# Stationary 2-type surfaces in a hypersphere 

In Memory of Professor Yozô Matsushima

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## 1. Introduction.

In terms of finite-type submanifolds [3], a well-known theorem of Takahashi [9] says that an $n$-dimensional, compact submanifold $M$ of $E^{m+1}$ is of 1-type if and only if $M$ is a minimal submanifold of a hypersphere $S^{m}$ of $E^{m+1}$. Such a submanifold is always mass-symmetric in $S^{m}$, i. e., the center of mass of $M$ is the center of $S^{m}$ in $E^{m+1}$. Thus, if one chooses the center of $S^{m}$ as the origin of $E^{m+1}$, then the position vector $x$ of $M$ has the following form:

$$
\begin{equation*}
x=x_{p}, \quad \Delta x_{p}=\lambda_{p} x_{p}, \tag{1.1}
\end{equation*}
$$

where $\lambda_{p}=n / r^{2}, r$ is the radius of $S^{m}$ and $\Delta$ is the Laplacian of $M$. Submanifolds of $S^{m}$ satisfying (1.1) are the simplest finite-type submanifolds. The study of such submanifolds has attracted many mathematicians for many years.

On the other hand, it was shown in [4] (see, also [3, p. 274]) that if $M$ is a compact hypersurface of $S^{m}$ such that $M$ is not a small hypersphere; then $M$ has constant mean curvature $\alpha^{\prime} \neq 0$ and constant scalar curvature $\tau$ if and only if $M$ is mass-symmetric and of 2-type. In this case, the position vector $x$ of $M$ in $E^{m+1}$ has the following form:

$$
\begin{equation*}
x=x_{p}+x_{q}, \quad \Delta x_{p}=\lambda_{p} x_{p} \quad \text { and } \quad \Delta x_{q}=\lambda_{q} x_{q} . \tag{1.2}
\end{equation*}
$$

Furthermore, $\alpha^{\prime}$ and $\tau$ are completely determined by $\left\{\lambda_{p}, \lambda_{q}\right\}$. Applying this result, we see that all non-minimal, isoparametric hypersurfaces of $S^{m}$ are masssymmetric and of 2-type.

Mass-symmetric, 2-type submanifolds of $S^{m}$ are the "simplest" submanifolds of $E^{m+1}$ next to minimal submanifolds of $S^{m}$. Many important submanifolds are known to be of 2 -type and are mass-symmetric (cf. [1, 3, 4, 7, 8]). For instance, it was shown in [8] that any compact, non-totally geodesic, parallel, Einstein, complex submanifold of complex projective space $\boldsymbol{C} P^{N}$ is of 2-type if

[^0]we regard $\boldsymbol{C} P^{N}$ as a submanifold of a Euclidean space by its first standard imbedding. The complete classification of mass-symmetric, 2 -type submanifolds of $S^{m}$ is formidably difficult. However, the case of surfaces in $S^{3}$ was done by the second author (cf. [3, p. 279]). In Section 4, we will solve this problem for surfaces in $S^{4}$.

Given an isometric immersion $f: M \rightarrow M^{\prime}$ of a surface $M$ into a Riemannian manifold $M^{\prime}$, one has the conformal total mean curvature $\tau(f)$ (cf. Section 5). Surfaces which are critical points of $\tau(f)$ are called stationary. Related to the Chen-Willmore problem, Weiner asked in [10] whether minimal surfaces of $S^{m}$ are the only stationary, mass-symmetric surfaces of $S^{m}$ ? N. Ejiri constructed in [5] a counter-example to Weiner's problem. It is easy to see that Ejiri's example is of 2-type.

In this paper, we will study stationary, mass-symmetric, 2-type surfaces of $S^{m}$ in detail. In particular, we will prove that such surfaces are in fact flat surfaces which lie fully in $S^{5}$ or in $S^{7}$. By completely determining the connection form of such surfaces, we show that such surfaces are obtained by some doubly-periodic isometric immersions of the Euclidean plane $\boldsymbol{R}^{2}$ into $S^{5}$ or $S^{7}$. In the case of $S^{5}$, a surprising phenomenon occurs. The connection form depends only on the eigenvalue $\lambda_{p}$ which satisfies $2 / 3<\lambda_{p}<2$. Furthermore, for each $\lambda_{p} \in(2 / 3,4 / 3]$, there is only one possibility for the connection form, while for each $\lambda_{p} \in(4 / 3,2)$ there are two possibilities. Moreover, for each such connection form, we can construct a stationary, mass-symmetric, 2 -type, flat torus in $S^{5}$. Although such a torus is not unique, it comes from a "unique" doublyperiodic immersion of $\boldsymbol{R}^{2}$ into $S^{5}$. We also show that the estimate on $\lambda_{p}$ is best possible.

In the case of $S^{7}$, the connection form depends on both $\lambda_{p}$ and $\lambda_{q}$ (and depends only on them). Such $\lambda_{p}$ and $\lambda_{q}$ must satisfy $0<\lambda_{p}<2<\lambda_{q}<\infty$. Furthermore, we can give some concrete examples for this case. More precisely, for any real number $d \in(2, \infty)$, there are a real number $c \in(0,2)$ and a stationary, mass-symmetric, 2-type flat torus in $S^{7}$ with $\left(\lambda_{p}, \lambda_{q}\right)=(c, d)$. For each such pair ( $c, d$ ), the flat torus in $S^{7}$ is obtained from a "unique" doubly-periodic immersion of $\boldsymbol{R}^{2}$ into $S^{7}$.

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## 2. Preliminaries.

Let $M$ be a compact Riemannian manifold and $\Delta$ the Laplacian of $M$ acting on differentiable functions in $C^{\infty}(M)$. Then $\Delta$ is an elliptic differential operator and it has an infinite sequence of eigenvalues: $0=\lambda_{0}<\lambda_{1}<\lambda_{2}<\cdots<\lambda_{k}<\cdots \uparrow \infty$. Let $V_{k}=\left\{f \in C^{\infty}(M) \mid \Delta f=\lambda_{k} f\right\}$ be the eigenspace of $\Delta$ with eigenvalue $\lambda_{k}$. Then each $V_{k}$ is finite-dimensional. If we define an inner product on $C^{\infty}(M)$ by $(f, g)=\int f g d V$, then the decomposition $\sum_{k \geq 0} V_{k}$ is orthogonal and dense in $C^{\infty}(M)$ (in $L^{2}$-sense). For each $f \in C^{\infty}(M)$, let $f_{t}$ be the projection of $f$ into $V_{t}$. Then we have the decomposition: $f=\sum_{t \geq 0} f_{t}$ (in $L^{2}$-sense).

For an isometric immersion $x: M \rightarrow E^{m+1}$ of a compact Riemannian manifold $M$ into the Euclidean ( $m+1$ )-space $E^{m+1}$, we put $x=\left(x_{1}, \cdots, x_{m+1}\right)$, where $x_{A}$ is the $A$-th Euclidean coordinate function of $M$. Thus, we may write $x=\Sigma_{t z 0} x_{t}$ (in $L^{2}$-sense) so that $\Delta x_{t}=\lambda_{t} x_{t}$ for each $t$. Since $M$ is compact, $x_{0}$ is a constant vector in $E^{m+1}$ and, moreover, there is a natural number $p$ such that $x_{p} \neq 0$ and $x=x_{0}+\sum_{t z p} x_{t}$. If there are infinitely many $x_{t}$ 's which are nonzero, we put $q=\infty$. Otherwise, there is an integer $q \geqq p$ such that

$$
x=x_{0}+\sum_{t=p}^{q} x_{t}, \quad x_{q} \neq 0 .
$$

In both cases, we have the following decomposition :

$$
\begin{equation*}
x=x_{0}+\sum_{t=p}^{q} x_{t} \quad\left(\text { in } L^{2} \text {-sense }\right) . \tag{2.1}
\end{equation*}
$$

The submanifold $M$ is said to be of finite type if $q$ is finite. Otherwise, $M$ is said to be of infinite type. The submanifold $M$ is said to be of $k$-type if there is exactly $k$ nonzero $x_{t}$ 's in the decomposition (2.1). The pair $[p, q]$ is called the order of the submanifold $M$ [3].

Let $M$ be an $n$-dimensional submanifold of an $m$-dimensional Riemannian manifold $M^{\prime}$. We choose a local field of orthonormal frames $e_{1}, \cdots, e_{n}, \xi_{n+1}$, $\cdots, \xi_{m}$ in $M^{\prime}$ such that $e_{1}, \cdots, e_{n}$ are tangent to $M$. Let $\omega^{1}, \cdots, \omega^{n}$ be the dual frame of $e_{1}, \cdots, e_{n}$. Denote by $\nabla^{\prime}$ the Riemannian connection of $M^{\prime}$. We put $\nabla^{\prime} e_{i}=\Sigma \omega_{i}{ }^{j} e_{j}+\Sigma \omega_{i}{ }^{r} \xi_{r}$ and $\nabla^{\prime} \xi_{r}=\Sigma \omega_{r}{ }^{i} \xi_{i}+\Sigma \omega_{r}{ }^{t} \xi_{t}, \quad i, j, k=1, \cdots, n ; r, s, t=n+1$, $\cdots, m$. By Cartan's Lemma, we have

$$
\begin{equation*}
\omega_{i}^{r}=\Sigma h_{i j}{ }^{r} \omega^{j}, \quad h_{i j}^{r}=h_{j i}^{r}, \tag{2.2}
\end{equation*}
$$

where $h_{i j}{ }^{r}$ are coefficients of the second fundamental form. The connection form of $M$ in $M^{\prime}$ is given by ( $\omega_{A}{ }^{B}$ ), $A, B=1, \cdots, m$.

Throughout this paper, we shall assume that the submanifold $M$ is compact unless mentioned otherwise.

## 3. 2-type submanifolds of hyperspheres.

Let $x: M \rightarrow E^{m+1}$ be an isometric immersion of a compact, $n$-dimensional Riemannian manifold $M$ into $E^{m+1}$. Denote by $\nabla$ and $\tilde{\nabla}$ the Riemannian connections of $M$ and $E^{m+1}$, respectively. And by $h, A$ and $D$ the second fundamental form, the Weingarten map and the normal connection of $M$ in $E^{m+1}$, respectively.

For a fixed vector $a$ in $E^{m+1}$ and vector fields $X, Y$ tangent to $M$, the formulas of Gauss and Weingarten give

$$
\begin{equation*}
Y X\langle H, a\rangle=\left\langle D_{Y} D_{X} H, a\right\rangle-\left\langle\nabla_{Y}\left(A_{H} X\right), a\right\rangle-\left\langle A_{D_{X} H} Y, a\right\rangle-\left\langle h\left(Y, A_{H} X\right), a\right\rangle \tag{3.1}
\end{equation*}
$$

where $H$ is the mean curvature vector of $M$ in $E^{m+1}$ and $\langle$,$\rangle the inner product$ of $E^{m+1}$. Let $e_{1}, \cdots, e_{n}$ be an orthonormal local frame field tangent to $M$. Equation (3.1) implies

$$
\begin{equation*}
\Delta H=\Delta^{D} H+\sum_{i=1}^{n}\left\{h\left(e_{i}, A_{H} e_{i}\right)+A_{D_{e_{i}} H_{e_{i}}}+\left(\nabla_{e_{i}} A_{H}\right) e_{i}\right\} \tag{3.2}
\end{equation*}
$$

where

$$
\begin{equation*}
\Delta^{D} H=\sum_{i=1}^{n}\left\{D_{\nabla_{e_{i}} i_{i}} H-D_{e_{i}} D_{e_{i}} H\right\} \tag{3.3}
\end{equation*}
$$

is the Laplacian of $H$ with respect to $D$. Regard $\nabla A_{H}$ and $A_{D H}$ as (1,2)-tensors on $M$ and we set $\bar{\nabla} A_{H}=\nabla A_{H}+A_{D H}$. Then we have

$$
\begin{equation*}
\operatorname{tr}\left(\bar{\nabla} A_{H}\right)=\sum_{i=1}^{n}\left\{\left(\nabla_{e_{i}} A_{H}\right) e_{i}+A_{D_{e_{i}} H_{e_{i}}}\right\} \tag{3.4}
\end{equation*}
$$

Let $U=\{u \in M \mid H \neq 0$ at $u\}$. Then $U$ is an open subset of $M$. On $U$ we choose an orthonormal local frame $\xi_{n+1}, \cdots, \xi_{m+1}$ normal to $M$ in $E^{m+1}$ so that $\xi_{n+1}$ is parallel to $H$. Then we have

$$
\begin{equation*}
\sum_{i=1}^{n} h\left(e_{i}, A_{H} e_{i}\right)=\left\|A_{n+1}\right\|^{2} H+\mathfrak{A}(H) \tag{3.5}
\end{equation*}
$$

where

$$
\begin{equation*}
A_{r}=A_{\xi_{r}}, \quad\left\|A_{n+1}\right\|^{2}=\operatorname{tr} A_{n+1}^{2} \tag{3.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathfrak{A}(H)=\sum_{r=n+2}^{m+1}\left(\operatorname{tr} A_{H} A_{r}\right) \xi_{r} \tag{3.7}
\end{equation*}
$$

on $U$. If $H=0$ at a point $u, \mathfrak{A}(H)$ is defined to be zero. It is clear that (3.5) and (3.6) hold trivially on $M-U$. Therefore, we have (3.5) and (3.7) on the whole submanifold $M$. The vector field $\mathfrak{A}(H)$ is a well-defined vector field perpendicular to $H$, which is called the allied mean curvature vector of $M$ in $E^{m+1}$. From (3.2), (3.4) and (3.5) we get

$$
\begin{equation*}
\Delta H=\Delta^{D} H+\left\|A_{n+1}\right\|^{2} H+\mathfrak{Y}(H)+\operatorname{tr}\left(\bar{\nabla} A_{H}\right) . \tag{3.8}
\end{equation*}
$$

Now, assume that $M$ is a submanifold of the unit hypersphere $S_{0}{ }^{m}(1)$ of $E^{m+1}$ centered at the origin 0 . Denote by $H^{\prime}, A^{\prime}$ and $D^{\prime}$ the mean curvature vector, the Weingarten map and the normal connection of $M$ in $S_{0}{ }^{m}(1)$, respectively. Then we have

$$
\begin{equation*}
H=H^{\prime}-x, \quad \Delta^{D} H=\Delta^{D^{\prime}} H^{\prime}, \quad D x=0 . \tag{3.9}
\end{equation*}
$$

Moreover, for any vector $\eta$ normal to $M$ in $S_{0}{ }^{m}(1)$, we have $A_{\eta}=A_{\eta}{ }^{\prime}$. Let $\xi$ be a unit vector parallel to $H^{\prime}$ with $H^{\prime}=\alpha^{\prime} \xi, \alpha^{\prime}=\left\|H^{\prime}\right\|$. (If $H^{\prime}=0$ at $u$, $\xi$ can be chosen to be an arbitrary unit normal vector of $M$ in $S_{0}{ }^{m}(1)$.) We have the following.

Lemma 1 ([3, p. 273]). Let $M$ be an $n$-dimensional submanifold of $S_{0}{ }^{m}(1)$ in $E^{m+1}$. Then we have

$$
\begin{equation*}
\Delta H=\Delta^{D^{\prime}} H^{\prime}+\mathfrak{A}^{\prime}\left(H^{\prime}\right)+\operatorname{tr}\left(\bar{\nabla} A_{H}\right)+\left(\left\|A_{\xi}\right\|^{2}+n\right) H^{\prime}-n(\alpha)^{2} x, \tag{3.10}
\end{equation*}
$$

where $\mathfrak{A}^{\prime}\left(H^{\prime}\right)$ is the allied mean curvature vector of $M$ in $S_{0}{ }^{m}(1)$ (which is zero on $\left\{u \in M \mid H^{\prime}=0\right.$ at $\left.u\right\}$ ).

We also need the following.
Lemma 2 ([3, p. 274]). If $M$ is a mass-symmetric, 2-type submanifold of $S_{0}{ }^{m}(1)$, then we have
(1) the mean curvature $\alpha^{\prime}$ is constant which is given by

$$
\begin{equation*}
\left(\alpha^{\prime}\right)^{2}=\left(1-\frac{\lambda_{p}}{n}\right)\left(\frac{\lambda_{q}}{n}-1\right) \neq 0 \tag{3.11}
\end{equation*}
$$

and
(2) $\operatorname{tr}\left(\bar{\nabla} A_{H}\right)=0,0<\lambda_{p}<n<\lambda_{q}<\infty$.

Lemma 3 ([4]). Let $M$ be an $n$-dimensional submanifold of $S_{0}{ }^{m}(1)$ in $E^{m+1}$. Then $\operatorname{tr}\left(\bar{\nabla} A_{H}\right)=0$ if and only if

$$
\begin{equation*}
n \operatorname{grad}\left(\alpha^{\prime}\right)^{2}+4 \operatorname{tr} A_{D H^{\prime}}=0 \tag{3.12}
\end{equation*}
$$

Lemma 2 implies that the mean curvature $\alpha^{\prime}=\left|H^{\prime}\right|$ of $M$ in $S_{0}{ }^{m}(1)$ is determined by the order. In the following, let $d H^{\prime}$ denote the $E^{m+1}$-valued 1-form defined by $\left(d H^{\prime}\right)(X)=\tilde{\nabla}_{X} H^{\prime}$, for $X$ tangent to $M$. Then we have $\left\|d H^{\prime}\right\|^{2}=$ $\left\|D^{\prime} H^{\prime}\right\|^{2}+\left\|A_{H^{\prime}}\right\|^{2}$. The following lemma shows that the length of $d H^{\prime}$ is also determined by the order of $M$.

Lemma 4. Let $M$ be a mass-symmetric, 2-type submanifold of $S_{0}{ }^{m}(1)$ in $E^{m+1}$. Then we have

$$
\begin{align*}
& \left\|d H^{\prime}\right\|^{2}=\left\{\lambda_{p}+\lambda_{q}-n\right\}\left\{n\left(\lambda_{p}+\lambda_{q}\right)-\lambda_{p} \lambda_{q}-n^{2}\right\} / n^{2},  \tag{3.13}\\
& \mathfrak{U}^{\prime}\left(H^{\prime}\right)=\left|H^{\prime}\right| \sum_{r=n+2}^{m}\left\{\operatorname{tr}\left(\nabla \omega_{n+1}^{r}\right)-\left\langle D^{\prime} \xi, D^{\prime} \xi_{r}\right\rangle\right\} \xi_{r} . \tag{3.14}
\end{align*}
$$

Proof. Let $\xi_{n+1}, \cdots, \xi_{m}$ be a local orthonormal normal basis of $M$ in $S_{0}{ }^{m}(1)$ such that $\xi_{n+1}=\xi$ is parallel to $H^{\prime}$ (this condition holds automatically on $\left\{u \in M \mid H^{\prime}=0\right.$ at $\left.u\right\}$ ). By using Lemma 2 and (3.3), we may find

$$
\begin{equation*}
\Delta^{D} H=\Delta^{D^{\prime}} H^{\prime}=\alpha^{\prime} \sum_{r=n+2}^{m}\left\{\left\langle D^{\prime} \xi, D^{\prime} \xi_{r}\right\rangle-\operatorname{tr}\left(\nabla \omega_{n+1}^{r}\right)\right\} \xi_{r}+\left\langle D^{\prime} \xi, D^{\prime} \xi\right\rangle H^{\prime}, \tag{3.15}
\end{equation*}
$$ where

$$
\begin{equation*}
\operatorname{tr}\left(\nabla \omega_{n+1}^{r}\right)=\sum_{i=1}^{n}\left(\nabla_{e_{i}} \omega_{n+1}^{r}\right)\left(e_{i}\right) . \tag{3.16}
\end{equation*}
$$

Since $M$ is mass-symmetric and of 2-type in $S_{0}{ }^{m}(1)$, we have (cf. [3, p. 256])

$$
\begin{equation*}
\Delta H=\left(\lambda_{p}+\lambda_{q}\right) H+\left(\lambda_{p} \lambda_{q} / n\right) x . \tag{3.17}
\end{equation*}
$$

Combining Lemma 1, Lemma 2, (3.15) and (3.17), we may obtain (3.13) and (3.14), (Q.E.D.)

Remark 1. By using Lemma 1, we may prove that there exist no masssymmetric, 3-type hypersurfaces with constant mean curvature in a hypersphere of $E^{n+2}$.

## 4. A non-existence theorem.

First, we mention the following [3, p. 279].
Theorem 1. Let $M$ be a mass-symmetric surface of $S_{0}{ }^{3}(1)$. Then $M$ is of 2-type if and only if $M$ is the product of two plane circles of different radii.

The Veronese surface in $S_{0}{ }^{4}(1)$ is a nice example of (mass-symmetric) 1-type surface which lies fully in $S_{0}{ }^{4}(1)$. In contrast with this, we give the following Non-existence Theorem.

Theorem 2. There exist no mass-symmetric, 2-type surfaces which lie fully in $S_{0}{ }^{4}(1)$.

Proof. Assume that $M$ is a mass-symmetric, 2-type surface which lies fully in $S_{0}{ }^{4}(1)$. Then we have $\left\langle D \xi, D \xi_{4}\right\rangle=0$. Thus, (3.14) reduces to

$$
\begin{equation*}
\mathfrak{U}^{\prime}\left(H^{\prime}\right)=\alpha^{\prime} \operatorname{tr}\left(\nabla \omega_{3}{ }^{4}\right) \xi_{4} . \tag{4.1}
\end{equation*}
$$

Combining (3.7) and (4.1), we obtain

$$
\begin{equation*}
\operatorname{tr}\left(A_{3} A_{4}\right)=\operatorname{tr}\left(\nabla \omega_{3}{ }^{4}\right) . \tag{4.2}
\end{equation*}
$$

On the other hand, by using constancy of $\alpha^{\prime}$ Lemma 2), Lemmas 2 and 3 imply $\operatorname{tr} A_{D \xi}=0$. Let $e_{1}, e_{2}$ be eigenvectors of $A_{4}$. Since $\operatorname{tr} A_{4}=0$, we may assume that $A_{4} e_{1}=\mu e_{1}$ and $A_{4} e_{2}=-\mu e_{2}$. Thus by using $\operatorname{tr} A_{D \xi}=0$, we find $\mu \omega_{3}{ }^{4}=0$, i. e., $A_{4} \omega_{3}{ }^{4}=0$. Combining this with (4.2), we obtain $\operatorname{tr}\left(A_{3} A_{4}\right)=0$.

Let $W=\left\{u \in M \mid A_{4} \neq 0\right.$ at $\left.u\right\}$. Assume that $W \neq \varnothing$ and $U$ is a connected component of $W$. Then $U$ is open and $D^{\prime} \xi_{3}=D^{\prime} \xi_{4}=0$. Let $e_{1}, e_{2}$ be an orthonormal tangent basis on $U$ such that, with respect to $e_{1}, e_{2}, A_{3}$ and $A_{4}$ are given by

$$
A_{3}=\left(\begin{array}{ll}
\beta & 0  \tag{4.3}\\
0 & \gamma
\end{array}\right), \quad A_{4}=\left(\begin{array}{rr}
c & b \\
b & -c
\end{array}\right) .
$$

Since $\operatorname{tr}\left(A_{3} A_{4}\right)=0$, (3.7) and (4.3) give $(\beta-\gamma) c=0$. On the other hand, Lemmas 2 and 4 imply $\left\|A_{3}\right\|^{2} \neq 2\left(\alpha^{\prime}\right)^{2}$. Thus $U$ is not pseudo-umbilical, i. e., $\beta \neq \gamma$. Consequently, we have $c=0$. Moreover, since $D^{\prime} \xi_{3}=D^{\prime} \xi_{4}=0$ on $U$, Ricci's equation gives $\left[A_{3}, A_{4}\right]=0$. Therefore, $b=0$ too. Hence, $W=\varnothing$. Thus, we have $A_{4}=0$ on $M$. This gives $\omega_{1}{ }^{4}=\omega_{2}{ }^{4}=0$. By taking exterior differentiation of these we obtain

$$
\begin{equation*}
\beta \omega^{1} \wedge \omega_{3}{ }^{4}=\gamma \omega^{2} \wedge \omega_{3}{ }^{4}=0 . \tag{4.4}
\end{equation*}
$$

Let $G$ denote the Gauss curvature of $M$. Then we have $G=1+\beta \gamma$. Let $V=$ $\{u \in M \mid G(u) \neq 1\}$. Then, on $V$, (4.4) implies $\omega_{3}{ }^{4}=0$, i. e., $D \xi_{3}=0$. Thus, Lemmas 2 and 4 imply that both $\beta$ and $\gamma$ are constant. Hence, by taking the exterior differentiation of $\omega_{1}{ }^{3}=\beta \omega^{1}$ and $\omega_{2}{ }^{3}=\gamma \omega^{2}$, we obtain $\omega_{1}{ }^{2}=0$. Thus, $G=0$. So, by the continuity of $G$ on $M$, we obtain $G \equiv 0$ or $G \equiv 1$. If $G \equiv 0$, then by $A_{4}=\omega_{3}{ }^{4}=0$, we conclude that $M$ is in fact a flat surface in a great hypersphere of $S_{0}{ }^{4}(1)$. This is a contradiction. Therefore, $G \equiv 1$ on $M$. Hence, $\beta \gamma=0$. Since $\beta+\gamma$ is constant, $\beta$ and $\gamma$ are both constant. Without loss of generality, we may assume that $\beta=0$. Since $M$ is of 2-type, $\gamma \neq 0$. Thus, we have $\omega_{1}{ }^{3}=0$ and $\omega_{2}{ }^{3}=\gamma \omega^{2}$, $\gamma \neq 0$. By taking exterior differentiation of these equations, we obtain $\omega_{2}{ }^{1}=0$ which implies $G=0$. This is a contradiction. (Q.E.D.)

## 5. Stationary, 2-type surfaces.

Let $f: M \rightarrow M^{\prime}$ be an isometric immersion of a surface $M$ into an $m$-dimensional Riemannian manifold $M^{\prime}$. We denote by $\alpha^{\prime}$ and $R^{\prime}$ the mean curvature of $f$ and the sectional curvature of $M^{\prime}$ with respect to the tangent space of $M$ and define $\tau(f)$ by

$$
\begin{equation*}
\tau(f)=\int_{M}\left(\left(\alpha^{\prime}\right)^{2}+R^{\prime}\right) d V \tag{5.1}
\end{equation*}
$$

It was proved in [2] that $\tau(f)$ is an invariant under conformal changes of the
metric of $M^{\prime}$ (cf. also [3, p. 207]). We call $\tau(f)$ the conformal total mean curvature. The variation of $\tau(f)$ was calculated in [10] (cf. also [3, pp. 213-225]). When $M^{\prime}$ is the unit hypersphere $S_{0}^{m}(1)$ of $E^{m+1}, f$ is a stationary point of $\tau$ if and only if

$$
\begin{equation*}
\Delta^{D^{\prime}} H^{\prime}=-2\left(\alpha^{\prime}\right)^{2} H^{\prime}+\left\|A_{\xi}\right\|^{2} H^{\prime}+\mathfrak{U}^{\prime}\left(H^{\prime}\right) \tag{5.2}
\end{equation*}
$$

where $H^{\prime}$ is the mean-curvature vector of $M$ in $S_{0}{ }^{m}(1)$ and $H^{\prime}=\alpha^{\prime} \xi, \alpha^{\prime}=\left|H^{\prime}\right|$.
In [5], Ejiri showed that the isometric immersion $f$ from the flat torus $S^{1}(1) \times S^{1}(\sqrt{1 / 3})$ into $S^{5}$ defined by

$$
f((x, y),(z, w))=(\sqrt{2 / 3} x, x z, x w, \sqrt{2 / 3} y, y z, y w)
$$

is a mass-symmetric, stationary, non-minimal surface in $S^{5}$.
It is easy to see that Ejiri's example is a 2-type surface in $S^{5}$. In the following, we want to classify stationary, 2-type, mass-symmetric surfaces in $S^{m}$. In particular, we shall obtain the following.

ThEOREM 3. Let $M$ be a stationary, mass-symmetric, 2-type surface in $S_{0}{ }^{m}(1)$. Then $M$ is a flat surface which lies fully in a totally geodesic $S_{0}{ }^{5}(1)$ or in a totally geodesic $S_{0}{ }^{7}(1)$ in $S_{0}{ }^{m}(1)$.

Proof. We need some lemmas.
LEMMA 5. Let $M$ be a stationary, mass-symmetric, 2-type surface in $S_{0}{ }^{m}(1)$ in $E^{m+1}$. Then we have
(1) $M$ is an $\mathfrak{H}$-surface,
(2) $\left|H^{\prime}\right|^{2}=\left(2-\lambda_{p}\right)\left(\lambda_{q}-2\right) / 4 \neq 0$,
(3) $\left\|A_{\xi}\right\|^{2}=\lambda_{p}+\lambda_{q}-2-\lambda_{p} \lambda_{q} / 4$,
(4) $\left\|D^{\prime} \xi\right\|^{2}=\lambda_{p} \lambda_{q} / 4$,
(5) $M$ is not pseudo-umbilical,
(6) $\operatorname{tr}\left(\bar{\nabla} A_{H^{\prime}}\right)=0$,
(7) $\operatorname{tr}\left(\nabla \omega_{3}{ }^{r}\right)=\left\langle D^{\prime} \xi, D^{\prime} \xi_{r}\right\rangle, r=4, \cdots, m$.

Conversely, if $M$ is an $\mathfrak{A}$-surface of $S_{0}{ }^{m}(1)$ satisfying (2), (3), (4), (6), and (7), then $M$ is a stationary, mass-symmetric 2-type surface in $S_{0}{ }^{m}(1)$.

Proof. If $M$ is mass-symmetric and of 2-type in $S_{0}{ }^{m}(1) \subset E^{m+1}, \alpha^{\prime}$ is a nonzero constant. So, there is a unit normal vector field $\xi$ on $M$ which is parallel to $H^{\prime}$. From Lemma 2 we obtain (2) and (6). Moreover, from Lemma 4 we find

$$
\begin{equation*}
\left\|A_{\xi}\right\|^{2}+\left\|D^{\prime} \xi\right\|^{2}=\lambda_{p}+\lambda_{q}-2 \tag{5.3}
\end{equation*}
$$

Since $M$ is stationary, (5.2) holds. Thus, by Lemmas 1 and 2, we find

$$
\begin{equation*}
\Delta H=2 \mathfrak{A}^{\prime}\left(H^{\prime}\right)+2\left(\left\|A_{\xi}\right\|^{2}-\left(\alpha^{\prime}\right)^{2}+1\right) H^{\prime}-2\left(\alpha^{\prime}\right)^{2} x . \tag{5.4}
\end{equation*}
$$

On the other hand, since $M$ is mass-symmetric and of 2 -type, we also have

$$
\begin{equation*}
\Delta H=\left(\lambda_{p}+\lambda_{q}\right) H^{\prime}+\left(\lambda_{p} \lambda_{q} / 2-\left(\lambda_{p}+\lambda_{q}\right)\right) x . \tag{5.5}
\end{equation*}
$$

Thus, by combining (5.4) and (5.5), we obtain (1), (2) and (3). Statement (4) follows from statement (3) and (5.3), Statement (5) follows from (2) and (3). Moreover, by applying Lemma 4 and (1), we obtain (7). The converse of this can be easily verified.

We choose $\left\{e_{1}, e_{2}\right\}$ which diagonalizes $A_{3}$. Then we have $h_{11}{ }^{4}=\cdots=h_{11}{ }^{m}=0$ because $M$ is not pseudo-umbilical and it is an $\mathfrak{A}$-surface. From Lemma 5, we also have $D^{\prime} \xi \neq 0$. Because $M$ is 2-dimensional, we may assume that $D^{\prime} \xi$ lies in the normal subspace spanned by $\xi_{4}$ and $\xi_{5}$. So, by a suitable choice of $\xi_{4}, \cdots, \xi_{m}$, we have

$$
\begin{gather*}
A_{3}=\left(\begin{array}{ll}
\beta & 0 \\
0 & \gamma
\end{array}\right), \quad A_{4}=\left(\begin{array}{ll}
0 & b \\
b & 0
\end{array}\right), \quad A_{6}=\left(\begin{array}{ll}
0 & c \\
c & 0
\end{array}\right), \\
A_{5}=A_{7}=\cdots=A_{m}=0, \quad D^{\prime} \xi_{3}=\omega_{3}{ }^{4} \xi_{4}+\omega_{3}{ }^{5} \xi_{5}, \tag{5.6}
\end{gather*}
$$

where $\beta$ and $\gamma$ are unequal constants.
Lemma 6. Under the hypothesis, $M$ is flat and $\omega_{1}{ }^{2}=b \omega_{3}{ }^{4}=0$.
Proof. Lemmas 2 and 3 imply $\operatorname{tr} A_{D^{\prime} \xi}=0$. Thus (5.6) gives $b \omega_{3}{ }^{4}=0$. So, by taking differentiation of $\omega_{1}{ }^{3}=\beta \omega^{1}$ and $\omega_{2}{ }^{3}=\gamma \omega^{2}$, and by using (5.6) and structure equations, we obtain $\omega_{1}{ }^{2}=0$. From $\omega_{1}{ }^{2}=0$, we see that $M$ is flat.

If $b \neq 0$, then Lemma 6 gives $\omega_{3}{ }^{4}=0$ and $D^{\prime} \xi_{3}$ being perpendicular to the first normal space. Thus, by choosing $\xi_{3}, \cdots, \xi_{m}$ such that the first normal space is spanned by $\xi_{3}$ and $\xi_{4}$, we obtain the following case (1). Otherwise, we have case (2):

Case (1). With respect to the frame field we have

$$
\begin{gathered}
A_{3}=\left(\begin{array}{ll}
\beta & 0 \\
0 & \gamma
\end{array}\right), \quad A_{4}=\left(\begin{array}{ll}
0 & b \\
b & 0
\end{array}\right), \quad A_{5}=\cdots=A_{m}=0, \\
D^{\prime} \xi_{3}=\omega_{3}{ }^{5} \xi_{5}, \quad \omega_{1}{ }^{2}=0, \quad b \neq 0,
\end{gathered}
$$

or
Case (2). With respect to the frame field, we have

$$
\begin{gathered}
A_{3}=\left(\begin{array}{ll}
\beta & 0 \\
0 & \gamma
\end{array}\right), \quad A_{4}=A_{5}=0, \quad A_{6}=\left(\begin{array}{ll}
0 & b \\
b & 0
\end{array}\right), A_{7}=\cdots=A_{m}=0, \\
D^{\prime} \xi_{3}=\omega_{3}{ }^{4} \xi_{4}+\omega_{3}{ }^{5} \xi_{5}, \quad \omega_{1}{ }^{2}=0 .
\end{gathered}
$$

In both cases, $\beta, \gamma$ and $b$ are constants with $b^{2}=1+\beta \gamma$.
We consider cases (1) and (2) separately.
If Case (1) holds, we have $\omega_{1}{ }^{3}=\beta \omega^{1}, \omega_{2}{ }^{3}=\gamma \omega^{2}, \omega_{1}{ }^{4}=b \omega^{2}, \omega_{2}{ }^{4}=b \omega^{1}, \omega_{i}{ }^{r}=0$ for $i=1,2 ; r=5, \cdots, m$. Taking differentiation of $\omega_{i}{ }^{r}=0$, we obtain $\omega_{4}{ }^{r}=0$, for $r=6, \cdots, m$. Thus,

$$
\begin{equation*}
D^{\prime} \xi_{4}=\omega_{4}{ }^{5} \xi_{5} . \tag{5.9}
\end{equation*}
$$

We put

$$
\begin{equation*}
\omega_{3}{ }^{5}=\mu_{1} \omega^{1}+\mu_{2} \omega^{2}, \quad \omega_{4}{ }^{5}=\eta_{1} \omega^{1}+\eta_{2} \omega^{2} \tag{5.10}
\end{equation*}
$$

Taking exterior differentiation of $\omega_{i}{ }^{5}=0, i=1,2$, we obtain $\omega_{i}{ }^{3} \wedge \omega_{3}{ }^{5}+\omega_{i}{ }^{4} \wedge \omega_{4}{ }^{5}=0$. Thus, by applying (5.7) and (5.10), we may obtain

$$
\begin{equation*}
\eta_{1}=\beta \mu_{2} / b, \quad \eta_{2}=\gamma \mu_{1} / b \tag{5.11}
\end{equation*}
$$

On the other hand, since $\omega_{3}{ }^{4}=0$, Lemma 5 and (5.10) imply

$$
\begin{equation*}
0=\operatorname{tr}\left(\nabla \omega_{3}{ }^{4}\right)=\mu_{1} \eta_{1}+\mu_{2} \eta_{2} \tag{5.12}
\end{equation*}
$$

Combining (5.11) and (5.12) we find $\mu_{1} \mu_{2} \doteq 0$. Since $\left(\mu_{1}\right)^{2}+\left(\mu_{2}\right)^{2}=\lambda_{p} \lambda_{q} / 4$ is a constant, we obtain $\mu_{1} \equiv 0$ or $\mu_{2} \equiv 0$. Without loss of generality, we may assume that $\mu_{2} \equiv 0$. Thus, we get

$$
\begin{equation*}
\omega_{3}{ }^{5}=\mu \omega^{1} \neq 0, \quad \mu^{2}=\lambda_{p} \lambda_{q} / 4, \quad \omega_{4}{ }^{5}=(\lambda \mu / b) \omega^{2} \tag{5.13}
\end{equation*}
$$

Now, since $\omega_{3}{ }^{r}=0$ for $r=6, \cdots, m$, Lemma 6 implies $\nabla \omega_{3}{ }^{r}=0$ where we use the definition of $\nabla \omega_{3}{ }^{r}$. Thus, by (7) of Lemma 5, (5.7) and (5.13), we obtain

$$
\begin{equation*}
\omega_{5}^{r}\left(e_{1}\right)=0, \quad r=6, \cdots, m . \tag{5.14}
\end{equation*}
$$

Moreover, by taking exterior differentiation of $\omega_{3}{ }^{r}=0$ and using (5.13), we may find $\omega^{1} \wedge \omega_{5}^{r}=0$. Combining this with (5.14) we find $\omega_{5}^{r}=0$ for $r=6, \cdots, m$. Since we already know that $\omega_{3}{ }^{r}=\omega_{4}{ }^{r}=0$ for $r=6, \cdots, m$, the normal subspace spanned by $\left\{\xi_{3}, \xi_{4}, \xi_{5}\right\}$ is parallel with respect to the normal connection $D^{\prime}$. Since the first normal subspace is spanned by $\left\{\xi_{3}, \xi_{4}\right\}$, Therefore, by a reduction theorem of submanifold [6], we may conclude that in fact $M$ lies in a totally geodesic $S_{0}{ }^{5}(1)$ of $S_{0}{ }^{m}(1)$. We summarize these as the following.

Lemma 7. Let $M$ be a stationary, mass-symmetric, 2-type surface in $S_{0}{ }^{m}(1)$. If Case (1) holds, then $M$ lies fully in a totally geodesic 5 -sphere $S_{0}{ }^{5}(1)$ of $S_{0}{ }^{m}(1)$. Moreover, with respect to a suitable orthonormal frame $\left\{e_{1}, e_{2}, \xi_{3}, \xi_{4}, \xi_{5}\right\}$ of $M$ in $S_{0}{ }^{5}(1)$, we have

$$
\begin{align*}
& A_{3}=\left(\begin{array}{ll}
\beta & 0 \\
0 & \gamma
\end{array}\right), \quad A_{4}=\left(\begin{array}{ll}
0 & b \\
b & 0
\end{array}\right), \quad A_{5}=0, \quad \omega_{1}^{2}=0, \quad D^{\prime} \xi_{3}=\mu \omega^{1} \xi_{5},  \tag{5.15}\\
& D^{\prime} \xi_{4}=(\gamma \mu / b) \omega^{2} \xi_{5}, \quad \beta \gamma+1=b^{2}, \quad \mu^{2}=\lambda_{p} \lambda_{q} / 4, \quad b^{2}(\beta-\gamma)+\gamma \mu^{2}=0 .
\end{align*}
$$

The last equation in (5.15) follows from the exterior differentiation of $\omega_{3}{ }^{4}=0$. Now, we shall consider Case (2). In this case, we have

$$
\begin{align*}
& \omega_{1}{ }^{2}=0, \quad \omega_{1}{ }^{3}=\beta \omega^{1}, \quad \omega_{2}{ }^{3}=\gamma \omega^{2}, \quad \omega_{i}{ }^{4}=\omega_{i}{ }^{5}=0, \quad \omega_{1}{ }^{6}=b \omega^{2}, \\
& \omega_{2}{ }^{6}=b \omega^{1}, \quad \omega_{i}^{r}=0, \quad \omega_{3}{ }^{6}=\cdots=\omega_{3}{ }^{m}=0, \quad i=1,2, \quad r=7, \cdots, m . \tag{5.16}
\end{align*}
$$

Moreover, we have

$$
\begin{equation*}
\beta \gamma+1=b^{2} \quad \text { and } \quad D^{\prime} \xi_{3}=\omega_{3}{ }^{4} \xi_{4}+\omega_{3}{ }^{5} \xi_{5} \neq 0 \tag{5.17}
\end{equation*}
$$

Taking exterior differentiation of $\omega_{i}{ }^{4}=\omega_{i}{ }^{5}=0$ and applying (5.17), we obtain

$$
\begin{equation*}
\beta \omega^{1} \wedge \omega_{3}{ }^{5}+b \omega^{2} \wedge \omega_{6}{ }^{5}=\gamma \omega^{2} \wedge \omega_{3}{ }^{5}+b \omega^{1} \wedge \omega_{6}{ }^{5}=0, \tag{5.18}
\end{equation*}
$$

$$
\begin{equation*}
\beta \omega^{1} \wedge \omega_{3}{ }^{4}+b \omega^{2} \wedge \omega_{6}{ }^{4}=\gamma \omega^{2} \wedge \omega_{3}{ }^{4}+b \omega^{1} \wedge \omega_{6}{ }^{4}=0 \tag{5.19}
\end{equation*}
$$

If $b=0$, (5.17) gives $\beta \gamma \neq 0$. Moreover, (5.18) and (5.19) imply $\omega_{3}{ }^{4}=\omega_{3}{ }^{5}=0$. This contradicts (5.17), Thus, we see that $b$ is a nonzero constant.

By taking exterior differentiation of $\omega_{i}^{r}=0, r=7, \cdots, m$, and applying (5.16), we get

$$
\begin{equation*}
\omega_{6}{ }^{r}=0, \quad r=7, \cdots, m . \tag{5.20}
\end{equation*}
$$

We put

$$
\begin{equation*}
\omega_{3}{ }^{4}=\alpha_{1} \omega^{1}+\alpha_{2} \omega^{2}, \quad \omega_{3}{ }^{5}=\delta_{1} \omega^{1}+\delta_{2} \omega^{2} . \tag{5.21}
\end{equation*}
$$

Then by (5.18) and (5.19) we find

$$
\begin{align*}
& \omega_{6}{ }^{4}=\left(\beta \alpha_{2} / b\right) \omega^{1}+\left(\gamma \alpha_{1} / b\right) \omega^{2},  \tag{5.22}\\
& \omega_{6}{ }^{5}=\left(\beta \delta_{2} / b\right) \omega^{1}+\left(\gamma \delta_{1} / b\right) \omega^{2} . \tag{5.23}
\end{align*}
$$

Because $\omega_{3}{ }^{6}=0$, Lemma 5 implies $0=\operatorname{tr}\left(\nabla \omega_{3}{ }^{6}\right)=\left\langle D^{\prime} \xi_{3}, D^{\prime} \xi_{6}\right\rangle=0$. Therefore, (5.21), (5.22) and (5.23) give

$$
\begin{equation*}
\alpha_{1} \alpha_{2}+\delta_{1} \delta_{2}=0 \tag{5.24}
\end{equation*}
$$

In the following, we may choose $\xi_{4}$ in such a way that

$$
\begin{equation*}
D_{e_{1}}^{\prime} \xi_{3}=\omega_{3}^{4}\left(e_{1}\right) \xi_{4}, \quad \delta_{1}=0 \tag{5.25}
\end{equation*}
$$

Therefore, we obtain

$$
\begin{equation*}
\omega_{8}{ }^{5}=\delta \omega^{2}, \quad \omega_{6}{ }^{5}=(\beta \delta / b) \omega^{1}, \tag{5.26}
\end{equation*}
$$

where $\delta=\delta_{2}$. Since $\delta_{1}=0$, (5.24) gives

$$
\begin{equation*}
\alpha_{1}=0 \quad \text { or } \quad \alpha_{2}=0 . \tag{5.27}
\end{equation*}
$$

If $\alpha_{1}=0$, we have

$$
\begin{equation*}
D_{e_{1}}^{\prime} \xi_{3}=0 \tag{5.28}
\end{equation*}
$$

In this case, we may choose $\xi_{4}$ in such a way that

$$
\begin{equation*}
D^{\prime} \xi_{3}=\omega_{3}{ }^{4} \xi_{4}, \quad \omega_{3}{ }^{4}=\alpha_{2} \omega^{2} \tag{5.29}
\end{equation*}
$$

Then we have $\omega_{3}{ }^{5}=0$. Therefore, by interchanging $\xi_{4}$ and $\xi_{6}$, we obtain Case (1). If $\delta=0$, the same argument holds. Consequently, we obtain the following.

Lemma 8. If $M$ is not the flat surface in a $S_{0}{ }^{5}(1)$ mentioned in Lemma 7, then, with respect to a suitable orthonormal frame $\left\{e_{1}, e_{2}, \xi_{3}, \cdots, \xi_{m}\right\}$, we have

$$
\begin{align*}
& A_{3}=\left(\begin{array}{ll}
\beta & 0 \\
0 & \gamma
\end{array}\right), \quad A_{4}=A_{5}=0, \quad A_{6}=\left(\begin{array}{ll}
0 & b \\
b & 0
\end{array}\right), \quad A_{7}=\cdots=A_{m}=0, \\
& \omega_{1}{ }^{2}=0, \quad \omega_{3}{ }^{4}=\alpha_{1} \omega^{1}, \quad \omega_{3}{ }^{5}=\delta \omega^{2}, \quad \omega_{3}{ }^{6}=\cdots=\omega_{3}{ }^{m}=0, \quad \omega_{4}{ }^{6}=-\left(\gamma \alpha_{1} / b\right) \omega^{2},  \tag{5.30}\\
& \omega_{5}{ }^{6}=-(\beta \delta / b) \omega^{1}, \quad \omega_{6}{ }^{7}=\cdots=\omega_{6}{ }^{m}=0, \quad \beta \gamma+1=b^{2} \neq 0, \quad \alpha_{1} \delta \neq 0 .
\end{align*}
$$

Now, we also need the following.
Lemma 9. Under the hypothesis of Lemma 8, we may choose the frame $\left\{e_{1}, e_{2}, \xi_{3}, \cdots, \xi_{m}\right\}$ in such a way that, in addition to (5.30), we also have

$$
\begin{equation*}
\omega_{4}{ }^{5}=0, \quad \omega_{4}^{r}=\omega_{5}^{r}=0 \quad \text { for } r=8, \cdots, m \tag{5.31}
\end{equation*}
$$

$$
\begin{equation*}
\alpha_{1} \text { and } \delta \text { are constant. } \tag{5.32}
\end{equation*}
$$

Proof. Since $A_{4}=0$, equation of Ricci implies

$$
0=\left\langle D_{e_{1}}^{\prime} D_{e_{2}}^{\prime} \xi_{6}, \xi_{4}\right\rangle-\left\langle D_{e_{2}}^{\prime} D_{e_{1}}^{\prime} \xi_{6}, \xi_{4}\right\rangle
$$

Thus, by (5.30) and constancy of $\beta, \gamma$ and $b$, we obtain

$$
\begin{equation*}
\gamma e_{1}\left(\alpha_{1}\right)=\beta \delta \omega_{5}{ }^{4}\left(e_{2}\right) . \tag{5.33}
\end{equation*}
$$

Similarly, by using [ $A_{5}, A_{6}$ ] $=0$ and equation of Ricci, we also have

$$
\begin{equation*}
\beta e_{2}(\delta)=\gamma \alpha_{1} \omega_{4}^{5}\left(e_{1}\right) . \tag{5.34}
\end{equation*}
$$

On the other hand, by (5.30), we have

$$
\begin{equation*}
\operatorname{tr}\left(\nabla \omega_{3}{ }^{4}\right)=e_{1}\left(\alpha_{1}\right), \quad \operatorname{tr}\left(\nabla \omega_{3}{ }^{5}\right)=e_{2}(\delta) \tag{5.35}
\end{equation*}
$$

Thus, by applying statement (7) of Lemma 5 and (5.30), we find

$$
\begin{equation*}
e_{1}\left(\alpha_{1}\right)=\delta \omega_{4}^{5}\left(e_{2}\right), \quad e_{2}(\delta)=\alpha_{1} \omega_{4}^{5}\left(e_{1}\right) . \tag{5.36}
\end{equation*}
$$

Since $M$ is not pseudo-umbilical (Lemma 5) and $\alpha_{1} \delta \neq 0$ (Lemma 8), (5.33), (5.34) and (5.36) imply

$$
\begin{equation*}
\omega_{4}{ }^{5}=0, \tag{5.37}
\end{equation*}
$$

$$
\begin{equation*}
e_{1}\left(\alpha_{1}\right)=e_{2}(\boldsymbol{\delta})=0 . \tag{5.38}
\end{equation*}
$$

On the other hand, since $\left[A_{3}, A_{4}\right]=\left[A_{3}, A_{5}\right]=0$, the equation of Ricci, (5.30) and (5.37) imply

$$
\begin{equation*}
e_{2}\left(\alpha_{1}\right)=e_{1}(\boldsymbol{\delta})=0 . \tag{5.39}
\end{equation*}
$$

Combining (5.38) and (5.39) and using the constancy of $\beta$ and $\gamma$, we see that $\gamma \alpha_{1}$ and $\beta \delta$ are constant.

Now, we want to prove that $\omega_{4}{ }^{r}=\omega_{5}{ }^{r}=0$ for $r=8, \cdots, m$. Since $\omega_{3}{ }^{r}=0$, statement (7) of Lemma 5 and (5.30) give

$$
\begin{equation*}
\alpha_{1} \omega_{4}^{\tau}\left(e_{1}\right)+\delta \omega_{5}^{\tau}\left(e_{2}\right)=0, \quad r=7, \cdots, m . \tag{5.40}
\end{equation*}
$$

On the other hand, from $\omega_{6}{ }^{r}=0, r=7, \cdots, m$, we find

$$
\begin{equation*}
-\gamma \alpha_{1} \omega_{4}^{\tau}\left(e_{1}\right)+\beta \delta \omega_{5}^{r}\left(e_{2}\right)=0, \quad r=7, \cdots, m . \tag{5.41}
\end{equation*}
$$

Since $(\beta+\gamma) \alpha_{1} \delta \neq 0$, (5.40) and (5.41) imply

$$
\begin{equation*}
\omega_{4}^{r}\left(e_{1}\right)=\omega_{5}^{r}\left(e_{2}\right)=0, \quad r=7, \cdots, m . \tag{5.42}
\end{equation*}
$$

Thus, we have

$$
\begin{equation*}
D_{e_{1}}^{\prime} \xi_{4}=-\alpha_{1} \xi_{3}, \quad D_{e_{2}}^{\prime} \xi_{5}=-\delta \xi_{3} . \tag{5.43}
\end{equation*}
$$

Now, since $D_{e_{2}} \xi_{4}$ has no component in $\operatorname{span}\left\{\xi_{3}, \xi_{4}, \xi_{5}\right\}$, we may choose $\xi_{7}$ in such a way that we have

$$
\begin{equation*}
D_{e_{2}}^{\prime} \xi_{4}=\omega_{4}{ }^{6}\left(e_{2}\right) \xi_{6}+\omega_{4}{ }^{7}\left(e_{2}\right) \xi_{7} \tag{5.44}
\end{equation*}
$$

In this way, we have $\omega_{4}{ }^{8}\left(e_{2}\right)=\cdots=\omega_{4}{ }^{m}\left(e_{2}\right)=0$. Combining this with (5.42), we obtain $\omega_{4}{ }^{8}=\cdots=\omega_{4}{ }^{m}=0$.

Taking exterior differentiation of $\omega_{3}{ }^{r}=0, r=7, \cdots, m$ and applying (5.30), we obtain

$$
\begin{equation*}
\alpha_{1} \omega_{4}^{r}\left(e_{2}\right)=\delta \omega_{5}^{r}\left(e_{1}\right), \quad r=7, \cdots, m . \tag{5.45}
\end{equation*}
$$

Combining this with $\omega_{4}^{r}=0$ for $r=8, \cdots, m$, and (5.40), we have $\omega_{5}^{r}=0$ for $r=$ $8, \cdots, m$. This proves the lemma.

From (5.42), we may put

$$
\begin{equation*}
\omega_{4}{ }^{7}=\mu_{1} \omega^{2}, \quad \omega_{5}{ }^{7}=\mu_{2} \omega^{1} . \tag{5.46}
\end{equation*}
$$

Taking exterior differentiation of (5.46) we may obtain

$$
\begin{equation*}
e_{1}\left(\mu_{1}\right)=e_{2}\left(\mu_{2}\right)=0 \tag{5.47}
\end{equation*}
$$

From (5.45) and (5.46) we get

$$
\begin{equation*}
\alpha_{1} \mu_{1}=\delta \mu_{2} . \tag{5.48}
\end{equation*}
$$

Since $\alpha_{1}$ and $\delta$ are nonzero constants, (5.47) and (5.48) show that $\mu_{1}$ and $\mu_{2}$ are constants, too. Since $\alpha_{1} \delta \neq 0$, (5.48) implies that either $\mu_{1}=\mu_{2}=0$ or $\mu_{1}$ and $\mu_{2}$ are nonzero constants satisfying $\alpha_{1} \mu_{1}=\delta \mu_{2}$. If $\mu_{1}=\mu_{2}=0$, we obtain $\omega_{4}{ }^{7}=\omega_{5}{ }^{7}=0$. Thus, by applying Lemmas 8 and 9 , and equation of Ricci, we find

$$
\beta \gamma \alpha_{1} \delta=b^{2}\left\langle D_{e_{1}}^{\prime} D_{e_{2}}^{\prime} \xi_{4}, \xi_{5}\right\rangle=b^{2}\left\langle D_{e_{2}}^{\prime} D_{e_{1}}^{\prime} \xi_{4}, \xi_{5}\right\rangle=-\alpha_{1} \delta b^{2} .
$$

Since $\beta \gamma+b^{2}=1$, we have $\alpha_{1} \delta=0$. This is a contradiction. Consequently, we conclude that $\mu_{1}$ and $\mu_{2}$ are nonzero constant. Now, taking exterior differentiation of $\omega_{4}{ }^{r}=0$ for $r=8, \cdots, m$ and applying Lemmas 11 and 12 and (5.46), we may obtain $\omega^{2} \wedge \omega_{7}{ }^{r}=0$ for $r=8, \cdots, m$. Similarly, by taking exterior differentiation of $\omega_{5}{ }^{r}=0, r=8, \cdots, m$, we may conclude that $\omega^{1} \wedge \omega_{7}{ }^{r}=0$. Consequently, we have $\omega_{7}{ }^{8}=\cdots=\omega_{7}{ }^{m}=0$. Combining these with (5.30) and (5.31), we see that the normal subspace $\nu=\operatorname{span}\left\{\xi_{3}, \cdots, \xi_{7}\right\}$ is parallel with respect to the normal connection $D^{\prime}$. Moreover, $\nu$ contains the first normal space $\operatorname{span}\left\{\xi_{3}, \xi_{6}\right\}$. Thus, by a reduction theorem of submanifolds [6], we conclude that $M$ is in fact contained in a totally geodesic 7 -sphere $S_{0}{ }^{7}(1)$ of $S_{0}{ }^{m}(1)$. Furthermore, from the connection form $\left(\omega_{A}{ }^{B}\right), A, B=1, \cdots, 8$, of $M$ in $S_{0}{ }^{m}(1)$, we may also conclude that $M$ lies fully in $S_{0}{ }^{7}(1)$. This completes the proof of the theorem.

## 6. Connection form.

Theorem 3 says that if $M$ is a stationary, mass-symmetric, 2-type surface in $S_{0}{ }^{m}(1)$, then $M$ is flat and it lies fully in a $S_{0}{ }^{5}(1)$ or $S_{0}{ }^{7}(1)$. In this section, we shall determine the connection form of such surfaces.

Theorem 4. If $M$ is a stationary, mass-symmetric, 2-type surface in $S_{0}{ }^{5}(1)$, then $M$ is flat and $2 / 3<\lambda_{p}<2$. Moreover, with respect to an adapted orthonormal frame field, the connection form is given by (6.1) if $2 / 3<\lambda_{p} \leqq 4 / 3$ and given by (6.1) or (6.2) if $4 / 3<\lambda_{p}<2$ :

$$
\left(\begin{array}{cc|ccc}
0 & 0 & \left.\begin{array}{ccc}
\frac{\sqrt{2}-\sqrt{2} c}{\sqrt{3 c-2}} \omega^{1} & \frac{\sqrt{c}}{\sqrt{3 c-2}} \omega^{2} & 0 \\
0 & 0 & \frac{\sqrt{c}}{\sqrt{3 c-2}} \omega^{2} \\
\frac{\sqrt{c}}{\sqrt{3 c-2}} \omega^{1} & 0 \\
\hline \frac{\sqrt{2} c-\sqrt{2}}{\sqrt{3 c-2}} \omega^{1} & \frac{-\sqrt{2}}{\sqrt{3 c-2}} \omega^{2} & 0 \\
0 & \frac{-c}{\sqrt{3 c-2}} \omega^{1} \\
\frac{-\sqrt{c}}{\sqrt{3 c-2}} \omega^{2} & \frac{-\sqrt{c}}{\sqrt{3 c-2}} \omega^{1} & 0 \\
0 & 0 & \frac{c}{\sqrt{3 c-2}} \omega^{1} \\
0 & \frac{\sqrt{2 c}}{\sqrt{3 c-2}} \omega^{2} & 0
\end{array}\right) . \begin{array}{c}
\sqrt{2 c-2} \omega^{2} \\
0
\end{array} & 0 & 0 \tag{6.1}
\end{array}\right)
$$

where $c=\lambda_{p}$ is a real number satisfying $2 / 3<c<2$, or

$$
\left(\begin{array}{cc|ccc}
0 & 0 & \frac{c-4}{2 \sqrt{3 c-4}} \omega^{1} & \frac{1}{2} \sqrt{c} \omega^{2} & 0  \tag{6.2}\\
0 & 0 & \frac{1}{2} \sqrt{3 c-4} \omega^{2} & \frac{1}{2} \sqrt{c} \omega^{1} & 0 \\
\hline \frac{4-c}{2 \sqrt{3 c-4}} \omega^{1} & -\frac{1}{2} \sqrt{3 c-4} \omega^{2} & 0 & 0 & \frac{-c}{\sqrt{6 c-8}} \omega^{1} \\
-\frac{1}{2} \sqrt{c} \omega^{2} & -\frac{1}{2} \sqrt{c} \omega^{1} & 0 & 0 & -\frac{\sqrt{c}}{\sqrt{2}} \omega^{2} \\
0 & 0 & \frac{c}{\sqrt{6 c-8}} \omega^{1} & \frac{\sqrt{c}}{\sqrt{2}} \omega^{2} & 0
\end{array}\right)
$$

where $c=\lambda_{p}$ is a real number satisfying $4 / 3<c<2$.
Proof. Under the hypothesis, Lemma 7 implies

$$
\begin{equation*}
\omega_{1}{ }^{2}=\omega_{1}{ }^{5}=\omega_{2}{ }^{5}=\omega_{3}^{4}=0, \quad \omega_{1}{ }^{3}=\beta \omega^{1}, \quad \omega_{2}{ }^{3}=\gamma \omega^{2}, \tag{6.3}
\end{equation*}
$$

$$
\begin{equation*}
\omega_{1}^{4}=b \omega^{2}, \quad \omega_{2}^{4}=b \omega^{1}, \quad \omega_{3}{ }^{5}=\mu \omega^{1}, \quad \omega_{4}{ }^{5}=\frac{\gamma \mu}{b} \omega^{2}, \tag{6.4}
\end{equation*}
$$

$$
\begin{equation*}
\beta \gamma+1=b^{2}, \quad \mu^{2}=\lambda_{p} \lambda_{q} / 4, \quad b^{2}(\beta-\gamma)+\gamma \mu^{2}=0 . \tag{6.5}
\end{equation*}
$$

Moreover, Lemmas 5 and 7 also imply

$$
\begin{equation*}
(\beta+\gamma)^{2}=\left(2-\lambda_{p}\right)\left(\lambda_{q}-2\right), \tag{6.6}
\end{equation*}
$$

$$
\begin{equation*}
\beta^{2}+\gamma^{2}=\lambda_{p}+\lambda_{q}-2-\lambda_{p} \lambda_{q} / 4 \tag{6.7}
\end{equation*}
$$

By using the second equation of (6.5), (6.6) and (6.7) we obtain

$$
\begin{equation*}
2 \mu^{2}=(\beta-\gamma)^{2} . \tag{6.8}
\end{equation*}
$$

Replacing $\xi_{5}$ by $-\xi_{5}$ if necessary, we may assume that

$$
\begin{equation*}
\sqrt{2} \mu=\beta-\gamma \tag{6.9}
\end{equation*}
$$

By using (6.5) and (6.9), we may obtain

$$
\begin{equation*}
\beta=\left(\gamma^{2}-2\right) / 3 \gamma, \quad \mu=-\sqrt{2}\left(1+\gamma^{2}\right) / 3 \gamma, \quad b^{2}=\left(1+\gamma^{2}\right) / 3 \tag{6.10}
\end{equation*}
$$

Replacing $\xi_{4}$ by $-\xi_{4}$ if necessary, we may assume that $b$ is positive, so we have

$$
\begin{equation*}
b=\left(1+\gamma^{2}\right)^{1 / 2} / \sqrt{3} . \tag{6.11}
\end{equation*}
$$

Substituting the first equation of (6.10) into (6.6) and (6.7) and then solving $\lambda_{q}$ in terms of $\lambda_{p}$, we may obtain

$$
\begin{equation*}
\lambda_{q}=2 \lambda_{p} /\left(3 \lambda_{p}-4\right) \quad \text { or } \quad \lambda_{q}=4 \lambda_{p} /\left(3 \lambda_{p}-2\right) \tag{6.12}
\end{equation*}
$$

On the other hand, since $M$ is mass-symmetric in $S_{0}{ }^{5}(1)$ and of 2-type, Theorem 9.1 of Chen [3, p. 307] gives $0<\lambda_{p}<2<\lambda_{q}$. Thus, by (6.12), we see that the first equation of (6.12) holds only when $4 / 3<\lambda_{p}<2$ and the second equation holds only when $2 / 3<\gamma_{p}<2$. Combining this with (6.3)-(6.7) and (6.9)-(6.11), we may obtain the theorem. (Q.E.D.)

Remark 2. Theorem 4 shows that the immersion is rigid.
Remark 3. In the next section, we will show that both cases of (6.1) and (6.2) occur and the estimates on $\lambda_{p}$ are best possible.

For a stationary, mass-symmetric, 2-type surface which lies fully in $S_{0}{ }^{7}(1)$, Lemmas 5, 8 and 9 and (5.46), (5.47) give

$$
\begin{gather*}
\omega_{1}{ }^{2}=\omega_{1}{ }^{4}=\omega_{2}{ }^{4}=\omega_{1}{ }^{5}=\omega_{2}{ }^{5}=\omega_{1}{ }^{7}=\omega_{2}{ }^{7}=\omega_{3}{ }^{6}=\omega_{3}{ }^{7}=\omega_{4}{ }^{5}=\omega_{6}{ }^{7}=0,  \tag{6.13}\\
\omega_{1}{ }^{3}=\beta \omega^{1}, \quad \omega_{2}{ }^{3}=\gamma \omega^{2}, \quad \omega_{1}{ }^{6}=b \omega^{2}, \quad \omega_{2}{ }^{6}=b \omega^{1},  \tag{6.14}\\
\omega_{3}{ }^{4}=\alpha_{1} \omega^{1}, \quad \omega_{3}{ }^{5}=\delta \omega^{2}, \quad \omega_{4}{ }^{6}=-\frac{\gamma \alpha_{1}}{b} \omega^{2}, \quad \omega_{5}{ }^{6}=-\frac{\beta \delta}{b} \omega^{1},  \tag{6.15}\\
\omega_{4}{ }^{7}=\mu_{1} \omega^{2}, \quad \omega_{5}{ }^{7}=\mu_{2} \omega^{1},  \tag{6.16}\\
\beta \gamma+1=b^{2}, \quad \alpha_{1} \mu_{1}=\delta \mu_{2}, \quad \beta \neq \gamma, \quad \beta+\gamma \neq 0,  \tag{6.17}\\
(\beta+\gamma)^{2}=\left(2-\lambda_{p}\right)\left(\lambda_{q}-2\right), \quad \beta^{2}+\gamma^{2}=\lambda_{p}+\lambda_{q}-\frac{1}{4} \lambda_{p} \lambda_{q},  \tag{6.18}\\
\alpha_{1}{ }^{2}+\delta^{2}=\frac{1}{4} \lambda_{p} \lambda_{q}, \quad \alpha_{1} \delta b \mu_{1} \mu_{2} \neq 0, \tag{6.19}
\end{gather*}
$$

where $\beta, \gamma, b, \alpha_{1}, \delta, \mu_{1}$ and $\mu_{2}$ are constants. Moreover, by taking differentiation of $\omega_{3}{ }^{6}=0$ and $\omega_{4}{ }^{5}=0$, we may also obtain

$$
\begin{equation*}
(\beta-\lambda) b^{2}=\beta \delta^{2}-\gamma \alpha_{1}^{2}, \quad \alpha_{1} \delta b^{2}=\beta \gamma \delta \alpha_{1}+b^{2} \mu_{1} \mu_{2} . \tag{6.20}
\end{equation*}
$$

Combining (6.17) and (6.20), we have

$$
\begin{equation*}
\left(\mu_{1}\right)^{2}=\left(\frac{\delta}{b}\right)^{2}, \quad\left(\mu_{2}\right)^{2}=\left(\frac{\alpha_{1}}{b}\right)^{2} . \tag{6.21}
\end{equation*}
$$

Without loss of generality, we may choose $\xi_{4}, \xi_{5}, \xi_{7}$ in such a way that $\alpha_{1}, \delta$ and $\mu_{1}$ are positive. Then we have $\mu_{1}=\delta / b$ and $\mu_{2}=\alpha_{1} / b$. From (6.18) we get $(\beta-\gamma)^{2}=\lambda_{p} \lambda_{q} / 2$. Furthermore, we have $0<\lambda_{p}<2<\lambda_{q}<\infty$ by Theorem 9.1 of [3, p. 307]. By interchanging $e_{1}$ and $e_{2}$ and replacing $\xi_{3}$ by $-\xi_{3}$ if necessary, we may assume that $\beta<\gamma$ and $\beta+\gamma>0$. In this case, we have $\beta+\gamma=\left[\left(2-\lambda_{p}\right)\left(\lambda_{q}-2\right)\right]^{1 / 2}$. From these we have the following.

Theorem 5. If $M$ is a stationary, mass-symmetric, 2-type surface which lies fully in $S_{0}{ }^{7}(1)$, then $M$ is a flat surface; moreover, with respect to an adapted orthonormal frame field, the connection form $\left(\omega_{A}{ }^{B}\right)$ is given by

$$
\left(\begin{array}{cc|ccccc}
0 & 0 & \beta \omega^{1} & 0 & 0 & b \omega^{2} & 0  \tag{6.22}\\
0 & 0 & \gamma \omega^{2} & 0 & 0 & b \omega^{1} & 0 \\
\hline-\beta \omega^{1} & -\gamma \omega^{2} & 0 & \alpha_{1} \omega^{1} & \delta \omega^{2} & 0 & 0 \\
0 & 0 & -\alpha_{1} \omega^{1} & 0 & 0 & -\frac{\gamma \alpha_{1}}{b} \omega^{2} & \frac{\delta}{b} \omega^{2} \\
0 & 0 & -\delta \omega^{2} & 0 & 0 & -\frac{\beta \delta}{b} \omega^{1} & \frac{\alpha_{1}}{b} \omega^{1} \\
-b \omega^{2} & -b \omega^{1} & 0 & \frac{\gamma \alpha_{1}}{b} \omega^{2} & \frac{\beta \delta}{b} \omega^{1} & 0 & 0 \\
0 & 0 & 0 & -\frac{\delta}{b} \omega^{2} & -\frac{\alpha_{1}}{b} \omega^{1} & 0 & 0
\end{array}\right)
$$

where $b, \beta, \gamma, \alpha_{1}, \delta$ are constants satisfying

$$
\begin{aligned}
& \beta=(1 / 2) \sqrt{c d / 2}+(1 / 2) \sqrt{(2-c)(d-2)}, \quad \gamma=-(1 / 2) \sqrt{c d / 2}+(1 / 2) \sqrt{(2-c)(d-2)}, \\
& \alpha_{1}=(\gamma-\beta)\left(2+3 \beta \gamma-\beta^{2}\right) / 2(\beta+\gamma), \quad \delta=(\beta-\gamma)\left(2+3 \beta \gamma-\gamma^{2}\right) / 2(\beta+\gamma), \\
& b=\sqrt{1+\beta \gamma},
\end{aligned}
$$

for some constants $c=\lambda_{p}$ and $d=\lambda_{q}$ so that $0<c<2<d<\infty$.
Remark 4. For any real numbers $c$ and $d$ with $0<c<2<d<\infty$, the connection form given in Theorem 5 satisfies the structure equations (or integrability condition). Thus, by Fundamental Theorem of Submanifolds, we see that there is a "unique" isometric immersion $y$ from $\boldsymbol{R}^{2}$ into $S_{0}{ }^{7}(1)$ whose connection form is given by (6.22). When $y$ is doubly-periodic, $y$ yields many stationary, masssymmetric, 2-type, flat surfaces in $S_{0}{ }^{7}(1)$ with $\lambda_{p}=c$ and $\lambda_{q}=d$. Theorem 5 also implies that all stationary, mass-symmetric, 2-type surfaces which lie fully in
$S_{0}{ }^{7}(1)$ are obtained in this way.

## 7. Examples.

Let $\boldsymbol{R}^{2}$ be the Euclidean plane with the Euclidean metric. Let $u, v$ and $w$ be real numbers with $u, v>0$. We define the lattice

$$
\begin{equation*}
\Lambda=\{(2 n \pi u, 2 m \pi v+2 n \pi w) \mid n, m \in \boldsymbol{Z}\} . \tag{7.1}
\end{equation*}
$$

The dual lattice of $\Lambda$ is given by

$$
\begin{equation*}
\Lambda^{*}=\left\{\left.\left(\frac{h}{2 \pi u}-\frac{k w}{2 \pi u v}, \frac{k}{2 \pi v}\right) \right\rvert\, h, k \in \boldsymbol{Z}\right\} . \tag{7.2}
\end{equation*}
$$

Let $T_{u v w}$ be the flat torus given by $\boldsymbol{R}^{2} / \Lambda$. Then the spectrum of $T_{u v w}$ is given by

$$
\begin{equation*}
\left\{\left.\left(\frac{h}{u}-\frac{k w}{u v}\right)^{2}+\frac{k^{2}}{v^{2}} \right\rvert\, h, k \in \boldsymbol{Z}\right\} . \tag{7.3}
\end{equation*}
$$

For any nonzero real number $\varepsilon$ and two natural numbers $h$ and $\bar{\varepsilon}$ satisfying

$$
\begin{equation*}
\varepsilon \neq 2 h \bar{\varepsilon}^{2} /\left(\bar{\varepsilon}^{2}-2 h^{2}\right), \tag{7.4}
\end{equation*}
$$

we put

$$
\begin{align*}
& u=\sqrt{3} \varepsilon \bar{\varepsilon} /\left(2 \varepsilon^{2}+\bar{\varepsilon}^{2}\right)^{1 / 2}, \quad v=\bar{\varepsilon} /\left(2 \varepsilon^{2}+\bar{\varepsilon}^{2}\right)^{1 / 2} \\
& w=(h-\varepsilon) \bar{\varepsilon} /\left(2 \varepsilon^{2}+\bar{\varepsilon}^{2}\right)^{1 / 2}, \quad e=\sqrt{2} \varepsilon /\left(2 \varepsilon^{2}+\bar{\varepsilon}^{2}\right)^{1 / 2} \tag{7.5}
\end{align*}
$$

and we define an isometric immersion $y$ from $\boldsymbol{R}^{2}$ into $S_{0}{ }^{5}(1) \subset E^{6}$ by

$$
\begin{align*}
y(s, t)= & \left(v \cos \frac{\varepsilon s}{u} \cos \frac{t}{v}, v \cos \frac{\varepsilon s}{u} \sin \frac{t}{v}, e \cos \frac{\bar{\varepsilon} s}{u}\right.  \tag{7.6}\\
& \left.v \sin \frac{\varepsilon s}{u} \cos \frac{t}{v}, v \sin \frac{\varepsilon s}{u} \sin \frac{t}{v}, e \sin \frac{\bar{\varepsilon} s}{u}\right) .
\end{align*}
$$

The immersion $y$ induces an isometric immersion from $T_{u v w}$ into $S_{0}{ }^{5}(1)$ which is denoted by $x$. Thus we have

$$
\begin{equation*}
x: T_{u v w} \longrightarrow S_{0}{ }^{5}(1) \subset E^{6} . \tag{7.7}
\end{equation*}
$$

It is easy to see that if $\varepsilon=\bar{\varepsilon}=h=1$, then (7.7) gives Ejiri's example mentioned in section 5 .

Proposition 1. For any nonzero real number $\varepsilon$ and two natural numbers $h$ and $\bar{\varepsilon}$ satisfying (7.4), the immersion $x: T_{u v w} \rightarrow S_{0}{ }^{5}(1)$ is a stationary, mass-symmetric, isometric immersion, where $u, v$ and $w$ are defined by (7.5). Furthermore, we have
(a) $x$ is of 1-type if and only if $\bar{\varepsilon}^{2}=4 \varepsilon^{2}$, in this case, $\lambda_{p}=2$.
(b) Otherwise, $x$ is of 2 -type with $\lambda_{p}$ and $\lambda_{q}$ given by

$$
\begin{equation*}
\left\{\lambda_{p}, \lambda_{q}\right\}=\left\{\left(\frac{\bar{\varepsilon}}{u}\right)^{2},\left(\frac{\varepsilon}{u}\right)^{2}+\left(\frac{1}{v}\right)^{2}\right\} . \tag{7.8}
\end{equation*}
$$

Proof (Outlined). From (7.6) we see that $x$ is an isometric immersion. The Laplacian of $T_{u v w}$ is given by $\Delta=-\partial^{2} / \partial s^{2}-\partial^{2} / \partial t^{2}$. Therefore, the coordinate functions of $x$ are eigenfunctions of $\Delta$ with eigenvalues given by (7.7), From (7.5) and (7.6) we know that $\lambda_{p}=\lambda_{q}$ if and only if $\bar{\varepsilon}^{2}=4 \varepsilon^{2}$. In this case, $x$ is of 1-type. Otherwise, $x$ is of 2-type.

By direct, long computation, we may prove that $T_{u v w}$ is an $\mathfrak{N}$-surface of $S_{0}{ }^{5}(1)$. Moreover, we may also prove that a mass-symmetric, 2-type, $\mathfrak{A}$-surface of $S_{0}{ }^{m}(1)$ is stationary if and only if $\Delta H=2\left(\left\|A_{\xi}\right\|^{2}-2\left(\alpha^{\prime}\right)^{2}\right) H^{\prime}+2 \alpha^{2} H$. So, by a long, straight-forward computation, we may in fact prove that the immersion $x$ satisfies this equation.

Remark 5. It is easy to check that $w$ satisfies $w^{2} \leqq\left(2 h^{2}+\bar{\varepsilon}^{2}\right) / 2$. If one chooses $w^{2} \in\left(0,\left(2 h^{2}+\bar{\varepsilon}^{2}\right) / 2\right)$, then one obtains two non-isometric tori. Otherwise, if $w=0$ or $w^{2}=\left(2 h^{2}+\bar{\varepsilon}^{2}\right) / 2$, one obtains only one torus. Moreover, if $w=0$, the torus is defined by a rectangular lattice.

Theorem 6. We have the following two statements.
(a) For each real number $c$ with $2 / 3<c \leqq 4 / 3$, there is a stationary, masssymmetric, 2-type, flat torus in $S_{0}{ }^{5}(1)$ whose connection form is given by (6.1).
(b) For each real number $c$ with $4 / 3<c<2$, there are two stationary, masssymmetric, 2-type, flat tori in $S_{0}{ }^{5}(1)$ whose connection forms are given by (6.1) and (6.2) respectively.

Proof. Consider the stationary, mass-symmetric, 2-type, flat torus in $S_{0}{ }^{5}(1)$ (with $\bar{\varepsilon}^{2} \neq 4 \varepsilon^{2}$ ) given by (7.7). According to (7.5) and (7.8), we have

$$
\begin{equation*}
\left\{\lambda_{p}, \lambda_{q}\right\}=\left\{\frac{2}{3}+\frac{1}{3 \varepsilon^{2}}, \frac{4}{3}+\frac{8}{3} \varepsilon^{2}\right\} \tag{7.9}
\end{equation*}
$$

Given a real number $c$ with $2 / 3<c<2$, we consider the following equation

$$
\begin{equation*}
c=\lambda_{p}=\frac{2}{3}+\frac{1}{3 \varepsilon^{2}} . \tag{7.10}
\end{equation*}
$$

The only thing we need to prove is that the range of $\varepsilon^{2}$ is $(1 / 4, \infty)$. One notices that for any fixed natural numbers $h$ and $\bar{\varepsilon}, \varepsilon$ depends on $w$ continuously over the domain. For instance, consider the case $h=\bar{\varepsilon}=1$, we have

$$
\begin{equation*}
\varepsilon=\left\{1+w\left(3-2 w^{2}\right)^{1 / 2}\right\} /\left(1-2 w^{2}\right), \quad \text { or } \quad \varepsilon=\left\{1-w\left(3-2 w^{2}\right)^{1 / 2}\right\} /\left(1-2 w^{2}\right) . \tag{7.11}
\end{equation*}
$$

In this case, the range of $w$ is $\left(-\frac{\sqrt{6}}{2}, \frac{\sqrt{6}}{2}\right)-\{1\}$ (with $w^{2}=3 / 2$ corresponding to the 1-type case). Now, it is not difficult to see that the range of $\varepsilon^{2}$ is $(1 / 4, \infty)$.

If $4 / 3<c<2$, we may also consider the equation

$$
\begin{equation*}
c=\lambda_{p}=\frac{4}{3}+\frac{8}{3} \varepsilon^{2} . \tag{7.12}
\end{equation*}
$$

By a similar argument we may see that the range of $\varepsilon^{2}$ is ( $0,1 / 4$ ). (Q.E.D.)
Remark 6. From the Fundamental Theorem of Submanifolds, we see that for any $c \in(2 / 3,4 / 3]$, there is an isometric immersion $y$ from $\boldsymbol{R}^{2}$ into $S_{0}{ }^{5}(1)$ whose connection form is given by (6.1). Such immersions are unique up to rigid motions. According to Theorem 6, for such $c$, there is a flat torus $T_{c}$ in $S_{0}{ }^{5}(1)$ whose connection form is also given by (6.1). Thus, if we lift the immersion $x$ of $T_{c}$ up to its universal covering $\boldsymbol{R}^{2}$, we obtain an isometric immersion $\bar{x}$ from $\boldsymbol{R}^{2}$ into $S_{0}{ }^{5}(1)$. Since $y$ and $\bar{x}$ have the same connection form, they only differ by a rigid motion. Consequently, the immersion $y$ is doubly-periodic. For $c \in(4 / 3,2)$, we have two isometric immersions $y_{1}, y_{2}$ from $\boldsymbol{R}^{2}$ into $S_{0}{ }^{5}(1)$ whose connection forms are given by (6.1) and (6.2), respectively. Theorem 6 implies that $y_{1}$ and $y_{2}$ are both doubly-periodic. From these, we conclude that all stationary, mass-symmetric, 2-type surfaces in $S_{0}{ }^{5}(1)$ are always obtained in this way.

In Remark 4, we know that a stationary, mass-symmetric, 2-type surface which lies fully in $S_{0}{ }^{7}(1)$ is obtained by a doubly-periodic isometric immersion of $\boldsymbol{R}^{2}$ into $S_{0}{ }^{7}(1)$ whose connection form is given by (6.22). In the following, we give some concrete examples of such surfaces in $S_{0}{ }^{7}(1)$.

Recall that for any real numbers $u, v$ and $w$ with $u, v>0$, we have a flat torus $T_{u v w}$. Given four natural numbers ( $n, m, \bar{n}, \bar{m}$ ), we put

$$
\begin{equation*}
\varepsilon=n-\frac{m w}{v}, \quad \bar{\varepsilon}=\bar{n}-\frac{\bar{m} w}{v} . \tag{7.13}
\end{equation*}
$$

We define an isometric immersion $y$ from $\boldsymbol{R}^{2}$ into $S_{0}{ }^{7}(1) \subset E^{8}$ by

$$
\begin{align*}
& y(s, t)=\left(c_{1} \cos \frac{\varepsilon s}{u} \cos \frac{m t}{v}, \quad c_{1} \cos \frac{\varepsilon s}{u} \sin \frac{m t}{v},\right.  \tag{7.14}\\
& c_{1} \sin \frac{\varepsilon s}{u} \cos \frac{m t}{v}, \quad c_{1} \sin \frac{\varepsilon s}{u} \sin \frac{m t}{v} \\
& c_{2} \cos \frac{\bar{\varepsilon} s}{u} \cos \frac{\bar{m} t}{v}, \quad c_{2} \cos \frac{\bar{\varepsilon} s}{u} \sin \frac{\bar{m} t}{v}, \\
& \left.c_{2} \sin \frac{\bar{\varepsilon} s}{u} \cos \frac{\bar{m} t}{v}, \quad c_{2} \sin \frac{\bar{\varepsilon} s}{u} \sin \frac{\bar{m} t}{v}\right),
\end{align*}
$$

where $c_{1}$ and $c_{2}$ are two real numbers satisfying

$$
\begin{equation*}
c_{1}^{2}+c_{2}^{2}=1, \quad c_{1}{ }^{2} \bar{\varepsilon}^{2}+c_{2}{ }^{2} \varepsilon^{2}=u^{2}, \quad c_{1}^{2} m^{2}+c_{2}^{2} \bar{m}^{2}=v^{2} \tag{7.15}
\end{equation*}
$$

The immersion $y$ induces an isometric immersion from $T_{u v w}$ into $S_{0}{ }^{7}(1) \subset E^{8}$ which is denoted by $x$. Thus we have

$$
\begin{equation*}
x: T_{u v w} \longrightarrow S_{0}{ }^{7}(1) \subset E^{8} . \tag{7.16}
\end{equation*}
$$

By using an argument similar to that of Proposition 1, we may prove the following.

PROPOSITION 2. If $v^{2}\left(\bar{\varepsilon}^{2}-\varepsilon^{2}\right) \neq u^{2}\left(m^{2}-\bar{m}^{2}\right)$, then the immersion $x: T_{u v w} \rightarrow$ $S_{0}{ }^{7}(1)$ is a mass-symmetric, 2-type, isometric immersion with

$$
\begin{equation*}
\left\{\lambda_{p}, \lambda_{q}\right\}=\left\{\left(\frac{\varepsilon}{u}\right)^{2}+\left(\frac{m}{v}\right)^{2}, \quad\left(\frac{\bar{\varepsilon}}{u}\right)^{2}+\left(\frac{\bar{m}}{v}\right)^{2}\right\} . \tag{7.17}
\end{equation*}
$$

Moreover, $T_{\text {uvw }}$ is an $\mathfrak{N}$-surface of $S_{0}{ }^{7}(1)$. Furthermore, the immersion $x$ is stationary if and only if the following equation holds:

$$
\begin{equation*}
2 c_{1}{ }^{2}\left[\left(\frac{\varepsilon}{u}\right)^{2}-\left(\frac{m}{v}\right)^{2}\right]^{2}+2 c_{2}{ }^{2}\left[\left(\frac{\bar{\varepsilon}}{u}\right)^{2}-\left(\frac{\bar{m}}{v}\right)^{2}\right]^{2}=\lambda_{p} \lambda_{q} . \tag{7.18}
\end{equation*}
$$

We also need the following.
Proposition 3. Let $[p, q]$ be the order of a stationary immersion given by (7.16). Then we have
(a) $2 / 3<\lambda_{p}<2$,
(b) if $2 / 3<\lambda_{p} \leqq 4 / 3$, then $\lambda_{q}>4 \lambda_{p} /\left(3 \lambda_{p}-2\right)$, and
(c) if $4 / 3<\lambda_{p}<2$, then $2 \lambda_{p} /\left(3 \lambda_{p}-4\right)>\lambda_{q}>4 \lambda_{p} /\left(3 \lambda_{p}-2\right)$.

Proof. Follows from (7.15), (7.18) and the fact $0<\lambda_{p}<2<\lambda_{q}$ of [3, p.307].
Given two real numbers $c$ and $d$ with $0<c<2<d<\infty$, we put

$$
\begin{align*}
& F(c, d)=c-\{c d(c-2) / 2(2-d)\}^{1 / 2} \\
& G(c, d)=d+\{c d(c-2) / 2(2-c)\}^{1 / 2} \tag{7.19}
\end{align*}
$$

Lemma 10. For any $d \in(2, \infty)$ and any rational number $r \neq 0$, there is a $c \in$ $(2 / 3,2)$ such that $G(c, d)=r^{2} F(c, d)$.

Proof. Under the hypothesis, it is easy to see that there is a $c \in(0,2)$ satisfying $G=r^{2} F$. Because $G$ is positive, $F$ is also positive. Thus, we obtain $c \in(2 / 3,2)$.

By using Lemma 10, (7.13), (7.16) and (7.20), we obtain the following.
Theorem 7. For any $d \in(2, \infty)$, there is a stationary, mass-symmetric,

2-type, flat torus in $S_{0}{ }^{7}(1)$ such that $\lambda_{q}=d$ and whose connection form is given by (6.22).

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