# Multiplicity of solutions for anisotropic quasilinear elliptic equations with variable exponents 

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#### Abstract

We study an anisotropic partial differential equation on a bounded domain $\Omega \subset \mathbb{R}^{N}$. We prove the existence of at least two nontrivial weak solutions using as main tools the mountain pass lemma and Ekeland's variational principle.


## 1 Introduction

Equations involving variable exponent growth conditions have been extensively studied in the last decade. We just remember the recent advances in [10, 12, 13, 1, $2,25,16,26,27,20,18,30,28,29]$. The large number of papers studying problems involving variable exponent growth conditions is motivated by the fact that this type of equations can serve as models in the theory of electrorhological fluids [17, 35, 36, 5, 1], image processing [4] or the theory of elasticity [40].

Typical models of elliptic equations with variable exponent growth conditions appeal to the so called $p(x)$-Laplace operator, i.e.

$$
\Delta_{p(x)} u:=\operatorname{div}\left(|\nabla u|^{p(x)-2} \nabla u\right),
$$

where $p(x)$ is a function satisfying $p(x)>1$ for each $x$. Recently, Mihăilescu-Pucci-Rădulescu extended in [24] the study involving the $p(x)$-Laplace operator

[^0]to the case of anisotropic equations with variable exponent growth conditions, where the differential operator considered has the form
\[

$$
\begin{equation*}
\sum_{i=1}^{N} \partial_{x_{i}}\left(\left|\partial_{x_{i}} u\right|^{p_{i}(x)-2} \partial_{x_{i}} u\right), \tag{1}
\end{equation*}
$$

\]

with $p_{i}(x)$ functions satisfying $\inf _{x} p_{i}(x)>1$ for each $i \in\{1, \ldots, N\}$. Undoubtedly, in the particular case when $p_{i}(x)=p(x)$ for each $i \in\{1, \ldots, N\}$ the above differential operator becomes $\sum_{i=1}^{N} \partial_{x_{i}}\left(\left|\partial_{x_{i}} u\right|^{p(x)-2} \partial_{x_{i}} u\right)$ and has similar properties with the $p(x)$-Laplace operator. On the other hand, the anisotropic equations with variable exponent growth conditions enable the study of equations with more complicated nonlinearities since the differential operator (1) allows a distinct behavior for partial derivatives in various directions.

Motivated by the above discussion, the goal of this paper is to investigate a problem of the type

$$
\begin{cases}-\sum_{i=1}^{N} \partial_{x_{i}}\left(\left|\partial_{x_{i}} u\right|^{p_{i}(x)-2} \partial_{x_{i}} u\right)=f(x, u), & \text { for } x \in \Omega  \tag{2}\\ u=0, & \text { for } x \in \partial \Omega\end{cases}
$$

where $\Omega \subset \mathbb{R}^{N}(N \geq 3)$ is a bounded domain with smooth boundary, $p_{i}$ are continuous functions on $\bar{\Omega}$ such that $2 \leq p_{i}(x)$ for any $x \in \bar{\Omega}$ and $i \in\{1, \ldots, N\}$. Our main result on problem (2) will supplement the results in [23, 24, 22, 21] obtained for similar anisotropic equations.

## 2 Preliminary results on variable exponent spaces

Assume $\Omega \subset \mathbb{R}^{N}$ is an open domain.
Set

$$
C_{+}(\bar{\Omega})=\{h ; h \in C(\bar{\Omega}), h(x)>1 \text { for all } x \in \bar{\Omega}\} .
$$

For any $p \in C_{+}(\bar{\Omega})$ we define

$$
p^{+}=\sup _{x \in \Omega} p(x) \quad \text { and } \quad p^{-}=\inf _{x \in \Omega} p(x) .
$$

For each $p \in C_{+}(\bar{\Omega})$, we recall the definition of the variable exponent Lebesgue space
$L^{p(\cdot)}(\Omega)=\{u ; u$ is a measurable real-valued function such that

$$
\left.\int_{\Omega}|u(x)|^{p(x)} d x<\infty\right\} .
$$

This space becomes a Banach space [19, Theorem 2.5] with respect to the Luxemburg norm, that is

$$
|u|_{p(\cdot)}=\inf \left\{\mu>0 ; \int_{\Omega}\left|\frac{u(x)}{\mu}\right|^{p(x)} d x \leq 1\right\}
$$

Moreover, $L^{p(\cdot)}(\Omega)$ is a reflexive space [19, Corollary 2.7] provided that $1<p^{-} \leq$ $p^{+}<\infty$. Furthermore, on such kind of spaces a Hölder type inequality is valid
[19, Theorem 2.1]. More exactly, denoting by $L^{q(\cdot)}(\Omega)$ the conjugate space of $L^{p(\cdot)}(\Omega)$, where $\frac{1}{p(x)}+\frac{1}{q(x)}=1$ for any $x \in \bar{\Omega}$, for each $u \in L^{p(\cdot)}(\Omega)$ and each $v \in L^{q(\cdot)}(\Omega)$ the Hölder type inequality reads as follows

$$
\begin{equation*}
\left|\int_{\Omega} u v d x\right| \leq\left(\frac{1}{p^{-}}+\frac{1}{q^{-}}\right)|u|_{p(\cdot)}|v|_{q(\cdot)} . \tag{3}
\end{equation*}
$$

An immediate consequence of Hölder's inequality is connected with some inclusions between various Lebesgue spaces involving variable exponent growth [19, Theorem 2.8]: if $0<|\Omega|<\infty$ and $p_{1}, p_{2}$ are variable exponents, so that $p_{1}(x) \leq p_{2}(x)$ almost everywhere in $\Omega$, then there exists the continuous embed$\operatorname{ding} L^{p_{2}(\cdot)}(\Omega) \hookrightarrow L^{p_{1}(\cdot)}(\Omega)$, whose norm does not exceed $|\Omega|+1$.

An important role in manipulating the generalized Lebesgue-Sobolev spaces is played by the modular of the $L^{p(\cdot)}(\Omega)$ space, which is the mapping $\rho_{p(\cdot)}$ : $L^{p(\cdot)}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\rho_{p(\cdot)}(u)=\int_{\Omega}|u|^{p(x)} d x
$$

provided that $p^{+}<\infty$. Spaces with $p^{+}=\infty$ have been studied by Edmunds, Lang and Nekvinda [6].

We point out some relations which can be established between the Luxemburg norm and the modular. If $\left(u_{n}\right), u \in L^{p(\cdot)}(\Omega)$ and $p^{+}<\infty$ then the following relations hold true

$$
\begin{align*}
|u|_{p(\cdot)}>1 & \Rightarrow|u|_{p(\cdot)}^{p^{-}} \leq \rho_{p(\cdot)}(u) \leq|u|_{p(\cdot)}^{p^{+}}  \tag{4}\\
|u|_{p(\cdot)}<1 & \Rightarrow \quad|u|_{p(\cdot)}^{p^{+}} \leq \rho_{p(\cdot)}(u) \leq|u|_{p(\cdot)}^{p^{-}}  \tag{5}\\
\left|u_{n}-u\right|_{p(\cdot)} & \rightarrow 0 \quad \Leftrightarrow \quad \rho_{p(\cdot)}\left(u_{n}-u\right) \rightarrow 0 . \tag{6}
\end{align*}
$$

Next, we define the variable exponent Sobolev space $W_{0}^{1, p(\cdot)}(\Omega)$ as the closure of $C_{0}^{\infty}(\Omega)$ under the norm

$$
\|u\|=|\nabla u|_{p(\cdot)} .
$$

The space $\left(W_{0}^{1, p(\cdot)}(\Omega),\|\cdot\|\right)$ is a separable and reflexive Banach space, provided that $1<p^{-} \leq p^{+}<\infty$. We recall that if $\Omega$ is a bounded, open domain in $\mathbb{R}^{N}$, $q \in C_{+}(\bar{\Omega})$ and $q(x)<p^{\star}(x)$ for all $x \in \bar{\Omega}$ then the embedding

$$
W_{0}^{1, p(\cdot)}(\Omega) \hookrightarrow L^{q(\cdot)}(\Omega)
$$

is compact and continuous, where $p^{\star}(x)=\frac{N p(x)}{N-p(x)}$ if $p(x)<N$ or $p^{\star}(x)=+\infty$ if $p(x) \geq N$. We refer to $[31,6,7,8,11,14,19]$ for further properties of variable exponent Lebesgue-Sobolev spaces.

Finally, we recall the definition and properties of the anisotropic variable exponent Sobolev spaces as they were introduced in [24]. With that end in view, we assume in the sequel that $\Omega$ is a bounded open domain in $\mathbb{R}^{N}$ and we denote
by $\vec{p}(\cdot): \bar{\Omega} \rightarrow \mathbb{R}^{N}$ the vectorial function $\vec{p}(\cdot)=\left(p_{1}(\cdot), \ldots, p_{N}(\cdot)\right)$. We define $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$, the anisotropic variable exponent Sobolev space, as the closure of $C_{0}^{\infty}(\Omega)$ with respect to the norm

$$
\|u\|_{\vec{p}(\cdot)}=\sum_{i=1}^{N}\left|\partial_{x_{i}} u\right|_{p_{i}(\cdot)}
$$

In the case when $p_{i}(\cdot) \in C_{+}(\bar{\Omega})$ are constant functions for any $i \in\{1, \ldots, N\}$ the resulting anisotropic Sobolev space is denoted by $W_{0}^{1, \vec{p}}(\Omega)$, where $\vec{p}$ is the constant vector $\left(p_{1}, \ldots, p_{N}\right)$. The theory of this type of spaces was developed in $[15,32,33,34,37,38]$. It was argued in [24] that $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is a reflexive Banach space.

On the other hand, in order to facilitate the manipulation of the space $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$, we introduced $\vec{P}_{+}, \vec{P}_{-}$in $\mathbb{R}^{N}$ as

$$
\vec{P}_{+}=\left(p_{1}^{+}, \ldots, p_{N}^{+}\right), \quad \vec{P}_{-}=\left(p_{1}^{-}, \ldots, p_{N}^{-}\right)
$$

and $P_{+}^{+}, P_{-}^{+}, P_{-}^{-} \in \mathbb{R}^{+}$as

$$
P_{+}^{+}=\max \left\{p_{1}^{+}, \ldots, p_{N}^{+}\right\}, \quad P_{-}^{+}=\max \left\{p_{1}^{-}, \ldots, p_{N}^{-}\right\}, \quad P_{-}^{-}=\min \left\{p_{1}^{-}, \ldots, p_{N}^{-}\right\} .
$$

Throughout this paper we assume that

$$
\begin{equation*}
\sum_{i=1}^{N} \frac{1}{p_{i}^{-}}>1 \tag{7}
\end{equation*}
$$

and define $P_{-}^{\star} \in \mathbb{R}^{+}$and $P_{-, \infty} \in \mathbb{R}^{+}$by

$$
P_{-}^{\star}=\frac{N}{\sum_{i=1}^{N} \frac{1}{p_{i}^{-}}-1}, P_{-, \infty}=\max \left\{P_{-}^{+}, P_{-}^{\star}\right\} .
$$

Finally, we recall a result regarding the compact embedding between $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ and variable exponent Lebesgue spaces (see, [24, Theorem 1]):

Theorem 1. Assume that $\Omega \subset \mathbb{R}^{N}(N \geq 3)$ is a bounded domain with smooth boundary. Assume relation (7) is fulfilled. For any $q \in C(\bar{\Omega})$ verifying

$$
\begin{equation*}
1<q(x)<P_{-, \infty} \text { for all } x \in \bar{\Omega} \tag{8}
\end{equation*}
$$

the embedding

$$
W_{0}^{1, \vec{p}(\cdot)}(\Omega) \hookrightarrow L^{q(\cdot)}(\Omega)
$$

is continuous and compact.

## 3 The main result

In this paper we study problem (2) in the particular case

$$
f(x, t)=A|t|^{a(x)-2} t+B|t|^{b(x)-2} t
$$

where $a: \bar{\Omega} \rightarrow \mathbb{R}, b: \bar{\Omega} \rightarrow \mathbb{R}$ are continuous functions such that

$$
\begin{equation*}
1<a^{-}<a^{+}<P_{-}^{-} \leq P_{+}^{+}<b^{-}<b^{+}<\min \left\{N, P_{-, \infty}\right\} \tag{9}
\end{equation*}
$$

and $A, B>0$. More precisely, we consider the following problem

$$
\begin{cases}-\sum_{i=1}^{N} \partial_{x_{i}}\left(\left|\partial_{x_{i}} u\right|^{p_{i}(x)-2} \partial_{x_{i}} u\right)=A|u|^{a(x)-2} u+B|u|^{b(x)-2} u, & \text { for } x \in \Omega  \tag{10}\\ u=0, & \text { for } x \in \partial \Omega\end{cases}
$$

We seek solutions for problem (10) belonging to the space $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ in the sense given below.
Definition 1. We say that $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is a weak solution for problem (10) if

$$
\int_{\Omega}\left\{\sum_{i=1}^{N}\left(\left|\partial_{x_{i}} u\right|^{p_{i}(x)-2} \partial_{x_{i}} u \partial_{x_{i}} v\right)-A|u|^{a(x)-2} u v-B|u|^{b(x)-2} u v\right\} d x=0
$$

for all $v \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$.
The main result of this paper is given by the following theorem.
Theorem 2. There exists $\mu>0$ such that, for any $A \in(0, \mu)$ and any $B \in(0, \mu)$, problem (10) has at least two distinct nontrivial weak solutions.

We point out that the result of Theorem 2 can be regarded as a generalization of Theorem 1 in [20], where a similar problem involving the $p(x)$-Laplace operator was studied.

## 4 Proof of Theorem 2

We start by introducing the energy functional corresponding to problem (10), that is $J: W_{0}^{1, \vec{p}(\cdot)}(\Omega) \rightarrow \mathbb{R}$,

$$
\begin{equation*}
J(u)=\int_{\Omega}\left\{\sum_{i=1}^{N} \frac{\left|\partial_{x_{i}} u\right|^{p_{i}(x)}}{p_{i}(x)}-A \frac{|u|^{a(x)}}{a(x)}-B \frac{|u|^{b(x)}}{b(x)}\right\} d x . \tag{11}
\end{equation*}
$$

Standard arguments assure that $J \in C^{1}\left(W_{0}^{1, \vec{p}(\cdot)}(\Omega), \mathbb{R}\right)$ and its Fréchet derivative is given by

$$
\begin{equation*}
\left\langle J^{\prime}(u), v\right\rangle=\int_{\Omega}\left\{\sum_{i=1}^{N}\left|\partial_{x_{i}} u\right|^{p_{i}(x)-2} \partial_{x_{i}} u \partial_{x_{i}} v-A|u|^{a(x)-2} u v-B|u|^{b(x)-2} u v\right\} d x, \tag{12}
\end{equation*}
$$

for all $u, v \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. Thus, the weak solutions of problem (10) are exactly the critical points of J. We shall prove that the functional J possesses two distinct critical points using as main tools the Mountain Pass Theorem (see, e.g. [3] or [39]) and Ekeland's Variational Principle (see, e.g. [9]).

The following lemma will be essential in proving our main result.
Lemma 1. The following assertions hold.
(i) There exists $\mu>0$ such that for any $A, B \in(0, \mu)$ we can find $\rho_{0}>0$ and $a>0$ such that

$$
J(u) \geq a>0, \quad \forall u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \text { with }\|u\|_{\vec{p}(\cdot)}=\rho_{0}
$$

(ii) There exists $\varphi \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that

$$
\lim _{t \rightarrow \infty} J(t \varphi)=-\infty
$$

(iii) There exists $\Phi \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that $\Phi \geq 0, \Phi \neq 0$ and

$$
J(t \Phi)<0
$$

for $t>0$ small enough.
Proof. (i) By condition (9) we have $1<a(x)<b(x)<P_{-, \infty}$, for all $x \in \bar{\Omega}$ and, consequently, Theorem 1 assures that $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is continuously and compactly embedded in $L^{a(x)}(\Omega)$ and $L^{b(x)}(\Omega)$.

The fact that $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is continuously and compactly embedded in $L^{a(x)}(\Omega)$ assures that there exists a positive constant $C_{1}$ such that

$$
\begin{equation*}
|u|_{a(\cdot)} \leq C_{1} \cdot\|u\|_{\vec{p}(\cdot)}, \quad \forall u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) . \tag{13}
\end{equation*}
$$

Similarly, $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is continuously and compactly embedded in $L^{b(x)}(\Omega)$ and this guarantees that there exists a positive constant $C_{2}$ such that

$$
\begin{equation*}
|u|_{b(\cdot)} \leq C_{2} \cdot\|u\|_{\vec{p}(\cdot)}, \quad \forall u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \tag{14}
\end{equation*}
$$

We fix $\rho_{0} \in(0,1)$ such that $\rho_{0}<\min \left\{\frac{1}{C_{1}}, \frac{1}{C_{2}}\right\}$. Then relations (13) and (14) imply

$$
|u|_{a(\cdot)}<1, \text { for all } u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \text { with }\|u\|_{\vec{p}(\cdot)}=\rho_{0}
$$

and

$$
|u|_{b(\cdot)}<1, \text { for all } u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \text { with }\|u\|_{\vec{p}(\cdot)}=\rho_{0}
$$

Furthermore, relation (5) yields

$$
\begin{equation*}
\int_{\Omega}|u|^{a(x)} d x \leq|u|_{a(\cdot)}^{a^{-}} \text {, for all } u \in W_{0}^{1, \vec{p}(\cdot)} \text { with }\|u\|_{\vec{p}(\cdot)}=\rho_{0} \tag{15}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{\Omega}|u|^{b(x)} d x \leq|u|_{b(\cdot)}^{b^{-}} \text {, for all } u \in W_{0}^{1, \vec{p}(\cdot)} \text { with }\|u\|_{\vec{p}(\cdot)}=\rho_{0} \tag{16}
\end{equation*}
$$

Relations (13) and (15) imply
$\int_{\Omega}|u|^{a(x)} d x \leq C_{1}^{a^{-}}\|u\|_{\vec{p}(\cdot)^{\prime}}^{a^{-}} \quad$ for all $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ with $\|u\|_{\vec{p}(\cdot)}=\rho_{0}$.
By relations (14) and (16) we get

$$
\begin{equation*}
\int_{\Omega}|u|^{b(x)} d x \leq C_{2}^{b^{-}}\|u\|_{\vec{p}(\cdot)^{\prime}}^{b^{-}} \text {for all } u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \text { with }\|u\|_{\vec{p}(\cdot)}=\rho_{0} \tag{18}
\end{equation*}
$$

Using relation (5), for all $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ with $\|u\|_{\vec{p}(\cdot)}<1$, we obtain

$$
\begin{align*}
\frac{\|u\|_{\vec{p}(\cdot)}^{P_{+}^{+}}}{N^{P_{+}^{+}-1}}=N\left(\frac{\sum_{i=1}^{N}\left|\partial_{x_{i}} u\right|_{p_{i}(\cdot)}}{N}\right)^{P_{+}^{+}} \leq & \sum_{i=1}^{N}\left|\partial_{x_{i}} u\right|_{p_{i}(\cdot)}^{P_{+}^{+}} \leq \\
& \sum_{i=1}^{N}\left|\partial_{x_{i}} u\right|_{p_{i}(\cdot)}^{p_{i}^{+}} \leq \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} u\right|^{p_{i}(x} d x \tag{19}
\end{align*}
$$

Relations (19), (18) and (17) show that for any $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ with $\|u\|_{\vec{p}(\cdot)}=$ $\rho_{0}$ we have

$$
\begin{aligned}
J(u) \geq & \frac{1}{P_{+}^{+} N^{P_{+}^{+}-1}}\|u\|_{+}^{P_{+}^{+}}(\cdot) \\
= & \frac{A}{a^{-}} C_{1}^{a^{-}}\|u\|_{\vec{p}(\cdot)}^{a^{-}}-\frac{B}{b^{-}} C_{2}^{b^{-}}\|u\|^{\frac{b}{p}}(\cdot) \\
= & \rho_{0}^{a_{+}^{+}-1}
\end{aligned} \rho_{0}^{P_{+}^{+}}-\frac{A}{a^{-}}\left(\frac{1}{2 P_{+}^{+} N^{P_{+}^{+}-1}} \rho_{1}^{a^{-}} \rho_{0}^{a_{+}^{+}}-\frac{B}{b^{-}} C_{2}^{b^{-}} \rho_{0}^{b^{-}}-\frac{A}{a^{-}} C_{1}^{a^{-}}\right)+.
$$

Defining

$$
\mu_{1}=\frac{1}{4 P_{+}^{+} N^{P_{+}^{+}-1}} \rho_{0}^{P_{+}^{+}-a^{-}} \frac{a^{-}}{C_{1}^{a^{-}}} \text {and } \mu_{2}=\frac{1}{4 P_{+}^{+} N^{P_{+}^{+}-1}} \rho_{0}^{P_{+}^{+}-b^{-}} \frac{b^{-}}{C_{2}^{b^{-}}}
$$

simple computations show that

$$
\rho_{0}^{a^{-}}\left(\frac{1}{2 P_{+}^{+} N^{P_{+}^{+}-1}} \rho_{0}^{P_{+}^{+}-a^{-}}-\frac{A}{a^{-}} C_{1}^{a^{-}}\right) \geq \frac{1}{4 P_{+}^{+} N^{P_{+}^{+}-1}} \rho_{0}^{P_{+}^{+}}, \quad \forall A \in\left(0, \mu_{1}\right)
$$

and

$$
\rho_{0}^{P_{+}^{+}}\left(\frac{1}{2 P_{+}^{+} N^{P_{+}^{+}-1}}-\frac{B}{b^{-}} C_{2}^{b^{-}} \rho_{0}^{b^{-}-P_{+}^{+}}\right) \geq \frac{1}{4 P_{+}^{+} N^{P_{+}^{+}-1}} \rho_{0}^{P_{+}^{+}}, \quad \forall B \in\left(0, \mu_{2}\right)
$$

Consequently, defining

$$
\begin{equation*}
\mu:=\min \left\{\mu_{1}, \mu_{2}\right\} \tag{20}
\end{equation*}
$$

and

$$
a:=\frac{1}{4 P_{+}^{+} N^{P_{+}^{+}-1}} \rho_{0}^{P_{+}^{+}}
$$

we conclude that for any $A \in(0, \mu)$ and any $B \in(0, \mu)$ we have

$$
J(u) \geq a>0
$$

for all $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ with $\|u\|_{\vec{p}(\cdot)}=\rho_{0}$, where $\rho_{0}$ was fixed such that $\rho_{0} \in$ $\left(0, \min \left\{1,1 / C_{1}, 1 / C_{2}\right\}\right)$ at the beginning of the proof of (i).
(ii) Let $\varphi \in C_{0}^{\infty}(\Omega), \varphi \geq 0, \varphi \neq 0$ and $t>1$. We have

$$
\begin{aligned}
J(t \varphi) & =\int_{\Omega}\left\{\sum_{i=1}^{N} \frac{t^{p_{i}(x)}}{p_{i}(x)}\left|\partial_{x_{i}} \varphi\right|^{p_{i}(x)}-A \frac{t^{a(x)}}{a(x)}|\varphi|^{a(x)}-B \frac{t^{b(x)}}{b(x)}|\varphi|^{b(x)}\right\} d x \\
& \leq \frac{t^{P_{+}^{+}}}{P_{-}^{-}} \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} \varphi\right|^{p_{i}(x)} d x-A \frac{t^{a^{-}}}{a^{+}} \int_{\Omega}|\varphi|^{a(x)} d x-B \frac{t^{b^{-}}}{b^{+}} \int_{\Omega}|\varphi|^{b(x)} d x \\
& \leq \frac{t^{P_{+}^{+}}}{P_{-}^{-}} \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} \varphi\right|^{p_{i}(x)} d x-B \frac{t^{b^{-}}}{b^{+}} \int_{\Omega}|\varphi|^{b(x)} d x .
\end{aligned}
$$

Since $b^{-}>P_{+}^{+}$, by (9) we deduce that $\lim _{t \rightarrow \infty} J(t \varphi)=-\infty$ and, thus, (ii) is proved.
(iii) Let $\Phi \in C_{0}^{\infty}(\Omega), \Phi \geq 0, \Phi \neq 0$ and $t \in(0,1)$. We conclude that

$$
\begin{aligned}
J(t \Phi) & =\int_{\Omega}\left\{\sum_{i=1}^{N} \frac{t^{p_{i}(x)}}{p_{i}(x)}\left|\partial_{x_{i}} \Phi\right|^{p_{i}(x)}-A \frac{t^{a(x)}}{a(x)}|\Phi|^{a(x)}-B \frac{t^{b(x)}}{b(x)}|\Phi|^{b(x)}\right\} d x \\
& \leq \frac{t^{P_{-}^{-}}}{P_{-}^{-}} \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} \Phi\right|^{p_{i}(x)} d x-A \frac{t^{a^{+}}}{a^{+}} \int_{\Omega}|\Phi|^{a(x)} d x-B \frac{t^{b^{+}}}{b^{+}} \int_{\Omega}|\Phi|^{b(x)} d x \\
& \leq \frac{t^{P_{-}^{-}}}{P_{-}^{-}} \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} \Phi\right|^{p_{i}(x)} d x-A \frac{t^{a^{+}}}{a^{+}} \int_{\Omega}|\Phi|^{a(x)} d x<0
\end{aligned}
$$

for $t<\delta^{1 /\left(P_{-}^{-}-a^{+}\right)}$with

$$
0<\delta<\min \left\{1, \frac{A P_{-}^{-} \int_{\Omega}|\Phi|^{a(x)} d x}{a^{+} \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} \Phi\right|^{p_{i}(x)} d x}\right\}
$$

It follows that (iii) is proved.
Thus, the proof of Lemma 1 is complete.
Proof of Theorem 2. Let $\mu>0$ be defined as in (20) and $A \in(0, \mu), B \in(0, \mu)$.
Using Lemma 1 (i) and (ii) and the Mountain Pass Theorem (see, e.g. [3]) we deduce that there exists a sequence $\left\{u_{n}\right\}$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that

$$
\begin{equation*}
J\left(u_{n}\right) \rightarrow \bar{c} \text { and } J^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in }\left(W_{0}^{1, \vec{p}(\cdot)}(\Omega)\right)^{\star} \tag{21}
\end{equation*}
$$

where $\left(W_{0}^{1, \vec{p}(\cdot)}(\Omega)\right)^{\star}$ is the dual space of $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$.

First, we show that $\left\{u_{n}\right\}$ is bounded in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. Assume by contradiction the contrary. Then, passing if necessary to a subsequence, still denoted by $\left\{u_{n}\right\}$, we may assume that $\left\|u_{n}\right\|_{\vec{p}(\cdot)} \rightarrow \infty$ as $n \rightarrow \infty$. Thus, we may consider that $\left\|u_{n}\right\|_{\vec{p}(\cdot)}>1$ for any integer $n$. Relations (21) and the above considerations imply that for $n$ large enough it holds that

$$
\begin{aligned}
1+\bar{c}+\left\|u_{n}\right\|_{\vec{p}(\cdot)} \geq & \geq\left(u_{n}\right)-\frac{1}{b^{-}}\left\langle J^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
\geq & \left(\frac{1}{P_{+}^{+}}-\frac{1}{b^{-}}\right) \int_{\Omega} \sum_{i=1}^{N}\left|\partial_{x_{i}} u_{n}\right|^{p_{i}(x)} d x+ \\
& A\left(\frac{1}{b^{-}}-\frac{1}{a^{-}}\right) \int_{\Omega}\left|u_{n}\right|^{a(x)} d x .
\end{aligned}
$$

For each $n$ and $i \in\{1, \ldots, N\}$ we define

$$
\xi_{n, i}= \begin{cases}P_{+}^{+}, & \text {if }\left|\partial_{x_{i}} u_{n}\right|_{p_{i}(\cdot)}<1 \\ P_{-}^{-}, & \text {if }\left|\partial_{x_{i}} u_{n}\right|_{p_{i}(\cdot)}>1\end{cases}
$$

Some elementary computations show that for all $u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ we have

$$
\begin{equation*}
\frac{\|u\|_{\vec{p}(\cdot)}^{P_{-}^{-}}}{N^{P_{-}^{-}-1}}=N\left(\frac{\sum_{i=1}^{N}\left|\partial_{x_{i}} u\right|_{p_{i}(\cdot)}}{N}\right)^{P_{-}^{-}} \leq \sum_{i=1}^{N}\left|\partial_{x_{i}} u\right|_{p_{i}(\cdot)}^{P_{-}^{-}} \tag{22}
\end{equation*}
$$

On the other hand, we point out that

$$
\begin{equation*}
\int_{\Omega}|u|^{a(x)} d x \leq|u|_{a(\cdot)}^{a^{+}}+|u|_{a(\cdot)}^{a^{-}}, \text {for all } u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \tag{23}
\end{equation*}
$$

Since, by Theorem 1, $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is compactly embedded in $L^{a(x)}(\Omega)$ and relation (23) holds true it follows that there exists a positive constant $C_{3}$ such that

$$
\begin{equation*}
\int_{\Omega}|u|^{a(x)} d x \leq C_{3}\left(\|u\|_{\vec{p}(\cdot)}^{a^{+}}+\|u\|_{\vec{p}(\cdot)}^{a^{-}}\right), \text {for all } u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \tag{24}
\end{equation*}
$$

Since $a^{-}<b^{-}$(see relation (9)), using relation (24) we find that there exists a positive constant $C_{4}$ such that

$$
\begin{equation*}
A\left(\frac{1}{b^{-}}-\frac{1}{a^{-}}\right) \int_{\Omega}|u|^{a(x)} d x \geq-C_{4}\left(\|u\|_{\vec{p}(\cdot)}^{a^{+}}+\|u\| \frac{a_{\vec{p}}^{-}}{\vec{p}(\cdot)}\right), \text { for all } u \in W_{0}^{1, \vec{p}(\cdot)}(\Omega) \tag{25}
\end{equation*}
$$

Using relations (4), (5), (22) and (25) we infer that for $n$ large enough we have

$$
\begin{aligned}
& 1+\bar{c}+\|u\|_{\vec{p}(\cdot)} \\
& \geq\left(\frac{1}{P_{+}^{+}}-\frac{1}{b^{-}}\right) \int_{\Omega} \sum_{i=1}^{N}\left|\partial_{x_{i}} u_{n}\right|^{p_{i}(x)} d x+A\left(\frac{1}{b^{-}}-\frac{1}{a^{-}}\right) \int_{\Omega}\left|u_{n}\right|^{a(x)} d x \\
& \geq\left(\frac{1}{P_{+}^{+}}-\frac{1}{b^{-}}\right) \sum_{i=1}^{N}\left|\partial_{x_{i}} u_{n}\right|_{p_{i}(\cdot)}^{\xi_{n, i}}-C_{4}\left(\|u\|_{\vec{p}(\cdot)}^{a^{+}}+\|u\|_{\vec{p}(\cdot)}^{a^{-}}\right) \\
& \geq\left(\frac{1}{P_{+}^{+}}-\frac{1}{b^{-}}\right) \sum_{i=1}^{N}\left|\partial_{x_{i}} u_{n}\right|_{p_{i}(\cdot)}^{P_{-}^{-}}-\left(\frac{1}{P_{+}^{+}}-\frac{1}{b^{-}}\right) \\
& \sum_{\left\{i ; \xi_{n, i}=P_{+}^{+}\right\}}\left(\left|\partial_{x_{i}} u_{n}\right|_{p_{i}(\cdot)}^{P_{-}^{-}}-\left|\partial_{x_{i}} u_{n}\right|_{p_{i}(\cdot)}^{P_{+}^{+}}\right)-C_{4}\left(\|u\| \frac{a_{\vec{p}}^{+}(\cdot)}{a^{+}}+\|u\| \frac{\vec{p}_{\vec{p}(\cdot)}}{a^{-}}\right) \\
& \geq\left(\frac{1}{P_{+}^{+}}-\frac{1}{b^{-}}\right) \frac{1}{N^{P_{-}^{-}-1}}\left\|u_{n}\right\|_{\vec{p}(\cdot)}^{P_{-}^{-}}-N\left(\frac{1}{P_{+}^{+}}-\frac{1}{b^{-}}\right) \\
& -C_{4}\left(\|u\|_{\vec{p}(\cdot)}^{a^{+}}+\|u\|_{\vec{p}(\cdot)}^{a^{-}}\right),
\end{aligned}
$$

where $C_{4}$ is a positive constant.
Taking into account that condition (9) holds true, dividing the above inequality by $\left\|u_{n}\right\|_{\vec{p}(\cdot)}^{P_{-}^{-}}$and passing to the limit as $n \rightarrow \infty$ we obtain a contradiction.

It follows that $\left\{u_{n}\right\}$ is bounded in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. This information and the fact that $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is a reflexive space implies that there exist a subsequence, still denoted by $\left\{u_{n}\right\}$, and $u_{1} \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that $\left\{u_{n}\right\}$ converges weakly to $u_{1}$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. Since, by Theorem 1 , the space $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is compactly embedded in $L^{a(x)}(\Omega)$ and $L^{b(x)}(\Omega)$, we conclude that $\left\{u_{n}\right\}$ converges strongly to $u_{1}$ in $L^{a(x)}(\Omega)$ and $L^{b(x)}(\Omega)$. Then, by inequality (3), we deduce

$$
\lim _{n \rightarrow \infty} \int_{\Omega}\left|u_{n}\right|^{a(x)-2} u_{n}\left(u_{n}-u_{1}\right) d x=0,
$$

and

$$
\lim _{n \rightarrow \infty} \int_{\Omega}\left|u_{n}\right|^{b(x)-2} u_{n}\left(u_{n}-u_{1}\right) d x=0
$$

On the other hand, by relation (21) we have

$$
\lim _{n \rightarrow \infty}\left\langle J^{\prime}\left(u_{n}\right), u_{n}-u_{1}\right\rangle=0
$$

Thus, by using the above equations, we get

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} u_{n}\right|^{p_{i}(x)-2} \partial_{x_{i}} u_{n}\left(\partial_{x_{i}} u_{n}-\partial_{x_{i}} u_{1}\right) d x=0 \tag{26}
\end{equation*}
$$

Relation (26) and the fact that $\left\{u_{n}\right\}$ converges weakly to $u_{1}$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ imply

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \sum_{i=1}^{N} \int_{\Omega}\left(\left|\partial_{x_{i}} u_{n}\right|^{p_{i}(x)-2} \partial_{x_{i}} u_{n}-\left|\partial_{x_{i}} u_{1}\right|^{p_{i}(x)-2} \partial_{x_{i}} u_{1}\right)\left(\partial_{x_{i}} u_{n}-\partial_{x_{i}} u_{1}\right) d x=0 \tag{27}
\end{equation*}
$$

Next, we recall that the following elementary inequality

$$
\begin{equation*}
\left(|\eta|^{t-2} \eta-|\varrho|^{t-2} \varrho\right)(\eta-\varrho) \geq 2^{-t}|\eta-\varrho|^{t}, \text { for all } \eta, \varrho \in \mathbb{R} \text {, } \tag{28}
\end{equation*}
$$

is valid for all $t \geq 2$.
Applying the above inequality in relation (27) we get

$$
\lim _{n \rightarrow \infty} \sum_{i=1}^{N} \int_{\Omega}\left|\partial_{x_{i}} u_{n}-\partial_{x_{i}} u_{1}\right|^{p_{i}(x)} d x=0
$$

and, consequently, $\left\{u_{n}\right\}$ converges strongly to $u_{1}$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$.
Then, by relation (21), we get

$$
J\left(u_{1}\right)=\bar{c}>0 \quad \text { and } \quad J^{\prime}\left(u_{1}\right)=0
$$

that is, $u_{1}$ is a nontrivial weak solution of problem (10).
Next, we prove that there exists a second weak solution $u_{2} \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that $u_{2} \neq u_{1}$.

By Lemma 1 (i), on the boundary of the ball centered at the origin and of radius $\rho_{0}$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$, denoted by $B_{\rho_{0}}(0)$, we have

$$
\inf _{\partial B_{\rho_{0}}(0)} J>0 .
$$

On the other hand, by Lemma 1 (iii), there exists $\Phi \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that $\Phi \geq 0$, $\Phi \neq 0$ and $J(t \Phi)<0$, for $t>0$ small enough. Moreover, by relations (17), (18), (19) and (5) we obtain that for any $u \in B_{\rho_{0}}(0)$, the inequality

$$
\begin{equation*}
J(u) \geq \frac{1}{P_{+}^{+} N^{P_{+}^{+}-1}}\|u\|_{\vec{p}(\cdot)}^{P_{+}^{+}}-\frac{A}{a^{-}} C_{1}^{a^{-}}\|u\|_{\vec{p}(\cdot)}^{a^{-}}-\frac{B}{b^{-}} C_{2}^{b^{-}}\|u\|_{\frac{b_{p}}{-}(\cdot)}^{b^{-}} \tag{29}
\end{equation*}
$$

holds true and it follows that

$$
\begin{equation*}
-\infty<\underline{c}:=\frac{\inf }{B_{\rho_{0}}(0)} J<0 . \tag{30}
\end{equation*}
$$

Particularly, we have found that J is bounded from below.
Let $\varepsilon$ such that

$$
\begin{equation*}
0<\varepsilon<\inf _{\partial B_{\rho_{0}}(0)} J-\inf _{B_{\rho_{0}}(0)} J . \tag{31}
\end{equation*}
$$

The same arguments as in the proof of Lemma 3.4 in [25] can be used in order to show that J is weakly lower semi-continuous on $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$.

Now, we concentrate our attention on functional $J: \overline{B_{\rho_{0}}(0)} \rightarrow \mathbb{R}$. Since $J \in C^{1}\left(\overline{B_{\rho_{0}}(0)}, \mathbb{R}\right), J$ is bounded from below in $\overline{B_{\rho_{0}}(0)}$ and $J$ is weakly lower semi-continuous on $\overline{B_{\rho_{0}}(0)}$, we can apply Ekeland's Variational Principle for $J$ (see [9]) in order to obtain that there exists $u_{\varepsilon} \in \overline{B_{\rho_{0}}(0)}$ such that:

1) $J\left(u_{\varepsilon}\right) \leq \frac{\inf }{B_{\rho_{0}}(0)} J+\varepsilon$
and
2) $J\left(u_{\varepsilon}\right)<J(u)+\varepsilon \cdot\left\|u-u_{\varepsilon}\right\|_{\vec{p}(\cdot)}, \forall u \in \overline{B_{\rho_{0}}(0)}$ with $u \neq u_{\varepsilon}$.

Actually, we have $J\left(u_{\varepsilon}\right) \leq \frac{\inf }{B_{\rho_{0}(0)}} J+\varepsilon \leq \inf _{B_{\rho_{0}}(0)} J+\varepsilon<\inf _{\partial B_{\rho_{0}}(0)} J$, since the last inequality holds true, then we get $u_{\varepsilon} \in B_{\rho_{0}}(0)$.

Now, we let $I: \overline{B_{\rho_{0}}(0)} \rightarrow \mathbb{R}$ defined by

$$
I(u)=J(u)+\varepsilon \cdot\left\|u-u_{\varepsilon}\right\|_{\vec{p}(\cdot)}, \quad \forall u \in \overline{B_{\rho_{0}}(0)} .
$$

It is clear that $u_{\varepsilon}$ is a minimum point of functional I and thus

$$
\begin{equation*}
\frac{I\left(u_{\varepsilon}+t \cdot v\right)-I\left(u_{\varepsilon}\right)}{t} \geq 0 \tag{32}
\end{equation*}
$$

for a small $t>0$ and $v \in B_{1}(0)$. By relation (32), we deduce that

$$
\frac{J\left(u_{\varepsilon}+t \cdot v\right)-J\left(u_{\varepsilon}\right)}{t}+\varepsilon\|v\|_{\vec{p}(\cdot)} \geq 0
$$

for a small $t>0$ and $v \in B_{1}(0)$.
Passing, in the above inequality, to the limit as $t \rightarrow 0$, it follows that $\left\langle J^{\prime}\left(u_{\varepsilon}\right), v\right\rangle+$ $\varepsilon\|v\|_{\vec{p}(\cdot)}>0$ and we infer that $\left\|J^{\prime}\left(u_{\varepsilon}\right)\right\| \leq \varepsilon$. This implies the existence of a sequence $\left\{z_{n}\right\}$ in $B_{\rho_{0}}(0)$ such that

$$
\begin{equation*}
J\left(z_{n}\right) \rightarrow \underline{c} \text { and } J^{\prime}\left(z_{n}\right) \rightarrow 0 \tag{33}
\end{equation*}
$$

It is obvious that $\left\{z_{n}\right\}$ is bounded in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. Since $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ is a reflexive space and $\left\{z_{n}\right\}$ is a bounded sequence, then there exists $u_{2} \in W_{0}^{1, \vec{p}(\cdot)}(\Omega)$ such that, up to a subsequence, $\left\{z_{n}\right\}$ converges weakly to $u_{2}$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$. Using similar arguments as in the case of the weak solution $u_{1}$, we can show that $\left\{z_{n}\right\}$ converges strongly to $u_{2}$ in $W_{0}^{1, \vec{p}(\cdot)}(\Omega)$.

Then, since $J \in C^{1}\left(W_{0}^{1, \vec{p}(\cdot)}(\Omega), \mathbb{R}\right)$ and relation (33) holds true, we obtain

$$
\begin{equation*}
J\left(u_{2}\right)=\underline{c}<0 \quad \text { and } \quad J^{\prime}\left(u_{2}\right)=0, \tag{34}
\end{equation*}
$$

that is, $u_{2}$ is a nontrivial weak solution for problem (10).
Finally, we conclude that $u_{1} \neq u_{2}$ since

$$
J\left(u_{1}\right)=\bar{c}>0>\underline{c}=J\left(u_{2}\right)
$$

Thus, Theorem 2 is completely proved.
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