# Hypersurfaces of $\boldsymbol{E}_{s}^{4}$ with proper mean curvature vector 

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#### Abstract

Submanifolds satisfying $\Delta \vec{H}=\lambda \vec{H}$ are named by B. Y. Chen submanifolds with proper mean curvature vector. We prove that a hypersurface of the pseudo-Euclidean space $E_{s}^{4}$ with $\Delta \vec{H}=\lambda \vec{H}$ and diagonalizable shape operator, has constant mean curvature.


## 1. Introduction.

Let $x: M^{m} \rightarrow E_{s}^{n}$ be an isometric immersion of an $n$-dimensional connected submanifold of a pseudo-Euclidean space $E_{s}^{m}$. If we denote by $\vec{x}, \vec{H}$, and $\Delta$ the position vector field, the mean curvature vector field, and the Laplace operator respectively of $M$, with respect to the induced metric of $M$, then it is well known that (e.g. [Ch1])

$$
\begin{equation*}
\Delta \vec{x}=-n \vec{H} . \tag{1}
\end{equation*}
$$

In particular, equation (1) shows that $M$ is a minimal submanifold of $E_{s}^{n}$ if and only if its coordinate functions are harmonic. We also observe that every minimal submanifold satisfies

$$
\begin{equation*}
\Delta \vec{H}=\overrightarrow{0} \tag{2}
\end{equation*}
$$

Submanifolds of $E_{s}^{n}$ which satisfy condition (2) are said to have harmonic mean curvature vector field. These submanifolds are often called biharmonic since, in view of (1), condition (2) is equivalent to $\Delta^{2} \vec{x}=\overrightarrow{0}$. Equation (2) is a special case of the equation

$$
\begin{equation*}
\Delta \vec{H}=\lambda \vec{H} \tag{3}
\end{equation*}
$$

Submanifolds of $E_{s}^{m}$ which satisfy condition (3) are said to have proper mean curvature vector field. Equations (2) and (3) can be related to the theory of harmonic and biharmonic maps as explained at the end of the present work (Section 3).

A conjecture of B. Y. Chen ([Ch2]) states that "the only biharmonic submanifolds of Euclidean spaces are the minimal submanifolds". For hypersurfaces in $E^{3}$ and $E^{4}$ the conjecture is supported by the work of several authors ([Ch1], [Di1], [Di2], [Ha-Vl], [De1]). However, it is not true in general for submanifolds in pseudo-Euclidean spaces $E_{s}^{m}$. Counterexamples were presented in $[\mathbf{C h}-\mathbf{I s} \mathbf{1}]$ and $[\mathbf{C h}-\mathbf{I s} \mathbf{2}]$. In contrast, there is

[^0]strong evidence that the conjecture is true for hypersurfaces in pseudo-Euclidean spaces. More precisely, in [Ch-Is1] is was shown that every biharmonic surface in $E_{s}^{3}$ is minimal, in [De-Ka-Pa] that every biharmonic hypersurface $M_{r}^{3}$ of $E_{s}^{4}$ whose shape operator is diagonal is minimal, and in [Ar-De-Ka-Pa] Chen's conjecture was proved for Lorentz hypersurfaces in $E_{1}^{4}$. Recently a "generalized Chen's conjecture" was posed in [Ca-Mo-On], stating that the only biharmonic submanifolds of a manifold with nonpositive sectional curvature are the minimal ones.

Equation (3) was first appeared in [Ch6] where surfaces in $E^{3}$ satisfying (3) were classified. Also, in $[\mathbf{C h} 7]$ it was shown that a submanifold $M$ of a Euclidean space satisfies (3) if and only if $M$ is biharmonic or of 1-type or of null 2-type. Hypersurfaces in $E^{4}$ satisfying (3) with the additional condition of conformal flatness were classified in $[\mathbf{G a}]$, and in $[\mathbf{D e} \mathbf{2}]$ it was proved that every hypersurface of $E^{4}$ satisfying (3) has constant mean curvature. For various other results about submanifolds satisfying (3) in Euclidean spaces, and more generally in space forms, contact, and Sasakian manifolds, we refer to [Ch4], [Ch5], [Ek-Ya], [In1], [In2].

The study of equation (3) for submanifolds in pseudo-Euclidean spaces was originally studied in $[\mathbf{F e}-\mathrm{Lu} 1]$, where the authors classifed surfaces $M_{r}^{2}(r=0,1)$ in the LorenzMinkowski space $E_{1}^{3}$. One of the possibilities for $M_{r}^{2}$ is that it is a submanifold of zero mean curvature $H$. The case of hypersurfaces $M_{r}^{n-1}(r=0,1)$ in $E_{1}^{n}$ satisfying (3) and such that the minimal polynomial of the shape operator is at most of degree two, was stydied in $[\mathbf{F e}-\mathrm{Lu} 2]$, showing that $M_{r}^{n-1}$ has constant mean curvature. Also in [Ch8] various classification theorems for submanifolds in a Minkowski space-time were presented.

The results of the previous paragraph suggest a further study of hypersurfaces of $E_{s}^{n}(0 \leq s \leq n)$ satisfying equation (3). Towards this direction we prove the following:

Theorem. Let $M_{r}^{3}(r=0,1,2,3)$ be a nondegenerate hypersurface of the pseudoEuclidean space $E_{s}^{4}$ with diagonal shape operator. If the mean curvature vector field $\vec{H}$ of $M_{r}^{3}$ satisfies $\Delta \vec{H}=\lambda \vec{H}$, then $M_{r}^{3}$ has constant mean curvature.

The idea of the proof is the following. Equation (3) reduces to the equations

$$
\begin{align*}
& S(\nabla H)=-\varepsilon \frac{3 H}{2}(\nabla H) \\
& \Delta H+\varepsilon H \operatorname{tr} S^{2}=\lambda H \tag{}
\end{align*}
$$

From the above equations together with Codazzi and Gauss equations we eliminate all derivatives. In this way we obtain a polynomial equation with constant coefficients which is satisfied by $H$, therefore $H$ must be constant.

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## 2. Preliminaries.

Hypersurfaces in $E_{s}^{4}$.
Consider the 4-dimensional vector space $\boldsymbol{R}^{4}$ with the standard basis $\left\{e_{1}, e_{2}, e_{3}, e_{4}\right\}$. Let $\langle$,$\rangle denote the indefinite inner product on \boldsymbol{R}^{4}$ whose matrix with respect to the standard basis is a diagonal matrix of index $s \in\{0,1,2,3,4\}$. The space $\boldsymbol{R}^{4}$ with one of these metrics is called the 4-dimensional pseudo-Euclidean space, and is denoted by $E_{s}^{4}$.

A vector $X \in E_{s}^{4}$ is called space-like, time-like, or light-like if $\langle X, X\rangle$ is positive, negative, or zero respectively. Let $x: M_{r}^{3} \rightarrow E_{s}^{4}$ be an isometric immersion of a hypersurface $M_{r}^{3}(r=0,1,2,3)$ in $E_{s}^{4}(s=0,1,2,3,4)$. The hypersurface $M_{r}^{3}$ can itself be endowed with a Riemannian or a pseudo-Riemannian metric structure, depending on whether the metric induced on $M_{r}^{3}$ from the pseudo-Riemannian metric on $E_{s}^{4}$, is positive-definite or indefinite.

Let $\vec{\xi}$ denote a unit normal vector field on $M_{r}^{3}$. Then $\langle\vec{\xi}, \vec{\xi}\rangle=\varepsilon$, where $\varepsilon=-1$ (time-like) or $\varepsilon=+1$ (space-like). The mean curvature vector $\vec{H}=H \vec{\xi}$ with $H=\frac{1}{3 \varepsilon} \operatorname{tr} S$ is a well-defined normal vector field of $M_{r}^{3}$ in $E_{s}^{4}$. Denote by $\nabla$ and $\tilde{\nabla}$ the Levi-Civita connections of $M_{r}^{3}$ and $E_{s}^{4}$ respectively. For any vector fields $X, Y$ tangent to $M_{r}^{3}$, the Gauss formula is given by

$$
\begin{equation*}
\tilde{\nabla}_{X} Y=\nabla_{X} Y+h(X, Y) \vec{\xi}, \tag{4}
\end{equation*}
$$

where $h$ is the second fundamental form. If $S$ is the shape operator of $M_{r}^{3}$ associated to $\vec{\xi}$, then the Weingarten formula is given by

$$
\begin{equation*}
\tilde{\nabla}_{X} \vec{\xi}=-S(X), \tag{5}
\end{equation*}
$$

where $\langle S(X), Y\rangle=\varepsilon h(X, Y)$. The Codazzi equation is given by

$$
\begin{equation*}
\left(\nabla_{X} S\right) Y=\left(\nabla_{Y} S\right) X \tag{6}
\end{equation*}
$$

and the Gauss equation by (cf. [ON])

$$
\begin{equation*}
R(X, Y) Z=\langle S(Y), Z\rangle S(X)-\langle S(X), Z\rangle S(Y) \tag{7}
\end{equation*}
$$

Our convention for the curvature tensor is

$$
R(X, Y) Z=\nabla_{X} \nabla_{Y} Z-\nabla_{Y} \nabla_{X} Z-\nabla_{[X, Y]} Z .
$$

The equation $\Delta \vec{H}=\lambda \vec{H}$.
We now consider a hypersurface $M_{r}^{3}$ of $E_{s}^{4}$ satisfying the condition

$$
\begin{equation*}
\Delta \vec{H}=\lambda \vec{H}, \quad \lambda \in \boldsymbol{R} \backslash\{0\} \tag{8}
\end{equation*}
$$

Here the Laplace operator $\Delta$ acting on a vector-valued function $\vec{V}$ is given by

$$
\Delta \vec{V}=\sum_{i=1}^{3}\left(\tilde{\nabla}_{\nabla_{e_{i} e_{i}}} \vec{V}-\tilde{\nabla}_{e_{i}} \tilde{\nabla}_{e_{i}} \vec{V}\right)
$$

with respect to a local orthonormal frame $\left\{e_{i}\right\}_{i=1}^{3}$.
The shape operator $S$ of a Riemannian hypersurface of $E_{s}^{4}$ is always diagonalizable, but for pseudo-Riemannian hypersurfaces there may be other forms for $S$ as well (e.g. [Ma]). In the present work we assume that the shape operator of the hypersurface $M_{r}^{3}$ in $E_{s}^{4}$ is diagonalizable, i.e. $S=\operatorname{diag}\left\{\lambda_{1}, \lambda_{2}, \lambda_{3}\right\}$. Here the $\lambda_{i}^{\prime} s$ are known as the principal curvature functions of $M_{r}^{3}$. The other possibilities for the shape operator need a separate investigation. The following proposition can be found in various forms (e.g. [Ch1], [Ch-Is1], $[\mathbf{F e}-\mathrm{Lu} 2]$ ), but we give a proof here adjusted to our problem.

Proposition 1. For an isometric immersion $x: M_{r}^{3} \rightarrow E_{s}^{4}$ with diagonal shape operator, the following formula holds:

$$
\Delta \vec{H}=\{2 S(\nabla H)+3 \varepsilon H(\nabla H)\}+\left\{\Delta H+\varepsilon H \operatorname{tr} S^{2}\right\} \vec{\xi}
$$

Proof. Let $\left\{X_{1}, X_{2}, X_{3}\right\}$ be an orthonormal frame such that $\nabla_{X_{i}} X_{j}(P)=0$, at some point $P \in M_{r}^{3}$. From the relation

$$
\tilde{\nabla}_{X_{i}} \tilde{\nabla}_{X_{i}} \vec{H}=X_{i} X_{i}(H) \vec{\xi}-2 X_{i}(H) S X_{i}-H\left(\nabla_{X_{i}} S\right) X_{i}-H h\left(S X_{i}, X_{i}\right) \vec{\xi}
$$

and summing with respect to $i$, we obtain that

$$
\Delta \vec{H}=\{2 S(\nabla H)+H \operatorname{tr} \nabla S\}+\left\{\Delta H+\varepsilon H t r S^{2}\right\} \vec{\xi}
$$

We need to find an expression for $\operatorname{tr} \nabla S$. If $\left\{e_{i}\right\}, i=1,2,3$ be an orthonormal basis of eigenvectors of the shape operator $S$ such that $S e_{i}=\lambda_{i} e_{i}$, then

$$
\begin{aligned}
\operatorname{tr} \nabla S= & \sum_{i=1}^{3} \epsilon_{i}\left(\nabla_{e_{i}} S\right) e_{i} \\
= & {\left[\epsilon_{1} e_{1}\left(\lambda_{1}\right)+\epsilon_{2}\left(\lambda_{2}-\lambda_{1}\right) \omega_{22}^{1}+\epsilon_{3}\left(\lambda_{3}-\lambda_{1}\right) \omega_{33}^{1}\right] e_{1} } \\
& +\left[\epsilon_{2} e_{2}\left(\lambda_{2}\right)+\epsilon_{1}\left(\lambda_{1}-\lambda_{2}\right) \omega_{11}^{2}+\epsilon_{3}\left(\lambda_{3}-\lambda_{2}\right) \omega_{33}^{2}\right] e_{2} \\
& +\left[\epsilon_{3} e_{3}\left(\lambda_{3}\right)+\epsilon_{1}\left(\lambda_{1}-\lambda_{3}\right) \omega_{11}^{3}+\epsilon_{2}\left(\lambda_{2}-\lambda_{3}\right) \omega_{22}^{3}\right] e_{3},
\end{aligned}
$$

where $\epsilon_{i}=\left\langle e_{i}, e_{i}\right\rangle= \pm 1$. From the Codazzi equation $\left(\nabla_{e_{1}} S\right) e_{2}=\left(\nabla_{e_{2}} S\right) e_{1}$ it follows that

$$
\epsilon_{i}\left(\lambda_{i}-\lambda_{j}\right) \omega_{i i}^{j}=\epsilon_{j} e_{j}\left(\lambda_{i}\right)
$$

with $i, j=1,2,3$. Therefore,

$$
\begin{aligned}
\operatorname{tr} \nabla S= & \epsilon_{1} e_{1}\left(\lambda_{1}+\lambda_{2}+\lambda_{3}\right) e_{1}+\epsilon_{2} e_{2}\left(\lambda_{1}+\lambda_{2}+\lambda_{3}\right) e_{2} \\
& +\epsilon_{3} e_{3}\left(\lambda_{1}+\lambda_{2}+\lambda_{3}\right) e_{3}=3 \varepsilon \nabla H,
\end{aligned}
$$

and this completes the proof.
Due to Proposition 1 condition (8) is equivalent to

$$
\begin{equation*}
\{2 S(\nabla H)+3 \varepsilon H(\nabla H)\}+\left\{\Delta H+\varepsilon H \operatorname{tr} S^{2}\right\} \vec{\xi}=\lambda H \vec{\xi} . \tag{9}
\end{equation*}
$$

Therefore we obtain the following necessary and sufficient conditions for a hypersurface $M_{r}^{3}$ of $E_{s}^{4}$ to satisfy $\Delta \vec{H}=\lambda \vec{H}$ :

$$
\begin{align*}
& S(\nabla H)=-\varepsilon \frac{3 H}{2}(\nabla H)  \tag{10}\\
& \Delta H+\varepsilon H \operatorname{tr} S^{2}=\lambda H \tag{11}
\end{align*}
$$

where the Laplace operator $\Delta$ acting on a scalar-valued function $f$ is given by (e.g. [Ch-Is1])

$$
\begin{equation*}
\Delta f=-\sum_{i=1}^{3} \epsilon_{i}\left(e_{i} e_{i} f-\nabla_{e_{i}} e_{i} f\right) \tag{12}
\end{equation*}
$$

Here $\left\{e_{1}, e_{2}, e_{3}\right\}$ is a local orthonormal frame of $T_{p} M_{r}^{3}$ with $\left\langle e_{i}, e_{i}\right\rangle=\epsilon_{i}= \pm 1$. Noting from equation (10) that $\nabla H$ is an eigenvector of the shape operator $S$, without loss of generality we can choose $e_{1}$ in the direction of $\nabla H$, and therefore the shape operator of $M_{r}^{3}$ takes the form

$$
S=\left(\begin{array}{ccc}
-\varepsilon \frac{3 H}{2} & &  \tag{13}\\
& \lambda_{2} & \\
& & \lambda_{3}
\end{array}\right)
$$

In the special case in which all principal curvatures are equal, using the relation $\operatorname{tr} S=3 \varepsilon H$, it follows immediately that $H=0$. Therefore it suffices to examine the case where all principal curvatures are different, and the case when the two principal curvatures are equal.

## 3. Proof of the theorem.

## Three mutually different principal curvatures.

We need to show that if $\Delta \vec{H}=\lambda \vec{H}$, and the shape operator has three mutually different principal curvatures, then $H$ is constant. Suppose on the contrary that $M_{r}^{3}$ $(r=0,1,2,3)$ does not have constant mean curvature $H$. Then $\nabla H \neq \overrightarrow{0}$, and (10) shows that $\nabla H$ is an eigenvector of $S$ with corresponding eigenvalue $\lambda_{1}=-\frac{3 \varepsilon H}{2}$. Expressing
$\nabla H$ as $\nabla H=e_{1}(H) e_{1}+e_{2}(H) e_{2}+e_{3}(H) e_{3}$, and since $e_{1}$ is in the direction of $\nabla H$ it follows that

$$
\begin{equation*}
e_{1}(H) \neq 0, e_{2}(H)=e_{3}(H)=0 \tag{14}
\end{equation*}
$$

By assumption, we have that $\varepsilon \frac{3 H}{2}+\lambda_{2} \neq 0, \varepsilon \frac{3 H}{2}+\lambda_{3} \neq 0, \lambda_{3}-\lambda_{2} \neq 0$.
We write $\nabla_{e_{i}} e_{j}=\sum_{k=1}^{3} \omega_{i j}^{k} e_{k}$, we take into account the action of $S$ on the basis $\left\{e_{1}, e_{2}, e_{3}\right\}$, and use the Codazzi equation (6). Then the relations

$$
\begin{array}{ll}
\left\langle\left(\nabla_{e_{1}} S\right) e_{2}, e_{1}\right\rangle=\left\langle\left(\nabla_{e_{2}} S\right) e_{1}, e_{1}\right\rangle & \left\langle\left(\nabla_{e_{2}} S\right) e_{3}, e_{3}\right\rangle=\left\langle\left(\nabla_{e_{3}} S\right) e_{2}, e_{3}\right\rangle \\
\left\langle\left(\nabla_{e_{1}} S e_{3}, e_{3}\right\rangle=\left\langle\left(\nabla_{e_{3}} S\right) e_{1}, e_{3}\right\rangle\right. & \left\langle\left(\nabla_{e_{2}} S\right) e_{3}, e_{2}\right\rangle=\left\langle\left(\nabla_{e_{3}} S\right) e_{2}, e_{2}\right\rangle \\
\left\langle\left(\nabla_{e_{1}} S\right) e_{2}, e_{2}\right\rangle=\left\langle\left(\nabla_{e_{2}} S\right) e_{1}, e_{2}\right\rangle & \left\langle\left(\nabla_{e_{1}} S\right) e_{3}, e_{3}\right\rangle=\left\langle\left(\nabla_{e_{3}} S\right) e_{1}, e_{3}\right\rangle
\end{array}
$$

imply that $\omega_{12}^{1}=\omega_{13}^{1}=0$, and that

$$
\begin{equation*}
\omega_{21}^{1}=\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}, \quad \omega_{31}^{3}=\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}, \quad \omega_{23}^{2}=\frac{e_{3}\left(\lambda_{2}\right)}{\lambda_{3}-\lambda_{2}}, \quad \omega_{32}^{3}=\frac{e_{2}\left(\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}} . \tag{15}
\end{equation*}
$$

Also, in view of (14) it follows that $\nabla_{e_{2}} e_{3}(H)-\nabla_{e_{3}} e_{2}(H)=\left[e_{2}, e_{3}\right](H)=0$. Thus, together with the Codazzi equations for $\left\langle\left(\nabla_{e_{1}} S\right) e_{2}, e_{3}\right\rangle=\left\langle\left(\nabla_{e_{2}} S\right) e_{1}, e_{3}\right\rangle$, and $\left\langle\left(\nabla_{e_{1}} S\right) e_{3}, e_{2}\right\rangle=\left\langle\left(\nabla_{e_{3}} S\right) e_{1}, e_{2}\right\rangle$ we obtain that

$$
\begin{equation*}
\omega_{13}^{2}=\omega_{21}^{3}=\omega_{32}^{1}=0 . \tag{16}
\end{equation*}
$$

We use Gauss equation (7) and the definition of the curvature tensor for $\left\langle R\left(e_{2}, e_{3}\right) e_{1}, e_{2}\right\rangle$, $\left\langle R\left(e_{2}, e_{3}\right) e_{1}, e_{3}\right\rangle$, and $\left\langle R\left(e_{3}, e_{1}\right) e_{2}, e_{3}\right\rangle$, to obtain that

$$
\begin{align*}
e_{3}\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right) & =\frac{e_{3}\left(\lambda_{2}\right)}{\lambda_{3}-\lambda_{2}}\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}-\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right)  \tag{17}\\
e_{2}\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right) & =\frac{e_{2}\left(\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}}\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}-\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right)  \tag{18}\\
e_{1}\left(\frac{e_{2}\left(\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}}\right) & =-\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}} \frac{e_{2}\left(\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}} . \tag{19}
\end{align*}
$$

If we combine relations (12) and (14), then equation (11) takes the form

$$
\begin{equation*}
\epsilon_{1} e_{1} e_{1}(H)+\epsilon_{1}\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}+\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right) e_{1}(H)-\varepsilon H\left(\frac{45 H^{2}}{2}-2 \lambda_{2} \lambda_{3}\right)+\lambda H=0 . \tag{20}
\end{equation*}
$$

Acting on (20) with $e_{2}$ and $e_{3}$ succesively, and combining the results with (17), (18) it follows that

$$
\begin{align*}
e_{2}\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right)= & -\frac{e_{2}\left(\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}}\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}-\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right) \\
& -\frac{2 \varepsilon H}{\epsilon_{1} e_{1}(H)}\left(\lambda_{2}-\lambda_{3}\right)^{2} \frac{e_{2}\left(\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}}  \tag{21}\\
e_{3}\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)= & -\frac{e_{3}\left(\lambda_{2}\right)}{\lambda_{3}-\lambda_{2}}\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}-\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right) \\
& -\frac{2 \varepsilon H}{\epsilon_{1} e_{1}(H)}\left(\lambda_{3}-\lambda_{2}\right)^{2} \frac{e_{3}\left(\lambda_{2}\right)}{\lambda_{3}-\lambda_{2}} . \tag{22}
\end{align*}
$$

Similarly, using Gauss equation for $\left\langle R\left(e_{1}, e_{2}\right) e_{1}, e_{2}\right\rangle$ and $\left\langle R\left(e_{3}, e_{1}\right) e_{1}, e_{3}\right\rangle$ we obtain that

$$
\begin{align*}
& e_{1}\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right)+\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right)^{2}=\epsilon_{1} \varepsilon \frac{3}{2} H \lambda_{2}  \tag{23}\\
& e_{2}\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)+\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)^{2}=\epsilon_{1} \varepsilon \frac{3}{2} H \lambda_{3} . \tag{24}
\end{align*}
$$

We will need the following:
Lemma 2. Let $M_{r}^{3}$ be a hypersurface of the pseudo-Euclidean space $E_{s}^{4}$ whose shape operator has the form (13), and three mutually different principal curvatures. Then $e_{2}\left(\lambda_{3}\right)=e_{3}\left(\lambda_{2}\right)=0$

Proof. Relations (15) and (16) imply that

$$
\begin{equation*}
\left[e_{1}, e_{2}\right]=\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}} e_{2} \tag{25}
\end{equation*}
$$

Applying both sides of (25) on $\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}$, and using (21), (23), (24), and (19) we deduce that

$$
\begin{aligned}
& {\left[\frac{\varepsilon H}{\epsilon_{1} e_{1}(H)}\left(\left(3 \frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}-\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)\left(\lambda_{2}-\lambda_{3}\right)^{2}+2\left(\lambda_{2}-\lambda_{3}\right) e_{1}\left(\lambda_{2}-\lambda_{3}\right)\right)\right.} \\
& \left.\quad+\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}-\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)^{2}+\frac{\varepsilon}{\epsilon_{1}} e_{1}\left(\frac{H}{e_{1}(H)}\right)\left(\lambda_{2}-\lambda_{3}\right)^{2}\right] \frac{e_{2}\left(\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}}=0 .
\end{aligned}
$$

We will show that if $e_{2}\left(\lambda_{3}\right) \neq 0$ we get a contradiction. Indeed, in that case we would have that

$$
\begin{aligned}
e_{1}\left(\frac{H}{e_{1}(H)}\right)= & -\frac{H}{e_{1}(H)}\left(\left(3 \frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}-\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)+2 \frac{e_{1}\left(\lambda_{2}-\lambda_{3}\right)}{\lambda_{2}-\lambda_{3}}\right) \\
& -\frac{\epsilon_{1}}{\varepsilon\left(\lambda_{2}-\lambda_{3}\right)^{2}}\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}-\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)^{2} .
\end{aligned}
$$

Acting with $e_{2}$ on both sides of the above equation, and in view of (18), (21), (25), we obtain that

$$
\begin{equation*}
2\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}-\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)=-\frac{\varepsilon H}{\epsilon_{1} e_{1}(H)}\left(\lambda_{2}-\lambda_{3}\right)^{2} . \tag{26}
\end{equation*}
$$

We apply $e_{2}$ on (26) and obtain

$$
\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}-\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)=-2 \frac{\varepsilon H}{\epsilon_{1} e_{1}(H)}\left(\lambda_{2}-\lambda_{3}\right)^{2} .
$$

From the above two equations it follows that $\lambda_{2}=\lambda_{3}$, which is a contradiction. Hence, we conclude that $e_{2}\left(\lambda_{3}\right)=0$. In an analogous manner, it can be shown that $e_{3}\left(\lambda_{2}\right)=0$.

Coming back to the proof of the Theorem, we use Lemma 2 and Gauss equation for $\left\langle R\left(e_{2}, e_{3}\right) e_{2}, e_{3}\right\rangle$ to obtain that

$$
\begin{equation*}
-\epsilon_{1}\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}\right)\left(\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)-\lambda_{2} \lambda_{3}=0 . \tag{27}
\end{equation*}
$$

Calculating $e_{1} e_{1}(H)$ from (23) and (24), and combining with (20) and (27) it follows that

$$
\begin{align*}
\left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}+\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right) e_{1}(H) & =-\frac{54+135 \varepsilon}{8 \epsilon_{1}} H^{3}+\frac{6+3 \varepsilon}{2 \epsilon_{1}} H \lambda_{2} \lambda_{3}+\frac{3}{4 \epsilon_{1}} \lambda H  \tag{28}\\
e_{1} e_{1}(H) & =\frac{54+315 \varepsilon}{8 \epsilon_{1}} H^{3}-\frac{6+7 \varepsilon}{2 \epsilon_{1}} H \lambda_{2} \lambda_{3}-\frac{7}{4 \epsilon_{1}} \lambda H \tag{29}
\end{align*}
$$

Acting with $e_{1}$ on both sides of (28) and using (23), (24), and (27) we deduce the expression

$$
\begin{align*}
& \left(\frac{e_{1}\left(\lambda_{2}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{2}}+\frac{e_{1}\left(\lambda_{3}\right)}{-\varepsilon \frac{3 H}{2}-\lambda_{3}}\right)\left(441 H^{2}-26 \lambda_{2} \lambda_{3}+10 \lambda\right) H \\
& \quad=\left(432 H^{2}-26 \lambda_{2} \lambda_{3}-3 \lambda\right) e_{1}(H) . \tag{30}
\end{align*}
$$

If we apply $e_{1}$ on (30) and use (28), (29), (30) we obtain the following algebraic relation between $H$ and $\lambda_{2} \lambda_{3}$ (notice that $H$ and $\lambda_{2} \lambda_{3}$ are real functions in general):

$$
\begin{align*}
0= & a_{80} H^{8}+a_{60} H^{6}+a_{40} H^{4}+a_{20} H^{2}+a_{00} \\
& +a_{61} H^{6}\left(\lambda_{2} \lambda_{3}\right)+a_{41} H^{4}\left(\lambda_{2} \lambda_{3}\right)+a_{21} H^{2}\left(\lambda_{2} \lambda_{3}\right)+a_{01}\left(\lambda_{2} \lambda_{3}\right)+a_{42} H^{4}\left(\lambda_{2} \lambda_{3}\right)^{2} \\
& +a_{22} H^{2}\left(\lambda_{2} \lambda_{3}\right)^{2}+a_{02}\left(\lambda_{2} \lambda_{3}\right)^{2}+a_{23} H^{2}\left(\lambda_{2} \lambda_{3}\right)^{3}+a_{03}\left(\lambda_{2} \lambda_{3}\right)^{3}+a_{04}\left(\lambda_{2} \lambda_{3}\right)^{4} \\
= & f\left(H, \lambda_{2} \lambda_{3}\right), \tag{31}
\end{align*}
$$

where $a_{i j}$ are known constants. Acting in (31) with $e_{1}$ twice, and using (28), (29), (30), we obtain another algebraic relation of $H$ and $\lambda_{2} \lambda_{3}$ of the form

$$
\begin{equation*}
g\left(H, \lambda_{2} \lambda_{3}\right)=0 \tag{32}
\end{equation*}
$$

Using a computer algebra program, we eliminate $\lambda_{2} \lambda_{3}$ between (31) and (32), so obtain an algebraic equation for $H$ with constant coefficients. Thus, we have concluded that the real function $H$ satisfies a polynomial equation $q(H)=0$ with constant coefficients, therefore it must be a constant. This contradicts our original assumption, so the Theorem is proved in this case.

## Two equal principal curvatures.

We need to show that if $\Delta \vec{H}=\lambda \vec{H}$, and the shape operator has two equal principal curvatures, then $H$ is constant. Assume the contrary, and try to get a contradiction. As in the previous case, $e_{1}$ can be chosen in the direction of $\nabla H$, yielding $\lambda_{1}=-\frac{3 \varepsilon H}{2}$ and

$$
e_{1}(H) \neq 0, e_{2}(H)=e_{3}(H)=0
$$

Then the shape operator of $M_{r}^{3}$ takes the form

$$
S=\left(\begin{array}{ccc}
-\varepsilon \frac{3 H}{2} & & \\
& \mu & \\
& & \mu
\end{array}\right)
$$

for some function $\mu$. From $\operatorname{tr} S=3 \varepsilon H$ if follows that $\mu=\varepsilon \frac{9 H}{4}$, and that $\operatorname{tr} S^{2}=\frac{99 H^{2}}{8}$. Applying the Codazzi equation (6) it follows that $\left\langle\left(\nabla_{e_{1}} S\right) e_{2}, e_{2}\right\rangle=\left\langle\left(\nabla_{e_{2}} S\right) e_{1}, e_{2}\right\rangle$ and that $\left\langle\left(\nabla_{e_{1}} S\right) e_{3}, e_{3}\right\rangle=\left\langle\left(\nabla_{e_{3}} S\right) e_{1}, e_{3}\right\rangle$, which in turn give that

$$
\begin{equation*}
\omega_{21}^{2}=\omega_{31}^{3}=-\frac{3}{5} \frac{e_{1}(H)}{H} \tag{33}
\end{equation*}
$$

The Gauss equation (7) applied to $\left\langle R\left(e_{1}, e_{2}\right) e_{1}, e_{2}\right\rangle$ implies that

$$
\begin{equation*}
e_{1}\left(\omega_{21}^{2}\right)=\epsilon_{1} \frac{27 H^{2}}{8}-\left(\omega_{21}^{2}\right)^{2} \tag{34}
\end{equation*}
$$

Equation (11) then reduces to

$$
\begin{equation*}
-\epsilon_{1} e_{1} e_{1}(H)-2 \epsilon_{1}\left(\omega_{21}^{2}\right)^{2} e_{1}(H)+\varepsilon \frac{99 H^{3}}{8}=\lambda H \tag{35}
\end{equation*}
$$

We act on (33) with $e_{1}$ and use (34) to obtain that

$$
e_{1} e_{1}(H)=\frac{40}{9} H\left(\omega_{21}^{2}\right)^{2}-\frac{45 H^{3}}{8} \epsilon_{1} .
$$

We substitute the above equation to (35) and get

$$
H\left[\epsilon_{1} \frac{10}{9}\left(\omega_{21}^{2}\right)^{2}+\lambda-\frac{45+99 \varepsilon}{8} H^{2}\right]=0
$$

and as $H \neq 0$ it follows that

$$
\epsilon_{1} \frac{10}{9}\left(\omega_{21}^{2}\right)^{2}+\lambda-\frac{45+99 \varepsilon}{8} H^{2}=0
$$

Acting by $e_{1}$ in the above equation and using (33) and (34), it follows that

$$
\epsilon_{1} \frac{10}{3}\left(\omega_{21}^{2}\right)^{2}-\frac{945+165 \varepsilon}{8} H^{2}=0
$$

If we eliminate the $\left(\omega_{21}^{2}\right)^{2}$ from the last two equations we obtain the relation

$$
\lambda+\frac{810-132 \varepsilon}{24} H^{2}=0
$$

that is $H$ is constant, which contradicts our assumption.

## 4. Relation with biharmonic maps.

In this section we describe the relation of equation (3) to the theory of harmonic and biharmonic maps. For relative background we refer to $[\mathbf{C a}-\mathbf{M o}-\mathbf{O n}]$, $[\mathbf{E e}-\mathrm{Le}]$, and $[\mathbf{U r}]$. Let $\left(M^{m}, g\right)$ and $\left(N^{n}, h\right)$ be Riemannian manifolds. A smooth map $\phi: M \rightarrow N$ is said to be harmonic if it is a critical point of the energy functional:

$$
E_{1}(\phi)=\frac{1}{2} \int_{M}|d \phi|^{2} d v_{g}
$$

Denote by $\nabla^{\phi}$ the connection of the vector bundle $\phi^{*} T N$ induced from the Levi-Civita connection $\nabla^{h}$ of $(N, h)$. The second fundamental form $\nabla d \phi$ is defined by

$$
\nabla d \phi(X, Y)=\nabla_{X}^{\phi} d \phi(Y)-d \phi\left(\nabla_{X} Y\right), \quad X, Y \in \Gamma(T M)
$$

where $\nabla$ is the Levi-Civita connection of $(M, g)$. The tension field $\tau(\phi)$ is a section of $\phi^{*} T N$ defined by

$$
\tau(\phi)=\operatorname{trace}(\nabla d \phi)
$$

It is well known that the map $\phi$ is harmonic if and only if its tension field vanishes. Now assume that $\phi: M \rightarrow N$ is an isometric immersion with mean curvature vector field $\vec{H}$. Then $m \vec{H}=\tau(\phi)$ (cf. [Ee-Sa, p. 119]), therefore the immersion $\phi$ is a harmonic map if and only if $M$ is a minimal submanifold of $N$.

A smooth map $\phi:(M, g) \rightarrow(N, h)$ is called biharmonic if it is a critical point of the bienergy functional:

$$
E_{2}(\phi)=\frac{1}{2} \int_{M}|\tau(\phi)|^{2} d v_{g} .
$$

This is a special case of a more general set-up suggested in [Ee-Sa] on studying polyharmonic maps. In [Ji1] and [Ji2] G. Y. Jiang derived the first variation formula of the bienergy showing that the Euler-Lagrange equation for $E_{2}$ is given by

$$
\tau_{2}(\phi)=-J_{\phi}(\tau(\phi))=0
$$

Here $J_{\phi}$ is the Jacobi operator of $\phi$ acting on sections $V \in \Gamma\left(\phi^{*} T N\right)$. It is defined by

$$
\begin{gathered}
J_{\phi}(V)=\bar{\Delta}_{\phi} V-R_{\phi}(V), \\
\bar{\Delta}_{\phi}=-\sum_{i=1}^{m}\left(\nabla_{e_{i}}^{\phi} \nabla_{e_{i}}^{\phi}-\nabla_{\nabla_{e_{i}} e_{i}}^{\phi}\right), \quad R_{\phi}(V)=\sum_{i=1}^{m} R^{N}\left(V, d \phi\left(e_{i}\right)\right) d \phi\left(e_{i}\right),
\end{gathered}
$$

where $R^{N}$ is the curvature tensor of $N$, and $\left\{e_{i}\right\}$ a local orthonormal frame field of $M$. If $x:\left(M^{m}, g\right) \rightarrow\left(E^{n}\right.$, canonical) is an isometric immersion, then $\bar{\Delta}_{x}$ is simply the Laplace operator $\Delta$ of $M$ with respect to the induced metric, thus

$$
\tau_{2}(x)=\Delta \tau(x)=\Delta(m \vec{H})=m \Delta \vec{H}
$$

Therefore, $M^{m}$ is a biharmonic submanifold of the Euclidean space $E^{n}$ with the canonical metric if and only if the immersion $x: M^{m} \rightarrow E^{n}$ is a biharmonic map. Finally, an isometric immersion $x: M \rightarrow N$ is called $\lambda$-biharmonic if it is a critical point of the functional

$$
E_{2, \lambda}(x)=E_{2}(x)+\lambda E(x), \quad \lambda \in \boldsymbol{R} .
$$

The Euler-Lagrange equation for $\lambda$-biharmonic immersions is

$$
\tau_{2}(x)=\lambda \tau(x)
$$

This is equivalent to the equation $\Delta \vec{H}=\lambda \vec{H}$, i.e. the submanifold $M$ has proper mean curvature vector field.

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