## PAPERS COMMUNICATED

## 65. Theory of Connections in the Generalized Finsler Manifold.

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Recently the theory of connections in the Finsler manifold $F_{n}$ has been developed somewhat by several authors, ${ }^{1}$ in which the parameters of the connections are in general functions not only of positions $x^{\nu}$ but also of directions $d x^{2}$. We shall consider some connections in the generalized Finsler manifold $K_{n}$ similarly as in the Finsler manifold, the parameters of the connections depending on $x^{\nu}, d x^{\nu}, d^{2} x^{\nu}, \ldots \ldots, d^{r} x^{\nu}$.

1. The base-connections. Let $X_{n}$ be a manifold of $n$ dimensions, in which a system of coordinates $x^{\nu}$ is taken, and we express for brevity $n$ differentials $d x^{\nu}$ of the coordinates by ${ }^{(1)} p^{\nu}: d x^{\nu}=\left({ }^{(1)}{ }^{\nu}\right.$. Introduce an arbitrary linear connection $\stackrel{(1)}{U}$, whose parameters $\stackrel{\Gamma}{\Gamma}_{\stackrel{(1)}{\nu}}^{(\alpha)}$ depend upon positions and directions, i.e. $x^{\nu}$ and $d x^{\nu}$, then we get covariant differentials of a line element $\stackrel{(1)}{p}$, which can be considered as a contravariant vector and expressed by the parameters ${\stackrel{(1)}{\Gamma}{ }_{\lambda \mu}}^{(1)}$ of the connection $\stackrel{(1)}{U}$ as follows:

$$
\begin{equation*}
\stackrel{(1)(1)}{\delta} p^{\nu}=d p^{(1)}+\stackrel{(1)}{\Gamma_{\lambda \mu}^{v}} \stackrel{(1)}{(1)} d x^{\nu} . \tag{1.1}
\end{equation*}
$$

We shall write $\stackrel{(2)}{p^{v}}$ for $\stackrel{(1)(1)(1)}{\delta} p^{\nu}$, which is function of $x^{\nu}$ as well as $\stackrel{(1)}{p^{\nu}}$. Let us call this contravariant vector ${ }^{(2)}{ }^{\nu}$ v the line element of the second kind.

We introduce moreover a new linear connection $\stackrel{(2)}{U}^{2}$ ) for $X_{n}$, whose
 nection $U$ is applied to the affinors, which depend on $x^{\nu}$ and the foregoing line elements $\stackrel{(1)}{p^{v}}$ and $\stackrel{(2)}{p}$, and by which we get the line element of the third kind
(1.2) $\quad \stackrel{(3)}{p^{v}}=\delta_{\delta}^{(2)(2)} p^{\nu}=d p^{(2)}+\Gamma_{\lambda_{\mu}}^{(2)} p^{(2)} d x^{\mu}$,

1) E. Noether: Göttinger Nachrichten, math.-phys. Kl. (1918), S. 37-44. L. Berwald: Jahresberichte d. Deutschen Math.-Ver. 34 (1925), S. 213-220. J. L. Synge: Trans. of Amer. Math. Soc. 27 (1925), p. 61-67. J. H. Taylor: Trans. of Amer. Math. Soc. 27 (1925), p. 246-264.
2) $\stackrel{(21}{U}$ is in general different from $\stackrel{(1)}{U}$, but we may also take same one with $\stackrel{(1)}{U}$ as $\stackrel{(2)}{U}$. It is completely analogous for following connections $\stackrel{(b)}{U}$.
$\stackrel{(3)}{p^{\nu}}$ s are functions of $x^{\nu}, \stackrel{(1)}{p}^{\nu}$ and $\stackrel{(2)}{p}^{\nu}$, as we can see easily. Introducing another linear connections $\stackrel{(3)}{U}, \stackrel{(4)}{U}, \ldots \ldots, \stackrel{(\underset{\sim}{U}}{U}$ and repeating this method we can get in general a line element of the $(k+1)$-th kind
 $\stackrel{(1)}{p^{v}}, \stackrel{(2)}{p^{v}}, \ldots \ldots, \stackrel{(i)}{p}{ }^{\text {v }}$. These connections $\stackrel{(1)}{U}, \stackrel{(2)}{U}, \ldots \ldots, \stackrel{(r)}{U}$ are called the base connections. The manifold $X_{n}$, to whose every point are associated every system of the line elements $\stackrel{(1)}{p}^{\nu}, \stackrel{(1)}{p}^{\nu}, \ldots \ldots, \stackrel{(r)}{p^{\nu}}$, is defined as the generalized Finsler manifold.
2. The connection $U$. Consider a contravariant or covariant vector, which depends on $x^{\nu}, d x^{\nu}, \ldots \ldots, d^{r} x^{\nu}$, then components of the vector will be considered as functions of $x^{v}, \stackrel{(1)}{p^{\nu}}, \ldots \ldots, \stackrel{(r)}{p^{v}}$, for we can substitute $\stackrel{(1)}{p^{\nu}}, \ldots \ldots, \stackrel{(r)}{p^{\nu}}$ for $d x^{\nu}, \ldots \ldots, d^{r} x^{\nu}$ by (1.1), (1.2) and (1.3). Now we define a connection $U$ for such contravariant or covariant vectors $v^{\nu}$ or $w_{\lambda}$ :

$$
\begin{align*}
& \delta v^{\nu}=\nabla_{\mu} v^{\nu} d x^{\mu}=d v^{\nu}+\Gamma_{\lambda \mu}^{\nu} v^{\lambda} d x^{\mu},  \tag{2.1}\\
& \delta w_{\lambda}=\nabla_{\mu} w_{\lambda} d x^{\mu}=d w_{\lambda}-\Gamma_{\lambda \mu}^{\prime \nu} w_{\nu} d x^{\mu}
\end{align*}
$$

where $\Gamma_{\lambda, ~}^{\nu}$ 's as well as $\Gamma_{\lambda \mu}^{\prime \prime \prime}$ 's depend upon $x^{\nu}, \stackrel{(1)}{p}_{p}^{\nu},{ }^{(2)}, \ldots \ldots, \stackrel{(r)}{p}^{\nu}$. Accordingly the covariant differential of any affinor can be defined from (2.1) and we can calculate the covariant derivatives

$$
\begin{align*}
& \nabla_{\mu} v^{\nu}=\frac{\partial v^{\nu}}{\partial x^{\mu}}+\Gamma_{\lambda \mu}^{\nu} v^{\lambda}+\sum_{i}^{1, r} \frac{\partial v^{\nu}}{\partial p^{(i)}}\left({ }^{(i)}{ }_{\mu}^{(i)} p^{(0)}-\stackrel{\Gamma}{\Gamma}_{\lambda \mu}^{(i)}{ }^{(i)} p^{\lambda}\right), \\
& \nabla_{\mu} w_{\lambda}=\frac{\partial w_{\lambda}}{\partial x^{\mu}}-\Gamma_{\lambda \mu}^{\prime \nu} w_{\nu}+\sum_{i}^{1, r} \frac{\partial w_{\lambda}}{\partial p^{(i)}}\left({ }_{\nabla}^{(i)}{ }_{\mu}^{(i)} p^{(i)}-{ }_{\Gamma}^{(i)}{ }_{\nu \mu}{ }^{(i)} p^{\nu}\right), \tag{2.2}
\end{align*}
$$

where ${\stackrel{(i)}{\nabla_{\mu}}{ }^{(i)} p^{(i)}}$ means covariant derivative of ${ }^{(i)}{ }^{(i)}$ referred to $\stackrel{(i)}{U}$.
3. Construction of the parameters $\Gamma_{\lambda \mu}^{\nu}$ and $\Gamma_{\lambda \mu}^{\prime \nu}$. The parameters $\Gamma_{\lambda \mu}^{\nu}$ and $\Gamma_{\lambda \mu}^{\prime \nu}$ are transformed in the same manner as that of the affine connection by any change of system of coordinates. Therefore

$$
\begin{equation*}
C_{\lambda \mu} \ddot{\nu}^{\nu}=\Gamma_{\lambda \mu}^{\nu}-\Gamma_{\lambda \mu}^{\prime \nu} \tag{3.1}
\end{equation*}
$$

$$
\begin{equation*}
S_{\dot{\lambda \mu}}^{\ddot{\nu}}=\frac{1}{2}\left(\Gamma_{\lambda \mu}^{\nu}-\Gamma_{\mu \lambda}^{\nu}\right), \quad S^{\prime} \ddot{\lambda \mu}{ }^{\nu}=\frac{1}{2}\left(\Gamma_{\lambda \mu}^{\prime \nu}-\Gamma_{\mu \lambda}^{\prime \nu}\right) \tag{3.2}
\end{equation*}
$$

are affinors. Let us consider a contravariant tensor $g^{\lambda \mu}$ and a corresponding covariant tensor $g_{\lambda_{\mu}}$, which depend upon $x^{\nu}, \stackrel{(1)}{p}^{\nu}, \ldots \ldots, \stackrel{(r)}{p^{\nu}}$ also, and put their covariant derivatives

$$
\begin{equation*}
\nabla_{\nu} g^{\lambda \mu}=Q_{\stackrel{\lambda}{\lambda \mu}}, \quad \nabla_{\nu} g_{\lambda \mu}=Q_{\lambda \mu \nu}^{\prime} \tag{3.3}
\end{equation*}
$$

Then it follows from (2.2)
(3. 4)
(3. 5)

$$
\Gamma_{\lambda \nu}^{\omega} g_{\omega j \nu}+\Gamma_{\mu \nu}^{\omega} g_{\omega \lambda}=\frac{\partial g_{\lambda \mu}}{\partial x^{\nu}}+Q_{\lambda \mu \nu}+\sum_{i}^{1, r} \frac{\partial g_{\lambda \mu}}{\partial^{(i)}}\left({ }^{(i)} \nabla_{\nu}^{(i)} p^{(i)}-\stackrel{(i)}{\tau \nu}_{(i)}^{(i)} p^{\tau}\right),
$$

where $Q_{\lambda \mu \nu}=g_{\lambda \omega} g_{\mu \tau} Q^{\omega \tau}{ }_{\nu \nu}$. We get in consequence of (3.5)
(3. 6) $\quad \Gamma_{\lambda \mu}^{\nu}=\left\{\begin{array}{c}\lambda \mu \\ \nu\end{array}\right\}+T_{\lambda \mu}^{\nu}{ }^{\nu}+W_{\dot{\mu} \mu}^{\nu}, \quad \Gamma_{\lambda \mu}^{\prime \nu}=\left\{\begin{array}{c}\lambda \mu \\ \nu\end{array}\right\}+T_{\dot{\lambda} \mu}^{\prime}{ }^{\nu}+W_{\dot{\lambda} \mu}{ }^{\nu}$, where we put
(3. 7)

$$
\left\{\begin{array}{l}
\lambda \mu \\
\nu
\end{array}\right\}=\frac{1}{2} g^{v \omega}\left(\frac{\partial g_{\lambda \omega}}{\partial x^{\mu}}+\frac{\partial g_{\omega \mu}}{\partial x^{\lambda}}-\frac{\partial g_{\lambda \mu}}{\partial x^{\omega}}\right)
$$

(3. 8)

$$
\begin{aligned}
& T_{i \mu}{ }^{\nu}=\frac{1}{2}\left(g_{\lambda \omega} Q_{i}^{;}{ }^{\omega \nu}+g_{\mu \omega} Q_{\lambda}{ }^{\omega \nu}-g^{\nu \sigma} g_{\mu \omega} g_{\lambda \tau} Q_{\sigma}{ }^{\cdot \omega \tau}\right),
\end{aligned}
$$

(3. 9)

$$
\begin{equation*}
T^{\prime} \ddot{\lambda} \dot{\mu}_{\nu}=T_{\dot{\lambda} \ddot{\mu}^{\nu}}-C_{\mu \ddot{\lambda}}^{\nu} \tag{3.10}
\end{equation*}
$$

4. The curvature tensor. We shall now find the curvature tensor. From (2. 2) it follows

$$
\begin{aligned}
& =\frac{\partial^{2} v^{\lambda}}{\partial x^{\nu} \partial x^{\nu}}+\frac{\partial \Gamma_{\rho v}^{\lambda}}{\partial x^{\underline{2}}} v^{\rho}+\Gamma_{\rho \nu}^{\lambda} \frac{\partial v^{\rho}}{\partial x^{u}}+\sum_{i} \frac{\partial^{2} v^{\lambda}}{\partial x^{\mu} \partial^{(i)} p^{(i)}}\left(\nabla_{\nu}^{(i)}{ }_{\nu}^{(i)} p^{\omega 0}-\Gamma_{\rho \nu}^{(i)} p^{(i)}\right) \\
& +\sum_{i} \frac{\partial v^{\lambda}}{\partial p^{(i)}}\left(\frac{\partial \nabla_{\nu}^{(i)}{ }_{\nu}^{(i)} p^{\omega}}{\partial x^{\mu}}-\frac{\partial \Gamma_{\Gamma_{p v}^{(i)}}^{(i)} p^{(i)}}{\partial x^{\mu}}\right)+\Gamma_{\omega ;}^{\lambda}\left(\frac{\partial v^{\omega}}{\partial x^{\nu}}+\Gamma_{p \nu}^{\omega} v^{\rho}\right. \\
& \left.+\sum_{i} \frac{\partial v^{(\omega)}}{\partial p^{(i)}}\left({ }_{\nabla}^{(i)}{ }_{\nu}^{(i)} p^{\sigma}-\stackrel{\Gamma}{\Gamma}_{\rho}^{(i)}{ }_{\nu}^{(i)} p^{\rho}\right)\right\}-\Gamma_{\lambda \mu}^{\prime / \omega} \nabla_{\omega} v^{\lambda}
\end{aligned}
$$

hence
(4. 1)

$$
\begin{aligned}
& \left.2 \nabla_{[\mu} \nabla_{\nu}\right\urcorner v^{\lambda}=R_{\lambda \cdots \rho}^{\cdots{ }_{\mu}} v^{p}-\sum_{i} R_{\lambda \mu \rho}^{(i)} \cdots \frac{\partial v^{\lambda(i)}}{\partial p^{(i)}} p^{\omega} \\
& +\sum_{i} \frac{\partial v^{\lambda}}{\partial_{p^{(i)}}^{(i)}} V_{\lambda \mu}^{\dot{\mu}}+2 S_{\mu \nu}^{\prime \cdots} \nabla_{\omega} v^{\lambda},
\end{aligned}
$$

where

$$
\begin{align*}
& R_{\lambda \mu \rho}^{\sim{ }_{\mu}^{\lambda}}=\frac{\partial \Gamma_{\rho \nu}^{\lambda}}{\partial x^{\nu}}-\frac{\partial \Gamma_{\rho \mu}^{\lambda}}{\partial x^{\nu}}+\Gamma_{\omega \mu}^{\lambda} \Gamma_{\rho \nu}^{\omega}-\Gamma_{\omega \nu}^{\lambda} \Gamma_{\rho \mu}^{\omega \omega}  \tag{4.2}\\
& +\sum_{j} \frac{\partial \Gamma_{\rho v}^{\lambda}}{\partial p^{(j)}}\left({ }_{\nabla}^{(j)}{ }_{\mu}^{(j)} p^{\omega}-{ }^{(j)}{ }_{\tau \mu}^{(j)} p^{(j)}\right)-\sum_{j} \frac{\partial \Gamma_{\rho \mu}^{\lambda}}{\partial p^{(j)}}\left({ }_{\nu}^{(j)}{ }_{\nu}^{(j)} p^{\omega \omega}-\Gamma_{\tau \nu}^{(j)} p^{(j)}\right), \tag{4.3}
\end{align*}
$$

We shall call $R_{\mu \nu \rho}^{\because \ddot{p}^{\lambda}}$ the curvature tensor of our connection $U$ and
 $\stackrel{(1)}{U}, \stackrel{(2)}{U}, \ldots \ldots, \stackrel{(r)}{U}$.
5. Special cases. When we put $r=1$, we get a connection in the general Finsler manifold. Let ${ }^{(1)}$ of $x^{\nu}$ be connected with that of $x^{\nu}+d x^{\nu}$,
 T. Hosokawa ${ }^{2}$ ) has studied. The connection of Berwald in the Finsler manifold $^{3)}$ is a more special one of this case, i.e. $\Gamma_{\lambda \mu}^{\nu}={\stackrel{(1)}{\Gamma_{\lambda \mu}}}^{\nu}$.

As the correspondence between line elements belonging to two consecutive points can be defined arbitrarily, we may put ${ }^{(6)}{ }_{\mu} p^{(i)}=0$,
 Berwald. Our connection contains that of $\mathrm{Craig}^{4}$ ) as a special case too, in which case $r$ is 2.

[^0]
[^0]:    1) This expression consists only formally, because ${ }^{(i)}{ }^{0}$ is not a vector field but a line element and $\stackrel{i}{(i)}_{\mu}^{(i)} p^{(i)}$ define a correspondence between two line elements belonging to a point $x^{\nu}$ and its consecutive point $x^{\nu}+d x^{\nu}$ respectively. Accordingly this expression does not give us the curvature tensor.
    2) T. Hosokawa: Science Reports, Tohoku Imp. University, series I, 19 (1930), p. 37-51.
    3) L. Berwald: Math. Zeitschrift, 25 (1926), S. 40-73.
    4) H. V. Craig: Trans. of Amer. Math. Soc., 33 (1931), p. 125-142.
