# EXOTIC NEGATIVELY CURVED STRUCTURES ON CAYLEY HYPERBOLIC MANIFOLDS 

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#### Abstract

We construct examples of closed negatively curved manifolds $M$ which are homeomorphic but not diffeomorphic to Cayley locally symmetric spaces. Given $\epsilon>0$, we can construct such an $M$ with sectional curvatures all in $[-4-\epsilon,-1]$.


## 1. Introduction

Margulis [16] discovered a strengthening of Mostow's strong rigidity theorem [17] to a phenomenon called Archimedean superrigidity valid for lattices in semisimple Lie groups $G$ of real rank bigger than or equal to two. (Here $G$ is assumed to be centerless and to contain no compact normal subgroup other than 1.) Later Corlette [5] proved a version of superrigidity for lattices in the automorphism groups of quarternion hyperbolic spaces or the Cayley hyperbolic plane. It is known that superrigidity fails for other real rank 1 situations; i.e., for lattices in the automorphism groups of the real or complex hyperbolic spaces. Stronger versions of Corlette superrigidity were later proven by Jost and Yau [13] and Mok, Siu and Yeung [19]. A consequence of these superrigidity theorems is that if $M$ and $N$ are homeomorphic closed negatively curved manifolds and the universal cover of $M$ is either a quaternionic hyperbolic space $\mathbb{H} \mathbf{H}^{n}, n \geq 2$, or the Cayley hyperbolic plane $\mathbb{O} \mathbf{H}^{2}$, then $M$ and $N$ are isometric up to a scaling of the metric on either of them by a

[^0]constant (the isometry being the unique harmonic map in the homotopy class of the homeomorphism [6], [10]) under any of the following three extra conditions on $N$ :

1. The curvature operator of $N$ is nonpositive [5].
2. The complexified sectional curvatures of $N$ are nonpositive [19].
3. The sectional curvatures of $N$ are pointwise $\frac{1}{4}$-pinched; i.e., lie in a closed interval $\left[-4 a_{x},-a_{x}\right]$ where $a_{x}>0$ and $x \in N$ (cf [11] and [24]).

In fact, each of Conditions 1 and 3 independently imply Condition 2.
In [9], homeomorphic pairs of closed negatively curved $n$-manifolds $M$ and $N$ are constructed where the universal cover $\widetilde{M}$ of $M$ is the complex hyperbolic space $\mathbb{C} \mathbf{H}^{m}$ (and $n=2 m$ ) but $M$ and $N$ are not diffeomorphic; indeed, given $\epsilon>0$, such pairs of $M$ and $N$ were constructed so that $\widetilde{M}=\mathbb{C} \mathbf{H}^{m}$ and the sectional curvatures of $N$ are "almost $\frac{1}{4}$-pinched", i.e., lie in $[-4-\epsilon,-1]$.

It was conjectured in [9] that such examples could be constructed where the universal cover $\widetilde{M}$ of $M$ is either the quaternionic hyperbolic plane $\mathbb{H} \mathbf{H}^{2}$ or $\mathbb{O} \mathbf{H}^{2}$. We prove here this conjecture for the case where $\widetilde{M}=\mathbb{O} \mathbf{H}^{2}$. The smooth manifolds $N$ are the connected sum $M \# \Sigma^{16}$ where $\Sigma^{16}$ is the unique smooth manifold homeomorphic but not diffeomorphic to the 16 -dimensional round sphere $S^{16}$. The case when $\widetilde{M}=\mathbb{H} \mathbf{H}^{n}$ is treated separately in [2] where we use a different technique to show that the manifolds $M \# \Sigma^{4 n}$ admit metrics of negative curvature but get a weaker result without the "almost $1 / 4$-pinched" conclusion. However, we believe that the method used in this paper could be used to get the pinching result for the case $\widetilde{M}=\mathbb{H} \mathbf{H}^{2}$. A corollary of our construction is that Condition 1 or 2 on $N$ in the superrigidity theorems mentioned above is optimal in a sense, i.e., neither of them can be replaced by the condition that the sectional curvatures of $N$ are nonpositive. On the other hand, in view of superrigidity under Condition 3 on $N$, we note that $\epsilon$ cannot be 0 in any of our examples.

We conclude this introduction with an outline of the paper. There are two problems that must be addressed: (1) How to put a negatively curved metric on $M \# \Sigma^{16}$. (2) How to show that $M \# \Sigma^{16}$ is not diffeomorphic to $M .\left(M \# \Sigma^{16}\right.$ is clearly homeomorphic to $M$.) We solve problem (1) in the next section and problem (2) in the final section of the paper. Broadly speaking, we follow the pattern established in
[8] and [9]. But the difficulties encountered are more formidable and require substantial modifications to the arguments in [9].

To solve the first problem, we construct a 1-parameter family $b_{\gamma}($, of Riemannian metrics on $\mathbb{R}^{16}$ indexed by $\gamma \in[e,+\infty)$ which satisfy the following properties:
(i) The sectional curvatures of $b_{\gamma}$ lie in the closed interval [ $-4-$ $\epsilon(\gamma),-1]$ where $\epsilon(\gamma)>0$ and $\epsilon(\gamma) \rightarrow 0$ as $\gamma \rightarrow+\infty$.
(ii) The ball of radius $\gamma$ about 0 in $\left(\mathbb{R}^{16}, b_{\gamma}\right)$ is isometric to a ball of radius $\gamma$ in real hyperbolic space $\mathbb{R} \mathbf{H}^{16}$.
(iii) The complement of the ball of radius $\gamma^{2}$ about 0 in $\left(\mathbb{R}^{16}, b_{\gamma}\right)$ is isometric to the complement of a ball of radius $\gamma^{2}$ in $\mathbb{O} \mathbf{H}^{2}$.

To construct these metrics we make use of the explicit description of the Riemannian curvature tensor for $\mathbb{O} \mathbf{H}^{2}$ given in [4]. We use this result together with [9, Lemma 3.18] to put an "almost $\frac{1}{4}$-pinched" negatively curved Riemannian metric on $M \# \Sigma^{16}$ provided $M$ has sufficiently large injectivity radius. Here $M$ is a closed, orientable Cayley hyperbolic manifold. This injectivity radius condition is satisfied when we pass to sufficiently large finite sheeted covers of $M$ since $\pi_{1}(M)$ is a residually finite group.

The second problem (i.e., to show that $M$ and $M \# \Sigma^{16}$ are not diffeomorphic) is reduced via Kirby-Siebenmann smoothing theory and using Mostow's strong rigidity theorem [17] together with its topological analogue [7] to showing that the group homomorphism

$$
\theta_{16}=\left[S^{16}, \mathrm{Top} / O\right] \underset{\phi^{*}}{\longrightarrow}[M, \operatorname{Top} / O]
$$

is monic where $\phi: M \rightarrow S^{16}$ is a degree 1 map. Now, a result of Okun [21] shows that $\phi^{*}$ is the initial map in a factoring of

$$
\theta_{16}=\left[S^{16}, \mathrm{Top} / O\right] \underset{\psi^{*}}{\longrightarrow}\left[\mathbb{O} \mathbf{P}^{2}, \mathrm{Top} / O\right]
$$

where $\psi: \mathbb{O} \mathbf{P}^{2} \rightarrow S^{16}$ is a degree 1 map and $\mathbb{O} \mathbf{P}^{2}$ is the Cayley projective plane. Hence it suffices to show that $\psi^{*}$ is monic. This is done by making delicate use of some calculations of Toda [22] on the stable homotopy groups of spheres.

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## 2. Tapering between $\mathbb{O} \mathbf{H}^{2}$ and $\mathbb{R} \mathbf{H}^{16}$

We begin with a brief description of the Cayley hyperbolic plane $0 \mathbf{H}^{2}$.

The Cayley numbers, denoted by $\mathbb{O}$, is an 8-dimensional non associative division algebra over the real numbers. It has a multiplicative identity 1 and a positive definite bilinear form $\langle$,$\rangle whose associated$ norm $\|\|$ is multiplicative, i.e., $\| u v\|=\| u\|\|v\|$. Every element $u \in \mathbb{O}$ can be written as $\alpha+u_{0}$ when $\alpha$ is real and $\left\langle\alpha, u_{0}\right\rangle=0$. The conjugation map $u \mapsto \bar{u}:=\alpha-u_{0}$ is an antiautomorphism, i.e., $\overline{(u v)}=\bar{v} \bar{u}$ for all $u, v \in \mathbb{O}$. Moreover, $u \bar{u}=\|u\|^{2}$ and one has the following identities which can be checked easily: $\langle u v, w\rangle=\langle\bar{v} \bar{u}, \bar{w}\rangle=\langle\bar{v}, \bar{w} u\rangle=\langle v, \bar{u} w\rangle$ for $u, v, w \in \mathbb{O}$.

On $\mathbb{O}^{2}=\mathbb{O} \times \mathbb{O}$, one has the positive definite bilinear form given by $\left\langle\left(u_{1}, u_{2}\right),\left(v_{1}, v_{2}\right)\right\rangle=\left\langle u_{1}, v_{1}\right\rangle+\left\langle u_{2}, v_{2}\right\rangle$ for $u_{1}, v_{1}, u_{2}, v_{2} \in \mathbb{O}$. The set $D=\left\{u \in \mathbb{O}^{2} \mid\langle u, u\rangle<1\right\}$ equipped with the metric given by formula (20.4) in [17, p. 144] is a model for the Cayley hyperbolic plane.

It is convenient for us to consider the set $\mathbb{O}^{2}$ itself as the underlying set for the Cayley hyperbolic plane $\mathbb{O} \mathbf{H}^{2}$ equipped with the metric gotten by scaling the above metric from $D$ to $\mathbb{O}^{2}$. This enables one to identify $\mathbb{O}^{2}$ with $T_{0}\left(\mathbb{O} \mathbf{H}^{2}\right)$, the tangent space to $\mathbb{O} \mathbf{H}^{2}$ at the origin $0 \in \mathbb{O}^{2}$.

The Riemannian metric on the distance sphere $S^{15}$ at distance $t$ from the origin can be described as follows. Firstly, one has the Hopf fibering of $S^{15}$ over $S^{8}$ with $S^{7}$ as fiber. This equips $S^{15}$ with complementary distributions $\eta_{1}, \eta_{2}$ where Whitney sum $\eta_{1} \oplus \eta_{2}$ equals the tangent bundle of $S^{15} . \eta_{1}$ is the 7 -dimensional distribution tangent to the $S^{7}$ fibers and $\eta_{2}$ is the 8-dimensional distribution perpendicular to $\eta_{1}$ (perpendicular with respect to the round metric on $S^{15}$ ). We call the subspace of the tangent space to $S^{15}$ belonging to the distribution $\eta_{1}$ "the vertical subspace" and the subspace belonging to $\eta_{2}$ "the horizontal subspace". The induced Riemannian metric $\langle$,$\rangle on S^{15}$ is then,

$$
\langle X, X\rangle=b^{2} X \cdot X, \quad\langle U, U\rangle=a^{2} U \cdot U \quad \text { and } \quad\langle X, U\rangle=0
$$

where $X \in \eta_{1}, U \in \eta_{2}, a=\sinh t, b=\sinh t \cosh t$ and "." is the inner
product with respect to the round metric on $S^{15}$. For brevity, we denote the distance sphere with the above metric by $S_{a, b}^{15}$.

Fix a smooth function $\phi:(0,+\infty) \rightarrow[0,+\infty)$ such that $\phi(t) \geq 0$ for all $t \in[0,+\infty)$. We put a Riemannian metric on $S^{15} \times(0,+\infty)$ using the function $\phi(t)$ as follows. The foliations $S^{15} \times t$ and $x \times(0,+\infty)$ are required to be perpendicular where $x \in S^{15}$ and $t \in(0,+\infty)$. We set $|N|=1$ where $N=\frac{\partial}{\partial t}$ and $t$ is the second coordinate variable in the product structure $S^{15} \times(0,+\infty)$. We require that the induced metric on $S^{15} \times t$ is $S_{a, b}^{15}$ where $a=\sinh t$ and $b=\sinh t \cosh \phi(t)$. This Riemannian manifold is denoted $S^{15} \times^{\phi}(0,+\infty)$. Notice that when $\phi(t)=t$ for all $t \in(0,+\infty), S^{15} \times^{\phi}(0,+\infty)$ is the punctured Cayley hyperbolic plane $\mathbb{O} \mathbf{H}^{2}-*$. Our object is to calculate the sectional curvatures of $S^{15} \times^{\phi}(0,+\infty)$.

Let $P$ be a real 2-plane tangent to $S^{15} \times^{\phi}(0,+\infty)$. Let $\tau$ denote the angle made by the plane $P$ with the distance sphere $S^{15}$. Then $P$ is spanned by vectors $\{u, \cos (\tau) v+\sin (\tau) N\}$ where vectors $\{u, v\}$ are tangent to the $S^{15} \times t$ foliation and satisfy $|u|=|v|=1$ and $(u \cdot v)=0$ where \| denotes the norm and $(\cdot)$ denotes the inner product with respect to the metric on $S^{15} \times{ }^{\phi}(0,+\infty)$. If $\bar{K}(P)$ denotes the sectional curvature of the plane $P$ in $S^{15} \times{ }^{\phi}(0,+\infty)$, then we have

$$
\begin{align*}
\bar{K}(P)= & \bar{K}(u, \cos \tau v+\sin \tau N)  \tag{1}\\
= & \cos ^{2} \tau \bar{K}(u, v)+\sin ^{2} \tau \bar{K}(u, N) \\
& +2 \sin \tau \cos \tau(\bar{R}(u, N) v \cdot u)
\end{align*}
$$

where $\bar{R}$ denotes the Riemann curvature tensor in $S^{15} \times^{\phi}(0,+\infty)$.
We shall denote the vectors tangent to $S^{15}$ and lying in the vertical subspace by symbols $X, Y$ and those lying in the horizontal subspace by symbols $U, V$. If $\sigma$ is the angle between the vector $u$ and the horizontal subspace and $\alpha$ is the angle between $v$ and the horizontal subspace, let $u=\sin \sigma X+\cos \sigma U$ and $v=\sin \alpha Y+\cos \alpha V$ where $|X|=|Y|=$ $|U|=|V|=1$. (Note $\sigma, \alpha \in[0, \pi / 2]$.) We calculate $\bar{K}(P)$ by explicitly calculating $\bar{K}(u, v), \bar{K}(u, N)$ and $(\bar{R}(u, N) v \cdot u)$ separately.

Before starting off to compute the above terms, we make the following important observations.

Firstly, recall that we identify $\mathbb{O}^{2}$ with the tangent space $T_{0}\left(\mathbb{O} \mathbf{H}^{2}\right)$ where $0=(0,0) \in \mathbb{O}^{2}$. Since $\mathbb{O} \mathbf{H}^{2}$ is a homogeneous space, the group $G$ of isometries of $\mathbb{O} \mathbf{H}^{2}$ acts transitively on $\mathbb{O} \mathbf{H}^{2}$. Further, $G_{0}$ - the
subgroup of $G$ fixing 0 - acts transitively on vectors of unit length in $T_{0}\left(\mathbb{O} \mathbf{H}^{2}\right)$. Therefore one can identify the tangent spaces at other points of $\mathbb{O} \mathbf{H}^{2}$ also with $\mathbb{O}^{2}$. In particular, we identify the vector $N$, normal to the distance spheres, with $(1,0) \in \mathbb{O}^{2}$. This would identify $0 \times \mathbb{O}$ with the horizontal subspace and the subspace of $\mathbb{O} \times 0$ perpendicular to $N$ with the vertical subspace. Thus the vectors $X, Y$ lying in the vertical subspace are purely imaginary Cayley numbers.

Secondly, we note that any $A \in O(16)$ with the property that it induces a permutation of the fibers of the Hopf fibration $S^{15} \rightarrow S^{8}$ determines an isometry $\bar{A}$ of $S^{15} \times{ }^{\phi}(0,+\infty)$ defined by $\bar{A}(x, t)=(A(x), t)$. All $A \in \operatorname{Spin}(9) \subseteq O(16)$ have this property. And since $\operatorname{Spin}(9)$ acts transitively on $S^{15}$, it acts transitively on the fibers of the Hopf fibration. In fact, for any leaf $L$ of Hopf fibration there exists an $A \in \operatorname{Spin}(9)$ such that $A^{2}=I$ and the fixed set of $A$ (acting on $S^{15}$ ) is $L$. To verify this, it suffices to verify it for $L=(\mathbb{O} \times 0) \cap S^{15}$. Here we can define $A(x, y)=(x,-y)(c f .[4, \S \S 3$ and 4$])$. Consequently, the submanifolds $L \times^{\phi}(0,+\infty)$ are totally geodesic in $S^{15} \times^{\phi}(0,+\infty)$. In particular, $L \times^{t}(0,+\infty)$ is totally geodesic in $S^{15} \times^{t}(0,+\infty)=\mathbb{O} \mathbf{H}^{2}-*$.

Thirdly, the Hopf submersion $S^{15} \rightarrow S^{8}$ is indeed a Riemannian submersion $S^{15}(1) \rightarrow S^{8}(r)$ where $S^{15}(1)$ is the sound sphere $S^{15}$ of radius 1 and $S^{8}(r)$ is the round sphere $S^{8}$ of radius $r$. Therefore, one has a Riemannian submersion from $S_{a, b}^{15} \rightarrow S_{a}^{8}$ and more generally a Riemannian submersion from $S^{15} \times^{\phi}(0,+\infty) \rightarrow S^{8} \times^{a}(0,+\infty)$ with fibers $S_{b}^{7}$ where $a, b$ are functions of $t$ described earlier and $S^{7}=S^{7}(1)$. Here $S^{8} \times^{a}(0,+\infty)$ denotes the product of $S^{8}(r)$ warped over $(0,+\infty)$ using $a(t)=\sinh (t)$ for the warping function.

Finally, since $\phi$ is a function from $(0,+\infty)$ to $[0,+\infty)$, the distance sphere $S_{a, b}^{15}$ where $a=\sinh \phi(t)$ and $b=\sinh \phi(t) \cosh \phi(t)$ is indeed the distance sphere in $\mathbb{O} \mathbf{H}^{2}$ at distance $\phi(t)$ from the origin. Scaling this metric on $S_{a, b}^{15}$ throughout by the factor $s=\frac{\sinh t}{\sinh \phi(t)}$, we get a sphere $S_{s a, s b}^{15}$ where $s a=\sinh t$ and $s b=\sinh t \cosh \phi(t)$. But this is the distance sphere in $S^{15} \times{ }^{\phi}(0,+\infty)$ at distance $t$ from the origin.

Since $S^{15}$ is a Riemannian hypersurface in $\mathbb{O} \mathbf{H}^{2}$ and also in $S^{15} \times^{\phi}$ $(0,+\infty)$, in the foregoing calculations, it is important for us to consider the shape operator $\mathcal{L}$ of $S^{15} \subset \mathbb{O} \mathbf{H}^{2}$ and the shape operator $L$ of $S^{15} \subset$ $S^{15} \times{ }^{\phi}(0,+\infty)$ corresponding to the normal vector field $N$ on $S^{15}$. The shape operator is a linear operator acting on each tangent space $T_{p} S^{15}$ at $p \in S^{15}$. We shall compute it by its action on vectors $X, U$ belonging
to the vertical and horizontal subspaces respectively. Using formulas from O'Neill [20], we get the following for $S^{15} \subset S^{15} \times{ }^{\phi}(0,+\infty)$ :

$$
L(X)=(\operatorname{coth} t+\tanh (\phi(t)) \phi(t)) X \text { and } L(U)=\operatorname{coth} t U .
$$

And for $S^{15} \subset \mathbb{O} \mathbf{H}^{2}$ we get,

$$
\mathcal{L}(X)=(\operatorname{coth} \phi(t)+\tanh \phi(t)) X \text { and } \mathcal{L}(U)=\operatorname{coth} \phi(t) U .
$$

Calculation of $\bar{K}(u, N)$.

$$
\begin{aligned}
\bar{K}(u, N)= & \bar{K}(\sin \sigma X+\cos \sigma U, N) \\
= & \sin ^{2} \sigma \bar{K}(X, N)+\cos ^{2} \sigma \bar{K}(U, N) \\
& +2 \sin \sigma \cos \sigma(\bar{R}(U, N) N \cdot X) .
\end{aligned}
$$

Now $X \in \eta_{1}$ and since the fibers $S^{7} \times{ }^{\phi}(0,+\infty)$ are totally geodesic in $S^{15} \times{ }^{\phi}(0,+\infty)$, the vector $\bar{R}(N, X) N \in \eta_{1}$. Therefore $(\bar{R}(U, N) N \cdot X)=$ $-(\bar{R}(N, U) N \cdot X)=-(\bar{R}(N, X) N \cdot U)=0$ since $U \in \eta_{2}$ and $\eta_{1}$ and $\eta_{2}$ are complementary distributions on $S^{15}$. Since $S^{7} \times{ }^{\phi}(0,+\infty)$ are totally geodesic in $S^{15} \times^{\phi}(0,+\infty)$ and since $X \in \eta_{1}$, the curvature $\bar{K}(X, N)$ in $S^{15} \times^{\phi}(0,+\infty)$ is the same as the curvature of the plane $\{X, N\}$ in $S^{7} \times^{b}(0,+\infty)$. Since the metric on $S^{7} \times^{b}(0,+\infty)$ is $d t^{2}+b^{2}\left(d S^{7}\right)^{2}$, the curvature of $\{X, N\}$ is $-\frac{\ddot{b}}{b}$ where $b=\sinh t \cosh \phi(t)$. Therefore,

$$
\bar{K}(X, N)=-\left(1+\tanh \phi(t) \ddot{\phi}(t)+(\phi(t))^{2}+2 \operatorname{coth} t \tanh \phi(t) \phi(t)\right) .
$$

Since $U$ is tangent to $S^{15},[U, N]=0$. And since $S^{15}(1) \times{ }^{\phi}(0 ;+\infty) \rightarrow$ $S^{8}(r) \times^{a}(0,+\infty)$ is a Riemannian submersion, by O'Neill's submersion formula [20], $\bar{K}(U, N)$ is the curvature of the plane spanned by $\{U, N\}$ in $S^{8}(r) \times^{a}(0,+\infty)$. Thus $\bar{K}(U, N)=-\frac{\ddot{a}}{a}=-\frac{\sinh t}{\sinh t}=-1$. Piecing together these components we get,

$$
\begin{align*}
\bar{K}(u, N)= & -1-\sin ^{2} \sigma\left(\ddot{\phi}(t) \tanh \phi(t)+(\dot{\phi}(t))^{2}\right.  \tag{2}\\
& +2 \dot{\phi}(t) \tanh \phi(t) \operatorname{coth} t) .
\end{align*}
$$

## Calculation of $\bar{K}(u, v)$.

Since the distance sphere in $S^{15} \times{ }^{\phi}(0,+\infty)$ is the distance sphere in $\mathbb{O} \mathbf{H}^{2}$ scaled by $s=\frac{\sinh t}{\sinh \phi(t)},\|s u\|=\|s v\|=1$ in $\mathbb{O} \mathbf{H}^{2}$. Using Gauss' equation for the submanifold $S^{15} \subset S^{15} \times{ }^{\phi}(0,+\infty)$,
(a) $\bar{K}(u, v)=K_{S^{15}}(u, v)-\left((L u \cdot u)(L v \cdot v)-(L u \cdot v)^{2}\right)$
where $K_{S^{15}}(u, v)$ is the curvature of $\{u, v\}$ in $S^{15} \subset S^{15} \times^{\phi}(0,+\infty)$. Similarly, using Gauss' equation for $S^{15} \subset \mathbb{O} \mathbf{H}^{2}$,
(b) $\quad \hat{K}(s u, s v)=\mathcal{K}_{S^{15}}(s u, s v)$

$$
-\left(\langle\mathcal{L}(s u), s u\rangle\langle\mathcal{L}(s v), s v\rangle-\langle\mathcal{L}(s u), s v\rangle^{2}\right) .
$$

where $\hat{K}$ is the curvature in $\mathbb{O} \mathbf{H}^{2}$ and $K_{S^{15}}$ is the curvature of $\{s u, s v\}$ in $S^{15} \subset \mathbb{O} \mathbf{H}^{2}$. Since the metrics on the distance spheres in $S^{15} \times^{\phi}(0,+\infty)$ and $\mathbb{O} \mathbf{H}^{2}$ differ by the scaling factor $s$, we have

$$
\frac{1}{s^{2}} \mathcal{K}_{S^{15}}(s u, s v)=K_{S^{15}}(u, v)
$$

Therefore $s^{2} \times(\mathrm{a})-(\mathrm{b})$ and rearranging gives,

$$
\begin{align*}
\bar{K}(u, v)= & \frac{1}{s^{2}}\left(\langle\mathcal{L}(s u), s u\rangle\langle\mathcal{L}(s v), s v\rangle-\langle\mathcal{L}(s u), s v\rangle^{2}\right)  \tag{3}\\
& -\left((L u \cdot u)(L v \cdot v)-(L u \cdot v)^{2}\right)+\frac{1}{s^{2}} \hat{K}(s u, s v) .
\end{align*}
$$

We now calculate the terms on the right-hand side.
A calculation yields,

$$
\begin{aligned}
& \frac{1}{s^{2}}\left(\langle\mathcal{L}(s u), s u\rangle\langle\mathcal{L}(s v), s v\rangle-\langle\mathcal{L}(s u), s v\rangle^{2}\right) \\
& =\frac{1}{s^{2}}\left(\operatorname{coth}^{2} \phi+\sin ^{2} \alpha+\sin ^{2} \sigma+\sin ^{2} \sigma \sin ^{2} \alpha \tanh ^{2} \phi\right. \\
& \quad-\sin ^{2} \sigma \sin ^{2} \alpha\left(\operatorname{coth}^{2} \phi+\tanh ^{2} \phi+2\right)\langle s X, s Y\rangle^{2} \\
& \quad-\cos ^{2} \sigma \cos ^{2} \alpha \operatorname{coth}^{2} \phi\langle s U, s V\rangle^{2} \\
& \left.\quad-2 \sin \sigma \cos \sigma \sin \alpha \cos \alpha\left(1+\operatorname{coth}^{2} \phi\right)\langle s X, s Y\rangle\langle s U, s V\rangle\right)
\end{aligned}
$$

and

$$
\begin{aligned}
&\left((L(u) \cdot u)(L(v) \cdot v)-(L(u) \cdot v)^{2}\right) \\
&= \operatorname{coth}^{2} t+\tanh ^{2} \phi(\dot{\phi})^{2} \sin ^{2} \sigma \sin ^{2} \alpha+\operatorname{coth} t \tanh \phi \dot{\phi}\left(\sin ^{2} \sigma+\sin ^{2} \alpha\right) \\
&-\sin ^{2} \sigma \sin ^{2} \alpha(\operatorname{coth} t+\tanh \phi \dot{\phi})^{2}(X \cdot Y)^{2} \\
&-\cos ^{2} \sigma \cos ^{2} \alpha \operatorname{coth}^{2} t(U \cdot V)^{2} \\
&-2 \sin \sigma \sin \alpha \cos \sigma \cos \alpha \operatorname{coth} t(\operatorname{coth} t+\tanh \phi \dot{\phi})(X \cdot Y)(U \cdot V) .
\end{aligned}
$$

Since $(u \cdot v)=0$, we get, $-\sin \alpha \sin \sigma(X \cdot Y)=\cos \alpha \cos \sigma(U \cdot V)$. Using this identity and the relations $\langle s X, s Y\rangle=(X \cdot Y),\langle s U, s V\rangle=(U \cdot V)$, the term

$$
\begin{aligned}
& \frac{1}{s^{2}}\left(\langle\mathcal{L}(s u), s u\rangle\langle\mathcal{L}(s v), s v\rangle-\langle\mathcal{L}(s u), s v\rangle^{2}\right) \\
& \quad\left((L(u) \cdot u)(L(v) \cdot v)-(L(u) \cdot v)^{2}\right)
\end{aligned}
$$

simplifies to

$$
\begin{align*}
\left(\frac{1}{s^{2}}-1\right)+ & \left(\sin ^{2} \sigma+\sin ^{2} \alpha\right)\left(\frac{1}{s^{2}}-\operatorname{coth} t \tanh \phi \dot{\phi}\right)  \tag{4}\\
& +\sin ^{2} \sigma \sin ^{2} \alpha \tanh ^{2} \phi\left(1-(X \cdot Y)^{2}\right)\left(\frac{1}{s^{2}}-(\dot{\phi})^{2}\right)
\end{align*}
$$

Now, to calculate $\hat{K}(s u, s v)$ we use the description of the Riemann curvature tensor $\hat{R}$ of the Cayley hyperbolic plane $\mathbb{O} \mathbf{H}^{2}$ in [4]. However, the action of the representation of $\operatorname{Spin}(9)$ in $[4]$ is different from the action described in [17]. Indeed, the map $(x, y) \mapsto(x, \bar{y})$ of $\mathbb{O} \times \mathbb{O} \rightarrow \mathbb{O} \times \mathbb{O}$ sends the $\operatorname{Spin}(9)$ action in [17] to that in [4]. In particular, the $\mathbb{O}$ lines in Mostow's description of $\mathbb{O} \mathbf{H}^{2}$ go to $\mathbb{O}$-lines, i.e., 8 -dimensional $\mathbb{R}$-subspaces $\mathcal{R}$ of the tangent space to $\mathbb{O} \mathbf{H}^{2}$ at 0 such that $\hat{K}(\mathcal{P})=-4$ for each 2-plane $\mathcal{P} \subset \mathcal{R}$, in the description in [4] under the above map. Applying this map and using the formula for sectional curvature $\hat{K}$ in [4] yields,

$$
\begin{equation*}
\hat{K}(s u, s v)=-1-3 \cos ^{2} \theta \tag{5}
\end{equation*}
$$

where $\theta$ is the angle between the vector $s v$ and the unique $\mathbb{O}$-line $\mathbb{O} u$ containing the vector su. Since $\mathbb{O} u$ is an 8 -dimensional subspace, it is important to note that

$$
\theta=\min _{\substack{w \in \mathbb{O} u \\\|w\|=1}} \Varangle(w, s v) .
$$

Hence the value $\cos \theta$ is the maximum for all such angles. Putting (4) and (5) into (3) we get
(6) $\bar{K}(u, v)=\left(\frac{1}{s^{2}}-1\right)$

$$
+\left(\sin ^{2} \sigma+\sin ^{2} \alpha\right)\left(\frac{1}{s^{2}}-\operatorname{coth} t \tanh \phi(t) \cdot \dot{\phi}(t)\right)
$$

$$
+\sin ^{2} \sigma \sin ^{2} \alpha \tanh ^{2} \phi(t)\left(1-(X \cdot Y)^{2}\right)\left(\frac{1}{s^{2}}-(\dot{\phi}(t))^{2}\right)
$$

$$
+\frac{1}{s^{2}}\left(-1-3 \cos ^{2} \theta\right)
$$

## Calculation of $(\bar{R}(u, N) v \cdot u)$.

Using the fact that $u=\sin \sigma X+\cos \sigma U$ and $v=\sin \alpha Y+\cos \alpha V$, we first expand out $(\bar{R}(u, N) v \cdot u)$ into 8 terms.

Claim 1. The terms $(\bar{R}(X, N) Y \cdot U),(\bar{R}(U, N) Y \cdot X)$ and $(\bar{R}(X, N) V \cdot X)$ are all zero.

Proof. The vectors $X, Y$ belong to the vertical subspace tangential to the fiber $S^{7}$ of the distance sphere $S^{15}$. And since $S^{7} \times^{\phi}(0,+\infty)$ is totally geodesic in $S^{15} \times{ }^{\phi}(0,+\infty)$ we conclude that the vectors $\bar{R}(X, N) Y$, $\bar{R}(Y, X) N$ and $\bar{R}(X, N) X$ are tangent to $S^{7} \times{ }^{\phi}(0,+\infty)$. Since the vectors $U$ and $V$ belong to the horizontal space, the terms $(\bar{R}(X, N) Y \cdot U)$, $(\bar{R}(Y, X) N \cdot U)$ and $(\bar{R}(X, N) X \cdot V)$ are zero and this completes the proof of Claim 1.
q.e.d.

To analyze the remaining terms we use the Codazzi-Mainardi equation for the submanifold $S^{15}$ of $S^{15} \times^{\phi}(0,+\infty)$. For this we recall the shape operator $L$ acting on the vectors tangent to $S^{15}$. We have $L(X)=(\operatorname{coth} t+\tanh \phi \dot{\phi}) X$ for vectors $X$ in vertical subspace and $L(U)=(\operatorname{coth} t) U$ for vectors $U$ in the horizontal subspace.

For vectors $\alpha, \beta, \gamma \in T_{p} S^{15}$, the Codazzi-Mainardi equation gives $(\bar{R}(\alpha, \beta) \gamma \cdot N)=-\left(\operatorname{Tor}_{L}(\alpha, \beta) \cdot \gamma\right)$ where $\operatorname{Tor}_{L}(\alpha, \beta)=\bar{D}_{\alpha} L(\beta)-$ $\bar{D}_{\beta} L(\alpha)-L([\alpha, \beta])$ where $\bar{D}$ is the Riemannian connection on $S^{15} \times^{\phi}$ $(0,+\infty)$. Applying this equation for the remaining 5 terms gives $(\bar{R}(X$, $N) X \cdot Y)=0,(\bar{R}(U, N) U \cdot V)=0$ and $(\bar{R}(U, N) U \cdot Y)=0$. For the remaining two terms we get $(\bar{R}(U, N) V \cdot X)=(\tanh \phi) \phi\left(\bar{D}_{V} U \cdot X\right)$ and $(\bar{R}(X, N) V \cdot U)=(\tanh \phi) \phi\left(\left(\bar{D}_{V} U-\bar{D}_{U} V\right) \cdot X\right)$. Thus $(\bar{R}(u, N) v$. $u)=\sin \sigma \cos \sigma \cos \alpha(\tanh \phi) \dot{\phi}\left(\left(2 \bar{D}_{V} U-\bar{D}_{U} V\right) \cdot X\right)$. Since the distance
spheres in $S^{15} \times{ }^{\phi}(0,+\infty)$ are gotten by scaling the metric on the distance spheres in $\mathbb{O} \mathbf{H}^{2}$ by a factor $s=\frac{\sinh t}{\sinh \phi(t)}$ they have the same affine connection. We therefore have, for vectors $U, V, X \in T_{p} S^{15}$, that tangential part of $\bar{D}_{V} U=$ tangential part of $\hat{D}_{V} U$ where $\hat{D}$ denotes the Riemannian connection on $\mathbb{O} \mathbf{H}^{2}$. Hence $\left(\bar{D}_{V} U \cdot X\right)=\left(\hat{D}_{V} U \cdot X\right)=$ $s^{2}\left\langle\hat{D}_{V} U, X\right\rangle$ where ( $\left.\cdot\right)$ and $\langle$,$\rangle denote the Riemannian metrics on$ $S^{15} \times^{\phi}(0,+\infty)$ and on $\mathbb{O} \mathbf{H}^{2}$ respectively.

On the other hand, proceeding exactly as above while simplifying $(\bar{R}(u, N) v \cdot u)$, we can show that

$$
\langle\hat{R}(s u, s N) s v, s u\rangle=\sin \sigma \cos \sigma \cos \alpha(\tanh \phi) s^{4}\left\langle 2 \hat{D}_{V} U-\hat{D}_{U} V, X\right\rangle
$$

Now, using relations $\left(\bar{D}_{V} U \cdot X\right)=s^{2}\left(\hat{D}_{V} U, X\right\rangle$ and $\left(\bar{D}_{U} V \cdot X\right)=$ $s^{2}\left\langle\hat{D}_{U} V, X\right\rangle$ deduced above we get the following:

$$
(\bar{R}(u, N) v \cdot u)=\frac{\dot{\phi}}{s^{2}}\langle\hat{R}(s u, s N) s v, s u\rangle .
$$

We then wish to calculate the term on the right-hand side using the formula for the curvature tensor for the Cayley hyperbolic plane $\mathbb{O} \mathbf{H}^{2}$ described in [4]. To be able to do so we must as before first transform our description of $\mathbb{O} \mathbf{H}^{2}$ to the description in [4] via the map $f:(x, y) \mapsto$ $(x, \bar{y})$ of $\mathbb{O}^{2} \rightarrow \mathbb{D}^{2}$. Also our curvature operator $\hat{R}$ is negative of that in [4]. Making these necessary changes and using the formula for the curvature operator $\hat{R}$ in [4, page 52], a calculation yields,

$$
\begin{aligned}
\frac{\dot{\phi}}{s^{2}}\langle\hat{R}(s u, s N) s v, s u\rangle & =-\frac{3 \dot{\phi}}{s} \sin \sigma \cos \sigma \cos \alpha\left\langle s^{2} X \bar{U}, s \bar{V}\right\rangle \\
& =-\frac{3 \dot{\phi}}{s} \sin \sigma \cos \sigma \cos \alpha\left\langle s^{2} U \bar{X}, s V\right\rangle .
\end{aligned}
$$

This together with the fact that

$$
\left\langle s^{2} u \bar{X}, s v\right\rangle=\cos \sigma \cos \alpha\left\langle s^{2} U \bar{X}, s V\right\rangle
$$

yields

$$
\begin{equation*}
(\bar{R}(u, N) v \cdot u)=-\frac{3 \dot{\phi}}{s} \sin \sigma\left\langle s^{2} u \bar{X}, s v\right\rangle=-\frac{3 \dot{\phi}}{s} \sin \sigma \cos \omega \tag{7}
\end{equation*}
$$

where $\omega$ is the angle between the unit length vectors $s^{2} u \bar{X}$ and $s v$. Finally, putting together the calculations (2), (6) and (7) into (1) gives,

$$
\begin{aligned}
\bar{K}(P)= & \cos ^{2} \tau\left(\frac{1}{s^{2}}-1\right. \\
& +\left(\sin ^{2} \sigma+\sin ^{2} \alpha\right)\left(\frac{1}{s^{2}}-\operatorname{coth}(t)(\tanh \phi(t)) \dot{\phi}(t)\right) \\
& +\sin ^{2} \sigma \sin ^{2} \alpha \tanh ^{2} \phi(t)\left(1-(X \cdot Y)^{2}\right)\left(\frac{1}{s^{2}}-\dot{\phi}(t)^{2}\right) \\
& \left.+\frac{1}{s^{2}}\left(-1-3 \cos ^{2} \theta\right)\right) \\
+ & \sin ^{2} \tau\left(-1-\sin ^{2} \sigma\left(\ddot{\phi}(t) \tanh \phi(t)+\dot{\phi}(t)^{2}\right.\right. \\
& +2 \dot{\phi}(t)(\tanh \phi(t)) \operatorname{coth}(t)) \\
- & 6 \sin \tau \cos \tau \sin \sigma \frac{\dot{\phi}}{s} \cos \omega .
\end{aligned}
$$

Combining and regrouping the above term we get
(8) $\bar{K}(P)$

$$
\begin{aligned}
= & -1-3\left(\frac{\cos \tau \cos \omega}{s}+\sin \tau \sin \sigma \dot{\phi}(t)\right)^{2} \\
& -\frac{3 \cos ^{2} \tau}{s^{2}}\left(\cos ^{2} \theta-\cos ^{2} \omega\right) \\
& +\cos ^{2} \tau\left(\sin ^{2} \sigma+\sin ^{2} \alpha\right)\left(\frac{1}{s^{2}}-\operatorname{coth}(t)(\tanh \phi(t)) \dot{\phi}(t)\right) \\
& -\sin ^{2} \tau \sin ^{2} \sigma \ddot{\phi}(t) \tanh \phi(t) \\
& -2 \sin ^{2} \tau \sin ^{2} \sigma \dot{\phi}(t)(\tanh \phi(t) \operatorname{coth}(t)-\dot{\phi}(t)) \\
& +\cos ^{2} \tau \sin ^{2} \sigma \sin ^{2} \alpha \tanh ^{2} \phi(t)\left(1-(X \cdot Y)^{2}\right)\left(\frac{1}{s^{2}}-\dot{\phi}(t)^{2}\right) .
\end{aligned}
$$

We now proceed to choose functions $\phi$ so that given an $\epsilon>0$, the curvature $\bar{K}(P)$ satisfies $-4-\epsilon \leq \bar{K}(P) \leq-1+\epsilon$ for all plane sections $P$ in $S^{15} \times{ }^{\phi}(0,+\infty)$.

Following [9] first fix a smooth function $\psi: \mathbb{R} \rightarrow[0,1]$ such that

$$
\begin{aligned}
\dot{\psi}(t) & \geq 0 \text { for all } t \in[1,2] \\
\psi^{-1}(0) & =(-\infty, 1) \text { and } \\
\psi^{-1}(1) & =[2,+\infty) .
\end{aligned}
$$

For each $c \geq 1$, let $\phi_{c}(t)=\psi\left(\frac{\ln t}{c}\right) t$ for all $t>0$. Therefore $\phi_{c}(t)=0$ for $t \in\left(0, e^{c}\right]$ and $\phi_{c}(t)=t$ for $t \in\left[e^{2 c},+\infty\right)$. As in [9] observe that the following limits hold uniformly in $t$ :

$$
\left\{\begin{array}{l}
\lim _{c \rightarrow+\infty}\left|\ddot{\phi}_{c}(t)\right|=0  \tag{9}\\
\lim \sup _{c \rightarrow+\infty} \dot{\phi}_{c}(t) \leq 1 \\
\lim \sup _{c \rightarrow+\infty}\left(\frac{1}{s^{2}}-(\dot{\phi})^{2}\right) \leq 0 \\
\lim _{c \rightarrow+\infty} \dot{\phi}_{c}(t)\left(\tanh \phi_{c}(t) \operatorname{coth} t-1\right)=0
\end{array}\right.
$$

(The 3rd inequality is a bit different from the corresponding inequality posited in $[9,(2.22)]$.) Now, for the angles $\theta$ and $\omega$ as in (8) we have the following:

Lemma 1. $|\cos \omega| \leq \cos \theta$.
Proof. Recall that $\omega$ is defined by the relation $\cos \omega=\left\langle s^{2} u \bar{X}, s v\right\rangle$. Consider the vector $u^{X}:=\left(s^{2} \sin \sigma, s^{2} \cos \sigma U \bar{X}\right)=s^{2} u \bar{X}$. It is easy to see that $\pm u^{X} \in \mathbb{O} u$ where $\mathbb{O} u$ is the unique $\mathbb{O}$-line containing the vector su. (Note that $X=-\bar{X}$ and hence $(U \bar{X}) X=U(\bar{X} X)=U$.) And obviously $\left\langle \pm u^{X}, s v\right\rangle= \pm \cos \omega$. Since $\theta=\min _{\substack{w \in \mathbb{O} u \\|w|=1}} \Varangle(w, s v)$, we conclude that $|\cos \omega| \leq \cos \theta$.
q.e.d.

Lemma 1 together with formulas (8) and (9) yield the following result when $\phi(t)$ is one of the functions $\phi_{c}(t)$.

Lemma 2. If $t \in\left(0, e^{c}\right]$ and $\phi=\phi_{c}$, then $\bar{K}(P)=-1$. Moreover, for all $t>0$, the following limit holds uniformly in $t$ :

$$
\limsup _{c \rightarrow+\infty} \bar{K}(P)=-1
$$

where $\phi=\phi_{c}$. Also $S^{15} \times^{0}(0,+\infty)$ is $\mathbb{R} \mathbf{H}^{16}$ less a point and $S^{15} \times^{t}$ $(0,+\infty)$ is $\mathbb{O} \mathbf{H}^{2}$ less a point. Hence $S^{15} \times^{\phi_{c}}\left(0, e^{c}\right]$ can be identified
with a closed ball of radius $e^{c}$ in $\mathbb{R} \mathbf{H}^{16}$ with its center deleted. And $S^{15} \times{ }^{\phi_{c}}\left[e^{2 c},+\infty\right)$ can be identified with $\mathbb{O} \mathbf{H}^{2}$ from which an open ball of radius $e^{2 c}$ is deleted.

To obtain a lower bound for $\bar{K}(P)$ we need the following lemma.

## Lemma 3.

(i) The maximum value of $|C \cos y \cos z+D \sin y \sin z|$ is $\max \{|C|$, $|D|\}$, as both $y$ and $z$ vary over $\mathbb{R}$.
(ii) The maximum value of

$$
\begin{aligned}
& B \cos ^{2} \tau \cos ^{2} \theta+(1-B) \cos ^{2} \tau \frac{a}{3} \\
&+2 \sqrt{B} \cos \tau \sin \tau \sin \sigma \cos \omega+\sin ^{2} \tau \sin ^{2} \sigma
\end{aligned}
$$

is 1 , where $\tau \in \mathbb{R}, B \in[0,1], \alpha, \sigma \in\left[0, \pi_{2}\right]$ and angles $\theta$ and $\omega$ are as in (8). And $a=\sin ^{2} \sigma+\sin ^{2} \alpha+\sin ^{2} \sigma \sin ^{2} \alpha$.

Proof. We skip the proof of (i) which can be proved by elementary calculus and proceed directly to prove (ii).
(ii) It is convenient to set

$$
\begin{aligned}
& f=B \cos ^{2} \tau \cos ^{2} \theta+(1-B) \cos ^{2} \tau \frac{a}{3} \\
&+2 \sqrt{B} \cos \tau \sin \tau \sin \sigma \cos \omega+\sin ^{2} \tau \sin ^{2} \sigma
\end{aligned}
$$

We first observe that $f$ is quadratic in $\sin \tau$ and $\cos \tau$. Hence, letting

$$
M=\left(\begin{array}{cc}
B \cos ^{2} \theta+(1-B) \frac{a}{3} & \sqrt{B} \sin \sigma \cos \omega \\
\sqrt{B} \sin \sigma \cos \omega & \sin ^{2} \sigma
\end{array}\right)
$$

we see that $f=(\cos \tau \sin \tau) M\binom{\cos \tau}{\sin \tau}$. By a linear algebra argument it follows that $f \leq 1$ for all $\tau \in \mathbb{R}$ if and only if the maximum eigenvalue of $M$ is less than or equal to 1 ; i.e., $f \leq 1$ for all $\tau \in \mathbb{R}$ if and only if $g=\operatorname{Trace}(M)-\operatorname{Determinant}(M) \leq 1$.

Now

$$
\begin{aligned}
& g=B \cos ^{2} \theta+(1-B) \frac{a}{3}+\sin ^{2} \sigma \\
& \quad-\left(B \cos ^{2} \theta+(1-B) \frac{a}{3}\right) \sin ^{2} \sigma+B \sin ^{2} \sigma \cos ^{2} \omega .
\end{aligned}
$$

Since $g$ is linear in $B$, thinking of $g$ as a function of $B$ and fixing $\theta, \omega$, $\sigma$ and $\alpha$, it is easy to see that the maximum value of $g$ occurs at either $B=0$ or at $B=1$. Therefore, to prove the lemma, it is sufficient to show that $\left.g\right|_{B=0} \leq 1$ and $\left.g\right|_{B=1} \leq 1$. It is easy to see that $\left.g\right|_{B=0} \leq 1$. To show $\left.g\right|_{B=1} \leq 1$, we show equivalently that for all $\tau \in \mathbb{R},\left.f\right|_{B=1} \leq 1$. This fact follows easily by observing that the formula for $\bar{K}(P)$ in (8) for values of $t \geq e^{2 c}$ (in which case $S^{15} \times{ }^{\phi_{c}}\left[e^{2},+\infty\right.$ ) is $\mathbb{O} \mathbf{H}^{2}$ less an open ball of radius $e^{2 c}$ ) reduces to $-\left.3 f\right|_{B=1}-1$ from which it follows that $\left.f\right|_{B=1} \leq 1$ for all $\tau \in \mathbb{R}$.
q.e.d.

Lemma 4. $\liminf _{c \rightarrow+\infty} \bar{K}(P)=-4$.
Proof. It suffices, because of Lemma 2, to show that $\liminf _{c \rightarrow+\infty} \bar{K}(P) \geq$ -4 . Because of (8) and (9), this is equivalent to showing that $\limsup _{c \rightarrow+\infty} v \leq$ 3 , where

$$
\begin{aligned}
v= & 3\left(\frac{\cos ^{2} \tau \cos ^{2} \theta}{s^{2}}+\sin ^{2} \tau \sin ^{2} \sigma\left(\dot{\phi}_{c}(t)\right)^{2}\right. \\
& \left.+\frac{2}{s} \sin \tau \cos \tau \sin \sigma \cos \omega \dot{\phi}_{c}(t)\right) \\
+ & \cos ^{2} \tau\left(\sin ^{2} \sigma+\sin ^{2} \alpha\right)\left(\operatorname{coth} t \tanh \phi_{c}(t) \dot{\phi}_{c}(t)-\frac{1}{s^{2}}\right) \\
+ & 2 \sin ^{2} \tau\left(\sin ^{2} \sigma\right) \dot{\phi}_{c}(t)\left(\tanh \phi_{c}(t) \operatorname{coth} t-\dot{\phi}_{c}(t)\right) \\
& +\cos ^{2} \tau \sin ^{2} \sigma \sin ^{2} \alpha \tanh ^{2} \phi_{c}(t)\left(1-(X \cdot Y)^{2}\right)\left(\left(\dot{\phi}_{c}(t)\right)^{2}-\frac{1}{s^{2}}\right)
\end{aligned}
$$

Let $B=\frac{1}{s^{2}}$ and $x=\psi\left(\frac{\ln t}{c}\right)$ and define $v_{1}$ by

$$
\begin{aligned}
v_{1}= & 3\left(B \cos ^{2} \tau \cos ^{2} \theta+\sin ^{2} \tau\left(\sin ^{2} \sigma\right) x^{2}\right. \\
& +2 \sqrt{B} \sin \tau \cos \tau \sin \sigma(\cos \omega) x) \\
& +\cos ^{2} \tau\left(\sin ^{2} \sigma+\sin ^{2} \alpha\right)(x-B)+2 \sin ^{2} \tau\left(\sin ^{2} \sigma\right) x(1-x) \\
& +\cos ^{2} \tau \sin \sigma \sin ^{2} \alpha \tanh ^{2} \phi_{c}(t)\left(1-(X \cdot Y)^{2}\right)\left(x^{2}-B\right)
\end{aligned}
$$

Using (9), it is easy to see that $v_{1}-v$ converges to 0 uniformly as $c \rightarrow+\infty$. Therefore, it suffices to show that $\limsup _{c \rightarrow+\infty} v_{1} \leq 3$. Since the
maximum values of $\tanh ^{2} \phi_{c}(t)$ and $\left(1-(X \cdot Y)^{2}\right)$ is 1 , it suffices to show that $\limsup v_{2} \leq 3$ where

$$
\begin{aligned}
& c \rightarrow+\infty \\
& v_{2}= 3\left(B \cos ^{2} \tau \cos ^{2} \theta+\sin ^{2} \tau\left(\sin ^{2} \sigma\right) x^{2}\right. \\
&+2 \sqrt{B} \sin \tau \cos \tau \sin \sigma(\cos \omega) x) \\
&+\cos ^{2} \tau\left(\sin ^{2} \sigma+\sin ^{2} \alpha\right)(x-B)+2 \sin ^{2} \tau\left(\sin ^{2} \sigma\right) x(1-x) \\
&+\cos ^{2} \tau \sin ^{2} \sigma \sin ^{2} \alpha\left(x^{2}-B\right)
\end{aligned}
$$

Now, define

$$
\begin{aligned}
v_{3}= & 3\left(B \cos ^{2} \tau \cos ^{2} \theta+\sin ^{2} \tau \sin ^{2} \sigma+2 \sqrt{B} \sin \tau \cos \tau \sin \sigma \cos \omega\right) \\
& +\cos ^{2} \tau\left(\sin ^{2} \sigma+\sin ^{2} \alpha+\sin ^{2} \sigma \sin ^{2} \alpha\right)(1-B) .
\end{aligned}
$$

Since $v_{2}$ is a continuous function of $B, \tau, \sigma, \alpha, \theta, \omega$ and $x$, we have $\lim _{x \rightarrow 1}\left(v_{2}-v_{3}\right)=0$ uniformly in $B \in[0,1], \tau, \omega \in \mathbb{R}$ and $\sigma, \alpha, \theta \in[0, \pi / 2]$.

Since $v_{3} \leq 3$ by Lemma 3(ii), it follows therefore that, given an $\epsilon>0$, there exists a $\delta>0$ such that $x>1-\delta$ implies $v_{2} \leq 3+\epsilon$. Hence to complete the proof of the Lemma, we may assume that $x \leq 1-\delta$. This, together with the fact that $\phi_{c}(t)=x t$ and the specific choice of functions $\phi_{c}(t)$ can be used to show that

$$
\begin{equation*}
\lim _{c \rightarrow+\infty} B=0 \text { uniformly in } t . \tag{10}
\end{equation*}
$$

Now (10), $x \leq 1-\delta$ and Lemma 3(i) together imply that $\limsup _{c \rightarrow+\infty} v_{2} \leq 3$ which completes the proof of the Lemma 4.
q.e.d.

## 3. Detecting exotic smooth structures

The purpose of this section is to prove the following result.
Theorem. Let $M^{16}$ be any closed locally Cayley hyperbolic manifold. Given $\epsilon>0$ then there exists a finite sheeted cover $\mathcal{N}^{16}$ of $M^{16}$ such that the following is true for any finite sheeted cover $N^{16}$ of $\mathcal{N}^{16}$.
(a) $N^{16}$ is not diffeomorphic to $N^{16} \# \Sigma^{16}$.
(b) $N^{16} \# \Sigma^{16}$ supports a negatively curved Riemannian metric whose sectional curvatures are contained in the interval $[-4-\epsilon,-1]$.

Here $\Sigma^{16}$ is the unique closed, oriented smooth 16 -dimensional manifold which is homeomorphic but not diffeomorphic to the sphere $S^{16}$. The existence and uniqueness of $\Sigma^{16}$ is a consequence of the following result which is implicit in [14]. However, for the reader's convenience, we derive it here from the Sullivan-Wall surgery exact sequence.

Proposition. The group of smooth homotopy spheres $\theta_{16}$ is cyclic of order 2.

Proof. We have the following surgery exact sequence from [23]:

$$
0 \rightarrow \theta_{16} \rightarrow \pi_{16}(F / O) \rightarrow L_{16}(O)=\mathbb{Z}
$$

This sequence together with the fact that $\theta_{16}$ is a finite group show that $\theta_{16}$ can be identified with the subgroup $S$ of $\pi_{16}(F / O)$ consisting of all elements having finite order. Next consider the exact sequence

$$
\pi_{16}(O) \xrightarrow{J} \pi_{16}(F) \rightarrow \pi_{16}(F / O) \rightarrow \pi_{15}(O)=\mathbb{Z} .
$$

This sequence and the fact that $\pi_{16}(F)=\pi_{16}^{s}$ is a finite group show that $S$ can be identified with cokernel of $J$. Recall now that Adams [1] proved that $J$ is monic. This result together with the facts that $\pi_{16}(O)=\mathbb{Z}_{2}$ and $\pi_{16}^{s}=\mathbb{Z}_{2} \oplus \mathbb{Z}_{2}$ (cf. [22]) show that $\theta_{16}=\mathbb{Z}_{2}$. q.e.d.

Now, the proof the theorem posited in the beginning of this section follows the pattern established in [8] and [9]. The main result of Okun's thesis [21, Theorem 5.1] gives a finite sheeted cover $\mathcal{N}^{16}$ of $M^{16}$ and a tangential map $f: \mathcal{N}^{16} \rightarrow \mathbb{O} \mathbf{P}^{2}$. And we can arrange that $\mathcal{N}^{16}$ has arbitrarily large preassigned injectivity radius $r$ by taking larger covers since $\pi_{1}\left(M^{16}\right)$ is residually finite. Once $r$ is determined, then this is the manifold $\mathcal{N}^{16}$ posited in the theorem. The argument in $[9$, pp. 69-70] is now easily adapted to yield the following lemma since the other ingredients - Mostow's strong rigidity theorem [17], its topological analogue [7] and Kirby-Siebenmann smoothing theory [14, pp. 25 and 194] remain valid.

Lemma 0. Let $N^{16}$ be any finite sheeted cover of $\mathcal{N}^{16}$. If $N^{16} \# \Sigma^{16}$ is diffeomorphic to $N^{16}$, then $\mathbb{O} \mathbf{P}^{2} \# \Sigma^{16}$ is concordant to $\mathbb{O} \mathbf{P}^{2}$.

The octave projective plane $\mathbb{O} \mathbf{P}^{2}$ is the mapping cone of the Hopf map $p: S^{15} \rightarrow S^{8}$. Let $\phi: \mathbb{O} \mathbf{P}^{2} \rightarrow S^{16}$ be the collapsing map obtained by identifying $S^{16}$ with $\mathbb{O} \mathbf{P}^{2} / S^{8}$ in an orientation preserving way.

Lemma 1. The homomorphism $\phi^{*}:\left[S^{16}, \mathrm{Top} / O\right] \rightarrow\left[\mathbb{O} \mathbf{P}^{2}, \mathrm{Top} / O\right]$ is monic.

Proof. Recall that [ $X$, Top $/ O$ ] is the zeroth cohomology group of $X$ in an extraordinary cohomology theory. By considering the long exact sequence in this theory determined by the pair $\left(\mathbb{O} \mathbf{P}^{2}, S^{8}\right)$ and using the identification of $\mathbb{O} \mathbf{P}^{2}$ with the mapping cone of $\phi$, it is seen that $\phi^{*}$ is monic if and only if

$$
(\Sigma p)^{*}:\left[S^{9}, \mathrm{Top} / O\right] \rightarrow\left[S^{16}, \mathrm{Top} / O\right]
$$

is the zero homomorphism. Here

$$
\Sigma p: S^{16} \rightarrow S^{9}
$$

denotes the suspension of $p$.
To show that $(\Sigma p)^{*}$ is the zero homomorphism consider the following commutative ladder of groups and homomorphisms:

$$
\begin{aligned}
& \theta_{16}=\left[S^{16}, \mathrm{Top} / O\right] \xrightarrow{\alpha}\left[S^{16}, F / O\right] \longleftarrow\left[S^{16}, F\right] \stackrel{J}{\longleftrightarrow}\left[S^{16}, O\right] \\
& (\Sigma p)^{*} \uparrow \quad \uparrow(\Sigma p)^{*} \quad \uparrow(\Sigma p)^{*} \\
& \theta_{9}=\left[S^{9}, \operatorname{Top} / O\right] \longrightarrow\left[S^{9}, F / O\right] \longleftarrow{ }_{\beta}\left[S^{9}, F\right] .
\end{aligned}
$$

The horizontal homomorphisms in this ladder are induced by the natural maps Top $/ O \rightarrow F / O, F \rightarrow F / O$, and $O \rightarrow F$. Now the following three facts used in conjunction with a simple "diagram chase" show that $(\Sigma p)^{*}$ is the zero homomorphism thus proving Lemma 1. q.e.d.

Fact 1. $\alpha$ is monic.
Fact 2. $\beta$ is an epimorphism.
Fact 3. Image $(\Sigma p)^{*} \subseteq$ Image $J$ where $(\Sigma p)^{*}:\left[S^{9}, F\right] \rightarrow\left[S^{16}, F\right]$ and $J:\left[S^{16}, O\right] \rightarrow\left[S^{16}, F\right]$ is the classical $J$-homomorphism.

It remains to verify these Facts. Fact 1 is due to Kervaire and Milnor [14]. A more modern proof is given by observing that $\alpha$ is a homomorphism in Sullivan's surgery exact sequence

$$
\cdots \longrightarrow L_{17}(0) \longrightarrow \theta_{16} \xrightarrow{\alpha} \pi_{16}(F / O) \longrightarrow \ldots
$$

and that $L_{17}(0)=0$. Fact 2 is due to Adams [1] who showed that

$$
J: \pi_{8}(O) \rightarrow \pi_{8}(F)
$$

is monic. (Now consider the homotopy exact sequence for the fibration $O \rightarrow F \rightarrow F / O$.) Fact 3 is a more special result which we proceed to prove.

During this proof we will use Toda's notation [22, p. 189] for special elements in the stable stems $G_{n}=\pi_{n}^{s}$. Recall that $G$ equal the direct sum of the $G_{n}$ is an anti-commutative graded ring with respect to composition as multiplication. First note that the homotopy class

$$
[\Sigma p]=a \sigma+x \in \pi_{7}^{s}
$$

where $a \in \mathbb{Z}$ and $x \in \pi_{7}^{s}$ has odd order. Using this together with the fact that $\pi_{16}^{s}$ has order 4 , we see that Image $(\Sigma p)^{*}$ is generated by the following three elements:

$$
v^{3} \circ a \sigma, \quad \mu \circ a \sigma, \quad \eta \circ \epsilon \circ a \sigma .
$$

And Theorem 14.1 (ii, iii) [22, p. 190] yields that

$$
\begin{aligned}
v^{3} \circ a \sigma & =a\left(v^{2} \circ(v \circ \sigma)\right)=0, \\
\eta \circ \epsilon \circ a \sigma & =a(\eta \circ(\sigma \circ \epsilon))=0, \quad \text { and } \\
\mu \circ a \sigma & =a(\mu \circ \sigma)=-a(\sigma \circ \mu)=-a(\eta \circ \rho)=a(\rho \circ \eta) .
\end{aligned}
$$

Consequently, Image $(\Sigma p)^{*}$ is contained in the subgroup of $\pi_{16}^{s}$ generated by $\rho \circ \eta$. Hence in order to complete the verification of Fact 3 it suffices to show that

$$
\rho \circ \eta \in \text { Image } J \text {. }
$$

To do this let $\eta_{15}: S^{16} \rightarrow S^{15}$ represent the element $\eta \in \pi_{1}^{S}$ and notice that the following digram commutes:

where $J^{\prime}$ denotes the $J$-homomorphism in dimension 15 . We recall that Kervaire and Milnor showed (cf. [18, p. 284]) that Image $J^{\prime}$ is a cyclic group of order 480. Using this fact together with Toda's calculation of $\pi_{15}^{s}$ in [22, p. 189] it is easily seen that

$$
\text { either } \rho \in \text { Image } J^{\prime} \text { or } \rho+\eta \circ k \in \text { Image } J^{\prime} \text {. }
$$

Consequently the above commutative diagram shows that

$$
\text { either } \rho \circ \eta \in \text { Image } J \text { or } \rho \circ \eta+\eta \circ k \circ \eta \in \text { Image } J \text {. }
$$

But Theorem 14.1(i) in [22, p. 190] yields that

$$
\eta \circ k \circ \eta=\eta^{2} \circ k=0
$$

and consequently $\rho \circ \eta \in \operatorname{Image} J$.
But Lemma 1 implies that $\mathbb{O} \mathbf{P}^{2} \# \Sigma^{16}$ is not concordant to $\mathbb{O} \mathbf{P}^{2}$ since the concordance classes of smooth structures on a smooth manifold $X$ are in bijective correspondence with $[X, \operatorname{Top} / O]$ provided $\operatorname{dim} X>4$. Thus assertion (a) of the Theorem is a direct consequence of Lemmas 0 and 1.

Now combining the construction of the previous section with [9, Lemma 3.18] it is seen that there exists a number $r_{16}>0$ (independent of $M^{16}$ ) such that if the injectivity radius of $\mathcal{N}^{16}$ is chosen to be larger than $r_{16}$, then assertion (b) of the Theorem is also true. Since Borel [3] has constructed closed Riemannian manifolds $M^{16}$ whose universal cover $\widetilde{M}^{16}$ is $\mathbb{O} \mathbf{H}^{2}$, Theorem produces the examples claimed in the Introduction.

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