\
& -D_{I}^{J} \sum_{k} \sum_{r, \eta}\left\{\sum_{s, \alpha}\left[D_{I}^{J}\right]_{\alpha}^{\beta_{s}} a_{\tau k} \bar{a}_{\eta k} S_{\alpha \beta_{s 7}}\right\}
\end{align*}
$$

or

$$
\begin{align*}
\frac{1}{2} \Delta u= & 2 \sum_{I, J} \sum_{k}\left|D_{I, k}^{J}\right|^{2}+\sum_{I, J}\left\{\bar{D}_{I}^{J} \sum_{s, j}\left[D_{I}^{J}\right]_{j}^{i_{s}} R_{j i_{s}}\right.  \tag{4.29}\\
& -\bar{D}_{I}^{J} \sum_{k, \gamma, \eta}\left(\sum_{s, \alpha}\left[D_{I}^{J}\right]_{a}^{\beta_{s}} a_{\tau k} \bar{a}_{\eta k} S_{\alpha \beta_{s} 7 \eta}\right) .
\end{align*}
$$

Letting

$$
\begin{gather*}
E=2 \sum_{I, J} \sum_{k}\left|D_{I, k}^{J}\right|^{2}  \tag{4.30}\\
R(\xi, \xi)=\sum_{T, J}\left\{\bar{D}_{I}^{J} \sum_{s, j}\left[D_{I}^{J}\right]_{j}^{i_{s}} R_{j i_{s}}\right\},  \tag{4.31}\\
S(\xi, \xi)=\sum_{I, J}\left\{\bar{D}_{I}^{J} \sum_{k, r, \eta}\left(\sum_{s, \alpha}\left[D_{I}^{J}\right]_{\alpha}^{\beta_{s}} a_{r k} \bar{a}_{\eta k} S_{\alpha \beta_{s, \eta}}\right)\right\}, \tag{4.32}
\end{gather*}
$$

we can then see that (4.29) gives exactly what we need.
If we define

$$
\begin{gather*}
D_{, k}^{J}=\sum_{i=1} \sum_{B} \varepsilon_{\alpha_{1} \ldots \alpha_{m}}^{\beta_{1} \ldots \beta_{m}} a_{\beta_{1} 1} \cdots\left(a_{\beta_{i} i k}\right) \cdots a_{\beta_{m} m}  \tag{4.33}\\
{\left[D^{J}\right]_{j}^{i}=\sum_{B} \varepsilon_{\alpha \alpha_{1} \ldots a_{m}}^{\beta_{1} \cdots \beta_{m}} a_{\beta_{11}} \cdots\left(a_{\beta_{i j}}\right) \cdots a_{\beta_{m} m}}  \tag{4.34}\\
{\left[D^{J}\right]_{\alpha}^{\beta_{i}}=\sum_{B} \varepsilon_{\alpha_{1} \cdots \alpha_{m}}^{\beta_{1} \ldots \beta_{m}} a_{\beta_{11}} \cdots\left(a_{\alpha j}\right) \cdots a_{\beta_{m} m}} \tag{4.35}
\end{gather*}
$$

then we have
Corollary 4.3. For the function $v$ defined in (3.18),

$$
\begin{equation*}
\frac{1}{2} \Delta v=\sum_{k} \sum_{J}\left|D_{, k}^{J}\right|^{2}+v R-\sum_{J} \bar{D}^{J}\left\{\sum_{k, \gamma, \eta}\left(\sum_{\alpha, i}\left[D^{J}\right]_{\alpha}^{\beta_{i}} S_{\alpha \beta_{i r \eta}}\right) a_{r k} \bar{a}_{\eta k}\right\} \tag{4.36}
\end{equation*}
$$

Proof. By looking at (4.29) we only have to explain the second term on the right hand side. In this case we can easily see that

$$
\begin{equation*}
\sum_{J} \bar{D}^{J} \sum_{i, j}\left[D^{J}\right]_{j}^{i} R_{j i}=\sum_{J} \bar{D}^{J} D^{J} R=v R . \tag{4.37}
\end{equation*}
$$

By substituting (4.37) in (4.29) and changing $D_{I}^{J} ' s$ into $D^{J \prime} s$, (4.36) is then obvious.

Corollary 4.4. For the function $v$, defined in (3.19), which is the ratio of volume elements of $M$ and $N$ in the case $m=n$,

$$
\begin{equation*}
\frac{1}{4} \Delta v=\sum_{k} D_{k} \bar{D}_{k}+\frac{v}{2}\left(R-\sum_{\alpha, \beta, k} a_{\alpha k} \bar{a}_{\beta k} S_{\alpha \beta}\right), \tag{4.38}
\end{equation*}
$$

where

$$
\begin{equation*}
D_{k}=\sum_{B} \varepsilon_{1, \cdots, n}^{i_{1} \cdots i_{n}} a_{i_{1} 1} \cdots\left(a_{i_{j} j k}\right) \cdots a_{i_{n} n} \tag{4.39}
\end{equation*}
$$

Proof. By the same reason as that for deriving (4.37),

$$
\begin{align*}
\bar{D} \sum_{k, \gamma, \eta} \sum_{\alpha, \beta}[D]_{\alpha}^{\beta} S_{\alpha \beta r \eta} a_{r k} \bar{a}_{\eta k} & =\sum_{r, \eta, k} \bar{D} D S_{r \eta} a_{r k} \bar{a}_{\eta k}  \tag{4.40}\\
& =v \sum_{r, \eta, k} S_{r \eta} a_{r k} \bar{a}_{\eta k} .
\end{align*}
$$

Then (4.38) is obvious, and is the formula proved by Chern in [5].

## 5. Applications

So far what we have derived is the general Laplacian formula for symmetric functions. In this section we use (4.36) and (4.12) to prove some geometrical conclusions. The first conclusion will be a generalization of a theorem in Chern's paper [5], and another one is a theorem, in a different form from that appearing in Kobayashi's recent paper [6], concerning distance-decreasing mappings.

Now we assume $f: \mathrm{M} \rightarrow N$ is a holomorphic mapping between hermitian manifolds with dimensions $m$ and $n$ respectively.

Definition 5.1. $f$ is degenerate at a point $p$ in $M$ if $u$ vanishes at $p$, where $u=\sum_{J}\left|D^{J}\right|^{2}$ is the ratio function of intermediate volume elements as in (3.18).

Geometrically, $f$ degenerate means that the induced linear map $f_{*}$ on the tangent space at $p$ is not univalent.

Definition 5.2. $f$ is totally degenerate if $u$ vanishes identically.
Definition 5.3. An Einstein manifold $M$ of dimension $m$ has $K$-exhaustion property if $M$ is exhausted by a sequence of open submanifolds

$$
\begin{equation*}
M_{1} \subset M_{2} \subset \cdots \subset M \tag{5.1}
\end{equation*}
$$

whose closures $\bar{M}_{a}$ are compact and such that
(1) to each $a=1,2, \cdots$ there is a $C^{\infty}$ function $v_{a} \geq 0$, defined in $M_{a}$, satisfying the inequality

$$
\begin{equation*}
\frac{1}{2} \Delta v_{a} \leq \frac{R}{m}+K \exp v_{a} \tag{5.2}
\end{equation*}
$$

where $K$ is a given positive constant;
(2) $v_{a}\left(p_{b}\right) \rightarrow \infty$ if $p_{b}$ is a divergent sequence of points in $M_{a}$ (an infinite sequence of points $p_{b}$ in $M_{a}$ is called divergent, if every compact open set of $M_{a}$ contains only a finite number of points of the sequence).

This kind of manifold is not uncommon because all bounded symmetric domains and especially the $m$-dimensional unit ball have this property.

Now we proceed to establish the main theorems of this section.
Theorem 5.1. For compact $M$, let $R$ be the scalar curvature of $M$, and $S_{\alpha \beta r y}$ the curvature tensor of $N$.
(1) If $R>0$ and the curvature transformations are non-positive on all kinds of tensors, then $f$ is totally degenerate.
(2) If $R<0$ and the curvature transformations are non-negative on all kinds of tensors, then $f$ is degenerate at some point of $M$.

Proof. (1) Since $M$ is compact, $u$ attains its maximum at a point $p_{0}$ in $M$. If $u$ is not identically zero, $u\left(p_{0}\right)>0$. Since $D^{J}$ is a function on $M$, covariant differentiation is the same as the usual one and hence $D^{J}{ }_{, k}$ is zero at a maximum point, i.e. $D^{J}{ }_{, k}\left(p_{0}\right)=0$. With our hypothesis and formula (4.36) it is clear that

$$
\begin{equation*}
(\Delta u)_{p_{0}}>0 . \tag{5.3}
\end{equation*}
$$

However, since $p_{0}$ is the maximum point of $u$, we have

$$
\begin{equation*}
(\Delta u)_{p_{0}} \leq 0 \tag{5.4}
\end{equation*}
$$

which is a contradiction.
Assertion (2) is proved similarly by the consideration of the minimum of $u$.
Theorem 5.2. Let $f: M \rightarrow N$ be a holomorphic mapping, where $M$ is an Einstein manifold of dimension $m$ with K-exhaustion property, and $N$ is an hermitian manifold of dimension $n$ with negative constant holomorphic sectional curvatures ( $=-K$, for some positive constant $K$ ). Then

$$
\begin{equation*}
v=\log u \leq v_{a} \tag{5.5}
\end{equation*}
$$

where $u=\sum_{\alpha, i} a_{a i} \bar{a}_{\alpha i}$ as in (3.16).
Proof. Let $E$ be the open set defined by $E=\left\{x \in M \mid v>v_{a}\right\}$. In $E$ we have $u=\sum_{a, i} a_{\alpha i} \bar{a}_{\alpha i} \geq 0$. From (4.31) we have

$$
\begin{equation*}
\frac{1}{2} \Delta \log u=\frac{1}{u}\left[\sum_{a, i, j} \bar{a}_{\alpha i} a_{\alpha j} R_{j i}-\sum_{i, k} \sum_{\alpha, \beta, r, \eta} \bar{a}_{\alpha i} a_{\beta i} a_{\tau k} \bar{a}_{\eta j k} S_{\alpha \beta T \eta}\right] \tag{5.6}
\end{equation*}
$$

The condition for $M$ to be Einsteinian is

$$
\begin{equation*}
R_{j i}=\frac{R}{m} \delta_{j i} \tag{5.7}
\end{equation*}
$$

in which case (5.6) reduces to

$$
\begin{equation*}
\frac{1}{2} \Delta v=\frac{R}{m}-\frac{1}{n} \sum_{i, k} \sum_{\alpha, \beta, \gamma, \eta} \bar{a}_{\alpha i} a_{\beta i} a_{\gamma k} \bar{a}_{\eta k} S_{\alpha \beta r \eta} \tag{5.8}
\end{equation*}
$$

The condition for $N$ to have negative constant holomorphic sectional curvature ( $=-K$ ) is

$$
\begin{aligned}
& \sum_{i, k} \sum_{\alpha, \beta, T, \eta} \bar{a}_{\alpha i} a_{\beta i} a_{\tau k} \bar{a}_{\eta k} S_{\alpha \beta r \eta} \\
& \quad=-K \sum_{i, k}\left(\sum_{\alpha} \bar{a}_{\alpha i} a_{\alpha k}\right)\left(\sum_{\beta} a_{\beta i} \bar{a}_{\beta k}\right)<-K u^{2}
\end{aligned}
$$

which reduces (5.8) to

$$
\begin{equation*}
\frac{1}{2} \Delta v \geq \frac{R}{m}+K \exp v \tag{5.9}
\end{equation*}
$$

From (5.2) and (5.9) we have

$$
\begin{equation*}
\frac{1}{2} \Delta\left(v-v_{a}\right) \geq K\left(\exp v-\exp v_{a}\right) \tag{5.10}
\end{equation*}
$$

Since the exponential function is increasing, $\Delta\left(v-v_{a}\right)>0$ in $E$, and $v-v_{a}$ cannot have a maximum in $E$. Hence $v-v_{a}$ must approach its least upper bound on a sequence of points $p_{1}, p_{2}, \ldots$ tending to the boundary of $E$. This sequence cannot have a limit point $p_{0}$ in $M_{a}$, for otherwise at the point $p_{0}, v-v_{a}>0$, and $p_{0}$ would belong to $E$ and be a maximum for $v-v_{a}$. It cannot be divergent either, for otherwise $v_{a} \rightarrow \infty$ and $v$ is bounded. These show that $E$ is empty, and hence $v \leq v_{a}$.
q.e.d.

In Chern's paper [5] we know that the unit ball $D_{m}$ is an Einstein manifold with $2 n(n+1)$-exhaustion property such that

$$
\begin{equation*}
v_{\rho}=\log \left(\frac{1-r^{2}}{\rho^{2}-r^{2}}\right)^{m+1} \tag{5.11}
\end{equation*}
$$

where $0<r<\rho<1$. As $\rho \rightarrow 1, v_{\rho} \rightarrow 0$ giving us
Theorem 5.3. Let $f: D_{m} \rightarrow N$ be a holomorphic mapping, where $D_{m}$ is the unit m-ball in $C^{m}$ with the standard Kaehler metric, and $N$ is a n-dimensional hermitian manifold with negative constant holomorphic sectional curvature $(=-2 m(m+1))$. Then $f$ is distance-decreasing.

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