# Normal approximation and smoothness for sums of means of lattice-valued random variables 

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#### Abstract

Motivated by a problem arising when analysing data from quarantine searches, we explore properties of distributions of sums of independent means of independent lattice-valued random variables. The aim is to determine the extent to which approximations to those sums require continuity corrections. We show that, in cases where there are only two different means, the main effects of distribution smoothness can be understood in terms of the ratio $\rho_{12}=\left(e_{2} n_{1}\right) /\left(e_{1} n_{2}\right)$, where $e_{1}$ and $e_{2}$ are the respective maximal lattice edge widths of the two populations, and $n_{1}$ and $n_{2}$ are the respective sample sizes used to compute the means. If $\rho_{12}$ converges to an irrational number, or converges sufficiently slowly to a rational number; and in a number of other cases too, for example those where $\rho_{12}$ does not converge; the effects of the discontinuity of lattice distributions are of smaller order than the effects of skewness. However, in other instances, for example where $\rho_{12}$ converges relatively quickly to a rational number, the effects of discontinuity and skewness are of the same size. We also treat higher-order properties, arguing that cases where $\rho_{12}$ converges to an algebraic irrational number can be less prone to suffer the effects of discontinuity than cases where the limiting irrational is transcendental. These results are extended to the case of three or more different means, and also to problems where distributions are estimated using the bootstrap. The results have practical interpretation in terms of the accuracy of inference for, among other quantities, the sum or difference of binomial proportions.


Keywords: algebraic irrational number; bootstrap; confidence interval; continuity correction; difference of binomial proportions; discontinuity; irrational number; sum of binomial proportions; transcendental number

## 1. Introduction

### 1.1. Background: The case of a single sample mean

Let $\hat{\theta}$ denote a statistical estimator of an unknown quantity $\theta$, and assume that $\hat{\theta}-\theta$ is asymptotically normally distributed with zero mean and variance $n^{-1} \sigma^{2}$, where $n$ is a measure of the sample size from which $\hat{\theta}$ was computed. In particular, the statistic $T=n^{1 / 2}(\hat{\theta}-\theta) / \sigma$ is asymptotically normal $\mathrm{N}(0,1)$. Under additional assumptions an Edgeworth expansion of the
distribution of $T$ can generally be formulated, having the form

$$
\begin{equation*}
P(T \leq x)=\Phi(x)+n^{-1 / 2} P(x) \phi(x)+o\left(n^{-1 / 2}\right) \tag{1.1}
\end{equation*}
$$

uniformly in $x$, where $\Phi$ and $\phi$ are the standard normal distribution and density functions, respectively, and $P$ is an even, quadratic polynomial.

For example, if $\hat{\theta}$ denotes the mean of a sample of size $n$ from a population with mean $\theta$, and if data from the population have finite third moment and a nonlattice distribution, then (1.1) holds with $P(x)=\frac{1}{6} \beta\left(1-x^{2}\right)$, where $\beta$ is the standardised skewness of the population. On the other hand, still in the case of the mean of a population with finite third moment, if the population is lattice then an extra, discontinuous term has to be added to (1.1).

This extra term reflects the discrete continuity correction that statisticians are often obliged to introduce when approximating a lattice distribution, for example the binomial distribution or the distribution of a sum of Poisson variates, using the smooth normal distribution:

$$
\begin{equation*}
P(T \leq x)=\Phi(x)+n^{-1 / 2} P(x) \phi(x)+n^{-1 / 2} d_{n}(x) \phi(x)+o\left(n^{-1 / 2}\right), \tag{1.2}
\end{equation*}
$$

where $d_{n}(x)=\left(e_{0} / \sigma\right) \psi_{n}(x)$ denotes the discontinuous term in the Edgeworth expansion, $e_{0}$ is the maximal span of the lattice, $\sigma^{2}$ is the population variance,

$$
\psi_{n}(x)=\psi\left\{\left(x-\xi_{n}\right) \sigma n^{1 / 2} / e_{0}\right\}, \quad \xi_{n}=\left(e_{0} / \sigma n^{1 / 2}\right)\left\{\frac{1}{2}-\psi\left(n x_{0} / e_{0}\right)\right\},
$$

$\psi(x)=\lfloor x\rfloor-x+\frac{1}{2},\lfloor x\rfloor$ is the largest integer not strictly exceeding $x$, and it is assumed that all points of support in the distribution (of which $\theta$ is the mean) have the form $x_{0}+v e_{0}$ for an integer $v$.

### 1.2. Contributions of this paper

We show that in multi-sample problems, where $\hat{\theta}$ is a sum of several independent means, the discontinuous term can be ignored if sample sizes are chosen judiciously. For example, if there are just two sample sizes (as in the case of a sum or difference of two binomial proportions), and if the lattice edge widths are identical (this simplifies our discussion here, but is not essential), then it is sufficient that the ratio of the two sample sizes converge to an irrational number, or converge sufficiently slowly to a rational number. These results are corollaries of Theorem 1 in Section 2.1, and they and other properties are discussed in Section 2.2.

We also show that the discontinuous term can be replaced by $O\left(n^{\delta-1}\right)$, for all $\delta>0$, provided that four moments are finite and the ratio of the two sample sizes converges sufficiently quickly to an irrational number of "type" 1. (See Section 2.3 for a definition of the type of an irrational number.) More generally, we explore the effect that type has on the size of the discontinuous term. Theorem 2 also gives an explicit formula for the discontinuous term, up to a remainder of order $n^{-1}$. Sections 2.5 and 2.6 show how to bound the discontinuous remainder term, for two different approaches to defining that quantity, and show how the effects of irrationals of different types can be teased from the remainder. Applications to the bootstrap are straightforward, and in fact Section 2.7 outlines a bootstrap version of Theorem 2 and discusses its implications.

We do not treat in any detail cases where the differences between two lattice distributions arise mainly in terms of their centres, rather than their lattice edge widths. For example, if two independent sample means $\bar{X}_{j}$, for $j=1,2$, are respectively averages of $n_{j}$ independent variables and are distributed on lattices $x_{j}+\nu n_{j}^{-1} e_{j}$; and if the difference $x_{1}-x_{2}$ between the lattice translations equals an irrational multiple of the ratio $\rho_{12}=\left(e_{2} n_{1}\right) /\left(e_{1} n_{2}\right)$; then the distribution of $\bar{X}_{1}+\bar{X}_{2}$ is non-lattice. While this problem and its implications are of mathematical interest, they do not enjoy the practical motivation of problems where, say, $x_{1}=x_{2}$ and $\rho_{12}$ can be almost arbitrary. For example, $x_{1}=x_{2}$ in the problem of constructing confidence intervals for the sum or difference of two binomial probabilities, based on samples of unequal size. Therefore, we address cases where the focus of attention is $\rho_{12}$ rather than $x_{1}-x_{2}$. Differences between lattice centres are permitted by our regularity conditions, but their role is not treated in detail.

### 1.3. Practical motivation

The extra term in (1.2), relative to (1.1), is of significant interest to a practitioner, since it causes significant inaccuracy when the central limit theorem is used to approximate the distribution of $T$. The presence of this extra term motivates the continuity correction, and also the fiducial approach taken by Clopper and Pearson [6] and Sterne [21] to estimating a binomial proportion, as well as a large, more recent literature discussing methodology for solving problems such as constructing confidence intervals for the difference or sum of two binomial proportions. See, for example, Hall [11], Duffy and Santer [7], Lee et al. [14], Agresti and Caffo [1], Brown et al. [4,5], Zhou et al. [23], Price and Bonnett [16], Brown and Li [3], Borkowf [2], Roths and Tebbs [19], Wang [22] and Zieliński [24].

The practical motivation for the work described in this paper came from data acquired during quarantine searches, where the construction of confidence intervals for the sum, rather than difference, of two binomial proportions was of interest. In detail, shipping containers arriving at a frontier contained a certain number, $N$ say, of consignments. Some of the consignments might be clean, but others could contain pests which needed to be detected and removed to prevent their introduction to the environment. To reduce the costs associated with inspection, quarantine services usually inspect only $n_{1}<N$ consignments. Consignments are assumed to be contaminated with probability $p_{1}$, and the number, $n_{1} \bar{X}_{1}$ say, of contaminated consignments found after routine (but incomplete) inspection of the items in each of $n_{1}$ consignments is assumed to follow a binomial distribution. Contaminated consignments are then "cleaned," and the members of a subsample of $n_{2}$ of them are reinspected. (The data gathered in this way comprise a "leakage survey.") The number of items, $n_{2} \bar{X}_{2}$, still found contaminated (for example, contaminated with a different kind of pest) are assumed to follow a binomial distribution with parameters $n_{2}$ and $p_{2}$, and typically it is argued that $\bar{X}_{1}$ and $\bar{X}_{2}$ are independent. An estimator of the proportion of items that pass through this inspection process, and are still contaminated, is given by

$$
\begin{equation*}
\bar{X}_{1}\left(1-\frac{n_{1}}{N}\right)+\bar{X}_{2}\left(1-\frac{n_{2}}{N}\right), \tag{1.3}
\end{equation*}
$$

which can be viewed as a sum of means of lattice-valued random variables where the lattice edge lengths are $e_{j}=1-N^{-1} n_{j}$ for $j=1,2$.

The quarantine inspection service aims to develop a strategy for choosing consignments, and items, to inspect. This reduces the associated costs, and minimises, to at least some extent, the number of contaminated items that cross the border. The performance of such a strategy is assessed, by the quarantine service, using a variety of statistics based on sums of binomials; (1.3) is just one example. Quarantine services are usually interested in providing confidence intervals as well as point estimators, and hence there is significant interest in estimating the distributions of statistics such as that at (1.3).

## 2. Main results

### 2.1. Edgeworth expansions with remainder equal to $o\left(n^{-1 / 2}\right)$

Let $X_{j i}$, for $1 \leq i \leq n_{j}$ and $j=1, \ldots, k$, denote independent random variables. Assume that each $X_{j i}$ has a nondegenerate lattice distribution, depending on $j$ but not on $i$ and with maximal lattice edge width $e_{j}$ and finite third moment. Suppose too that $k \geq 2$. Put $\bar{X}_{j}=n_{j}^{-1} \sum_{i} X_{j i}$, $\mu_{j}=E\left(X_{j i}\right), \sigma_{j}^{2}=\operatorname{var}\left(X_{j i}\right)$ and

$$
\begin{equation*}
S=\sum_{j=1}^{k} \bar{X}_{j} . \tag{2.1}
\end{equation*}
$$

The model (2.1) includes cases of apparently greater generality, for example where signed weights are incorporated in the series, since the absolute values of the weights can be incorporated into (2.1) by modifying the lattice edge widths, and negative signs can be addressed by reflecting the summand distributions.

Since third moments are finite then, if the distributions of $X_{11}, \ldots, X_{k 1}$ were to satisfy a smoothness condition, such as that of Cramér, we could express the distribution of $S$ in a oneterm Edgeworth expansion:

$$
\begin{equation*}
P\left\{\frac{S-E(S)}{(\operatorname{var} S)^{1 / 2}} \leq x\right\}=\Phi(x)+n^{-1 / 2} \frac{1}{6} \beta\left(1-x^{2}\right) \phi(x)+o\left(n^{-1 / 2}\right), \tag{2.2}
\end{equation*}
$$

where we take $n=n_{1}+\cdots+n_{k}$ to be the asymptotic parameter, and

$$
\begin{equation*}
\beta=\beta(n)=\frac{n^{1 / 2} E(S-E S)^{3}}{(\operatorname{var} S)^{3 / 2}}=\frac{n^{1 / 2} \sum_{j} n_{j}^{-2} E\left(X_{j 1}-E X_{j 1}\right)^{3}}{\left(\sum_{j} n_{j}^{-1} \operatorname{var} X_{j 1}\right)^{3 / 2}} \tag{2.3}
\end{equation*}
$$

is a measure of standardised skewness and, under our assumptions, is bounded as $n \rightarrow \infty$. Result (2.2) is a version of (1.1) in a particular case.

However, in general (2.2) does not hold in the lattice-valued case that we are considering. For example, if $k \geq 1$ and the $X_{j i} \mathrm{~s}$, for all $i$ and $j$, have a common lattice distribution, then, as was made clear by Esseen [9], any expansion of the distribution of $X$ has to include a discontinuous term of size $n^{-1 / 2}$ (specifically, the term $n^{-1 / 2} d_{n}(x) \phi(x)$ in (1.2)) that reflects the "continuity
correction" needed to approximate the discontinuous distribution of $T=\{S-E(S)\} /(\operatorname{var} S)^{1 / 2}$ by a continuous normal distribution.

When exploring this problem, we suppose that the sequence of values of $n$ is strictly increasing. Further, we assume that

$$
\begin{equation*}
\min _{1 \leq j \leq k} \liminf _{n \rightarrow \infty}\left(n_{j} / n\right)>0 \tag{2.4}
\end{equation*}
$$

In Theorem 1, below, we fix both $k$ and the distributions of $X_{j 1}$, for $1 \leq j \leq k$. This means that $e_{1}, \ldots, e_{k}$ are fixed too. However, for each $n$ we consider there to be a potentially new sequence of values $n_{1}, \ldots, n_{k}$. In particular, the ratios $n_{j_{1}} / n_{j_{2}}$ can change considerably from one choice of $n$ to another, although, in view of (2.4), $n_{j_{1}} / n_{j_{2}}$ is bounded away from zero and infinity as $n \rightarrow \infty$.

In the first part of Theorem 1, below, we also impose the following condition on at least one of the ratios $\rho_{j_{1} j_{2}}=\left(e_{j_{2}} n_{j_{1}}\right) /\left(e_{j_{1}} n_{j_{2}}\right)$ :

$$
\begin{equation*}
\text { for each integer } \ell \geq 1, \quad n^{1 / 2}\left|\sin \left(\ell \rho_{j_{1} j_{2}} \pi\right)\right| \rightarrow \infty \tag{2.5}
\end{equation*}
$$

as $n \rightarrow \infty$.
Theorem 1. Assume that $E\left|X_{j 1}\right|^{3}<\infty$ for $j=1, \ldots, k$; that $X_{j 1}$ is distributed on a lattice $x_{j}+v e_{j}$, for integers $v$, where $e_{j}$ is the maximal lattice edge width; and that (2.4) holds. (i) If, for some pair $j_{1}, j_{2}$ with $1 \leq j_{1}<j_{2} \leq k, \rho_{j_{1} j_{2}}$ satisfies (2.5), then the one-term Edgeworth expansion at (2.2) holds uniformly in $x$. (ii) However, if $\rho_{j_{1} j_{2}}$ equals a fixed rational number (not depending on $n$ ) for each pair $j_{1}, j_{2}$, and if the points $x_{j}$ can all be taken equal, then the expansion at (2.2) fails to hold because it does not include an appropriate discontinuous term of size $n^{-1 / 2}$.

### 2.2. Circumstances where (2.5) holds

If $\rho_{0}$ is irrational, then $\left|\sin \left(\ell \rho_{0} \pi\right)\right|>0$ for all integers $\ell$. Therefore, (2.5) holds if $\rho_{j_{1} j_{2}}$ converges to an irrational number as $n \rightarrow \infty$.

However, in many cases (2.5) holds without the sequence $\rho_{j_{1} j_{2}}$ converging. For example, assume for simplicity that the lattice edge widths $e_{j}$ are all identical, let $\rho_{1}$ and $\rho_{2}$ be two distinct irrational numbers, and let the sequence of values of the ratio $n_{j_{1}} / n_{j_{2}}$ be a sequence of convergents of $\rho_{1}$ and $\rho_{2}$, chosen so that an infinite number of convergents come from each $\rho_{j}$. (For a definition of convergents of irrational numbers, see, e.g., Leveque [15], p. 70.) Then (2.5) holds, although the sequence $\rho_{j_{1} j_{2}}$ does not converge.

Importantly, (2.5) also holds in many cases where each $\rho_{j_{1} j_{2}}$ is close to a rational number, indeed where each $\rho_{j_{1} j_{2}}$ converges to a rational number. For example, we claim that (2.5) obtains if $\rho_{j_{1} j_{2}}=1+\varepsilon_{j_{1} j_{2}}$, where $\varepsilon_{j_{1} j_{2}}=\varepsilon_{j_{1} j_{2}}(n)$, which can be either positive or negative, converges to zero strictly more slowly than $n^{1 / 2}$ :

$$
\begin{equation*}
\varepsilon_{j_{1} j_{2}} \rightarrow 0, \quad n^{1 / 2}\left|\varepsilon_{j_{1} j_{2}}\right| \rightarrow \infty \tag{2.6}
\end{equation*}
$$

In this case, for each fixed, positive integer $\ell$,

$$
\sin \left(\ell \rho_{j_{1} j_{2}} \pi\right)=\sin (\ell \pi)+\ell \pi \varepsilon_{j_{1} j_{2}} \cos (\ell \pi)+O\left(\varepsilon_{j_{1} j_{2}}^{2}\right)
$$

from which it follows that

$$
\left|\sin \left(\ell \rho_{j_{1} j_{2}} \pi\right)\right| \sim \begin{cases}\ell \pi\left|\varepsilon_{j_{1} j_{2}}\right| & \text { if } \ell \text { is an even integer }  \tag{2.7}\\ 1 & \text { if } \ell \text { is an odd integer }\end{cases}
$$

where $a_{n} \sim b_{n}$ means that the ratio $a_{n} / b_{n}$ converges to 1 . Assumption (2.5) follows from (2.6) and (2.7).

A similar argument can be used to prove that if $\rho_{j_{1} j_{2}}=\rho_{0}+\varepsilon_{j_{1} j_{2}}$, where $\rho_{0}$ is a fixed rational number and $\varepsilon_{j_{1} j_{2}}=\varepsilon_{j_{1} j_{2}}(n)$ satisfies (2.6), then (2.5) is true. (The case $\rho_{0}=0$ is excluded by (2.4).) These examples make it clear that there is not a great deal of latitude in the assumption, imposed in part (ii) of Theorem 1, that each $\rho_{j_{1} j_{2}}$ should equal a fixed rational number. In particular, for (2.5) to fail it is not sufficient that each $\rho_{j_{1} j_{2}}$ converge to a rational.

### 2.3. Refinement of bound on remainder term in Edgeworth expansions

In Section 2.1, we showed that, if (2.5) holds, the discontinuous term of size $n^{-1 / 2}$, in expansions such as (1.2), is actually of smaller order than $n^{-1 / 2}$. To obtain a more concise bound on the discontinuous term, we shall investigate in detail cases where one or more of the ratios $\rho_{j_{1} j_{2}}$ converge to an irrational number as $n$ diverges. However, this treatment requires a definition of the "type" of an irrational, and we give that next.

If $x$ is a real number, let $\langle x\rangle$ denote the distance from $x$ to the nearest integer. (In particular, if $\lfloor x\rfloor$ is the integer part function, $\langle x\rangle=\min \{x-\lfloor x\rfloor, 1-(x-\lfloor x\rfloor)\}$.) We say that the irrational number $\rho$ is of type $\eta$ if $\eta$ equals the supremum of all $\zeta$ such that $\liminf _{p \rightarrow \infty} p^{\zeta}\langle p \rho\rangle=0$, where $p \rightarrow \infty$ through integer values. Properties of convergents of irrational numbers (specifically, Dirichlet's Theorem) can be used to prove that the type of any given irrational number always satisfies $\eta \geq 1$. It follows from Roth's Theorem (Roth [18]) that all algebraic irrationals (that is, all irrational numbers that are roots of polynomials with rational coefficients) are of minimal type, i.e., $\eta=1$, which is one of the cases we consider below.

More generally, if a number is chosen randomly, for example as the value of a random variable having a continuous distribution on the real line, then with probability 1 it is an irrational of type 1. Irrationals that are not algebraic are said to be transcendental, and can have type strictly greater than 1. (However, the transcendental number $e$ is of type 1.) Known upper bounds to the types of $\pi, \pi^{2}$ and $\log 2$ are $6.61,4.45$ and 2.58 , respectively. Liouville numbers have type $\eta=\infty$. The type of an irrational number is one less than its irrationality measure (or equivalently, one less than its approximation exponent or Liouville-Roth constant). We refer the reader to Ribenboim [17] for more information about types of irrational numbers.

Next, we introduce notation which helps us to define an approximation to the discontinuous term, an analogue of $d_{n}(x)$ in (1.2), when $k=2$. (Here, $k$ is as in (2.1).) Assuming that the lattice, on which the distribution of $X_{j i}$ is supported, consists of points $x_{j}+v e_{j}$ for integers $v$, define
$\xi_{j n}=e_{j}\left(\sigma_{j} n_{j}^{1 / 2}\right)^{-1}\left\{\left(n_{j} x_{j} / e_{j}\right)-\left\lfloor n_{j} x_{j} / e_{j}\right\rfloor\right\}$ and

$$
\begin{equation*}
\xi_{n}(x)=\left\{x-\left(c_{1} \xi_{1 n}+c_{2} \xi_{2 n}\right)\right\} \frac{\sigma_{1} n_{1}^{1 / 2}}{c_{1} e_{1}} \tag{2.8}
\end{equation*}
$$

where, recalling that $\sigma_{j}^{2}=\operatorname{var}\left(X_{j i}\right)$, we define $c_{j}$ for $j=1$ and 2 by

$$
\begin{equation*}
c_{j}=\left(\frac{n_{j}^{-1} \sigma_{j}^{2}}{n_{1}^{-1} \sigma_{1}^{2}+n_{2}^{-1} \sigma_{2}^{2}}\right)^{1 / 2} \tag{2.9}
\end{equation*}
$$

Let $\alpha \in\left(0, \frac{1}{2}\right)$ and partition the set of all integers into adjacent blocks each comprised of $2\left\lfloor n^{\alpha}\right\rfloor+1$ consecutive integers. Write $\bar{\nu}_{\ell}$ for the central integer in the $\ell$ th block, which we denote by $\mathcal{N}_{\ell}$ where $-\infty<\ell<\infty$ and $\mathcal{N}_{\ell+1}$ is immediately to the right of $\mathcal{N}_{\ell}$ on the number line. Given $v \in \mathcal{N}_{\ell}$, put $\nu_{\ell}=v-\bar{\nu}_{\ell}$.

Let $c_{3}=e_{2} n_{1} / \sigma_{1} n_{2}$ and $c_{4}=\left(e_{1} / \sigma_{2}\right)\left(n_{1} / n_{2}\right)^{1 / 2}$, and note that $c_{1}, \ldots, c_{4}$ are strictly positive functions of $n$ and are bounded away from zero and infinity as $n$ diverges. Put $\gamma=$ $\prod_{j=1,2}\left(e_{j} / \sigma_{j}\right)$, and, given an integer $r_{0} \geq 1$, define

$$
\begin{align*}
\phi(u, x) & =\phi\left\{\left(x / c_{1}\right)-c_{3} u\right\} \phi\left(c_{4} u\right), \quad \phi_{r}(u, x)=(\partial / \partial u)^{r} \phi(u, x), \\
K_{n}(x) & =\gamma \sum_{r=0}^{r_{0}} \sum_{-\infty<\ell<\infty} \frac{\phi_{r}\left(\bar{v}_{\ell} / n_{1}^{1 / 2}, x\right)}{r!n_{1}^{r / 2}} \sum_{v \in \mathcal{N}_{\ell}} v_{\ell}^{r} \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\}, \tag{2.10}
\end{align*}
$$

where, as in Section 1.1, $\psi(x)=\lfloor x\rfloor-x+\frac{1}{2}$.
We claim that the infinite series in the definition of $K_{n}(x)$ is absolutely convergent, uniformly in $x$. To appreciate why, note that

$$
\begin{equation*}
\sup _{-\infty<x<\infty}\left|\phi_{r}\left(\bar{v}_{\ell} / n_{1}^{1 / 2}, x\right)\right| \leq C_{1}(r) \phi\left(c_{4} \bar{v}_{\ell} / n_{1}^{1 / 2}\right) \tag{2.11}
\end{equation*}
$$

where, here and below, the notation $C_{j}(r)$ will denote a constant depending on $r$ but not on $n$. Using (2.4) and (2.11), we deduce that

$$
\begin{equation*}
\sum_{-\infty<\ell<\infty}\left\{\sup _{-\infty<x<\infty}\left|\phi_{r}\left(\bar{v}_{\ell} / n_{1}^{1 / 2}, x\right)\right|\right\} \leq C_{2}(r) n^{(1 / 2)-\alpha} \tag{2.12}
\end{equation*}
$$

(In more detail, without loss of generality the block $\mathcal{N}_{0}$ is centred at 0 , in which case, when bounding the series on the left-hand side of (2.12), $\bar{\nu}_{\ell}$ can be interpreted as $2 \ell n^{\alpha}$. Consequently the left-hand side of (2.12) is bounded by a constant multiple of $C_{1}(r) \int \phi\left(2 u n^{\alpha} / n_{1}^{1 / 2}\right) d u$, and (2.12) follows.)

More simply, since (a) $\left|\nu_{\ell}\right| \leq n^{\alpha}$, (b) $\left|\mathcal{N}_{\ell}\right| \leq\left(2 n^{\alpha}+1\right)$ (where we define $\left|\mathcal{N}_{\ell}\right|=\# \mathcal{N}_{\ell}$ ), and (c) $|\psi| \leq \frac{1}{2}$, then

$$
\begin{equation*}
\sup _{-\infty<x<\infty}\left|\sum_{v \in \mathcal{N}_{\ell}} v_{\ell}^{r} \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\}\right| \leq C_{3}(r) n^{(r+1) \alpha} \tag{2.13}
\end{equation*}
$$

Combining (2.10), (2.12) and (2.13), and replacing each summand on the right-hand side of (2.10) by its absolute value, we obtain the bound: $n^{-1 / 2}\left|K_{n}(x)\right| \leq C_{4}\left(r_{0}\right)$, uniformly in $x$. This inequality demonstrates the claimed absolute convergence of the series in (2.10).

Recall the definition of $S$ at (2.1), and that $\rho_{j_{1} j_{2}}=\left(e_{j_{2}} n_{j_{1}}\right) /\left(e_{j_{1}} n_{j_{2}}\right)$. Part (i) of Theorem 2, below, captures the analogue of the discontinuous term, $d_{n}(x)$, in a multisample version of (1.2), and part (ii) gives conditions under which the net contribution of that term equals $O\left(n^{\delta-(1 / 2)-(1 / 2 \eta)}\right)$, for all $\delta>0$ when some $\rho_{j_{1} j_{2}}$ is sufficiently close to an irrational number of type $\eta$.

Theorem 2. Assume that $E\left|X_{j 1}\right|^{4}<\infty$ for $j=1, \ldots, k$; that $X_{j 1}$ is distributed on a lattice $x_{j}+v e_{j}$,for integers $v$, where $e_{j}$ is the maximal lattice edge width; and that (2.4) holds. Choose $r_{0} \geq 4 \alpha /(1-2 \alpha)$ in (2.10). (i) If $k=2$ and $K_{n}$ is as defined at (2.10), then

$$
\begin{equation*}
P\left\{\frac{S-E(S)}{(\operatorname{var} S)^{1 / 2}} \leq x\right\}=\Phi(x)+n^{-1 / 2} \frac{1}{6} \beta\left(1-x^{2}\right) \phi(x)+\left(n_{1} n_{2}\right)^{-1 / 2} K_{n}(x)+O\left(n^{-1}\right), \tag{2.14}
\end{equation*}
$$

uniformly in $x$. (ii) If, for some pair $j_{1}, j_{2}$ with $1 \leq j_{1}<j_{2} \leq k$, the ratio $\rho_{j_{1} j_{2}}=\left(e_{j_{2}} n_{j_{1}}\right) /$ $\left(e_{j_{1}} n_{j_{2}}\right)$ satisfies

$$
\begin{equation*}
\left|\rho_{j_{1} j_{2}}-\rho_{0}\right|=O\left(n^{-(1 / 2)\{1+(1 / \eta)+\delta\}}\right) \tag{2.15}
\end{equation*}
$$

for some $\delta>0$, where $\rho_{0}$ is an irrational number of type $\eta$, then, for each $\delta>0$,

$$
\begin{equation*}
P\left\{\frac{S-E(S)}{(\operatorname{var} S)^{1 / 2}} \leq x\right\}=\Phi(x)+n^{-1 / 2} \frac{1}{6} \beta\left(1-x^{2}\right) \phi(x)+O\left(n^{\delta-(1 / 2)-(1 / 2 \eta)}\right) \tag{2.16}
\end{equation*}
$$

uniformly in $x$.
Result (2.16) is of particular interest in the case $\eta=1$, which encompasses almost all irrational numbers (with respect to Lebesgue measure), including all the algebraic irrationals and some transcendental numbers. When $\eta=1$,

$$
\begin{equation*}
P\left\{\frac{S-E(S)}{(\operatorname{var} S)^{1 / 2}} \leq x\right\}=\Phi(x)+n^{-1 / 2} \frac{1}{6} \beta\left(1-x^{2}\right) \phi(x)+O\left(n^{\delta-1}\right) \tag{2.17}
\end{equation*}
$$

uniformly in $x$ for each $\delta>0$. Result (2.17) implies that the lattice nature of the distribution of $X_{j i}$ can be ignored, almost up to terms of second order in Edgeworth expansions, when considering the impact of latticeness on the accuracy of normal approximations.

### 2.4. Practical choice of $\boldsymbol{n}_{\mathbf{1}}$ and $\boldsymbol{n}_{\mathbf{2}}$

In practice it is not difficult to choose $n_{1}$ and $n_{2}$ so that (2.15) holds. To see how, assume for simplicity that the lattice edge widths $e_{1}$ and $e_{2}$ are identical, as they would be if (for example) $S$ were equal to a sum or difference of estimators of binomial proportions. If $\rho_{0}$ is an irrational number then the convergents $m_{1} / m_{2}$ of $\rho_{0}$ satisfy

$$
\begin{equation*}
\left|\left(m_{1} / m_{2}\right)-\rho_{0}\right| \leq m_{2}^{-2} . \tag{2.18}
\end{equation*}
$$

(See e.g. Leveque [15], equation (29), p. 180.) Therefore, if $n_{1}$ and $n_{2}$ are relatively prime and $n_{1} / n_{2}$ is a convergent of $\rho_{0}$, then (2.15), for each $\delta \in(0,3-(1 / \eta)]$, follows from (2.18). The most difficult case, as far as (2.15) is concerned, is the one where the convergence rate in (2.15) is fastest, and arises when $\eta=1$. There we need to ensure that

$$
\begin{equation*}
\left|\rho_{j_{1} j_{2}}-\rho_{0}\right|=O\left(n^{-1-\delta}\right) \tag{2.19}
\end{equation*}
$$

for some $\delta>0$. Now, (2.19) holds whenever $n_{1} / n_{2}$ is a convergent of $\rho_{0}$, and the KhinchinLévy Theorem (see, e.g., pp. 82-83 of Einsiedler and Ward [8]) implies that the convergents are reasonably closely spaced; the numerators and denominators generally increase by factors of only $\pi^{2} /(12 \log 2) \approx 1.87$. Moreover, there are many ratios $n_{1} / n_{2}$ on either side of convergents for which (2.19) holds.

The pair $\left(n_{1}, n_{2}\right)$ can be chosen from tables of, or formulae for, convergents for commonly arising irrationals of type 1 . See, for example, Griffiths [10] and references therein, and note that $e$ and any algebraic irrational is of type 1 .

### 2.5. Alternative formula for $K_{n}$, and derivation of (2.16) from (2.14) when $\eta=1$

Part (i) of Theorem 2 can be stated for a version of $K_{n}(x)$ simpler than that at (2.10):

$$
\begin{equation*}
K_{n}(x)=\gamma \sum_{\nu} \phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} v\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v\right\} \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} \tag{2.20}
\end{equation*}
$$

Indeed, the $K_{n}(x)$ at (2.20) is just $\gamma I_{4}(x)$, where $I_{4}(x)$ is as defined at (4.17) in the proof of Theorem 1, and in fact that formula provides a convenient point of access to a proof of (2.14) with $K_{n}(x)$ as at (2.20). However, in the case $\eta>1$ it is not straightforward to pass from (2.20) to (2.16), and that is why we used the definition of $K_{n}(x)$ at (2.10).

To appreciate that (2.16) follows from (2.20) when $\eta=1$, note that the definition of $K_{n}(x)$ at (2.20) is equivalent to:

$$
\begin{equation*}
\gamma^{-1} K_{n}(x)=\sum_{\nu} \Psi(x, \nu) \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} \tag{2.21}
\end{equation*}
$$

where

$$
\Psi(x, v)=\phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} v\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v\right\} .
$$

If (2.15) holds with $\eta=1$ then a standard argument for bounding discrepancies of sequences (see p. 123 of Kuipers and Niederreiter [13]) can be used to prove that for all $\delta>0$,

$$
\begin{equation*}
\sup _{-\infty<z<\infty}\left|\sum_{\nu=1}^{N} \psi\left(z-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right)\right|=O\left(N^{\delta}\right) \tag{2.22}
\end{equation*}
$$

Note too that

$$
\begin{equation*}
\sup _{v \geq 1} \sup _{-\infty<x<\infty}|\Psi(x, v+1)-\Psi(x, \nu)| \leq C n^{-1 / 2} \tag{2.23}
\end{equation*}
$$

Taking $a_{v}=\Psi(x, \nu)$ and $b_{\nu}=\psi\left\{\xi_{n}(x)-\left(e_{2} n_{1} / e_{1} n_{2}\right) \nu\right\}$, and employing Abel's method of summation, we can write:

$$
\sum_{\nu=1}^{N} a_{v} b_{v}=a_{N} \sum_{\nu=1}^{N} b_{v}-\sum_{\nu=1}^{N-1}\left(a_{v+1}-a_{v}\right) \sum_{j=1}^{\nu} b_{j},
$$

which in company with (2.22) and (2.23) allows us to prove that, provided $N=O\left(n^{C}\right)$ for some $C>0$,

$$
\begin{equation*}
\sup _{-\infty<x<\infty}\left|\sum_{v=1}^{N} \Psi(x, v) \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\}\right|=O\left(N^{\delta}\right) \tag{2.24}
\end{equation*}
$$

for all $\delta>0$. More simply, if $N \geq n^{2}$ then

$$
\begin{align*}
\sup _{-\infty<x<\infty}\left|\sum_{\nu=N+1}^{\infty} \Psi(x, v) \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\}\right| \\
\leq \frac{1}{2} \sup _{-\infty<x<\infty} \sum_{\nu=N+1}^{\infty} \Psi(x, v)=O(1) \tag{2.25}
\end{align*}
$$

Combining (2.24) and (2.25), using a similar argument to treat series where $v \leq 0$, and noting the definition of $K_{n}(x)$ at (2.21), we deduce that $\sup _{x}\left|K_{n}(x)\right|=O\left(n^{\delta}\right)$ for all $\delta>0$. In the case $k=2$, and for $\eta=1$, this gives (2.16) as a corollary of (2.14).

### 2.6. Derivation of (2.16) from (2.14) when $\eta \geq 1$

Our proof of (2.16), in Section 4.2, will proceed by deriving implicitly a version of (2.14) in the case $k \geq 2$, and showing that, if (2.15) holds, then that version of (2.14) entails (2.16). The relative complexity of a form of (2.14) for general $k$ discouraged us from including it in Theorem 2, but it is nevertheless instructive to show how, when $k=2$, one can obtain (2.16) from (2.14). We outline the proof below, highlighting the properties of irrational numbers, particularly the differences between the case of irrationals of type $\eta=1$ and the case of those of larger type, that determine the bound for the remainder in (2.16).

Note that if $q$ is a polynomial function then, applying Koksma's inequality (see, e.g., Theorem 5.1, p. 143 of [13]) and the Erdős-Turán inequality (see, e.g., formula (2.42), p. 114 of [13]), it can be shown that

$$
\begin{equation*}
\chi(N, q, \tau) \equiv \sup _{-\infty<z<\infty}\left|\sum_{i=1}^{N} q(i / N) \psi(z-\tau i)\right| \leq C_{1}(q)\left\{\frac{N}{m}+\sum_{\ell=1}^{m} \frac{1}{\ell|\sin (\ell \tau \pi)|}\right\} \tag{2.26}
\end{equation*}
$$

for all integers $m \geq 1$. Here $\tau>0$ is permitted to vary with $N$, and the constant $C_{1}(q)$ depends on the degree and the coefficients of $q$ but not on the positive integer $N$ or on $m, z$ or $\tau$.

We shall take $\tau=\rho_{12}$, a function of $n$, in which case, since

$$
\left||\sin (\ell \tau \pi)|-\left|\sin \left(\ell \rho_{0} \pi\right)\right|\right| \leq \ell \pi\left|\rho_{12}-\rho_{0}\right|
$$

we have:

$$
\begin{equation*}
|\sin (\ell \tau \pi)|+\ell \pi\left|\rho_{12}-\rho_{0}\right| \geq\left|\sin \left(\ell \rho_{0} \pi\right)\right|=\sin \left(\pi\left\langle\ell \rho_{0}\right\rangle\right) \geq 2\left\langle\ell \rho_{0}\right\rangle \geq 2 C_{2}(\delta) \ell^{-\eta-\delta} \tag{2.27}
\end{equation*}
$$

for any given $\delta>0$ and all $\ell \geq 1$, where $C_{2}(\delta)>0$ depends on $\delta$ but not on $\ell$, the last inequality in (2.27) follows from the assumption that $\rho_{0}$ is of type $\eta$, and the second-last inequality comes from the fact that $0 \leq\langle x\rangle \leq \frac{1}{2}$ for all real numbers $x$, and $\sin (\pi x) \geq 2 x$ whenever $0 \leq x \leq \frac{1}{2}$. If

$$
\begin{equation*}
\left|\rho_{12}-\rho_{0}\right| \leq C_{2}(\delta) \pi^{-1} m^{-(1+\eta+\delta)} \tag{2.28}
\end{equation*}
$$

then it follows from (2.27) that $|\sin (\ell \tau \pi)| \geq\left\langle\ell \rho_{0}\right\rangle$ for $1 \leq \ell \leq m$, and so

$$
\begin{equation*}
\sum_{\ell=1}^{m} \frac{1}{\ell|\sin (\ell \tau \pi)|} \leq \sum_{\ell=1}^{m} \frac{1}{\ell\left\langle\ell \rho_{0}\right\rangle} \tag{2.29}
\end{equation*}
$$

A standard argument for bounding the discrepancy of a sequence (see, e.g., p. 123 of [13]) can be used to show that, since $\rho_{0}$ is an irrational number of type $\eta$,

$$
\begin{equation*}
\sum_{\ell=1}^{m} \frac{1}{\ell\left\langle\ell \rho_{0}\right\rangle}=O\left(m^{\eta-1+\delta}\right) \tag{2.30}
\end{equation*}
$$

for all $\delta>0$. Therefore, provided that (2.28) holds, we can deduce from (2.26) and (2.29) that

$$
\begin{equation*}
\chi\left(N, q, \rho_{12}\right) \leq C_{3}(q) m^{-1}\left(N+m^{\eta+\delta}\right)=O\left(N^{1-(1 / \eta)+\delta_{1}}\right) \tag{2.31}
\end{equation*}
$$

where $\delta_{1}>0$ and the inequality holds for all $m$ and the identity is true if $m / N^{1 / \eta}$ is bounded away from zero and infinity as $N \rightarrow \infty$. When $m$ has the latter property, (2.28) is satisfied, for all sufficiently large $N$, if

$$
\begin{equation*}
\left|\rho_{12}-\rho_{0}\right|=O\left(N^{-\left\{1+(1 / \eta)+\delta_{2}\right\}}\right) \tag{2.32}
\end{equation*}
$$

for some $\delta_{2}>\delta_{1}$.
Note that it is at (2.30) that the type, $\eta$, of the irrational number $\rho_{0}$ enters into consideration. In the case $\eta=1$ the exponent $\delta$ in (2.30) could not be removed or reduced, perhaps by replacing the implicit factor $m^{\delta}$ in (2.30) by $(\log m)^{C}$ for some $C>0$, without an analogous strengthening of Roth's Theorem. Formula (2.30) also marks the step at which it becomes apparent that a poorer bound will be obtained for an irrational number of type 1, relative to one of type $\eta>1$.

Applying the bound (2.31), for several versions of the polynomial $q$, in the case $N=2\left\lfloor n^{\alpha}\right\rfloor+$ 1 , we deduce that

$$
\begin{equation*}
n^{-r / 2} \sup _{-\infty<z<\infty}\left|\sum_{v \in \mathcal{N}_{\ell}} v_{\ell}^{r} \psi\left(z-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right)\right|=O\left(n^{r\{\alpha-(1 / 2)\}+\alpha\left\{1-(1 / \eta)+\delta_{1}\right\}}\right), \tag{2.33}
\end{equation*}
$$

provided that (2.32) holds, i.e., $\left|\rho_{12}-\rho_{0}\right|=O\left(n^{-\alpha\left\{1+(1 / \eta)+\delta_{2}\right\}}\right)$. Now, the only constraint on $\alpha$ is $0<\alpha<\frac{1}{2}$, and so we can choose $\alpha$ as close to $\frac{1}{2}$, but less than $\frac{1}{2}$, as we desire. In particular, if $\delta_{3}>0$ is given, and we choose $\alpha=\frac{1}{2}-\delta_{4}$ where $\delta_{4}>0$ is sufficiently small, then by taking $\delta_{1}$ in (2.33) to be small we obtain:

$$
\begin{equation*}
\max _{1 \leq r \leq r_{0}} n^{-r / 2} \sup _{-\infty<z<\infty}\left|\sum_{v \in \mathcal{N}_{\ell}} v_{\ell}^{r} \psi\left(z-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right)\right|=O\left(n^{(1 / 2)\{1-(1 / \eta)\}+\delta_{3}}\right), \tag{2.34}
\end{equation*}
$$

provided that

$$
\begin{equation*}
\left|\rho_{12}-\rho_{0}\right|=O\left(n^{-(1 / 2)\left\{1+(1 / \eta)+\delta_{5}\right\}}\right), \tag{2.35}
\end{equation*}
$$

where $\delta_{5}>0$ can be made as small as we like simply by choosing $\delta_{4}$ small. Now, (2.35) follows from (2.15). It therefore follows from (2.34), and the definition of $K_{n}(x)$ at (2.10), that if (2.15) holds for some $\delta>0$ then

$$
\begin{equation*}
\sup _{-\infty<x<\infty}\left(n_{1} n_{2}\right)^{-1 / 2}\left|K_{n}(x)\right|=O\left(n^{\delta-(1 / 2)\{1+(1 / \eta)\}}\right) \tag{2.36}
\end{equation*}
$$

for all $\delta>0$. Results (2.14) and (2.36) imply (2.16), as had to be shown.

### 2.7. Expansions relating to the bootstrap

In this section we show that, despite the potential for problems arising from discreteness, the bootstrap (including the double bootstrap) applied to inference based on the distribution of $\{S-$ $E(S)\} /(\operatorname{var} S)^{1 / 2}$, generally (when (2.15) holds and $\rho_{0}$ is of type 1) produces confidence regions and hypothesis tests with the same orders of magnitude of coverage or level accuracy, up to terms of size $n^{\delta-1}$ for all $\delta>0$, as it would in the case of smooth sampling distributions. This result is of practical importance, since standard percentile bootstrap methods applied to lattice distributions are frustrated by the effects of discontinuities; see, e.g., Singh [20] and Hall [12].

For brevity, when establishing this property we treat only the context of Theorem 2. We begin by stating an analogue of (2.14) there, valid when $k=2$. The arguments used to prove part (i) of Theorem 2 can be employed to show that

$$
\begin{align*}
P\left[\left.\frac{S^{*}-E\left(S^{*} \mid \mathcal{X}\right)}{\{\operatorname{var}(S \mid \mathcal{X})\}^{1 / 2}} \leq x \right\rvert\, \mathcal{X}\right]=\Phi(x)+n^{-1 / 2} & \frac{1}{6} \hat{\beta}\left(1-x^{2}\right) \phi(x) \\
& +\left(n_{1} n_{2}\right)^{-1 / 2} \widehat{K}_{n}(x)+n^{-1} \Delta_{1}(x) \tag{2.37}
\end{align*}
$$

where, analogously to the definitions in Section 2.1, $S^{*}=\sum_{j} \bar{X}_{j}^{*} ; \bar{X}_{j}^{*}=n_{j}^{-1} \sum_{i} X_{j i}^{*}$ and $X_{j 1}^{*}, \ldots, X_{j n_{j}}^{*}$ are drawn by sampling randomly, with replacement, from $\mathcal{X}_{j}=\left(X_{j 1}, \ldots, X_{j n_{j}}\right)$; $\mathcal{X}=\left(\mathcal{X}_{1}, \ldots, \mathcal{X}_{k}\right) ;$

$$
\hat{\beta}=\frac{n^{1 / 2} E\left[\left\{S^{*}-E\left(S^{*} \mid \mathcal{X}\right)\right\}^{3} \mid \mathcal{X}\right]}{\left\{\operatorname{var}\left(S^{*} \mid \mathcal{X}\right)\right\}^{3 / 2}} ;
$$

using (2.10) or (2.20), respectively, as the model for $K_{n}(x)$,

$$
\begin{aligned}
& \widehat{K}_{n}(x)=\hat{\gamma} \sum_{r=0}^{r_{0}} \sum_{-\infty<\ell<\infty} \frac{\hat{\phi}_{r}\left(\bar{v}_{\ell} / n_{1}^{1 / 2}, x\right)}{r!n_{1}^{r / 2}} \sum_{v \in \mathcal{N}_{\ell}} v_{\ell}^{r} \psi\left\{\hat{\xi}_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\}, \\
& \widehat{K}_{n}(x)=\hat{\gamma} \sum_{v} \phi\left(\frac{x}{\hat{c}_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\hat{\sigma}_{1} n_{2}} v\right) \phi\left\{e_{2}\left(\hat{\sigma}_{2} n_{2}^{1 / 2}\right)^{-1} v\right\} \psi\left\{\hat{\xi}_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\},
\end{aligned}
$$

where $\hat{\gamma}=\prod_{j=1,2}\left(e_{j} / \hat{\sigma}_{j}\right), \hat{\phi}_{r}(u, x)=(\partial / \partial u)^{r} \hat{\phi}(u, x), \hat{\phi}(u, x)=\phi\left\{\left(x / \hat{c}_{1}\right)-\hat{c}_{3} u\right\} \times \phi\left(\hat{c}_{4} u\right)$,

$$
\hat{c}_{j}=\left(\frac{n_{j}^{-1} \hat{\sigma}_{j}^{2}}{n_{1}^{-1} \hat{\sigma}_{1}^{2}+n_{2}^{-1} \hat{\sigma}_{2}^{2}}\right)^{1 / 2}
$$

for $j=1$ and $2, \hat{c}_{3}=e_{2} n_{1} / \hat{\sigma}_{1} n_{2}, \hat{c}_{4}=\left(e_{1} / \hat{\sigma}_{2}\right)\left(n_{1} / n_{2}\right)^{1 / 2}$, and $\hat{\xi}_{n}(x)$ is defined using the empirical analogue of (2.8); and, for $C_{1}>0$ sufficiently large and for some $C_{2}>0$,

$$
\begin{equation*}
P\left\{\sup _{-\infty<x<\infty}\left|\Delta_{1}(x)\right|>C_{1} n^{-1}\right\}=O\left(n^{-C_{2}}\right) \tag{2.38}
\end{equation*}
$$

The assumptions needed for (2.37) are those imposed for part (i) of Theorem 2. The size of $C_{2}$ in (2.38) depends to some extent on the distributions of $X_{1 i}$ and $X_{2 i}$ (recall that at this point we are assuming that $k=2$ ), but for distributions such as the Bernoulli or Poisson, which have all moments finite, $C_{2}$ can be taken arbitrarily large if $C_{1}$ is sufficiently large. The connection to moments here arises because the $O\left(n^{-C_{2}}\right)$ bound in (2.38) is derived using a method related to Markov's inequality, which can be applied at a higher order if more moments are finite.

The methods used in Sections 2.5 and 2.6 to derive uniform bounds to $K_{n}$ can also be employed to bound $\widehat{K}_{n}$, giving

$$
\begin{equation*}
P\left\{\sup _{-\infty<x<\infty}\left|\widehat{K}_{n}(x)\right|>n^{\delta+(1 / 2)\{1-(1 / \eta)\}}\right\}=O\left(n^{-C_{3}}\right) \tag{2.39}
\end{equation*}
$$

for all $\delta>0$ and some $C_{3}>0$, provided that (2.15) holds. In (2.39), $\eta$ denotes the type of the irrational number $\rho_{0}$ appearing in (2.15), and for sampling distributions such as the Bernoulli or Poisson (with all moments finite), $C_{3}$ can be taken arbitrarily large. Therefore, treating the case of irrationals of type 1 , we deduce from (2.37)-(2.39) that

$$
\begin{equation*}
P\left[\left.\frac{S^{*}-E\left(S^{*} \mid \mathcal{X}\right)}{\left\{\operatorname{var}\left(S^{*} \mid \mathcal{X}\right)\right\}^{1 / 2}} \leq x \right\rvert\, \mathcal{X}\right]=\Phi(x)+n^{-1 / 2} \frac{1}{6} \hat{\beta}\left(1-x^{2}\right) \phi(x)+n^{-1} \Delta_{2}(x) \tag{2.40}
\end{equation*}
$$

where

$$
P\left\{\sup _{-\infty<x<\infty}\left|\Delta_{2}(x)\right|>n^{\delta}\right\}=O\left(n^{-C_{2}}\right) \quad \text { for all } C_{2}, \delta>0
$$

A similar argument, employing the methods introduced in Section 2.6, can be used to prove that (2.40) continues to hold if $k \geq 2$, provided that the assumptions imposed in part (ii) of Theorem 2 hold. Therefore, the properties stated in the first paragraph of this section hold.

## 3. Numerical properties

Throughout this section, we take $k=2$ and let $X_{j i}$ be a Bernoulli random variable satisfying $P\left(X_{j i}=0\right)=1-P\left(X_{j i}=1\right)=p_{j}$ for $j=1,2$, where $p_{1}=0.4$ and $p_{2}=0.6$. Thus, $\rho_{12}=$ $e_{2} n_{1} /\left(e_{1} n_{2}\right)=n_{1} / n_{2}$, where $n_{1}$ and $n_{2}$ are the two sample sizes. We take $n_{2}$ to be the integer nearest to $\rho_{0} n_{1}$, and vary $n_{1}$ between 10 and $80 ; n_{1}$ is plotted on the horizontal axes of each of our graphs. The probability

$$
\begin{equation*}
P(x)=P\left[\{S-E(S)\} /(\operatorname{var} S)^{1 / 2} \leq x\right] \tag{3.1}
\end{equation*}
$$

was approximated by averaging over the results of $10^{5}$ Monte Carlo simulations.
To illustrate the influence of $\rho_{12}$ on the oscillatory behaviour of $P(x)$, and in particular on the accuracy of the normal approximation, each panel in Figure 1 plots $P(x)$ against $n_{1}$ for $x=\Phi^{-1}(\alpha)=z_{\alpha}$ and $\alpha=0.95,0.85$ and 0.75 . The top left panel of Figure 1 shows results for $\rho_{0}=1$ (indicated by the lines with circles) and $\rho_{0}=2$ (lines with dots), and it is clear that in both cases there is significant oscillatory behaviour, arising principally from the term in $K_{n}(x)$ in (2.14). The top right panel of Figure 1 shows that these oscillations decline markedly, and the accuracy of the normal approximation improves considerably, if $\rho_{0}=2^{1 / 2}$. This property reflects the results reported in Section 2.

Of course, $\rho_{0}=2^{1 / 2}$ is an algebraic irrational. The bottom left panel of Figure 1 shows that broadly similar values of $P(x)$, although with somewhat more oscillation (reflecting the relatively low upper bounds given in Theorem 1), are obtained for $\rho_{0}=\pi / 2$, a transcendental irrational whose type is bounded above by 6.61. The bottom right panel of Figure 1 addresses one of the results reported in Section 2.2, specifically that there may be less oscillatory behaviour when $\rho_{12}$ converges slowly to a rational number than when it converges quickly. We consider the cases $n_{2}=n_{1}+\left[n_{1}^{1 / 5}\right]$ and $n_{2}=n_{1}+\left[n_{1}^{3 / 5}\right]$, where $[x]$ denotes the integer nearest to $x$. In the first case, $\rho_{12}$ converges relatively quickly to 1 , and in the second case the convergence is relatively slow. Figure 1 demonstrates that, as anticipated, the oscillatory behaviour is less pronounced, and the normal approximation better, in the "slow" case.

Finally, Figure 2 shows that broadly similar results are obtained for coverage probabilities of percentile bootstrap confidence intervals for $E(S)$. Let $s_{\alpha}$ denote the $\alpha$-level quantile of the distribution of $S-E(S)$, and let $\hat{s}_{\alpha}$, the parametric bootstrap estimator of $s_{\alpha}$, be the $\alpha$-level quantile of the distribution of $S^{*}-S$ given $\mathcal{X}$, i.e. $\hat{s}_{\alpha}=\inf \left\{s: P\left(S^{*}-S \leq s \mid \mathcal{X}\right) \geq \alpha\right\}$. A naive $\alpha$-level one-sided percentile-bootstrap confidence interval for $E(S)$, with nominal coverage probability $1-\alpha$, is given by

$$
\begin{equation*}
\mathcal{I}_{\alpha}=\left(-\infty, S-\hat{s}_{\alpha}\right] . \tag{3.2}
\end{equation*}
$$

In the figure, we give plots of estimates of the coverage probability $P\left\{E(S) \in \mathcal{I}_{\alpha}\right\}$ of $\mathcal{I}_{\alpha}$ against $n_{1}$, estimated using $10^{5}$ Monte-Carlo simulations, for $\alpha=0.95$. We used $B=9999$ simulations in each bootstrap step. Each panel depicts the case $\rho_{0}=1$, and successive panels also give results when $\rho_{0}=3^{1 / 2}, 5^{1 / 2}, e$ and $\phi=\left(1+5^{1 / 2}\right) / 2$, respectively. Each of these values of $\rho_{0}$ is an irrational of type 1 , and in each instance the oscillations are markedly less, and the normal approximation markedly improved, relative to the case $\rho_{0}=1$.


Figure 1. Plots of $P(x)$ against $n_{1}$. Plots are given for $x=\Phi^{-1}(\alpha)=z_{\alpha}$ and $\alpha=0.95,0.85$, and 0.75 , and for $n_{2}$ equal to the nearest integer to $\rho_{0} n_{1}$, with $\rho_{0}=1$ or 2 (top left), $\rho_{0}=1$ or $2^{1 / 2}$ (top right), $\rho_{0}=1$ or $\pi / 2$ (bottom left) and $\rho_{0}$ converges to 1 rapidly or slowly (bottom right; see text for details).

## 4. Proofs

### 4.1. Proof of Theorem 1

### 4.1.1. Proof of part (i) of Theorem 1

Here we show that if (2.5) holds for some $\rho_{j_{1} j_{2}}$, where $j_{1} \neq j_{2}$, then (2.2) obtains. Some of the asymptotic expansions in our argument are taken a little further than is necessary for (2.2); the extra detail will be used in the proof of Theorem 2.


Figure 2. Plots of estimates of $P\left\{E(S) \in \mathcal{I}_{\alpha}\right\}$, against $n_{1}$; see text for details. Each panel shows the case $\rho_{0}=1$ and also, in respective panels, the cases $\rho_{0}=3^{1 / 2}, \rho_{0}=5^{1 / 2}, \rho_{0}=e$ and $\rho_{0}=\left(1+5^{1 / 2}\right) / 2$. Throughout, $x=\Phi^{-1}(\alpha)$ where $\alpha=0.95$.

Step 1: Proof that it is sufficient to consider the case $k=2$. Without loss of generality, (2.5) holds for $\rho_{12}$, and in this case we write $S-E(S)=S_{1}+S_{2}$, where $S_{1}=(1-E)\left(\bar{X}_{1}+\bar{X}_{2}\right)$ and $S_{2}=(1-E)\left(\bar{X}_{3}+\cdots+\bar{X}_{k}\right)$, where $E$ denotes the expectation operator. Recall that $S_{2}$ is independent of $\bar{X}_{1}$ and $\bar{X}_{2}$. Suppose we can prove that, analogously to (2.2),

$$
\begin{equation*}
P\left\{\frac{S_{1}}{\left(\operatorname{var} S_{1}\right)^{1 / 2}} \leq x\right\}=\Phi(x)+n^{-1 / 2} \frac{1}{6} \beta_{1}\left(1-x^{2}\right) \phi(x)+o\left(n^{-1 / 2}\right) \tag{4.1}
\end{equation*}
$$

uniformly in $x$, where, reflecting (2.3),

$$
\beta_{1}=\beta_{1}(n)=\frac{n^{1 / 2} E\left(S_{1}^{3}\right)}{\left(\operatorname{var} S_{1}\right)^{3 / 2}}
$$

If we prove that (2.2), in the case of general $k$, follows from (4.1), we shall have shown that it is sufficient to derive Theorem 1 the case $k=2$.

Since $P(S \leq x)=E\left\{P\left(S_{1} \leq x-S_{2} \mid S_{2}\right)\right\}$ then we can deduce from (4.1) that

$$
\begin{align*}
P(S \leq x)=E\left(\Phi\left\{\frac{x-S_{2}}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}+\frac{\beta_{1}}{6 n^{1 / 2}}[1\right. & \left.-\left\{\frac{x-S_{2}}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}^{2}\right] \\
& \left.\times \phi\left\{\frac{x-S_{2}}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}\right)+o\left(n^{-1 / 2}\right), \tag{4.2}
\end{align*}
$$

uniformly in $x$. Let $R=S_{2} /\left(\operatorname{var} S_{1}\right)^{1 / 2}$, and put $\tau_{1}^{2}=\operatorname{var}(R)$, which is bounded away from zero and infinity as $n \rightarrow \infty$. It is straightforward to prove that, if $N$ denotes a normally distributed random variable with the same mean (i.e., zero mean) and variance as $S_{2}$, then

$$
\begin{align*}
& E\left(\left[1-\left\{\frac{x-S_{2}}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}^{2}\right] \phi\left\{\frac{x-S_{2}}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}\right) \\
& \quad=\int\left[1-\left\{\frac{x}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}-t\right\}^{2}\right] \phi\left\{\frac{x}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}-t\right\} d P(R \leq t)  \tag{4.3}\\
& \quad=\int\left[1-\left\{\frac{x}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}-t\right\}^{2}\right] \phi\left\{\frac{x}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}-t\right\} \frac{1}{\tau_{1}} \phi\left(\frac{t}{\tau_{1}}\right) d t+O\left(n^{-1 / 2}\right)  \tag{4.4}\\
& \quad=E\left(\left[1-\left\{\frac{x-N}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}^{2}\right] \phi\left\{\frac{x-N}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}\right)+O\left(n^{-1 / 2}\right) \tag{4.5}
\end{align*}
$$

uniformly in $x$. The passage from (4.3) to (4.4) can be accomplished by integrating by parts in (4.3), then using an Edgeworth expansion of the distribution of $R$, then separating out the term in $n^{-1 / 2}$ in that expansion, and finally, undoing the integration by parts as it applies to the leading term in the Edgeworth expansion.

Let $\tau_{2}^{2}=\operatorname{var} S_{2}$ and $\beta_{2}=\beta_{2}(n)=n^{1 / 2} E\left(S_{2}^{3}\right) / \tau_{2}^{3}$. If

$$
\Phi\left(r / \tau_{2}\right)+n^{-1 / 2} \frac{1}{6} \beta_{2}\left\{1-\left(r / \tau_{2}\right)^{2}\right\} \phi\left(r / \tau_{2}\right)
$$

represents the two-term Edgeworth approximation to $P\left(S_{2} \leq r\right)$ that would be employed if the distribution of $S_{2}$ were continuous, then it can be proved that, uniformly in $x$,

$$
\begin{align*}
E\left[\Phi\left\{\frac{x-S_{2}}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}\right]= & \int \Phi\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\} d_{r}\left\{\Phi\left(x / \tau_{2}\right)\right. \\
& \left.+n^{-1 / 2} \frac{1}{6} \beta_{2}\left\{1-\left(r / \tau_{2}\right)^{2}\right\} \phi\left(r / \tau_{2}\right)\right\} \\
& + \begin{cases}o\left(n^{-1 / 2}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{3}<\infty \\
O\left(n^{-1}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{4}<\infty\end{cases} \tag{4.6}
\end{align*}
$$

To derive (4.6), first integrate by parts on the left-hand side, writing it as

$$
\begin{align*}
\frac{1}{\left(\operatorname{var} S_{1}\right)^{1 / 2}} \int \phi\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\} & P\left(S_{2} \leq r\right) d r \\
& =\int \phi\left\{\frac{x}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}-t\right\} P(R \leq t) d t \tag{4.7}
\end{align*}
$$

Next, write down an Edgeworth expansion, (E) say, for the joint distribution of $\bar{X}_{3}, \ldots, \bar{X}_{k}$, up to terms of $o\left(n^{-1 / 2}\right)$ when $\max _{j} E\left|X_{j 1}\right|^{3}<\infty$ and $O\left(n^{-1}\right)$ when $\max _{j} E\left|X_{j 1}\right|^{4}<\infty$. The expansion will include the conventional discontinuous terms of size $n^{-1 / 2}$. Use (E) to the derive discontinuous term $n^{-1 / 2} D$, say, up to a remainder of smaller order $n^{-1 / 2}$, in an Edgeworth expansion of the distribution of $R$. Since the function $\phi$ is smooth, the impact of $n^{-1 / 2} D$ on the right-hand side of (4.7) equals $O\left(n^{-1}\right)$, this being obtained by multiplying together the factor $n^{-1 / 2}$ and another term of order $n^{-1 / 2}$ that results from integrating $D$ against a smooth function. Therefore, (4.6) holds.

Combining (4.2), (4.5) and (4.6), we deduce that

$$
\begin{align*}
P(S \leq x)= & \int\left(\Phi\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}+\frac{\beta_{1}}{6 n^{1 / 2}}\left[1-\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}^{2}\right] \phi\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}\right) \\
& \times d_{r}\left\{\Phi\left(x / \tau_{2}\right)+n^{-1 / 2} \frac{1}{6} \beta_{2}\left\{1-\left(r / \tau_{2}\right)^{2}\right\} \phi\left(r / \tau_{2}\right)\right\}+o\left(n^{-1 / 2}\right), \tag{4.8}
\end{align*}
$$

uniformly in $x$. Result (4.8) is equivalent to (2.2), and so (4.2), representing (2.2) in the case $k=2$, implies (2.2) for general $k \geq 2$, as had to be shown.

Step 2: Proof of (2.2) when $k=2$. In this section, we shall show that, if $k=2$ and (2.5) holds for $\rho_{12}=e_{2} n_{1} /\left(e_{1} n_{2}\right)$, then (2.2) holds.

To this end, define

$$
T=(S-E S) /(\operatorname{var} S)^{1 / 2}=\frac{\bar{X}_{1}+\bar{X}_{2}-\mu_{1}-\mu_{2}}{\left(n_{1}^{-1} \sigma_{1}^{2}+n_{2}^{-1} \sigma_{2}^{2}\right)^{1 / 2}}=c_{1} T_{1}+c_{2} T_{2},
$$

where $T_{j}=\left(\bar{X}_{j}-\mu_{j}\right) /\left(n_{j}^{-1} \sigma_{j}^{2}\right)^{1 / 2}$ and $c_{1}$ and $c_{2}$ are defined as at (2.9). In this notation,

$$
\begin{align*}
P(T \leq x)= & P\left(c_{1} T_{1}+c_{2} T_{2} \leq x\right)=E\left\{P\left(c_{1} T_{1} \leq x-c_{2} T_{2} \mid T_{2}\right)\right\} \\
= & E\left\{\Phi\left(\frac{x-c_{2} T_{2}}{c_{1}}\right)+n_{1}^{-1 / 2} A_{1}\left(\frac{x-c_{2} T_{2}}{c_{1}}\right)+n_{1}^{-1 / 2} D_{1}\left(\frac{x-c_{2} T_{2}}{c_{1}}\right)\right\} \\
& + \begin{cases}o\left(n^{-1 / 2}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{3}<\infty \\
O\left(n^{-1}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{4}<\infty,\end{cases} \tag{4.9}
\end{align*}
$$

where $A_{j}$ and $D_{j}$ will refer to the smooth and discontinuous terms, respectively, in the $n_{j}^{-1 / 2}$ component of an Edgeworth expansion of the distribution of $T_{j}$ for $j=1,2$. In particular,
$n_{j}^{-1 / 2} A_{j}$ and $n_{j}^{-1 / 2} D_{j}$ are the counterparts of the second and third terms, respectively, on the right-hand side of formula (35) p. 56 of Esseen [9].

Writing $B$ for either $\Phi$ or $A_{1}$, appearing on the right-hand side of (4.9), we have:

$$
\begin{aligned}
E\left\{B\left(\frac{x-c_{2} T_{2}}{c_{1}}\right)\right\} & =\int B\left(\frac{x-c_{2} u}{c_{1}}\right) d P\left(T_{2} \leq u\right) \\
& =\frac{c_{2}}{c_{1}} \int B^{\prime}\left(\frac{x-c_{2} u}{c_{1}}\right) P\left(T_{2} \leq u\right) d u
\end{aligned}
$$

As in the argument leading to (4.6) it can be shown that the discontinuous term $n_{2}^{-1 / 2} D_{2}$, in the Edgeworth expansion of $P\left(T_{2} \leq x\right)$, contributes only $O\left(n^{-1}\right)$. Therefore, if we write $\mathcal{E}_{2}(u)$ for the Edgeworth approximation to $P\left(T_{2} \leq u\right)$ that includes the leading Gaussian term, plus the continuous part of the component of order $n_{2}^{-1 / 2}$, and neglects everything else, we deduce from (4.9) that

$$
\begin{align*}
P(T \leq x)= & \int\left\{\Phi\left(\frac{x-c_{2} u}{c_{1}}\right)+n_{1}^{-1 / 2} A_{1}\left(\frac{x-c_{2} u}{c_{1}}\right)\right\} d_{u} \mathcal{E}_{2}(u) \\
& +n_{1}^{-1 / 2} E\left\{D_{1}\left(\frac{x-c_{2} T_{2}}{c_{1}}\right)\right\}+ \begin{cases}o\left(n^{-1 / 2}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{3}<\infty \\
O\left(n^{-1}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{4}<\infty\end{cases} \tag{4.10}
\end{align*}
$$

Now we turn our attention to:

$$
\begin{align*}
E\left\{D_{1}\left(\frac{x-c_{2} T_{2}}{c_{1}}\right)\right\} & =\int D_{1}\left(\frac{x-c_{2} u}{c_{1}}\right) d P\left(T_{2} \leq u\right) \\
& =-\int P\left(T_{2} \leq u\right) d_{u} D_{1}\left(\frac{x-c_{2} u}{c_{1}}\right)  \tag{4.11}\\
& =I_{1}(x)+n_{2}^{-1 / 2} I_{2}(x)+ \begin{cases}o(1) & {\text { if } \max _{j} E\left|X_{j 1}\right|^{3}<\infty}_{O\left(n^{-1 / 2}\right)} \\
\text { if } \max _{j} E\left|X_{j 1}\right|^{4}<\infty,\end{cases} \\
& =n_{2}^{-1 / 2} I_{2}(x)+ \begin{cases}o(1) & \text { if } \max _{j} E\left|X_{j 1}\right|^{3}<\infty \\
O\left(n^{-1 / 2}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{4}<\infty,\end{cases} \tag{4.12}
\end{align*}
$$

where

$$
\begin{equation*}
I_{1}(x)=\int D_{1}\left(\frac{x-c_{2} u}{c_{1}}\right) \phi(u) d u, \quad I_{2}(x)=\int D_{1}\left(\frac{x-c_{2} u}{c_{1}}\right) d D_{2}(u) \tag{4.13}
\end{equation*}
$$

To obtain the third identity in the string of formulae leading to (4.12), we used the integration by parts step at (4.11), a short Taylor expansion of $P\left(T_{2} \leq u\right)$ with a remainder of $o\left(n^{-1 / 2}\right)$ if $\max _{j} E\left|X_{j 1}\right|^{3}<\infty$ and $O\left(n^{-1}\right)$ if $\max _{j} E\left|X_{j 1}\right|^{4}$, and the fact that $\int\left|d D_{1}\right|=O\left(n^{1 / 2}\right)$ uniformly in $x$. (This can be deduced either directly or by making use of (4.14) below.) Finally, it can be shown, arguing as in the proof of (4.6), that $I_{1}(x)=O\left(n^{-1 / 2}\right)$, from which (4.12) follows.

Note too that, with $\sigma_{j}$ defined as immediately above (2.1),

$$
\begin{align*}
D_{j}(x) & =\frac{e_{j}}{\sigma_{j}} \psi\left\{\frac{\left(x-\xi_{1 n}\right) \sigma_{j} n_{j}^{1 / 2}}{e_{j}}\right\} \phi(x) \\
& =\frac{e_{j}}{\sigma_{j}} \psi\left[\frac{\sigma_{j} n_{j}^{1 / 2} x}{e_{j}}-\left\{\left(n_{j} x_{j} / e_{j}\right)-\left\lfloor n_{j} x_{j} / e_{j}\right\rfloor\right\}\right] \phi(x), \tag{4.14}
\end{align*}
$$

where, as in Sections 1 and $2, \psi(x)=\lfloor x\rfloor-x+\frac{1}{2},\lfloor x\rfloor$ is the largest integer not strictly exceeding $x$, and $\xi_{j n}=e_{j}\left(\sigma_{j} n_{j}^{1 / 2}\right)^{-1}\left\{\left(n_{j} x_{j} / e_{j}\right)-\left\lfloor n_{j} x_{j} / e_{j}\right\rfloor\right\}$ if the lattice is located at points $x_{j}+v e_{j}$ for integers $v$, see Esseen [9], (29), (31) and (35) pp. 55/56. Defining $\gamma=\left(e_{1} e_{2} / \sigma_{1} \sigma_{2}\right)$, as in Section 2.3; putting

$$
\psi_{j}(x)=\psi\left\{\frac{\sigma_{j} n_{j}^{1 / 2}}{e_{j}}\left(x-\xi_{j n}\right)\right\}
$$

and noting that, by (4.14), $D_{j}(x)=\left(e_{j} / \sigma_{j}\right) \psi_{j}(x) \phi(x)$; we deduce that

$$
\begin{aligned}
I_{2}(x) / \gamma & =\frac{1}{\gamma} \int D_{1}\left(\frac{x-c_{2} u}{c_{1}}\right) d D_{2}(u)=\int\left(\psi_{1} \phi\right)\left(\frac{x-c_{2} u}{c_{1}}\right) d\left\{\psi_{2}(u) \phi(u)\right\} \\
& =\int\left(\psi_{1} \phi\right)\left(\frac{x-c_{2} u}{c_{1}}\right)\left\{\phi(u) d \psi_{2}(u)+\psi_{2}(u) d \phi(u)\right\} \\
& =\int\left(\psi_{1} \phi\right)\left(\frac{x-c_{2} u}{c_{1}}\right) \phi(u) d \psi_{2}(u)+\frac{1}{2} \int\left(\psi_{1} \phi\right)\left(\frac{x-c_{2} u}{c_{1}}\right) \psi_{2}(u) d \phi(u) .
\end{aligned}
$$

The last-written integral equals $O(1)$, uniformly in $x$, and so, with $I_{2}$ as at (4.13),

$$
\begin{equation*}
I_{2}(x)=\gamma I_{3}(x)+O(1) \tag{4.15}
\end{equation*}
$$

uniformly in $x$, where

$$
I_{3}(x)=\int\left(\psi_{1} \phi\right)\left(\frac{x-c_{2} u}{c_{1}}\right) \phi(u) d \psi_{2}(u) .
$$

Since $\psi_{2}$ has jumps of size +1 at points $u$ where $\left(u-\xi_{2 n}\right) \sigma_{2} n_{2}^{1 / 2} / e_{2}$ is an integer, i.e. $u=$ $u_{\nu} \equiv \xi_{2 n}+e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v$ for an integer $v$, then

$$
\begin{align*}
I_{3}(x)= & \sum_{v}\left(\psi_{1} \phi\right)\left(\frac{x-c_{2} u_{v}}{c_{1}}\right) \phi\left(u_{v}\right) \\
= & \sum_{v} \phi\left(\frac{x}{c_{1}}-\frac{c_{2}}{c_{1}}\left\{\xi_{2 n}+e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v\right\}\right) \phi\left\{\xi_{2 n}+e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v\right\} \\
& \times \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} \tag{4.16}
\end{align*}
$$

where $\xi_{n}$ is as at (2.8) and we have used the fact that

$$
\begin{aligned}
\psi_{1}\left(\frac{x-c_{2} u_{v}}{c_{1}}\right) & =\psi\left[\frac{\sigma_{1} n_{1}^{1 / 2}}{e_{1}}\left\{\left(x-c_{2} u_{v}\right) c_{1}^{-1}-\xi_{1 n}\right\}\right] \\
& =\psi\left\{\frac{\left(x-c_{2} u_{v}-c_{1} \xi_{1 n}\right) \sigma_{1} n_{1}^{1 / 2}}{c_{1} e_{1}}\right\} \\
& =\psi\left[\frac{\left\{x-\left(c_{1} \xi_{1 n}+c_{2} \xi_{2 n}\right)-c_{2} e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v\right\} \sigma_{1} n_{1}^{1 / 2}}{c_{1} e_{1}}\right] \\
& =\psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\}
\end{aligned}
$$

with

$$
\xi_{n}(x)=\left\{x-\left(c_{1} \xi_{1 n}+c_{2} \xi_{2 n}\right)\right\} \frac{\sigma_{1} n_{1}^{1 / 2}}{c_{1} e_{1}}
$$

Recall that $\xi_{j n}=e_{j}\left(\sigma_{j} n_{j}^{1 / 2}\right)^{-1}\left\{\left(n_{j} x_{j} / e_{j}\right)-\left\lfloor n_{j} x_{j} / e_{j}\right\rfloor\right\}$ if the lattice is located at points $x_{j}+$ $v e_{j}$ for integers $v$. In particular, $\xi_{j n}=O\left(n^{-1 / 2}\right)$ for $j=1,2$. Therefore, Taylor expanding the arguments of the functions $\phi$ at (4.16), and defining

$$
\begin{equation*}
I_{4}(x)=\sum_{v} \phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} v\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v\right\} \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} \tag{4.17}
\end{equation*}
$$

we deduce that

$$
\begin{equation*}
I_{3}(x)=I_{4}(x)+O(1) \tag{4.18}
\end{equation*}
$$

uniformly in $x$. Combining (4.10), (4.12), (4.15) and (4.18), we deduce that

$$
\begin{align*}
P(T \leq x)= & \int\left\{\Phi\left(\frac{x-c_{2} u}{c_{1}}\right)+n_{1}^{-1 / 2} A_{1}\left(\frac{x-c_{2} u}{c_{1}}\right)\right\} d_{u} \mathcal{E}_{2}(u) \\
& +\left(n_{1} n_{2}\right)^{-1 / 2} \gamma I_{4}(x)+ \begin{cases}o\left(n^{-1 / 2}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{3}<\infty \\
O\left(n^{-1}\right) & \text { if } \max _{j} E\left|X_{j 1}\right|^{4}<\infty\end{cases} \tag{4.19}
\end{align*}
$$

If we can show that

$$
\begin{equation*}
\sup _{-\infty<x<\infty}\left|I_{4}(x)\right|=o\left(n^{1 / 2}\right) \tag{4.20}
\end{equation*}
$$

then it will follow from (4.19), in cases where $\max _{j} E\left|X_{j 1}\right|^{3}<\infty$, that

$$
\begin{equation*}
P(T \leq x)=\int\left\{\Phi\left(\frac{x-c_{2} u}{c_{1}}\right)+n_{1}^{-1 / 2} A_{1}\left(\frac{x-c_{2} u}{c_{1}}\right)\right\} d_{u} \mathcal{E}_{2}(u)+o\left(n^{-1 / 2}\right) \tag{4.21}
\end{equation*}
$$

The right-hand side here is Edgeworth expansion we would expect the distribution of $T$ to enjoy if we were able to ignore the latticeness of the distributions of $X_{j 1}$ for $j=1,2$. That is, (4.21) is just (2.2) in the particular case $k=2$. Therefore, provided (4.20) holds then we shall have shown that (2.2) holds whenever $k=2$. It remains to derive (4.20).

Step 3: Proof of (4.20). Given $\varepsilon>0$, partition the set of all integers into adjacent blocks $\mathcal{N}_{\ell}$, for $-\infty<\ell<\infty$, where each block consists of just $2\left\lfloor n^{1 / 2} \varepsilon\right\rfloor+1$ consecutive integers, and the central integer is denoted by $\bar{\nu}_{\ell}$. Recalling the definition of $I_{4}(x)$ at (4.17), we deduce that

$$
\begin{equation*}
I_{4}=\sum_{-\infty<\ell<\infty} J_{1, \ell}, \tag{4.22}
\end{equation*}
$$

where

$$
\begin{equation*}
J_{1, \ell}(x)=\sum_{\nu \in \mathcal{N}_{\ell}} \phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} v\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} v\right\} \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} \tag{4.23}
\end{equation*}
$$

Now,

$$
\begin{equation*}
J_{1, \ell}=J_{2, \ell}+R_{\ell}, \tag{4.24}
\end{equation*}
$$

where

$$
\begin{equation*}
J_{2, \ell}(x)=\phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} \bar{v}_{\ell}\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} \bar{v}_{\ell}\right\} \sum_{v \in \mathcal{N}_{\ell}} \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} \tag{4.25}
\end{equation*}
$$

and $R_{\ell}$ is defined naively by (4.24). Given an integer $r$, let $\ell(r)$ denote the unique value of $\ell$ such that $r \in \mathcal{N}_{\ell}$. Then, since $|\psi| \leq 1$,

$$
\begin{align*}
\left|\sum_{-\infty<\ell<\infty} R_{\ell}\right| \leq & \sum_{r} \left\lvert\, \phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} r\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} r\right\}\right. \\
& \left.-\phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} \bar{v}_{\ell(r)}\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} \bar{v}_{\ell(r)}\right\} \right\rvert\, \\
\leq & C_{1} \varepsilon n^{1 / 2}, \tag{4.26}
\end{align*}
$$

where the constant $C_{1}$ does not depend on $\varepsilon$ or $n$.
Let $\rho=e_{2} n_{1} /\left(e_{1} n_{2}\right)$, and define

$$
\chi_{\mathcal{N}}(z, \rho) \equiv \sum_{\nu \in \mathcal{N}} \psi(z-\rho \nu) .
$$

In this notation,

$$
\begin{equation*}
J_{2, \ell}=\phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} \bar{v}_{\ell}\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} \bar{\nu}_{\ell}\right\} \chi_{\mathcal{N}_{\ell}}\left\{\xi_{n}(x), \rho\right\} . \tag{4.27}
\end{equation*}
$$

If we can prove that, whenever the set $\mathcal{N}$ consists of $|\mathcal{N}|$ consecutive integers and $C_{2}<C_{3}$ are positive constants,

$$
\begin{equation*}
\sup _{C_{2} n^{1 / 2} \leq|\mathcal{N}| \leq C_{3} n^{1 / 2}-\infty<z<\infty} \sup _{\chi_{\mathcal{N}}(z, \rho) \mid=o\left(n^{1 / 2}\right), ~}^{\text {and }} \tag{4.28}
\end{equation*}
$$

as $|\mathcal{N}| \rightarrow \infty$, then it will follow from (4.27) that

$$
\begin{align*}
\left|\sum_{-\infty<\ell<\infty} J_{2, \ell}\right| & =o\left[n^{1 / 2} \sum_{-\infty<\ell<\infty} \phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} \bar{v}_{\ell}\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} \bar{v}_{\ell}\right\}\right] \\
& =o\left(n^{1 / 2}\right) \tag{4.29}
\end{align*}
$$

for each $\varepsilon>0$, since the series on the first right-hand side of (4.29) is bounded uniformly in $n$. (To appreciate why, observe that $\bar{v}_{\ell}$ is approximately an integer multiple of $n^{1 / 2}$, plus a constant.) Note that, since the left-hand side of (4.28) involves the supremum over $z$, then that quantity does not depend on the location of the set $\mathcal{N}$ on the line, only on the number of consecutive integers it contains.

The desired result (4.20) follows from (4.22), (4.24), the fact that (4.26) holds for each $\varepsilon>0$, and (4.29). To complete the proof of (4.20), we shall derive (4.28). Specifically, we shall prove that, in cases where (2.5) is satisfied for $\rho_{12}=\rho=e_{2} n_{1} /\left(e_{1} n_{2}\right)$, (4.28) obtains.

Assume that $\mathcal{N}$ consists of $p$ consecutive integers, where $C_{2} n^{1 / 2} \leq p \leq C_{3} n^{1 / 2}$. Koksma's inequality (see, e.g., Theorems 1.3 and 5.1, pp. 91 and 143 of [13]), and the Erdős-Turán inequality (see, e.g., formula (2.42), p. 114 of [13]), can be combined to prove that, for all integers $m \geq 1$,

$$
\begin{align*}
\sup _{z}\left|\chi_{\mathcal{N}}\left(z, \rho_{12}\right)\right| & \leq C_{4}\left\{\frac{p}{m}+\sum_{\ell=1}^{m} \frac{1}{\ell} \sup _{z}\left|\sum_{r=1}^{p} \exp \left(2 \pi i \ell r \rho_{12}\right)\right|\right\} \\
& \leq C_{4}\left\{\frac{p}{m}+\sum_{\ell=1}^{m} \frac{1}{\ell\left|\sin \left(\ell \rho_{12} \pi\right)\right|}\right\} \tag{4.30}
\end{align*}
$$

where $C_{4}$ is an absolute constant. Since (2.5) is assumed to hold with $\left(j_{1}, j_{2}\right)=(1,2)$ then, for each fixed $m$,

$$
\max _{1 \leq \ell \leq m}\left|\sin \left(\ell \rho_{12} \pi\right)\right|^{-1}=o\left(n^{1 / 2}\right)
$$

Hence, by (4.30),

$$
\begin{equation*}
\sup _{z}\left|\chi_{\mathcal{N}}\left(z, \rho_{12}\right)\right| \leq \frac{C_{3} C_{4} n^{1 / 2}}{m}+o\left(n^{1 / 2}\right) \tag{4.31}
\end{equation*}
$$

where the $o\left(n^{1 / 2}\right)$ term is of that order uniformly in $\mathcal{N}$ such that $C_{2} n^{1 / 2} \leq|\mathcal{N}| \leq C_{3} n^{1 / 2}$. However, $m$ can be taken arbitrarily large, and none of $C_{2}, C_{3}$ and $C_{4}$ depends on $m$ or $n$. Therefore, (4.31) implies (4.28).

### 4.1.2. Proof of part (ii) of Theorem 1

We can write

$$
\bar{X}_{1}+\cdots+\bar{X}_{k}=\frac{e_{1}}{n_{1}}\left(Y_{1}+\cdots+Y_{k}\right)+\mu
$$

where $\mu$ is deterministic and, for each $j, Y_{j}$ is the sum of $n_{j}$ random variables $Y_{j 1}, \ldots, Y_{j n_{j}}$, each having a lattice distribution (not depending on $n$ ) supported on the set of points $\rho_{1 j} \ell$ for $\ell \in \mathbb{Z}$, and with the $Y_{j i}$ s being totally independent. Of course, $\rho_{11}=1$. Since each $\rho_{j_{1} j_{2}}$ equals a rational number, not depending on $n$, then the set $\bigcup_{j}\left\{\rho_{1 j} \ell, \ell \in \mathbb{Z}\right\}$ can itself be represented as a maximal lattice, $\mathcal{L}$ say, not depending on $n$. The distribution of

$$
Y_{1}+\cdots+Y_{k}=\sum_{j=1}^{k} \sum_{i=1}^{n_{j}} Y_{j i}
$$

can be viewed as the distribution of the sum of $n=n_{1}+\cdots+n_{k}$ independent and identically distributed random variables each having a mixture distribution, $D_{n}$ say, with support confined to $\mathcal{L}$. Although $D_{n}$ depends on $n$, since it is always supported on the same lattice, standard methods can be used to derive an Edgeworth expansion of the distribution of $Y_{1}+\cdots+Y_{k}$, from which it can be seen that there is a nonvanishing discontinuous term, not present in (2.2).

### 4.2. Proof of Theorem 2

Step 1: Proof that it is sufficient to consider the case $k=2$. We give the argument only in outline, since it parallels that in step 1 of the derivation of Theorem 1. Suppose it is possible to derive the version of (4.1) where the remainder $o\left(n^{-1 / 2}\right)$ is replaced by $O\left(n^{\xi-1}\right)$, for all $\xi>0$. Then, as in the earlier proof, we have (4.2) where the remainder term is $O\left(n^{\xi-1}\right)$, for all $\xi>0$, instead of $o\left(n^{-1 / 2}\right)$. The string of arguments leading to (4.5) holds without change, as too does (4.6). Combining the revised (4.2) with the old (4.5) and (4.6) we deduce the following version of (4.8):

$$
\begin{aligned}
P(S \leq x)= & \int\left(\Phi\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}+\frac{\beta_{1}}{6 n^{1 / 2}}\left[1-\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}^{2}\right] \phi\left\{\frac{x-r}{\left(\operatorname{var} S_{1}\right)^{1 / 2}}\right\}\right) \\
& \times d\left\{\Phi\left(x / \tau_{2}\right)+n^{-1 / 2} \frac{1}{6} \beta_{2}\left\{1-\left(r / \tau_{2}\right)^{2}\right\} \phi\left(r / \tau_{2}\right)\right\}+O\left(n^{\xi-1}\right),
\end{aligned}
$$

uniformly in $x$ and for all $\xi>0$. This formula is equivalent to (2.2), with the remainder there replaced by $O\left(n^{\xi-1}\right)$, and so we have shown that it suffices to consider $k=2$.

Step 2: Completion of proof of Theorem 2. Combining (4.10) and (4.12) in the case $\max _{j} E\left|X_{j 1}\right|^{4}<\infty$, and noting (4.15) and (4.18), we deduce the version of (4.19) when $\max _{j} E\left|X_{j 1}\right|^{4}<\infty$.

Next we reintroduce the notation noted below (2.8), where $\alpha \in\left(0, \frac{1}{2}\right), \mathcal{N}_{\ell}$ (for $-\infty<\ell<\infty$ ) is a partition of the set of all integers into adjacent blocks each containing $2\left\lfloor n^{\alpha}\right\rfloor+1$ consecutive integers, $\bar{\nu}_{\ell}$ is the central integer in $\mathcal{N}_{\ell}$, and $\nu_{\ell}=v-\bar{\nu}_{\ell}$ for $v \in \mathcal{N}_{\ell}$. Property (4.22) continues to
hold, with $J_{1, \ell}$ still given by (4.23). Again we define $R_{\ell}$ and $J_{2, \ell}$ by (4.24) and (4.25). However, this time we give an expansion for, rather than an upper bound to, $R_{\ell}$. As a first step, note that

$$
\begin{aligned}
R_{\ell}(x)= & J_{1, \ell}(x)-J_{2, \ell}(x) \\
= & \sum_{v \in \mathcal{N}_{\ell}}\left[\phi\left\{\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}}\left(\bar{v}_{\ell}+v_{\ell}\right)\right\} \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1}\left(\bar{v}_{\ell}+v_{\ell}\right)\right\}\right. \\
& \left.-\phi\left(\frac{x}{c_{1}}-\frac{e_{2} n_{1}^{1 / 2}}{\sigma_{1} n_{2}} \bar{v}_{\ell}\right) \phi\left\{e_{2}\left(\sigma_{2} n_{2}^{1 / 2}\right)^{-1} \bar{v}_{\ell}\right\}\right] \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} .
\end{aligned}
$$

Taylor-expanding, and using the argument in the paragraph immediately below that containing (2.10), we deduce that

$$
\begin{array}{r}
\sum_{-\infty<\ell<\infty} R_{\ell}(x)=\sum_{r=1}^{r_{0}} \sum_{-\infty<\ell<\infty} \frac{\phi_{r}\left(\bar{v}_{\ell} / n_{1}^{1 / 2}, x\right)}{r!n_{1}^{r / 2}} \sum_{v \in \mathcal{N}_{\ell}} v_{\ell}^{r} \psi\left\{\xi_{n}(x)-\frac{e_{2} n_{1}}{e_{1} n_{2}} v\right\} \\
+O\left(n^{\alpha} \cdot n^{1 / 2} \cdot n^{\left(r_{0}+1\right)\{\alpha-(1 / 2)\}}\right) \tag{4.32}
\end{array}
$$

uniformly in $x$. Adding $\sum_{\ell} J_{2, \ell}$ to either side of (4.32) has the effect, on the right-hand side, of changing the range of summation of the first series to $0 \leq r \leq r_{0}$. Therefore,

$$
\begin{equation*}
\sum_{-\infty<\ell<\infty}\left\{J_{2, \ell}(x)+R_{\ell}(x)\right\}=\gamma^{-1} K_{n}(x)+O\left(n^{\alpha} \cdot n^{1 / 2} \cdot n^{\left(r_{0}+1\right)\{\alpha-(1 / 2)\}}\right) \tag{4.33}
\end{equation*}
$$

uniformly in $x$, where $\gamma=\prod_{j=1,2}\left(e_{j} / \sigma_{j}\right)$ and $K_{n}$ is at (2.10). If $r_{0} \geq 4 \alpha /(1-2 \alpha)$, as stipulated in Theorem 2, then the " $O$ " remainder in (4.33) is just $O(1)$. In this case,

$$
\begin{equation*}
\gamma I_{4}(x)=\gamma \sum_{-\infty<\ell<\infty} J_{1, \ell}(x)=\gamma \sum_{-\infty<\ell<\infty}\left\{J_{2, \ell}(x)+R_{\ell}(x)\right\}=K_{n}(x)+O(1) \tag{4.34}
\end{equation*}
$$

uniformly in $x$. Part (i) of Theorem 2, which addresses only the case $k=2$, follows from (4.19) and (4.34). Part (ii) of Theorem 2, in the case $k=2$, follows from (4.34) and (2.36). In view of Part 1 of the proof of Theorem 2, this is sufficient to complete the proof of the theorem.

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