# Invariant nuclear space of a second quantization operator 

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#### Abstract

Let $S^{\prime}(\mathbf{R})$ be the dual of the Schwartz space $S(\mathbf{R})$, A a self-adjoint operator in $L^{2}(\mathbf{R})$ and $d \Gamma(\mathbf{A})^{*}$ the adjoint operator of $d \Gamma(\mathbf{A})$ which is the second quantization operator of $\mathbf{A}$. It is proven that under a suitable condition on $\mathbf{A}$ there exists a nuclear subspace $S$ of a fundamental space $S_{\mathbf{A}}$ of Hida's type on $S^{\prime}(\mathbf{R})$ such that $d \Gamma(\mathbf{A}) S \subset S$ and $e^{-t a \Gamma(\mathbf{A})} S \subset S$, which enables us to show that a stochastic differential equation arising from the central limit theorem for spatially extended neurons:


$$
d X(t)=d W(t)-d \Gamma(\mathbf{A})^{*} X(t) d t
$$

has a unique solution on the dual space $S^{\prime}$ of $S$, where $W(t)$ is an $S^{\prime}$-valued Wiener process.

## 1. Introduction

Concerning with infinite dimensional geometry and analysis, several types of fundamental spaces on infinite dimensional topological vector spaces have attracted several authors ([1], [4], [9], [11], [13], [14]). As it has been known by [5], the nuclearity of the space gives us the regularization theorem which guarantees the existence of a strong solution of the stochastic differential equation. However, [7] tried to construct a unique weak solution of a Segal-Langevin type stochastic differential equation on a suitable space of infinite dimensional generalized functionals which is not nuclear, and the fundamental spaces used in the Malliavin calculus are known not to be nuclear [2]. With this background, we consider spaces of Hida's type which are nuclear.

Let $\left(S_{\mathbf{A}}\right)$ be a fundamental space of Hida's type and $d \Gamma(\mathbf{A})$ the second quantization operator. Inspired by the works [11], [12], we construct a fundamental space which is invariant under the semi-group $e^{-t d \Gamma(\mathbf{A})}$ and is nuclear and smaller than $\left(S_{\mathbf{A}}\right)$ even if $\left(S_{\mathbf{A}}\right)$ is not nuclear. This enables us to obtain a unique strong solution of the stochastic differential equation

$$
\begin{equation*}
d X(t)=d W(t)-d \Gamma(\mathbf{A})^{*} X(t) d t \tag{1.1}
\end{equation*}
$$

which is a special case of the types considered in [7].
First we begin with giving some notations and explanations. Let $E$ be a real locally convex topological vector space and $E^{\prime}$ the topological dual space of $E$. We denote by $\langle$,$\rangle the pairing of E$ and $E^{\prime}$, and by $|\cdot|_{E}$ the norm of $E$ when $E$ is a Hilbert space. Let $\mathscr{H}$ be a separable real Hilbert space densely and continuously embedded in $E$. Then identifying $\mathscr{H}^{\prime}$ with $\mathscr{H}$, we have

$$
\begin{equation*}
E^{\prime} \subset \mathscr{H} \subset E . \tag{1.2}
\end{equation*}
$$

Let $\mu$ be the countably additive Gaussian measure on $E$ whose characteristic functional is given by

$$
\begin{equation*}
\int_{E} \exp [i\langle x, \xi\rangle] d \mu(x)=\exp \left[-\frac{1}{2}|\xi|_{\mathscr{H}}^{2}\right], \quad \xi \in E^{\prime} \tag{1.3}
\end{equation*}
$$

If we replace $E$ by $S^{\prime}(\mathbf{R}),(E, \mu)$ is called the white noise space [11].
Now we state our main result. Let $\mathbf{A}$ be a self-adjoint operator in Hilbert space $\mathscr{H}$ and $L^{2}(E, \mu)$ the space of square integrable functions with respect to $\mu$. Further we denote by $\left(S_{\mathrm{A}}\right)$ a fundamental space of Hida's type determined by $\mathbf{A}$ and denote by $d \Gamma(\mathbf{A})$ the second quantization operator of A, which will be precisely defined later. From now on we denote the domain of a linear operator $\mathbf{T}$ defined densely in $\mathscr{H}$ by $\mathscr{D}(\mathbf{T})$ and set $C^{\infty}(\mathbf{T})=$ $\bigcap_{n=1}^{\infty} \mathscr{D}\left(\mathbf{T}^{n}\right)$. We always consider $\mathscr{D}\left(\mathbf{T}^{n}\right)$ as a Hilbert space equipped with the inner product $\left(\mathbf{T}^{n} \cdot, \mathbf{T}^{n} \cdot\right)_{\mathscr{H}}$. We mean by $\mathbf{A} \geq \lambda, \lambda \in \mathbf{R}$, that $(\mathbf{A} f, f)_{\mathscr{H}} \geq \lambda(f, f)_{\mathscr{H}}$ for all $f \in \mathscr{D}(\mathbf{A})$.

Theorem 1.1. Let A be a self-adjoint operator in $\mathscr{H}$. Suppose that $\mathbf{A} \geq 1+\varepsilon$, for some $\varepsilon>0$ and there exist a self-adjoint operator $\mathbf{B}$ in $\mathscr{H}$ and natural numbers $p$ and $q$, satisfying the following conditions;
(1) $\mathscr{D}\left(\mathbf{B}^{p}\right) \subset \mathscr{D}(\mathbf{A})$,
(2) the identity map of $\mathscr{D}\left(\mathbf{B}^{q}\right)$ into $\mathscr{H}$ is a Hilbert Schmidt operator,
(3) $\mathbf{A} C^{\infty}(\mathbf{B}) \subset C^{\infty}(\mathbf{B})$.

Then there exists a nuclear subspace $S$ of $L^{2}(E, \mu)$ such that

$$
d \Gamma(\mathbf{A}) S \subset S
$$

Further, suppose that
(4) $e^{-t \boldsymbol{A}} C^{\infty}(\mathbf{B}) \subset C^{\infty}(\mathbf{B})$. Then

$$
e^{-t d \Gamma(\mathbf{A})} S \subset S
$$

If there exists a positive self-adjoint operator $\mathbf{B}$ and $p, q \in \mathbf{N}$ satisfying the
conditions (1) $\sim(4)$ of Theorem 1.1, those conditions (1) $\sim(4)$ still hold when we replace $\mathbf{B}$ by $\mathbf{B}+2 \mathbf{I}$. Then, since $\mathbf{B}+2 \mathbf{I} \geq 1+\varepsilon$ for some $\varepsilon>0$, the desired space $S$ in Theorem 1.1 will be given by $\left(S_{\mathbf{B}+21}\right)$.

## 2. Fundamental space of Hida's type

Before defining a fundamental space of Hida's type, we introduce the following notation. Let $\mathscr{H}$ be a separable Hilbert space. For $f_{i} \in \mathscr{H}$, $i=1,2, \cdots, n$, we denote the tensor product of them by

$$
\begin{equation*}
f_{1} \otimes f_{2} \otimes \cdots \otimes f_{n} \tag{2.1}
\end{equation*}
$$

and define the symmetric tensor product of them by

$$
\begin{equation*}
f_{1} \hat{\otimes} f_{2} \hat{\otimes} \cdots \hat{\otimes} f_{n}=\frac{1}{n!} \sum_{\sigma \in E_{n}} f_{\sigma(1)} \otimes f_{\sigma(2)} \otimes \cdots \otimes f_{\sigma(n)} \tag{2.2}
\end{equation*}
$$

where $\Xi_{n}$ is the symmetric group of degree $n$.
Let $\mathscr{G}$ and $\mathscr{F}$ be the sets of finite linear combinations of terms of (2.1) and (2.2) types, respectively. For $f_{i}, g_{i} \in \mathscr{H}, i=1,2, \cdots, n$, we first set

$$
\begin{equation*}
\left(f_{1} \otimes f_{2} \otimes \cdots \otimes f_{n}, g_{1} \otimes g_{2} \otimes \cdots \otimes g_{n}\right)_{\mathscr{H}^{\otimes n}}=\left(f_{1}, g_{1}\right)_{\mathscr{H}}\left(f_{2}, g_{2}\right)_{\mathscr{H}} \cdots\left(f_{n}, g_{n}\right)_{\mathscr{H}} \tag{2.3}
\end{equation*}
$$

and

$$
\begin{align*}
& \left(f_{1} \hat{\otimes} f_{2} \hat{\otimes} \cdots \hat{\otimes} f_{n}, g_{1} \hat{\otimes} g_{2} \hat{\otimes} \cdots \hat{\otimes} g_{n}\right)_{\mathscr{H}^{2 n}} \\
& \quad=\left(\frac{1}{n!}\right)^{2} \sum_{\sigma, \tau \in \Xi_{n}}\left(f_{\sigma(1)}, g_{\tau(1)}\right)_{\mathscr{H}}\left(f_{\sigma(2)}, g_{\tau(2)}\right)_{\mathscr{H}} \cdots\left(f_{\sigma(n)}, g_{\tau(n)}\right)_{\mathscr{H}} \tag{2.4}
\end{align*}
$$

Then the inner products $(\cdot, \cdot)_{\mathscr{H}^{\infty n}}$ on $\mathscr{G}$ and $(\cdot, \cdot)_{\mathscr{H}^{\otimes n}}$ on $\mathscr{F}$ are naturally extended for the linear combinations. Let $\mathscr{H}^{\otimes n}$ and $\mathscr{H}^{\otimes \theta}$ be the completions of $\mathscr{G}$ and $\mathscr{F}$ with respect to the inner products $(\cdot, \cdot)_{\mathscr{P}^{\otimes n}}$ and $(\cdot, \cdot)_{\mathscr{P}^{8 n}}$. Clearly

$$
\begin{equation*}
\mathscr{H}^{\hat{\otimes} n} \subset \mathscr{H}^{\otimes n} \tag{2.5}
\end{equation*}
$$

and

$$
\begin{align*}
& \left(f_{1} \hat{\otimes} f_{2} \hat{\otimes} \cdots \hat{\otimes} f_{n}, g_{1} \hat{\otimes} g_{2} \hat{\otimes} \cdots \hat{\otimes} g_{n}\right)_{\mathscr{H}^{\circ n}} \\
& \quad=\left(f_{1} \hat{\otimes} f_{2} \hat{\otimes} \cdots \hat{\otimes} f_{n}, g_{1} \hat{\otimes} g_{2} \hat{\otimes} \cdots \hat{\otimes} g_{n}\right)_{\varkappa^{8 n}} \tag{2.6}
\end{align*}
$$

We define the Wick ordering, denoted by $: x^{\otimes n}:$, for $x \in E$, according to the case where $E=S^{\prime}(\mathbf{R})$. First the Wick product : $\left\langle x, \xi_{1}\right\rangle\left\langle x, \xi_{2}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle$ : of random variables $\left\langle x, \xi_{k}\right\rangle, x \in E, \xi_{k} \in E^{\prime}, k=1,2, \cdots, n$, with respect to the probability space $(E, \mu)$ is defined by the following recursion relation [6];

$$
\begin{aligned}
&::\left\langle x, \xi_{1}\right\rangle:=\left\langle x, \xi_{1}\right\rangle \\
&:\left\langle x, \xi_{1}\right\rangle\left\langle x, \xi_{2}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle:=\left\langle x, \xi_{1}\right\rangle:\left\langle x, \xi_{2}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle: \\
&-\sum_{k=2}^{n} \int_{E}\left\langle x, \xi_{1}\right\rangle\left\langle x, \xi_{k}\right\rangle d \mu(x):\left\langle x, \xi_{2}\right\rangle \cdots\left\langle x,{ }^{`} \xi_{k}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle:, \quad n \geq 2
\end{aligned}
$$

where $\left\langle x,{ }^{`} \xi_{k}\right\rangle$ means that the term $\left\langle x, \xi_{k}\right\rangle$ is excluded in the product. Using the Wick product, we define the Wick ordering $: x^{\otimes n}$ : by

$$
\begin{equation*}
\left\langle: x^{\otimes n}:, \xi_{1} \hat{\otimes} \xi_{2} \hat{\otimes} \cdots \hat{\otimes} \xi_{n}\right\rangle \equiv:\left\langle x, \xi_{1}\right\rangle\left\langle x, \xi_{2}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle: . \tag{2.7}
\end{equation*}
$$

Let $\left\{e_{i}: i=0,1,2, \cdots\right\}$ be a complete orthonormal system in $\mathscr{H}$ taken from $E^{\prime}$. The well known Wiener-Ito theorem states that the space $L^{2}(E, \mu)$ has the following orthogonal decomposition

$$
L^{2}(E, \mu)=\bigoplus_{n=0}^{\infty} \mathbf{K}_{n}
$$

where $\mathbf{K}_{n}$ consists of $n$-homogeneous chaos, i.e. each $\varphi$ in $\mathbf{K}_{n}$ has the formal expression

$$
\varphi(x)=\left\langle: x^{\otimes n}:, \hat{f}_{n}\right\rangle, \quad \hat{f}_{n} \in \mathscr{H}^{\hat{\otimes} n} .
$$

In fact if $\hat{f}_{n}$ is represented by $\sum_{i_{1}, i_{2}, \ldots, i_{n}=0}^{\infty} a_{i_{1}, i_{2}, \ldots, i_{n}} e_{i_{1}} \hat{\otimes} e_{i_{2}} \hat{\otimes} \cdots \hat{\otimes} e_{i_{n}}$, then the right hand side of the above expression is given by

$$
\sum_{i_{1}, i_{2}, \ldots, i_{n}=0}^{\infty} a_{i_{1}, i_{2}, \ldots, i_{n}}\left\langle: x^{\otimes n}:, e_{i_{1}} \hat{\otimes} e_{i_{2}} \hat{\otimes} \cdots \hat{\otimes} e_{i_{n}}\right\rangle .
$$

Thus each $\psi \in L^{2}(E, \mu)$ can be represented uniquely in the following way;

$$
\begin{equation*}
\psi(x)=\sum_{n=0}^{\infty}\left\langle: x^{\otimes n}:, \hat{f}_{n}\right\rangle, \quad \mu-\text { a.e. } x \in E . \tag{2.8}
\end{equation*}
$$

Moreover, we have [11]

$$
\begin{equation*}
|\psi|_{L^{2}(E, \mu)}^{2}=\sum_{n=0}^{\infty} n!\left|\hat{f}_{n}\right|_{\mathscr{F}^{8} n}^{2} \tag{2.9}
\end{equation*}
$$

Let $\mathbf{A}$ be a positive self-adjoint operator in $\mathscr{H}$. Then there exists a unique positive self-adjoint operator $\Gamma(\mathbf{A})$ in $L^{2}(E, \mu)$ such that (see [12])

$$
\Gamma(\mathbf{A}) 1=1
$$

and for $\xi_{i} \in \mathscr{D}(\mathbf{A}), i=1,2, \cdots, n$,

$$
\begin{align*}
& \Gamma(\mathbf{A}):\left\langle x, \xi_{1}\right\rangle\left\langle x, \xi_{2}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle: \\
& \quad=:\left\langle x, \mathbf{A} \xi_{1}\right\rangle\left\langle x, \mathbf{A} \xi_{2}\right\rangle \cdots\left\langle x, \mathbf{A} \xi_{n}\right\rangle: . \tag{2.10}
\end{align*}
$$

We denote by $\mathscr{P}_{\mathbf{A}}$ the collection of all polynomials of the form

$$
\omega(x)=P\left(\left\langle x, \xi_{1}\right\rangle,\left\langle x, \xi_{2}\right\rangle, \cdots,\left\langle x, \xi_{m}\right\rangle\right), \quad \xi_{i} \in C^{\infty}(\mathbf{A}),
$$

where $P\left(t_{1}, \cdots, t_{m}\right)$ is a polynomial of $\left(t_{1}, \cdots, t_{m}\right)$. For each $p \in \mathbf{R}$ we define a semi-norm ${ }_{\mathrm{A}}\|\cdot\|_{2, p}$ by

$$
\begin{equation*}
\mathbf{A}\|\omega\|_{2, p}^{2}=\int_{E}\left|\Gamma(\mathbf{A})^{p} \omega(x)\right|^{2} d \mu(x) . \tag{2.11}
\end{equation*}
$$

It is not difficult to see that $\Gamma(\mathbf{A})^{p}=\Gamma\left(\mathbf{A}^{p}\right)$. By (2.7), each $\omega$ in $\mathscr{P}_{\mathbf{A}}$ has the following expression;

$$
\omega(x)=\sum_{n=0}^{\infty}\left\langle: x^{\otimes n}:, \hat{g}_{n}\right\rangle, \quad \hat{g}_{n} \in C^{\infty}(\mathbf{A})^{\hat{\otimes} n}
$$

where $C^{\infty}(\mathbf{A})^{\hat{\otimes} n}=\overbrace{C^{\infty}(\mathbf{A}) \hat{\otimes} C^{\infty}(\mathbf{A}) \hat{\otimes} \cdots \hat{\otimes} C^{\infty}(\mathbf{A})}^{n \text { times }}$ is the set of finite linear combinations of the form $\xi_{1} \hat{\otimes} \xi_{2} \hat{\otimes} \cdots \hat{\otimes} \xi_{n}$ with $\xi_{i} \in C^{\infty}(\mathbf{A}), i=1,2, \cdots, n$. We note that there exists a natural number $k(\omega)$ such that $\hat{g}_{n}=0$ for $n \geq k(\omega)$. Since

$$
\hat{g}_{n}=\sum_{i=1}^{m(n)} a_{i}(n) \xi_{i_{1}} \hat{\otimes} \xi_{i_{2}} \hat{\otimes} \cdots \hat{\otimes} \xi_{i_{n}}, \quad \xi_{i_{k}} \in C^{\infty}(\mathbf{A}), k=1,2, \cdots, n,
$$

by (2.9), $A\|\cdot\|_{2, p}^{2}$ can be also represented as

$$
\begin{equation*}
\mathbf{A}\|\omega\|_{2, p}^{2}=\sum_{n=0}^{\infty} n!\left|\left(\mathbf{A}^{p}\right)^{\otimes n} \hat{g}_{n}\right|_{\varkappa^{8 n}}^{2}, \tag{2.12}
\end{equation*}
$$

where

$$
\left(\mathbf{A}^{p}\right)^{\otimes n}=\mathbf{A}^{p} \otimes \mathbf{A}^{p} \otimes \cdots \otimes \mathbf{A}^{p}
$$

For $p \geq 0,\left(S_{\mathbf{A}}\right)_{p}$ is the completion of $\mathscr{P}_{\mathbf{A}}$ with respect to the semi-norm ${ }_{A}\|\cdot\|_{2, p}$. We define the fundamental space $\left(S_{\mathbf{A}}\right)$ of the Hida distributions on $E$ by

$$
\begin{equation*}
\left(S_{\mathrm{A}}\right)=\bigcap_{p \geq 0}\left(S_{\mathrm{A}}\right)_{p} \tag{2.13}
\end{equation*}
$$

Let $\left(S_{\mathbf{A}}\right)_{-p}$ be the topological dual space of $\left(S_{\mathbf{A}}\right)_{p}$. Then we have

$$
\begin{equation*}
\left(S_{\mathrm{A}}\right)^{\prime}=\bigcup_{p \geq 0}\left(S_{\mathbf{A}}\right)_{-p} \tag{2.14}
\end{equation*}
$$

There are several criteria for nuclearity of a fundamental space of Hida's type, such as ([2], [10]). Here we show a sufficient condition.

Proposition 2.1. Let $\mathbf{Z}$ be a positive self-adjoint operator in $\mathscr{H}$ with eigenvalues $\lambda_{k}>1, k=0,1, \cdots$, such that $\sum_{k=0}^{\infty} \lambda_{k}^{-\gamma}<+\infty$ for some $\gamma>0$. Then $\left(S_{\mathbf{Z}}\right)$ is a nuclear space.

Proof. Here we mimic a proof from [11]. It is sufficient to show that for any $p \geq 0$, there exist an $s>0$ such that the inclusion map $i:\left(S_{\mathbf{z}}\right)_{p+2 s} \rightarrow$ $\left(S_{\mathbf{z}}\right)_{p}$ is a nuclear operator. For the notational simplicity, we prove the assertion in the case where $p=0$. This is equivalent to show that the inclusion map $t: L^{2}(E, \mu) \rightarrow\left(S_{\mathbf{Z}}\right)_{-s}$ is a Hilbert-Schmidt operator. Let $\mathbf{N}_{0}=\{0,1,2, \cdots\}$ and let $\mathbf{I}_{n}$ be the set of all ordered $n$-tuples $\alpha=\left(\alpha_{1}, \ldots, \alpha_{n}\right), \alpha_{1} \leq \cdots \leq \alpha_{n}$ in $\mathbf{N}_{0}^{n}$. For $\alpha \in \mathbf{I}_{n}$, define

$$
n_{k}(\alpha)=\#\left\{j: \alpha_{j}=k\right\}, \quad n(\alpha)!=\prod_{k=0}^{\infty} n_{k}(\alpha)!
$$

Let $\left\{e_{k}: k=0,1,2, \cdots\right\}$ be a complete orthonormal system associated with eigenfunctions of $\mathbf{Z}$ such that $\mathbf{Z} e_{k}=\lambda_{k} e_{k}, k=0,1,2, \cdots$. For each $\alpha \in \mathbf{I}_{n}$ we set

$$
\begin{aligned}
\mathrm{H}_{\alpha}(x) & =1 \quad \text { for } n=0, \\
\mathrm{H}_{\alpha}(x) & =(n(\alpha)!)^{-\frac{1}{2}} \prod_{k=0}^{\infty}:\left\langle x, e_{k}\right\rangle^{n_{k}(\alpha)}: \quad \text { for } n \neq 0 \\
& =(n(\alpha)!)^{-\frac{1}{2}} \prod_{k=0}^{\infty}\left\langle: x^{\otimes n_{k}(\alpha)}:, e_{k}^{\hat{\otimes}_{n_{k}(\alpha)}}\right\rangle .
\end{aligned}
$$

Moreover we have

$$
\begin{aligned}
\Gamma(\mathbf{Z})^{r} \mathbf{H}_{\alpha}(x) & =(n(\alpha)!)^{-\frac{1}{2}} \prod_{k=0}^{\infty}\left\langle: x^{\otimes n_{k}(\alpha)}:,\left(\mathbf{Z}^{r}\right)^{\otimes n_{k}(\alpha)} e_{k}^{\hat{\otimes}_{k}(\alpha)}\right\rangle \\
& =\prod_{k=0}^{\infty}\left(\lambda_{k}^{r}\right)^{n_{k}(\alpha)} \mathbf{H}_{\alpha}(x), \quad \text { for } r \in \mathbf{R}
\end{aligned}
$$

and hence

$$
\begin{aligned}
\sum_{\alpha} \mathbf{z}\left\|l \mathbf{H}_{\alpha}\right\|_{2,-s}^{2} & =\sum_{\alpha}\left|\Gamma(\mathbf{Z})^{-s} \mathrm{H}_{\alpha}\right|_{L^{2}(E, \mu)}^{2} \\
& =\sum_{\alpha}\left\{\frac{1}{\prod_{k=0}^{\infty} \lambda_{k}^{n_{k}(\alpha)}}\right\}^{2 s} \\
& \leq \sum_{n=0}^{\infty}\left[\sum_{k_{1}, \ldots, k_{n}=0}^{\infty}\left\{\frac{1}{\lambda_{k_{1}} \cdots \lambda_{k_{n}}}\right\}^{2 s}\right] \\
& =\sum_{n=0}^{\infty}\left\{\sum_{k=0}^{\infty} \lambda_{k}^{-2 s}\right\}^{n}<\infty,
\end{aligned}
$$

provided that $\sum_{k=0}^{\infty} \lambda_{k}^{-2 s}<1$, which is valid for sufficiently large $s$ by the assumptions $\sum_{k=0}^{\infty} \lambda_{k}^{-\gamma}<+\infty$ and $\lambda_{k}>1$. This, together with the fact that $\left\{\mathrm{H}_{\alpha}(x): \alpha \in \mathbf{I}_{n}, n=0,1,2, \cdots\right\}$ forms a complete orthonormal system of $L^{2}(E, \mu)$, implies the inclusion map $i:\left(S_{\mathbf{z}}\right)_{p+2 s} \rightarrow\left(S_{\mathbf{z}}\right)_{p}$ is a Hilbert-Schmidt operator.

Example 2.1 ([4], [11]). Let $\mathbf{Z}$ be the operator

$$
\begin{equation*}
\mathbf{Z}=-\left(\frac{d}{d x}\right)^{2}+x^{2}+1 \tag{2.15}
\end{equation*}
$$

Then

$$
\mathbf{Z} e_{n}=(2 n+2) e_{n}, n=0,1,2, \cdots
$$

where $\left\{e_{n}: n=0,1,2, \cdots\right\}$ is the complete orthonormal system consisting of Hermite functions in $L^{2}(\mathbf{R})$. If we take $E=S^{\prime}(\mathbf{R})$, then $\left(S_{\mathbf{z}}\right)$ becomes a nuclear space by Proposition 2.1 and originally it is called the fundamental space of the Hida distributions.

## 3. Proof of Theorem 1.1

Before proving Theorem 1.1, we define the second quantization operator $d \Gamma(\mathbf{A})$ of a self-adjoint operator $\mathbf{A}$ as

$$
\begin{align*}
& d \Gamma(\mathbf{A}):\left\langle x, \xi_{1}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle: \\
= & \sum_{i=1}^{n}:\left\langle x, \xi_{1}\right\rangle \cdots\left\langle x, \mathbf{A} \xi_{i}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle: \\
= & \left\langle: x^{\otimes n}:, \mathbf{N}_{\mathbf{A}}\left(\xi_{1} \hat{\otimes} \cdots \hat{\otimes} \xi_{n}\right)\right\rangle, \tag{3.1}
\end{align*}
$$

where

$$
\mathbf{N}_{\mathbf{A}}=\mathbf{A} \otimes I \otimes \cdots \otimes I+I \otimes \mathbf{A} \otimes I \otimes \cdots \otimes I+\cdots+I \otimes \cdots \otimes I \otimes \mathbf{A}
$$

Let $\mathbf{B}$ be a self-adjoint operator which satisfies the conditions in Theorem 1.1 and set $\widetilde{\mathbf{B}}=\mathbf{B}+2$ I. Consider the fundamental space $S=\left(S_{\tilde{\mathbf{B}}}\right)$ of the Hida distributions on $E$. Take any $\omega(x) \in \mathscr{P}_{\tilde{\mathbf{B}}}$ then the following expression holds,

$$
\begin{equation*}
\omega(x)=\sum_{n=0}^{\infty}\left\langle: x^{\otimes n}:, \hat{h}_{n}\right\rangle, \quad \hat{h}_{n} \in C^{\infty}(\tilde{\mathbf{B}})^{\hat{\otimes}^{n}} . \tag{3.2}
\end{equation*}
$$

By the assumption (1) of Theorem 1.1, $\omega(x) \in \mathscr{D}(d \Gamma(\mathbf{A}))$, so that

$$
\begin{equation*}
d \Gamma(\mathbf{A}) \omega(x)=\sum_{n=0}^{\infty}\left\langle: x^{\otimes n}:, \mathbf{N}_{\mathbf{A}} \hat{h}_{n}\right\rangle \tag{3.3}
\end{equation*}
$$

Define the Hilbert space $\mathscr{H}_{\tilde{\mathbf{B}}^{\text {l }}}$ for any natural number $l$ by

$$
\mathscr{H}_{\tilde{\mathbf{B}}^{l}}=\mathscr{D}\left(\tilde{\mathbf{B}}^{l}\right) \equiv\left\{h \in \mathscr{H}:\left|\tilde{B}^{l} h\right|_{\mathscr{H}}<\infty\right\}
$$

where the inner product of the Hilbert space $\mathscr{H}_{\tilde{\mathbf{B}}^{2}}$ is given by $(\cdot, \cdot)_{\mathscr{H} \tilde{\mathbf{B}}^{I}}=$ $\left(\widetilde{\mathbf{B}}^{l} \cdot, \widetilde{\mathbf{B}}^{l} \cdot\right)_{\mathscr{H}}$.

Noticing that $C^{\infty}(\mathbf{B})=\bigcap_{n=1}^{\infty} \mathscr{D}\left(\mathbf{B}^{n}\right)$ is a complete metric space equipped with a countable system of semi norms $|\cdot|_{n}=\left|\mathbf{B}^{n} \cdot\right|_{\mathscr{H}}, n=1,2, \cdots$, for any natural number $k$, we have, by Baire's category theorem, a natural number $l_{1}(\geq k)$ and a constant $c_{1} \geq 1$ such that

$$
\begin{equation*}
\left|\tilde{\mathbf{B}}^{\tilde{}} \mathbf{A} h\right|_{\mathscr{H}} \leq c_{1}\left|\tilde{\mathbf{B}}^{l_{1}} h\right|_{\mathscr{H}}, \quad h \in C^{\infty}(\tilde{\mathbf{B}}) . \tag{3.4}
\end{equation*}
$$

For any $1 \leq n \leq k$ and any $\xi \in C^{\infty}(\mathbf{B})$ such that $|\xi|_{\mathscr{H}} \leq 1$, the assumption (2) of Theorem 1.1 implies that $\left(\mathbf{B}^{n} \mathbf{A} h, \xi\right)_{\mathscr{\not}}=\left(h, \mathbf{A B}^{n} \xi\right)_{\mathscr{H}}$ is continuous in $h$ on $C^{\infty}(\mathbf{B})$, so that noticing that

$$
\left|\mathbf{B}^{k} \mathbf{A} h\right|_{\mathscr{H}}=\sup \left\{\left(\mathbf{B}^{k} \mathbf{A} h, \xi\right)_{\mathscr{H}} ; \xi \in C^{\infty}(\mathbf{B}),|\xi|_{\mathscr{H}} \leq 1\right\}
$$

we see that the norm $\left|\mathbf{B}^{n} \mathbf{A} h\right|_{\mathscr{E}}$ is lower-semicontinuous in $h$ on $C^{\infty}(\mathbf{B})$. The assumption (3) of Theorem 1.1 yields that

$$
\left\{h \in C^{\infty}(\mathbf{B}) ;\left|\mathbf{B}^{n} \mathbf{A} h\right|_{\mathscr{H}}<\infty\right\}=\bigcup_{m=0}^{\infty}\left\{h \in C^{\infty}(\mathbf{B}) ;\left|\mathbf{B}^{n} \mathbf{A} h\right|_{\mathscr{H}} \leq m\right\} \supset C^{\infty}(\mathbf{B}) .
$$

Baire's category theorem then asserts that for some $a>0$ the closed set $\Delta_{a}=\left\{h \in C^{\infty}(\mathbf{B}) ;\left|\mathbf{B}^{n} \mathbf{A} h\right|_{\mathscr{H}} \leq a\right\}$ contains a ball of the form

$$
h_{0}+\left\{h \in C^{\infty}(\mathbf{B}) ;|h|_{(n)} \leq b\right\}
$$

with $h_{0} \in C^{\infty}(\mathbf{B})$ and $b>0$. If $|h|_{l(n)} \leq b, h_{0}-h \in \Delta_{a}$. Since $\left|\mathbf{B}^{n} \mathbf{A} h\right|_{\mathscr{H}}$ is symmetric, so $-h_{0}+h \in \Delta_{a}$. Since $h_{0}+h \in \Delta_{a}$, we have

$$
\left|\mathbf{B}^{n} \mathbf{A}(2 h)\right|_{\mathscr{H}}=\left|\mathbf{B}^{n} \mathbf{A}\left(-h_{0}+h+h_{0}+h\right)\right|_{\mathscr{H}} \leq 2 a .
$$

Thus we get for $1 \leq n \leq k$,

$$
\left|\mathbf{B}^{n} \mathbf{A} h\right|_{\mathscr{H}} \leq a b^{-1}\left|\mathbf{B}^{l(n)} h\right|_{\mathscr{H}}, \quad h \in C^{\infty}(\mathbf{B}) .
$$

For $p, q$ given in Theorem 1.1, set $p \vee q=\max \{p, q\}$. Then the lowersemicontinuity of $|\mathbf{A} h|_{\mathscr{H}}$ on $\mathscr{D}\left(\mathbf{B}^{p \vee q}\right)$ is proved similarly by the assumptions (1) and (2) of Theorem 1.1, so that by Baire's category theorem again we get

$$
|\mathbf{A} h|_{\mathscr{H}} \leq \text { const. }\left|\mathbf{B}^{p \vee q} h\right|_{\mathscr{H}}, \quad h \in \mathscr{D}\left(\mathbf{B}^{p \vee q}\right) .
$$

Since $C^{\infty}(\tilde{\mathbf{B}}) \subset C^{\infty}(\mathbf{B}) \subset \mathscr{D}\left(\mathbf{B}^{p \vee q}\right)$, the two inequalities above hold for $h \in C^{\infty}(\tilde{\mathbf{B}})$, which together with

$$
\left|\tilde{\mathbf{B}}^{k} \mathbf{A} h\right|_{\mathscr{H}} \leq \sum_{n=0}^{k} \frac{k!}{n!(k-n)!} 2^{k-n}\left|\mathbf{B}^{n} \mathbf{A} h\right|_{\mathscr{H}}
$$

completes the proof of (3.4)
Here we prepare a lemma concerning with the operator norm of tensor product of linear operators on Hilbert space, which will be used later. Given Hilbert spaces $\mathscr{H}_{1}$ and $\mathscr{H}_{2}$ and a linear operator T from $\mathscr{H}_{1}$ to $\mathscr{H}_{2}$, we denote the operator norm of T by

$$
\|\mathrm{T}\|_{\mathscr{H}_{1} \rightarrow \mathscr{H}_{2}}=\sup _{x \in \mathscr{H}_{1}} \frac{|\mathrm{~T} x|_{\mathscr{H}_{2}}}{|x|_{\mathscr{H}_{1}}} .
$$

We note that (3.4) implies

$$
\begin{equation*}
\|\mathbf{A}\|_{\mathscr{H}_{\mathbf{B}_{1} 1} \rightarrow \mathscr{H}_{\mathbf{B}^{k}}} \leq c_{1} . \tag{3.5}
\end{equation*}
$$

For simplicity we use the notation $\mathscr{H}_{\tilde{\mathbf{B}}^{\hat{1}} n}^{\hat{n}}$ instead of $\left(\mathscr{H}_{\tilde{\mathbf{B}}}\right)^{\hat{\otimes} n}$.
Let $U_{i}, i=1,2, \cdots, n$ be bounded linear operators from $\mathscr{H}_{\tilde{\mathbf{B}}^{2}}$ to $\mathscr{H}_{\tilde{\mathbf{B}}^{k}}$ for any natural number $l$ and $k$ such that $l \geq k$. We have the following lemma, which is an extended version of the proposition of [12] on p. 299.

Lemma 3.1. Let $U_{i}, i=1,2, \cdots, n$ be bounded linear operators from $\mathscr{H}_{\widetilde{\mathbf{B}} 1}$ to $\mathscr{H}_{\tilde{\mathbf{B}}^{k}}$. Then

Proof. Let $\left\{\theta_{k}: k=0,1,2, \cdots\right\}$ be a complete orthonormal basis of $\mathscr{H}_{\tilde{\mathbf{B}}^{2}},\left\{\zeta_{k}: k=0,1,2, \cdots\right\}$ a complete orthonormal basis of $\mathscr{H}_{\tilde{\mathbf{B}}^{k}}$ and $\Sigma c_{k_{1} \ldots k_{n}}\left(\theta_{k_{1}}\right.$ $\otimes \cdots \otimes \theta_{k_{i}} \otimes \zeta_{k_{i+1}} \otimes \cdots \otimes \zeta_{k_{n}}$ ) a finite sum in the space $\mathscr{H}_{\mathbf{B}^{i}}^{\otimes i} \otimes \mathscr{H}_{\mathbf{B}^{i}}^{\otimes^{n-i}}$. We get

$$
\begin{aligned}
& \mid \mathrm{I} \otimes \cdots \otimes \mathrm{I} \otimes U_{i} \otimes \mathrm{I} \otimes \cdots \otimes \mathrm{I} \\
& \left.\left(\sum c_{k_{1} \ldots k_{n}} \theta_{k_{1}} \otimes \cdots \otimes \theta_{k_{i}} \otimes \zeta_{k_{i}+1} \otimes \cdots \otimes \zeta_{k_{n}}\right)\right|_{\mathscr{H}_{\frac{\mathbf{B}^{i}}{}{ }^{-1-1} \otimes \mathscr{H}_{\mathbf{B}^{k}}^{\otimes n-i+1}}^{2}} \\
& =\sum \sum c_{k_{1} \ldots k_{n}} c_{k_{1} \ldots k_{i-1} l_{i} k_{i+1} \ldots k_{n}}\left(U_{i} \theta_{k_{i}}, U_{i} \theta_{l_{i}}\right)_{\mathscr{Z}_{\widetilde{\mathbf{B}}}} \\
& =\sum_{k_{1} \ldots k_{i-1} k_{i+1} \ldots k_{n}}\left|\sum_{l_{i}} c_{k_{1} \ldots k_{i-1} l_{i} k_{i+1} \ldots k_{n}} U_{i} \theta_{l_{i}}\right|_{\mathscr{F}_{\tilde{\mathbf{B}}^{k}}}^{2} \\
& \leq\left\|U_{i}\right\|_{\mathscr{H}_{\tilde{\mathbf{B}}} 1}^{2} \rightarrow \mathscr{H}_{\tilde{\mathbf{B}}^{2}} \sum_{k_{1} \ldots k_{i-1} k_{i+1} \ldots k_{n}}\left|\sum_{l_{i}} c_{k_{1} \ldots k_{i-1} l_{i} k_{i+1} \ldots k_{n}} \theta_{l_{i}}\right|_{\mathscr{H}_{\tilde{\mathbf{B}}}}^{2}
\end{aligned}
$$

so that

## Hence

On the other hand, by making use of (2.5) and (2.6) we have

$$
\begin{aligned}
& =\sup _{\hat{f} \in \mathscr{H}_{\mathbb{B}^{8}}^{8 n}}\left(\left(U_{1} \otimes U_{2} \otimes \cdots \otimes U_{n}\right) \hat{f},\left(U_{1} \otimes U_{2} \otimes \cdots \otimes U_{n}\right) \hat{f}\right)_{\mathscr{H}_{\mathbf{B}^{k}}^{\otimes n}} /(\hat{f}, \hat{f})_{\mathscr{H}_{\mathbb{B}^{8}}^{\theta n}}
\end{aligned}
$$

$$
\begin{align*}
& =\left\|U_{1} \otimes U_{2} \otimes \cdots \otimes U_{n}\right\|_{\mathscr{E}_{\mathbf{B}^{-1}}^{\infty} \rightarrow \mathscr{H}_{\mathbf{B}^{k}}^{\infty}}^{2} . \tag{3.7}
\end{align*}
$$

Combining (3.6) and (3.7) we obtain that

Now we return to the proof of Theorem 1.1. Since $\mathrm{N}_{\boldsymbol{A}} \hat{h}_{n}$ is also symmetric, we have by (2.6),

$$
\begin{aligned}
& \left|\Gamma(\tilde{\mathbf{B}})^{k} d \Gamma(\mathbf{A})\left\langle: x^{\otimes n}:, \hat{h}_{n}\right\rangle\right|_{L^{2}(\mathbf{E}, d \mu)}^{2}=n!\left|\left(\left(\tilde{\mathbf{B}}^{k}\right)^{\otimes n} \mathbf{N}_{\mathbf{A}}\right) \hat{h}_{n}\right|_{\mathscr{e}^{8 n}}^{2}
\end{aligned}
$$

$$
\begin{aligned}
& \leq n!\left(\sum_{i=1}^{n}\left|(I \otimes \cdots \otimes I \otimes \stackrel{i}{\mathbf{A}} \otimes I \otimes \cdots \otimes I) \hat{n}_{n}\right|_{\left.\mathscr{E} \otimes \mathbf{B}_{\mathbf{K}}\right)^{2}}{ }^{2}\right.
\end{aligned}
$$

$$
\begin{aligned}
& \leq n!n^{2}\|\mathbf{A}\|_{\mathscr{H}_{\tilde{\mathbf{B}}^{1}+1}^{2}} \rightarrow \mathscr{H}_{\widetilde{\mathbf{B}}^{k}}\left|\left(\widetilde{\mathbf{B}}^{t_{1}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{P}^{8}{ }^{8 n}} \quad \text { (by Lemma 3.1) } \\
& \leq\left. n!n^{2} c_{1}^{2}| |\left(\tilde{\mathbf{B}}^{l_{1}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{H}^{8 n}} ^{2} \quad \text { (by (3.5)) }
\end{aligned}
$$

and hence

$$
\begin{aligned}
\tilde{\mathbf{B}}\|d \Gamma(\mathbf{A}) \omega\|_{2, k}^{2} & =\left|\Gamma(\tilde{\mathbf{B}})^{k} d \Gamma(\mathbf{A}) \omega\right|_{L^{2}(E, d \mu)}^{2} \\
& =\sum_{n=0}^{\infty}\left|\Gamma(\tilde{\mathbf{B}})^{k} d \Gamma(\mathbf{A})\left\langle: x^{\otimes n}:, \hat{h}_{n}\right\rangle\right|_{L^{2}(E, d \mu)}^{2}
\end{aligned}
$$

$$
\begin{equation*}
\leq \sum_{n=0}^{\infty} n!n^{2} c_{1}^{2}\left|\left(\tilde{\mathbf{B}}^{l_{1}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{H}^{8 n}}^{2} \tag{3.8}
\end{equation*}
$$

Since for the natural number $q$ given in the assumption (2) of Theorem 1.1, $\tilde{\mathbf{B}}^{q} \geq 2^{q}$, $\left(\tilde{\mathbf{B}}^{q}\right)^{-1} \leq 2^{-q}$, we choose a natural number $m_{1}$ such that $\left(\tilde{\mathbf{B}}^{q}\right)^{-m_{1}} \leq$ $\left(2 c_{1}\right)^{-1}$. Then the right hand side of (3.8) is dominated by

$$
\begin{aligned}
\sum_{n=0}^{\infty} n!n^{2} c_{1}^{2}\left(\frac{1}{2 c_{1}}\right)^{2 n}\left|\left(\tilde{\mathbf{B}}^{l_{1}+q m_{1}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{H}^{8 n}}^{2} & \leq \sum_{n=0}^{\infty} n!\left|\left(\tilde{\mathbf{B}}^{l_{1}+q m_{1}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{H}^{8 n}}^{2} \\
& =\left|\Gamma(\tilde{\mathbf{B}})^{l_{1}+q m_{1}} \omega\right|_{L^{2}(\mathbf{E}, d \mu)}^{2}=\tilde{\mathbf{B}}\|\omega\|_{2, l_{1}+q m_{1}}^{2}
\end{aligned}
$$

Therefore we have $d \Gamma(\mathbf{A}) S \subset S$. As noted in Proposition 2.1, $S$ is a nuclear space. Thus the proof of the first half of Theorem 1.1 is completed.

Now we will prove the second half of the theorem. Note that $e^{-t d \Gamma(\mathbf{A})}$ is the self-adjoint operator defined on $L^{2}(E, \mu)$ such that

$$
\begin{equation*}
e^{-t d \Gamma(\mathbf{A})}=\Gamma\left(e^{-t \mathbf{A}}\right) \tag{3.9}
\end{equation*}
$$

This follows from the identity

$$
\frac{d}{d t} \Gamma\left(e^{-t \mathbf{A}}\right):\left\langle x, \xi_{1}{ }^{1}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle:=-d \Gamma(\mathbf{A}) \Gamma\left(e^{-t \mathbf{A}}\right):\left\langle x, \xi_{1}\right\rangle \cdots\left\langle x, \xi_{n}\right\rangle:
$$

which is verified by the induction on $n$. By (2.10), (3.2) and (3.9), we have the following expression;

$$
\begin{equation*}
e^{-t d \Gamma(\mathbf{A})} \omega(x)=\sum_{n=0}^{\infty}\left\langle: x^{\otimes n}:,\left(e^{-t \mathbf{A}}\right)^{\otimes n} \hat{h}_{n}\right\rangle . \tag{3.10}
\end{equation*}
$$

For any natural number $k$, by the assumption (4) of Theorem 1.1 and the manner similar to that in (3.4) we find a natural number $l_{2}(\geq k)$ and a constant $c_{2} \geq 1$ such that

$$
\begin{equation*}
\left|\tilde{\mathbf{B}}^{k} e^{-t \mathbf{A}} h\right|_{\mathscr{H}} \leq c_{2}\left|\tilde{\mathbf{B}}^{t_{2}} h\right|_{\mathscr{H}}, \quad h \in C^{\infty}(\tilde{\mathbf{B}}) \tag{3.11}
\end{equation*}
$$

Then we have

$$
\begin{aligned}
& \left|\Gamma(\tilde{\mathbf{B}})^{k} e^{-t d \Gamma(\mathbf{A})}\left\langle: x^{\otimes n}:, \hat{h}_{n}\right\rangle\right|_{L^{2}(E, d \mu)}^{2} \\
& =\left.n!| |\left(\tilde{\mathbf{B}}^{k} e^{-\boldsymbol{t} \mathbf{A}}\right)^{\otimes n} \hat{h}_{n}\right|_{\boldsymbol{e}^{8 n}} ^{2} \\
& =n!\left|\left(e^{-t \mathbf{\lambda}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{*} \Psi_{\mathbf{B}^{\star}}^{\infty}}^{2}
\end{aligned}
$$

$$
\begin{aligned}
& \leq n!\left\|e^{-t \mathbf{A}}\right\|_{\mathscr{A}_{\tilde{\mathbf{B}}^{l_{2}}}^{2 n} \rightarrow \mathscr{H}_{\tilde{\mathbf{B}}^{k}}}\left|\left(\tilde{\mathbf{B}}^{\tilde{l}_{2}}\right)^{\otimes n} \hat{h}^{n}\right|_{\mathscr{E}^{8 n}}^{2} \quad \text { (by Lemma 3.1) }
\end{aligned}
$$

$$
\leq n!c_{2}^{2 n}\left|\left(\tilde{\mathbf{B}}^{l_{2}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{F}^{8 n}}^{2} \quad \text { (by (3.11)) }
$$

and

$$
\begin{align*}
\tilde{\mathbf{B}}\left\|e^{-t d \Gamma(\mathbf{A})} \omega\right\|_{2, k}^{2} & =\left|\Gamma(\tilde{\mathbf{B}})^{k} e^{-t d \Gamma(\mathbf{A})} \omega\right|_{L^{2}(\mathbf{E}, d \mu)}^{2} \\
& =\sum_{n=0}^{\infty}\left|\Gamma(\tilde{\mathbf{B}})^{k} e^{-t d \Gamma(\mathbf{A})}\left\langle: x^{\otimes n}:, \hat{h}_{n}\right\rangle\right|_{L^{2}(E, d \mu)}^{2} \\
& \leq \sum_{n=0}^{\infty} n!c_{2}^{2 n}\left|\left(\tilde{\mathbf{B}}^{l_{2}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{H}^{8 n}}^{2} . \tag{3.12}
\end{align*}
$$

In the same manner as the proof of the first half, we choose a natural number $m_{2}$ such that $\left(\tilde{\mathbf{B}}^{q}\right)^{-m_{2}} \leq c_{2}^{-1}$. Then the right hand side of (3.12) is dominated by

$$
\sum_{n=0}^{\infty} n!c_{2}^{2 n}\left(\frac{1}{c_{2}}\right)^{2 n}\left|\left(\tilde{\mathbf{B}}^{l_{2}+q m_{2}}\right)^{\otimes n} \hat{h}_{n}\right|_{\mathscr{H}^{8 n}}^{2}=\tilde{\mathbf{B}}|\omega|_{2, l_{2}+q m_{2}}^{2}
$$

which completes the proof of the second half of Theorem 1.1.

## 4. Inveriant nuclear space

In this section we discuss the conditions of Theorem 1.1. Especially the conditions (3) and (4) have been examined by several authors such as J. Fröhlich [3]. Let A and B be positive self-adjoint operators in the separable Hilbert space $\mathscr{H}$. In the sequel we denote by $c_{i}, i=3,4, \cdots$ positive constants.

### 4.1. Case where $A$ and $B$ are non-commutative.

Lemma 4.1. Let $\mathbf{D}$ and $\mathbf{B}$ be positive self-adjoint operators in $\mathscr{H}$ such that $C^{\infty}(\mathbf{B}) \subset C^{\infty}(\mathbf{D})$ and $\mathbf{D} \geq 1+\varepsilon$ for some $\varepsilon>0$. Suppose that $\mathbf{B}$ has a bounded inverse and for any natural number $n$, there exists a constant $c_{3}$ such that

$$
\left|\mathbf{B}^{n} \mathbf{D} f\right|_{\mathscr{H}} \leq c_{3}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}, \quad f \in C^{\infty}(\mathbf{B})
$$

Then, for any natural number $n$, there exists a constant $c_{4}$ such that

$$
\left|\mathbf{B}^{n} e^{-t \mathbf{D}} f\right|_{\mathscr{H}} \leq c_{4}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}, \quad f \in C^{\infty}(\mathbf{B})
$$

Proof. Since $\mathbf{B}^{n} \mathbf{D B}^{-n}$ is a bounded operator, we have

$$
\begin{aligned}
\left|\mathbf{B}^{n} e^{-t \mathbf{D}} f\right|_{\mathscr{H}} & =\left|\mathbf{B}^{n} e^{-t \mathbf{D}} \mathbf{B}^{-n} \mathbf{B}^{n} f\right|_{\mathscr{H}} \\
& =\left|e^{-t \mathbf{B}^{n} \mathbf{D B}-n} \mathbf{B}^{n} f\right|_{\mathscr{H}} \\
& \leq c_{4}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}
\end{aligned}
$$

Next, we havve the following proposition which is implied by Lemma 4.1.

Suppose that $\mathbf{D}$ is a positive self-adjoint operator in $\mathscr{H}$ and $\mathbf{D} \geq 1+\varepsilon$ for some $\varepsilon>0$. We choose and fix a complete orthonormal system $\left\{e_{i}: i=0,1,2, \cdots\right\}$ of the Hilbert space $\mathscr{H}$ taken from $\mathscr{D}(\mathbf{D})$. Given $\left\{\lambda_{i}\right\}$ such that $\lambda_{i} \geq 1+\varepsilon$, $\varepsilon>0, \lambda_{i} \uparrow \infty$, define a positive self-adjoint operator $\mathbf{B}$ in $\mathscr{H}$ by $\mathbf{B} e_{i}=\lambda_{i} e_{i}$. Let $\left\{d_{i, j}\right\}_{i, j=0}^{\infty}$ be the infinite dimensional matrix such that

$$
\mathbf{D} e_{i}=\sum_{j=0}^{\infty} d_{i, j} e_{j} .
$$

Proposition 4.1. Suppose that

$$
\begin{equation*}
M_{n} \equiv \sum_{i=0}^{\infty} \frac{\sum_{j=0}^{\infty} d_{i, j}^{2} \lambda_{j}^{2 n}}{\lambda_{i}^{2 n}}<+\infty, \quad \text { for all } n \in \mathbf{N} \tag{4.1}
\end{equation*}
$$

Then the conditions (1), (2), (3), (4) of Theorem 1.1 are satisfied for $\mathbf{A}=\mathbf{D}$ and $\mathbf{B}$.
Proof. To prove (1), let $f=\sum a_{i} e_{i}$ be a finite sum. Then $f \in \mathscr{D}(\mathbf{D})$ and hence

$$
\mathbf{D} f=\sum a_{i} \mathbf{D} e_{i} .
$$

Therefore

$$
\begin{aligned}
|\mathbf{D} f|_{\mathscr{H}}^{2} & =\left|\sum a_{i} \mathbf{D} e_{i}\right|_{\mathscr{H}}^{2} \\
& =\left|\sum a_{i} \lambda_{i} \lambda_{i}^{-1} \mathbf{D} e_{i}\right|_{\mathscr{H}}^{2} \\
& \leq\left(\sum a_{i}^{2} \lambda_{i}^{2}\right)\left(\sum \lambda_{i}^{-2}\left|\mathbf{D} e_{i}\right|_{\mathscr{H}}^{2}\right) \\
& \leq c_{5}|\mathbf{B} f|_{\mathscr{H}}^{2},
\end{aligned}
$$

where $c_{5}=\sum \lambda_{i}^{-2}\left|\mathbf{D} e_{i}\right|_{\neq H}^{2}=\sum \lambda_{i}^{-2} \sum_{j=0}^{\infty} d_{i, j}^{2} \leq M_{1}$.
Since $\left\{\lambda_{i}^{-1} e_{i}: i=0,1,2, \cdots\right\}$ forms a complete orthonormal system of $\mathscr{D}(\mathbf{B})$ and

$$
\sum_{i=0}^{\infty}\left|\frac{1}{\lambda_{i}} e_{i}\right|_{\mathscr{H}}^{2}=\sum_{i=0}^{\infty} \frac{1}{\lambda_{i}^{2}} \leq \sum_{i=0}^{\infty} \frac{\left|\mathbf{D} e_{i}\right|_{\mathscr{H}}^{2}}{\lambda_{i}^{2}}<+\infty,
$$

then the identity map $\mathscr{D}(\mathbf{B}) \rightarrow \mathscr{H}$ becomes a Hilbert-Schamit operator. Of course $\mathbf{B}$ has a bounded inverse $\mathbf{B}^{-1}$ such that $\mathbf{B}^{-1} e_{i}=\lambda_{i}^{-1} e_{i}$.

Since for any integer $i$ and natural number $n$,

$$
\sum_{j=0}^{\infty} d_{i, j}^{2} \lambda_{j}^{2 n}<\infty
$$

from (4.1), we have $\mathbf{D} e_{i} \in \mathscr{D}\left(\mathbf{B}^{n}\right)$, so that for any $f \in C^{\infty}(\mathbf{B})$ and for any fixed natural number N ,

$$
\begin{aligned}
\left|\mathbf{B}^{n} \sum_{i=0}^{\mathrm{N}}\left(f, e_{i}\right)_{\mathscr{H}} \mathbf{D} e_{i}\right|_{\mathscr{H}}^{2} & =\left|\sum_{i=0}^{\mathrm{N}}\left(f, e_{i}\right)_{\mathscr{H}} \mathbf{B}^{n} \mathbf{D} e_{i}\right|_{\mathscr{H}}^{2} \\
& =\left|\sum_{i=0}^{\mathrm{N}}\left(f, e_{i}\right)_{\mathscr{H}} \lambda_{i}^{n} \lambda_{i}^{-n} \mathbf{B}^{n} \mathbf{D} e_{i}\right|_{\mathscr{H}}^{2} \\
& \leq\left(\sum_{i=0}^{\mathrm{N}}\left(f, e_{i}\right)_{\mathscr{H}}^{2} \lambda_{i}^{2 n}\right)\left(\sum_{i=0}^{\mathrm{N}} \lambda_{i}^{-2 n}\left|\mathbf{B}^{n} \mathbf{D} e_{i}\right|_{\mathscr{H}}^{2}\right) \\
& \leq \mathbf{M}_{n}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}^{2} .
\end{aligned}
$$

This implies the condition (3) of Theorem 1.1.
Since B has a bounded inverse, Lemma 4.1 yields the condition (4) of Theorem 1.1.

Proposition 4.2. Suppose that $\mathbf{D}$ and $\mathbf{B}$ satisfy the assumptions of Proposition 4.1 and $\mathbf{C}$ is a positive self-adjoint operator in $\mathscr{H}$ such that $\mathbf{C}$ and $\mathbf{B}$ are commutative and $\mathscr{D}(\mathbf{B}) \subset \mathscr{D}(\mathbf{C})$. Then the conditions (1), (2), (3), (4) of Theorem 1.1 are satisfied for $\mathbf{A}=\mathbf{C}+\mathbf{D}$ and $\mathbf{B}$.

Proof. The condition (1) is obvious from $\mathscr{D}(\mathbf{B}) \subset \mathscr{D}(\mathbf{D})$ and $\mathscr{D}(\mathbf{B}) \subset \mathscr{D}(\mathbf{C})$. It has been already proved in Proposition 4.1 that the condition (2) is satisfied.

Since $\mathscr{D}(\mathbf{B}) \subset \mathscr{D}(\mathbf{C})$, there exists a constant $c_{6}$ such that

$$
\begin{equation*}
|\mathbf{C} f|_{\mathscr{H}} \leq c_{6}|\mathbf{B} f|_{\mathscr{H}}, \quad f \in C^{\infty}(\mathbf{B}) \tag{4.2}
\end{equation*}
$$

By the commutativity of $\mathbf{C}$ and $\mathbf{B}$ and (4.2), for any natural number $n$, we have

$$
\begin{aligned}
\left|\mathbf{B}^{n} \mathbf{A} f\right|_{\mathscr{H}}^{2} & =\left|\mathbf{B}^{n}(\mathbf{C}+\mathbf{D}) f\right|_{\mathscr{H}}^{2} \\
& \leq 2\left\{\left|\mathbf{B}^{n} \mathbf{C} f\right|_{\mathscr{H}}^{2}+|\mathbf{B D} f|_{\mathscr{H}}^{2}\right\} \\
& \leq 2\left\{\left|\mathbf{C B}^{n} f\right|_{\mathscr{H}}^{2}+\mathbf{M}_{n}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}^{2}\right\} \\
& \leq 2\left\{c_{6}\left|\mathbf{B}^{n+1} f\right|_{\mathscr{H}}^{2}+\mathbf{M}_{n}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}^{2}\right\} \\
& \leq c_{7}\left|\mathbf{B}^{n+1} f\right|_{\mathscr{H}}^{2} .
\end{aligned}
$$

To prove (4), it suffices to show that for some constant $0<c_{8}<\infty$

$$
\left|\mathbf{B}^{n} e^{-t \mathbf{A}} f\right|_{\mathscr{H}} \leq c_{8}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}, \quad f \in C^{\infty}(\mathbf{B})
$$

By Theorem 1.19 of Chap. IX in [8], we have an integral equation

$$
e^{-t \mathbf{A}} f=e^{-t \mathbf{C}} f+\int_{0}^{t} e^{-(t-\tau) \mathbf{C}} \mathbf{D} e^{-\tau \mathbf{A}} f d \tau
$$

so that operating $\mathbf{B}^{n}$ on both sides of the above equality and using the commutativity of $\mathbf{C}$ and B, the assumption of Proposition 4.1, we get

$$
\begin{aligned}
\left|\mathbf{B}^{n} e^{-t \mathbf{A}} f\right|_{\mathscr{H}} & \leq\left|e^{-t \mathbf{C}} \mathbf{B}^{n} f\right|_{\mathscr{H}}+\int_{0}^{t}\left|e^{-(t-\tau) \mathbf{C}} \mathbf{B}^{n} \mathbf{D} e^{-\tau \mathbf{A}} f\right|_{\mathscr{H}} d \tau \\
& \leq\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}+\int_{0}^{t}\left|\mathbf{B}^{n} \mathbf{D} e^{-\tau \mathbf{A}} f\right|_{\mathscr{H}} d \tau \\
& \leq\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}+\sqrt{\mathbf{M}_{n}} \int_{0}^{t}\left|\mathbf{B}^{n} e^{-\tau \mathbf{A}} f\right|_{\mathscr{H}} d \tau
\end{aligned}
$$

Gronwall's lamma then yields

$$
\left|\mathbf{B}^{n} e^{-t \mathbf{A}} f\right|_{\mathscr{H}} \leq e^{t \sqrt{M_{n}}}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}=: c_{8}\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}
$$

This completes the proof of Proposition 4.2.
We have a sufficient condition for (4) of Theorem 1.1. (See [3].)
Proposition 4.3. Let $\mathbf{A}$ and $\mathbf{B}$ be positive self-adjoint operators in $\mathscr{H}$ such that $\mathbf{A} \geq 1+\varepsilon$, for some $\varepsilon>0$ and the condition (3) of Theorem 1.1 is satisfied. Suppose that for any natural numbers $n$ and $k$, there exist a natural number $\alpha(n) \geq n$ and a constant $C_{k}^{n}>0$ such that

$$
\left|\left(\mathbf{A}^{k} \mathbf{B}^{n}-\mathbf{B}^{n} \mathbf{A}^{k}\right) f\right|_{\mathscr{H}} \leq C_{k}^{n}\left|\mathbf{B}^{\alpha(n)} f\right|_{\mathscr{H}}, \quad f \in C^{\infty}(\mathbf{B})
$$

and

$$
\sum_{k=0}^{\infty} \frac{t^{k}}{k!} C_{k}^{n}<\infty, \quad \text { for any } t>0
$$

Then for any natural number $n$, there exists a constant $C^{n}(t)>0$ such that

$$
\left|\mathbf{B}^{n} e^{-t \mathbf{A}} f\right|_{\mathscr{H}} \leq C^{n}(t)\left|\mathbf{B}^{\alpha(n)} f\right|_{\mathscr{H}}
$$

Proof. For any natural number $n$ and $f \in C^{\infty}(\mathbf{B})$, we have

$$
\begin{aligned}
\left|\mathbf{B}^{n} \sum_{k=0}^{\mathrm{N}} \frac{(-t)^{k}}{k!} \mathbf{A}^{k} f\right|_{\mathscr{H}} & =\left|\sum_{k=0}^{\mathrm{N}} \frac{(-t)^{k}}{k!} \mathbf{B}^{n} \mathbf{A}^{k} f\right|_{\mathscr{H}} \\
& \leq\left|\sum_{k=0}^{\mathrm{N}} \frac{(-t)^{k}}{k!} \mathbf{A}^{k} \mathbf{B}^{n} f\right|_{\mathscr{H}}+\left|\sum_{k=0}^{\mathrm{N}} \frac{(-t)^{k}}{k!}\left(\mathbf{A}^{k} \mathbf{B}^{n}-\mathbf{B}^{n} \mathbf{A}^{k}\right) f\right|_{\mathscr{H}} \\
& \leq\left|\sum_{k=0}^{\mathbf{N}} \frac{(-t)^{k}}{k!} \mathbf{A}^{k} \mathbf{B}^{n} f\right|_{\mathscr{H}}+\sum_{k=0}^{\mathrm{N}} \frac{t^{k}}{k!} C_{k}^{n}\left|\mathbf{B}^{\alpha(n)} f\right|_{\mathscr{H}} .
\end{aligned}
$$

From the above estimate it is seen that

$$
\left|\mathbf{B}^{n} e^{-t \mathbf{A}} f\right|_{\mathscr{H}} \leq\left|e^{-t \mathbf{A}} \mathbf{B}^{n} f\right|_{\mathscr{H}}+\sum_{k=0}^{\infty} \frac{t^{k}}{k!} C_{k}^{n}\left|\mathbf{B}^{\alpha(n)} f\right|_{\mathscr{H}}
$$

$$
\leq\left(1+\sum_{k=0}^{\infty} \frac{t^{k}}{k!} C_{k}^{n}\right)\left|\mathbf{B}^{\alpha(n)} f\right|_{\ngtr} \quad(\text { by } A>1)
$$

Remark 4.1. If $\mathbf{D}$ is a bounded operator in Lemma 4.1, then the conclusion is derived from Proposition 4.3 without $\mathbf{B}$ having a bounded inverse.

### 4.2. Case where $A$ and $B$ are commutative.

Here we assume that $\mathbf{A}$ and $\mathbf{B}$ are positive self-adjoint operators defined in the separable Hilbert space $\mathscr{H}$.

By the spectral theorem we have the following spectral representations;

$$
\mathbf{A}=\int_{0}^{\infty} v d E(v), \quad \mathbf{B}=\int_{0}^{\infty} \lambda d F(\lambda) .
$$

We say that $\mathbf{A}$ and $\mathbf{B}$ are commutative if $E(v)$ and $F(\lambda)$ are commutative for all $v, \lambda \geq 0$.

Proposition 4.4. Suppose that $\mathbf{A}$ and $\mathbf{B}$ are commutative and there exists the inverse $\mathbf{B}^{-1}$ of $\mathbf{B}$ such that $\mathbf{B}^{-1}$ is compact on $\mathscr{H}$. Then $\mathbf{A}$ has countable eigenvectors which form a complete orthonormal system of $\mathscr{H}$.

Proof. Since $\mathbf{B}^{-1}$ is compact and $\operatorname{ker} \mathbf{B}^{-1}=\{0\}$, we have

$$
\mathbf{B}^{-1}=\sum_{n=1}^{\infty} \gamma_{n} P_{n},
$$

where $\left\{P_{n}\right\}$ are orthogonal projections on $\mathscr{H}$ satisfying

$$
\operatorname{dim} P_{n} \mathscr{H}<+\infty
$$

and

$$
\mathscr{H}=\bigoplus_{n=1}^{\infty} P_{n} \mathscr{H} .
$$

By the commutativity of $E(v)$ and $P_{n}$, A maps $P_{n} \mathscr{H}$ into $P_{n} \mathscr{H}$, which asserts Proposition 4.4.

By Proposition 4.4, A has countable eigenvalues $\left\{v_{i}\right\}$ such that $\mathbf{A} e_{i}=v_{i} e_{i}$, $i=0,1,2, \cdots$, where $\left\{e_{i}\right\}$ forms a complete orthonormal system of $\mathscr{H}$. Setting $\mathbf{B} e_{i}=\lambda_{i} e_{i}, i=0,1,2, \cdots$, we have

Corollary 4.1. Let $\mathbf{A}$ and $\mathbf{B}$ be as above. Supose that $\mathbf{A} \geq 1+\varepsilon$ for some $\varepsilon>0$. If for some natural number $k$

$$
\sum_{i=0}^{\infty} \frac{v_{i}^{2}}{\lambda_{i}^{2 k}}=c_{9}<+\infty,
$$

then the conditions (1), (2), (3) and (4) of Theorem 1.1 are satisfied for $\mathbf{A}$ and $\mathbf{B}$.
Proof. Following the same argument as in the proof of Proposition 4.1, we see that for any finite sum $f=\sum a_{i} e_{i}$ and any integer $n \geq 0$,

$$
\left|\mathbf{B}^{n} \mathbf{A} f\right|_{\mathscr{H}}^{2} \leq c_{9}\left|\mathbf{B}^{n+k} f\right|_{\mathscr{H}}^{2},
$$

which implies the conditions (1) and (3) of Theorem 1.1. Since $v_{i}>1$, $i=0,1,2, \cdots,(2)$ is proved by the same argument as in Proposition 4.1. Since $A$ and $B$ are commutative, we have

$$
\left|\mathbf{B}^{n} e^{-t \boldsymbol{A} \boldsymbol{A}}\right|_{\mathscr{H}}^{2} \leq\left|\mathbf{B}^{n} f\right|_{\mathscr{H}}^{2},
$$

which completes the proof of Corollary 4.1.

## 5. Strong solution for a Segal-Langevin type equation

In [7], they discussed a fluctuation phenomena for interacting, spatially extended neurons and as a limit equation, they found a suitable fundamental space $\mathscr{D}_{E}$ of functionals on $E$ and studied Segal-Langevin type stochastic differential equations including a class of the weak version of (1.1):

$$
\begin{equation*}
d X_{F}(t)=d W_{F}(t)+X_{-d \Gamma(\mathbf{A}) F}(t) d t, \quad F \in \mathscr{D}_{E} . \tag{5.1}
\end{equation*}
$$

A stochastic process $X_{F}(t)$ indexed by elements in $\mathscr{D}_{E}$ is called a continuous $L\left(\mathscr{D}_{E}\right)$-process if for any fixed $F \in \mathscr{D}_{E}, X_{F}(t)$ is a real continuous process and $X_{\alpha F+\beta G}(t)=\alpha X_{F}(t)+\beta X_{G}(t)$ almost surely for real numbers $\alpha, \beta$ and elements $F, G \in \mathscr{D}_{E}$ and further $E\left[X_{F}(t)^{2}\right]$ is continuous on $\mathscr{D}_{E} . W_{F}(t)$ is an $L\left(\mathscr{D}_{E}\right)$-Wiener process such that for any fixed $F \in \mathscr{D}_{E}, W_{F}(t)$ is a real Wiener process.

Although the above $\mathscr{D}_{E}$ is not nuclear, appealing to the results in [7], we get a unique continuous $L\left(\mathscr{D}_{E}\right)$-process satisfying (5.1).

We consider the case where for $\mathbf{A}$ in (5.1) there exists a self-adjoint operator B satisfying all the conditions of Theorem 1.1. In this case, by Theorem 1.1, there is a nuclear space $S$ invariant under both $d \Gamma(\mathbf{A})$ and a strong continuous semigroup $T(t)=e^{-t d \Gamma(\mathbf{A})}$. If we replace $\mathscr{D}_{E}$ by $S$ in (5.1), then by the regularization theorem [5] there exists an $S^{\prime}$-valued Wiener process $W(t)$ such that $\langle W(t), F\rangle=W_{F}(t)$ almost surely and the strong form of the equation with $\mathscr{D}_{E}$ replaced by $S$ in (5.1) is the following stochastic differential equation on $S^{\prime}$ :

$$
d X(t)=d W(t)-d \Gamma(\mathbf{A})^{*} X(t) d t
$$

Let $T(t)^{*}$ be the adjoint operator of $T(t)$. Since $S$ is nuclear, again by the regularization theorem, the stochastic integral $\int_{0}^{t} T(t-s)^{*} d W(s)$ is well defined
from the weak form such that

$$
\left\langle\int_{0}^{t} T(t-s)^{*} d W(s), F\right\rangle=\int_{0}^{t}\langle d W(s), T(t-s) F\rangle
$$

Since $T(t-s) F=F+\int_{s}^{t} T(\tau-s)(-d \Gamma(\mathbf{A})) F d \tau$, we get

$$
\int_{0}^{t} T(t-s)^{*} d W(s)=W(t)+\int_{0}^{t}\left(-d \Gamma(\mathbf{A})^{*}\right)\left(\int_{0}^{\tau} T(\tau-s)^{*} d W(s)\right) d \tau .
$$

Noticing that

$$
\int_{0}^{t}\left(-d \Gamma(\mathbf{A})^{*}\right) T(\tau)^{*} X(0) d \tau=T(t)^{*} X(0)-X(0)
$$

we see that

$$
X(t)=T(t)^{*} X(0)+\int_{0}^{t} T(t-s)^{*} d W(s)
$$

is a unique strong solution of (1.1) on $S^{\prime}$.

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