PAPERS COMMUNICATED

65. Theory of Connections in the Generalized Finsler Manifold.

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Recently the theory of connections in the Finsler manifold F_n has been developed somewhat by several authors,¹⁾ in which the parameters of the connections are in general functions not only of positions x^{\vee} but also of directions dx^{\vee} . We shall consider some connections in the generalized Finsler manifold K_n similarly as in the Finsler manifold, the parameters of the connections depending on x^{\vee} , dx^{\vee} , d^2x^{\vee} ,, d^rx^{\vee} .

1. The base-connections. Let X_n be a manifold of n dimensions, in which a system of coordinates x^{ν} is taken, and we express for brevity n differentials dx^{ν} of the coordinates by $p^{\nu}: dx^{\nu} = p^{\nu}$. Introduce an arbitrary linear connection $\stackrel{(1)}{U}$, whose parameters $\stackrel{(1)}{\Gamma}_{\lambda\mu}^{\nu}$ depend upon positions and directions, i.e. x^{ν} and dx^{ν} , then we get covariant differentials of a line element $\stackrel{(1)}{p^{\nu}}$, which can be considered as a contravariant vector and expressed by the parameters $\stackrel{(1)}{\Gamma}_{\lambda\mu}^{\nu}$ of the connection $\stackrel{(1)}{U}$ as follows:

(1. 1) $\begin{array}{c} \delta p^{\nu} = dp^{\nu} + \Gamma^{\nu}_{\lambda\mu} p^{\lambda} dx^{\nu} . \end{array}$

We shall write p^{ν} for δp^{ν} , which is function of x^{ν} as well as p^{ν} . Let us call this contravariant vector p^{ν} the line element of the second kind.

We introduce moreover a new linear connection $U^{(2)}$ for X_n , whose parameters $\Gamma^{(2)}_{\lambda\mu}$ are functions of x^{ν} as well as p^{ν} and p^{ν} . This connection $U^{(2)}$ is applied to the affinors, which depend on x^{ν} and the foregoing line elements $p^{(1)}_{\nu}$ and p^{ν} , and by which we get the line element of the third kind

(1. 2)
$$p^{(3)} = \delta^{(2)} p^{(2)} = dp^{(2)} + \Gamma^{(3)}_{\lambda\mu} p^{\lambda} dx^{\mu} ,$$

¹⁾ E. Noether: Göttinger Nachrichten, math.-phys. Kl. (1918), S. 37-44. L. Berwald: Jahresberichte d. Deutschen Math.-Ver. **34** (1925), S. 213-220. J. L. Synge: Trans. of Amer. Math. Soc. **27** (1925), p. 61-67. J. H. Taylor: Trans. of Amer. Math. Soc. **27** (1925), p. 246-264.

²⁾ $\overset{(2)}{U}$ is in general different from $\overset{(1)}{U}$, but we may also take same one with $\overset{(1)}{U}$ as $\overset{(2)}{U}$. It is completely analogous for following connections $\overset{(k)}{U}$.

⁽³⁾ p^{ν} 's are functions of x^{ν} , p^{ν} and p^{ν} , as we can see easily. Introducing another linear connections U, U, \dots, U and repeating this method we can get in general *a line element of the* (k+1)-th kind

(1. 3)
$$p^{\nu} = \delta p^{\nu} = dp^{\nu} + \Gamma^{(k)}_{\lambda\mu} p^{\lambda} dx^{\nu}$$
,

where $\Gamma_{\lambda \mu}^{(k)}$'s are parameters of the connection $\stackrel{(k)}{U}$ and functions of x^{ν} , p^{ν} , p^{ν} , p^{ν} ,, p^{ν} . These connections $\stackrel{(1)}{U}$, $\stackrel{(2)}{U}$,, $\stackrel{(r)}{U}$ are called the base connections. The manifold X_n , to whose every point are associated every system of the line elements p^{ν} , p^{ν} ,, p^{ν} , is defined as the generalized Finsler manifold.

2. The connection U. Consider a contravariant or covariant vector, which depends on x^{ν} , dx^{ν} ,, $d^{r}x^{\nu}$, then components of the vector will be considered as functions of x^{ν} , p^{ν} ,, p^{ν} , for we can substitute $p^{(1)}_{\nu}$,, $p^{(r)}_{\nu}$ for dx^{ν} ,, $d^{r}x^{\nu}$ by (1. 1), (1. 2) and (1. 3). Now we define a connection U for such contravariant or covariant vectors v^{ν} or w_{λ} :

(2. 1)
$$\begin{aligned} \delta v^{\nu} = \nabla_{\mu} v^{\nu} dx^{\mu} = dv^{\nu} + \Gamma^{\nu}_{\lambda \mu} v^{\lambda} dx^{\mu} , \\ \delta w_{\lambda} = \nabla_{\mu} w_{\lambda} dx^{\mu} = dw_{\lambda} - \Gamma^{\prime}_{\lambda \mu} w_{\nu} dx^{\mu} , \end{aligned}$$

where $\Gamma_{\lambda\mu}^{\nu}$'s as well as $\Gamma_{\lambda\mu}^{\prime\nu}$'s depend upon x^{ν} , p^{ν} , p^{ν} , p^{ν} ,, p^{ν} . Accordingly the covariant differential of any affinor can be defined from (2.1) and we can calculate the covariant derivatives

(2. 2)

$$\nabla_{\mu}v^{\nu} = \frac{\partial v^{\nu}}{\partial x^{\mu}} + \Gamma^{\nu}_{\lambda\mu}v^{\lambda} + \sum_{i}^{1,r} \frac{\partial v^{\nu}}{\partial p^{\omega}} (\stackrel{(i)}{\nabla}_{\mu}p^{\nu} - \Gamma^{(i)}_{\lambda\mu}p^{\lambda}),$$

$$\nabla_{\mu}w_{\lambda} = \frac{\partial w_{\lambda}}{\partial x^{\mu}} - \Gamma^{\prime\nu}_{\lambda\mu}w_{\nu} + \sum_{i}^{1,r} \frac{\partial w_{\lambda}}{\partial p^{\omega}} (\stackrel{(i)}{\nabla}_{\mu}p^{\nu} - \Gamma^{(i)}_{\nu\mu}p^{\nu}),$$

where $\bigtriangledown_{\mu}^{(*)} p^{\nu}$ means covariant derivative of p^{ν} referred to U.

3. Construction of the parameters $\Gamma^{\nu}_{\lambda\mu}$ and $\Gamma^{\nu}_{\lambda\mu}$. The parameters $\Gamma^{\nu}_{\lambda\mu}$ and $\Gamma^{\nu}_{\lambda\mu}$ are transformed in the same manner as that of the affine connection by any change of system of coordinates. Therefore

$$(3. 1) C_{\lambda\mu}^{\prime\nu} = \Gamma_{\lambda\mu}^{\nu} - \Gamma_{\lambda\mu}^{\prime\nu}$$

(3. 2)
$$S_{\lambda\mu}^{\prime\nu} = \frac{1}{2} (\Gamma_{\lambda\mu}^{\nu} - \Gamma_{\mu\lambda}^{\nu}), \qquad S_{\lambda\mu}^{\prime\nu} = \frac{1}{2} (\Gamma_{\lambda\mu}^{\prime\nu} - \Gamma_{\mu\lambda}^{\prime\nu})$$

are affinors. Let us consider a contravariant tensor $g^{\lambda\mu}$ and a corresponding covariant tensor $g_{\lambda\mu}$, which depend upon x^{ν} , p^{ν} ,, p^{ν} also, and put their covariant derivatives

(3. 3) $\nabla_{\nu}g^{\lambda\mu} = Q^{\lambda\mu}_{\ldots\nu}, \quad \nabla_{\nu}g_{\lambda\mu} = Q'_{\lambda\mu\nu}.$

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Then it follows from (2.2)

(3. 4)
$$Q_{\mu\nu}^{\lambda\mu} = \frac{\partial g^{\lambda\mu}}{\partial x^{\nu}} + \Gamma_{\mu\nu}^{\lambda} g^{\mu\nu} + \Gamma_{\mu\nu}^{\mu} g^{\lambda\nu} + \sum_{i}^{1,r} \frac{\partial g^{\lambda\mu}}{\partial p^{\nu}} (\stackrel{(i)}{\bigtriangledown} \stackrel{(i)}{\bigtriangledown} \stackrel{(i)}{p^{\nu}} - \stackrel{(i)}{\Gamma} \stackrel{(i)}{\bigtriangledown} \stackrel{(i)}{p^{\tau}}),$$

(3. 5)
$$\Gamma_{\lambda\nu}^{\omega} g_{\mu\mu} + \Gamma_{\mu\nu}^{\omega} g_{\mu\lambda} = \frac{\partial g_{\lambda\mu}}{\partial x^{\nu}} + Q_{\lambda\mu\nu} + \sum_{i}^{1,r} \frac{\partial g_{\lambda\mu}}{\partial p^{\nu}} (\stackrel{(i)}{\bigtriangledown} \stackrel{(i)}{\bigtriangledown} \stackrel{(i)}{p^{\nu}} - \stackrel{(i)}{\Gamma} \stackrel{(i)}{\underset{\nu}{}} \stackrel{(i)}{p^{\tau}}),$$

where $Q_{\lambda\mu\nu} = g_{\lambda\omega}g_{\mu\tau}Q^{\omega\tau}_{...\nu}$. We get in consequence of (3.5)

(3. 6) $\Gamma^{\nu}_{\lambda\mu} = \{^{\lambda\mu}_{\nu}\} + T^{\cdot,\nu}_{\lambda\mu} + W^{\cdot,\nu}_{\lambda\mu}, \qquad \Gamma^{\prime\nu}_{\lambda\mu} = \{^{\lambda\mu}_{\nu}\} + T^{\prime,\nu}_{\lambda\mu} + W^{\cdot,\nu}_{\lambda\mu},$ where we put

(3. 7)
$$\{ {}^{\lambda\mu}_{\nu} \} = \frac{1}{2} g^{\nu\omega} \left(\frac{\partial g_{\lambda\omega}}{\partial x^{\mu}} + \frac{\partial g_{\omega\mu}}{\partial x^{\lambda}} - \frac{\partial g_{\lambda\mu}}{\partial x^{\omega}} \right),$$

(3. 8)
$$T_{\lambda\mu}^{\cdot,\nu} = \frac{1}{2} (g_{\lambda\omega} Q_{\mu}^{\cdot,\nu\nu} + g_{\mu\omega} Q_{\lambda}^{\cdot,\nu\nu} - g^{\nu\sigma} g_{\mu\omega} g_{\lambda\tau} Q_{\sigma}^{\cdot,\nu\tau}), + S_{\lambda\mu}^{\cdot,\nu} - g^{\nu\omega} (g_{\lambda\tau} S_{\omega\mu}^{\cdot,\nu\tau} + g_{\mu\tau} S_{\omega\lambda}^{\cdot,\tau}),$$

$$(3. 9) T'_{\lambda\mu} = T_{\lambda\mu} - C_{\mu\lambda}^{\mu\nu},$$

(3. 10)
$$W_{\lambda\mu}^{*,\nu} = \frac{1}{2} g^{\nu\omega} \sum_{i}^{i,r} \left(\frac{\partial g_{\lambda\omega}}{\partial p^{\tau}} \bigtriangledown_{\mu} p^{\tau} + \frac{\partial g_{\omega\mu}}{\partial p^{\tau}} \bigtriangledown_{\lambda} p^{\tau} - \frac{\partial g_{\lambda\mu}}{\partial p^{\tau}} \bigtriangledown_{\omega} p^{\tau} \right) \\ - \frac{1}{2} g^{\nu\omega} \sum_{i}^{i,r} \left(\frac{\partial g_{\lambda\omega}}{\partial p^{\tau}} \varGamma_{\sigma\mu}^{(i)} + \frac{\partial g_{\omega\mu}}{\partial p^{\tau}} \varGamma_{\sigma\lambda}^{(i)} - \frac{\partial g_{\lambda\mu}}{\partial p^{\tau}} \varGamma_{\sigma\omega}^{(i)} \right) p^{\sigma}.$$

4. The curvature tensor. We shall now find the curvature tensor. From (2.2) it follows

$$\begin{split} \nabla_{\mu}\nabla_{\nu}v^{\lambda} &= \frac{\partial}{\partial x^{\mu}} \nabla_{\nu}v^{\lambda} + \Gamma^{\lambda}_{\omega\mu} \nabla_{\nu}v^{\omega} - \Gamma'^{\omega}_{\nu\mu} \nabla_{\omega}v^{\lambda} + \sum_{i} \frac{\partial \nabla_{\nu}v^{\lambda}}{\partial p^{\omega}} (\stackrel{(i)}{\nabla} \stackrel{(i)}{v} \stackrel{(i)}{p^{\omega}} - \stackrel{(i)}{\Gamma} \stackrel{(i)}{v_{\mu}} \stackrel{(i)}{p^{\tau}}) \\ &= \frac{\partial^{2}v^{\lambda}}{\partial x^{\mu} \partial x^{\nu}} + \frac{\partial \Gamma^{\lambda}_{\mu\nu}}{\partial x^{\mu}} v^{\rho} + \Gamma^{\lambda}_{\rho\nu} \frac{\partial v^{\rho}}{\partial x^{\mu}} + \sum_{i} \frac{\partial^{2}v^{\lambda}}{\partial x^{\mu} \partial p^{\omega}} (\stackrel{(i)}{\nabla} \stackrel{(i)}{v} \stackrel{(i)}{p^{\omega}} - \stackrel{(i)}{\Gamma} \stackrel{(i)}{\rho_{\nu}} \stackrel{(i)}{p^{\tau}}) \\ &+ \sum_{i} \frac{\partial v^{\lambda}}{\partial p^{\omega}} \left(\frac{\partial \stackrel{(i)}{\nabla} \stackrel{(i)}{v} - \frac{\partial \stackrel{(i)}{\rho} \stackrel{(i)}{v}}{\partial x^{\mu}} \stackrel{(i)}{p^{\rho}} \right) + \Gamma^{\lambda}_{\omega\mu} \left\{ \frac{\partial v^{\omega}}{\partial x^{\nu}} + \Gamma^{\omega}_{\rho\nu} v^{\rho} \\ &+ \sum_{i} \frac{\partial v^{\omega}}{\partial p^{\sigma}} (\stackrel{(i)}{\nabla} \stackrel{(i)}{v} \stackrel{(i)}{p^{\sigma}} - \stackrel{(i)}{\Gamma} \stackrel{(i)}{\rho} \stackrel{(i)}{v} \stackrel{(i)}{p^{\rho}} \right) - \Gamma'^{\lambda}_{\lambda\mu} \nabla_{\omega} v^{\lambda} \\ &+ \sum_{j} \left\{ \frac{\partial^{2}v^{\lambda}}{\partial p^{\omega} \partial x^{\lambda}} + \frac{\partial \Gamma^{\lambda}_{\rho\nu}}{\partial p^{\omega}} v^{\rho} + \Gamma^{\lambda}_{\rho\nu} \frac{\partial v^{\rho}}{\partial p^{\omega}} + \sum_{i} \frac{\partial^{2}v^{\lambda}}{\partial p^{\omega} \partial p^{\tau}} (\stackrel{(i)}{\nabla} \stackrel{(i)}{v} \stackrel{(i)}{p^{\tau}} - \stackrel{(i)}{\Gamma} \stackrel{(i)}{\rho} \stackrel{(i)}{p^{\rho}} \right) \\ &+ \sum_{i} \frac{\partial v^{\lambda}}{\partial p^{\tau}} (\frac{\partial \stackrel{(i)}{\nabla} \stackrel{(i)}{v} - \stackrel{(i)}{\partial p^{\mu}} \stackrel{(i)}{v} \stackrel{(i)}{p^{\mu}} + \frac{\partial \Gamma^{\lambda}_{\rho\nu}}{\partial p^{\mu}} \stackrel{(i)}{p^{\mu}} \right) - \frac{\partial v^{\lambda}}{\partial p^{\tau}} \stackrel{(i)}{\Gamma} \stackrel{(i)}{v} \left\{ \stackrel{(i)}{\nabla} \stackrel{(j)}{v} \stackrel{(j)}{p^{\tau}} - \stackrel{(i)}{\Gamma} \stackrel{(j)}{\rho} \stackrel{(j)}{p^{\rho}} \right) , \end{split}$$

hence

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(4. 1) $2 \nabla_{(\mu} \nabla_{\nu)} v^{\lambda} = R_{\lambda \mu \rho}^{\dots \lambda} v^{\rho} - \sum_{i} \stackrel{(i)}{R_{\lambda \mu \rho}^{\dots \omega}} \frac{\partial v^{\lambda}}{\partial p^{\omega}} p^{i}$ $+ \sum_{i} \frac{\partial v^{\lambda}}{\partial n^{\omega}} V^{(i)}_{\lambda \mu} + 2S'_{\mu \nu}^{\dots \omega} \nabla_{\omega} v^{\lambda},$

where

We shall call $R_{\mu\nu\rho}^{\dots\lambda}$ the curvature tensor of our connection U and $\stackrel{(i)}{R_{\mu\nu\rho}}$ can be considered as the curvature tensors belonging to $\stackrel{(i)}{U}$, $\stackrel{(i)}{U}$, $\stackrel{(i)}{U}$, $\stackrel{(i)}{U}$.

5. Special cases. When we put r=1, we get a connection in the general Finsler manifold. Let p^{ν} of x^{ν} be connected with that of $x^{\nu} + dx^{\nu}$, so that $\delta p^{\nu} = 0$, i.e. $dp^{\nu} = -\Gamma_{\lambda\mu}^{(1)} p^{\lambda} dx^{\mu}$, then we get the case, which T. Hosokawa² has studied. The connection of Berwald in the Finsler manifold³ is a more special one of this case, i.e. $\Gamma_{\lambda\mu}^{\nu} = \Gamma_{\lambda\mu}^{(1)}$. As the correspondence between line elements belonging to two

As the correspondence between line elements belonging to two consecutive points can be defined arbitrarily, we may put $\bigtriangledown_{\mu}^{(i)} p^{\nu} = 0$, then $V_{\lambda\mu}^{(i)} = 0$ and we get a direct generalisation of the connection of Berwald. Our connection contains that of Craig⁴⁾ as a special case too, in which case r is 2.

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¹⁾ This expression consists only formally, because $p^{(i)}$ is not a vector field but a line element and $\nabla_{\mu}p^{\nu}$ define a correspondence between two line elements belonging to a point x^{ν} and its consecutive point $x^{\nu} + dx^{\nu}$ respectively. Accordingly this expression does not give us the curvature tensor.

²⁾ T. Hosokawa: Science Reports, Tohoku Imp. University, series I, **19** (1930), p. 37-51.

³⁾ L. Berwald: Math. Zeitschrift, 25 (1926), S. 40-73.

⁴⁾ H. V. Craig: Trans. of Amer. Math. Soc., 33 (1931), p. 125-142.