# ON DIFFERENTIABLE MANIFOLDS WITH $(\phi, \psi)$-STRUCTURES 

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(Received September 5, 1960)

1. Introduction. In a previous paper, ${ }^{1)}$ we have defined for some odd dimensional manifolds two kinds of structures which we have called ( $\phi, \xi, \eta$ )structure and $(\phi, \xi, \eta, g)$-structure. The latter is a $(\phi, \xi, \eta)$-structure with a positive definite Riemannian metric $g$ which stands in a notable relation with the ( $\phi, \xi, \eta$ )structure. These structures are remarkable in the sense that any differentiable manifold with $(\phi, \xi, \eta)$-structure is an almost contact manifold and any almost contact manifold admits ( $\phi, \xi, \eta, g$ )-structure.

In this paper, we shall study two kinds of structures for differentiable manifolds of any dimension, the first one ( $(\phi, \psi)$-structure) may be regarded as generalizations of almost complex structure, almost product structure and ( $\phi, \xi, \eta)$ structure, and the second one ( $\phi, \psi, g$ )-structure) may be regarded as generalizasions of almost Hermitian structure, almost product metric structure and ( $\phi, \xi$, $\eta, g)$-structure. We shall confine ourselves only to algebraic considerations, analytic considerations will be published in later papers.

## 2. $(\phi, \psi)$-structures.

$1^{\circ}$. Let $M^{n}$ be a differentiable manifold of dimension $n$. Suppose first that there exist over $M^{n}$ two tensor fields $\phi_{j}^{\prime}$ and $\psi_{j}^{(2)}$ of type ( 1,1 ) which satisfy the following conditions:

$$
\begin{aligned}
\operatorname{rank}\left|\boldsymbol{\phi}_{j}^{i}\right| & =l, \\
\operatorname{rank}\left|\boldsymbol{\psi}_{j}^{i}\right| & =m, \\
\phi_{j}^{i} \psi_{k}^{j} & =0, \\
\boldsymbol{\psi}_{j}^{i} \phi_{k}^{\prime} & =0,
\end{aligned}
$$

1) S. Sasaki, On differentiable manifolds with certain structures which are closely related to almost contact structure I, Tòhoku Math. Journ. 12(1960) pp. 459-476.
2) We assume, unless otherwise stated, that the indices run the following range of integers:

$$
\begin{aligned}
& i, j, k, a, \beta, \gamma=1,2, \ldots \ldots, n(=l+m), \\
& a, b, c=1,2, \ldots \ldots, l, \\
& p, q, r=l+1, \ldots \ldots, n \\
& A=1,2, \ldots \ldots, l^{\prime}, \quad A^{*}=l^{\prime}+A, \\
& E, F=1, \ldots \ldots, l_{1}, \quad H=l_{1}+1, \ldots \ldots, l\left(=l_{1}+l_{2}\right) \\
& L=l+1, \ldots \ldots, l+m^{\prime}, \quad L *=l+m^{\prime}+L, \\
& M, N=l+1, \ldots \ldots, l+m_{1}, \\
& S, T=l+m_{1}+1, \ldots \ldots, n .
\end{aligned}
$$

$$
\begin{equation*}
\varepsilon \phi_{j}^{\prime} \phi_{k}^{j}+\varepsilon^{\prime} \psi_{i}^{i} \psi_{k}^{j}=\delta_{k}^{i}, \tag{2.5}
\end{equation*}
$$

where $l, m$ are non negative integers such that

$$
\begin{equation*}
l+m=n \tag{2.6}
\end{equation*}
$$

and $\varepsilon, \varepsilon^{\prime}$ are +1 or -1 . In such case we say that the manifold $M^{n}$ in consideration has a $(\phi, \psi)$-structure of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$ or $M^{n}$ is a differentiable manifold with a $(\phi, \psi)$-structure of type ( $\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}$ ). From our definition we see that if a differentiable manifold $M^{n}$ admits a ( $\phi, \psi$ ) -structure of type (sgn $\varepsilon$, sgn $\varepsilon^{\prime}$ ), then it admits also a $(-\phi, \psi)$-structure, a $(\phi,-\psi)$-structure and a $(-\phi,-\psi)$ structure, all of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$. Hence, we identify all of these structures.
$2^{\circ}$. First we shall prove the following
THEOREM 1. Suppose $M^{n}$ be a differentiable manifold with a ( $\phi, \psi$ ). structure of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$. Then in every sufficiently small coordinate neighborhood $U$ of $M^{n}$, we can find frames $\left(\xi_{a}^{i}, \xi_{p}^{l}\right)(a=1, \ldots \ldots, l ; p=l+1, \ldots$ $\ldots, n$ ) such that

$$
\left\{\begin{array}{l}
\phi_{j}^{i}=\lambda_{b}^{a} \xi_{a}^{i} \eta_{\eta}^{b},  \tag{2.7}\\
\psi_{j}^{i}=\mu_{q}^{p} \xi_{p}^{i} \eta_{j}^{\eta},
\end{array}\right.
$$

where $\left(\eta_{j}^{a}, \eta_{j}^{p}\right)$ are the inverse matrix of $\left(\xi_{a}^{i}, \xi_{p}^{i}\right)$ and $\lambda_{b}^{a}, \mu_{q}^{p}$ are scalars such that

$$
\left\{\begin{array}{l}
\lambda_{b}^{a} \lambda_{c}^{b}=\varepsilon \delta_{l}^{a},  \tag{2.8}\\
\mu_{q}^{p} \mu_{r}^{q}=\varepsilon^{\prime} \delta_{q}^{p}
\end{array}\right.
$$

PROOF. As the rank of $\left|\phi_{j}^{i}\right|$ is equal to $l$, there exist $m$ linearly independent vector fields over $U$ which are solutions of the equation

$$
\begin{equation*}
\phi_{j}^{\prime} \xi^{j}=0 \tag{2.9}
\end{equation*}
$$

Let us denote any such vector fields by $\xi_{p}^{i}$ and take $n$ vector fields $\xi_{a}^{i}, \xi_{p}^{i}$ over $U$ so that they are linearly independent. If we put the inverse matrix of ( $\xi_{a}^{i}, \xi_{p}^{\prime}$ ) by ( $\eta_{i}^{a}, \eta_{j}^{p}$ ), then $\phi_{j}^{i}$ can be written as

$$
\phi_{j}^{i}=\lambda_{\beta}^{\alpha} \xi_{a}^{i} \eta_{j}^{\beta} .
$$

By virtue of the construction, we have

$$
\begin{equation*}
\phi_{j}^{\prime} \xi_{p}^{\prime}=0 \tag{2.10}
\end{equation*}
$$

Therefore, we can easily see that $\lambda_{p}^{\alpha}=0$ and hence we get

$$
\begin{equation*}
\phi_{j}^{i}=\lambda_{b}^{a} \xi_{a}^{i} \eta_{j}^{b}+\lambda_{b}^{p} \xi_{p}^{i} \eta_{j}^{b} . \tag{2.11}
\end{equation*}
$$

Next, $\psi_{j}^{i}$ can be written also as

$$
\psi_{j}^{l}=\mu_{\beta}^{\alpha} \xi_{a}^{l} \eta_{j}^{\beta} .
$$

If we put this into (2.3), we get

$$
\phi_{j}^{i} \psi_{k}^{j}=\mu_{\beta}^{a}\left(\phi_{i}^{i} \xi_{a}^{j}\right) \eta_{k}^{\beta}=0
$$

However, $\phi_{j}^{\prime} \xi_{a}^{\prime}$ 's do not vanish for any value of $a$, so we get $\mu_{\beta}^{a} \eta_{k}^{\beta}=0$. Hence, we get $\mu_{\beta}^{a}=0$. Therefore, $\psi_{j}^{i}$ has the following form:

$$
\begin{equation*}
\psi_{j}^{i}=\mu_{q}^{p} \xi_{p}^{i} \eta_{j}^{q}+\mu_{b}^{p} \xi_{p}^{i} \eta_{j}^{b} . \tag{2.12}
\end{equation*}
$$

Thirdly, putting (2.11) and (2.12) into (2.4), we get

$$
\psi_{j}^{\prime} \phi_{k}^{j}=\left(\mu_{a}^{p} \lambda_{b}^{a}+\mu_{a}^{p} \lambda_{b}^{q}\right) \xi_{p}^{i} \eta_{t}^{b}=0 .
$$

Hence, we see that the relation

$$
\begin{equation*}
\mu_{a}^{p} \lambda_{b}^{a}+\mu_{q}^{p} \lambda_{b}^{q}=0 \tag{2.13}
\end{equation*}
$$

holds good.
Finally, putting (2.11) and (2.12) into the left hand side of (2.5), we get

$$
\varepsilon \phi_{j}^{\prime} \phi_{k}^{j}+\varepsilon^{\prime} \psi_{j}^{i} \psi_{k}^{j}=\varepsilon \lambda_{b}^{\alpha} \lambda_{c}^{b} \xi_{a}^{i} \eta_{k}^{c}+\varepsilon^{\prime} \mu_{q}^{p} \mu_{\xi}^{q} \xi_{p}^{i} \eta_{k}^{r} .
$$

Comparing this with

$$
\delta_{k}^{i}=\xi_{\alpha}^{l} \eta_{k}^{\alpha}
$$

we see, by virtue of (2.5), that the relations

$$
\left\{\begin{align*}
\varepsilon \lambda_{b}^{a} \lambda_{c}^{b} & =\delta_{c}^{a},  \tag{2.14}\\
\varepsilon^{\prime} \mu_{q}^{p} \mu_{r}^{a} & =\delta_{r}^{p}, \\
\varepsilon \lambda_{b}^{p} \lambda_{c}^{b}+\varepsilon^{\prime} \mu_{q}^{p} \mu_{c}^{a} & =0
\end{align*}\right.
$$

hold good.
Now, we take another frame $\overline{\xi_{\alpha}^{\bar{c}}}$ which are given by

$$
\left\{\begin{array}{l}
\overline{\xi_{b}^{i}}=\lambda_{b}^{a} \xi_{a}^{i}+\lambda_{b}^{p} \xi_{\nu}^{i},  \tag{2.15}\\
\bar{\xi}_{q}^{i}=\quad \varepsilon^{\prime} \mu_{q}^{p} \xi_{p}^{i} .
\end{array}\right.
$$

Then, the inverse matrix $\overline{\boldsymbol{\eta}}_{j}^{\alpha}$ of $\overline{\boldsymbol{\xi}_{\alpha}^{i}}$ is easily seen to be

$$
\left\{\begin{array}{l}
\bar{\eta}_{j}^{a}=\varepsilon \lambda_{b}^{a} \eta_{j}^{b},  \tag{2.16}\\
\overline{\eta_{j}^{p}}=\mu_{b}^{p} \eta_{j}^{b}+\mu_{q}^{p} \eta_{j}^{q} .
\end{array}\right.
$$

The equation (2.15) can be solved with respect to $\xi_{\alpha}^{l}$ giving
and the equation (2.16) can be solved with respect to $\eta_{j}^{\alpha}$ giving

$$
\left\{\begin{array}{l}
\eta_{j}^{a}=\lambda_{b}^{a} \overline{\eta_{j}^{j}}  \tag{2.18}\\
\eta_{j}^{p}=-\varepsilon^{\prime} \mu_{g}^{p} \mu_{b}^{a} \lambda_{\lambda_{b}^{b}}^{\eta_{j}^{a}}+\varepsilon^{\prime} \mu_{g}^{p} \overline{\eta_{j}} .
\end{array}\right.
$$

If we put (2.17) and (2.18) into (2.11) and (2.12), we can easily see that the relations

$$
\left\{\begin{array}{c}
\phi_{j}^{l}=\lambda_{b}^{a \overline{\xi_{i}}} \overline{\eta_{j}^{\bar{p}}},  \tag{2.19}\\
\psi_{j}^{l}=\mu_{u}^{p} \overline{\xi_{p}^{i}} \eta_{j}^{\bar{q}}
\end{array}\right.
$$

hold good, where $\lambda_{b}^{a}, \mu_{l}^{p}$ are scalars over $U$ such that

$$
\left\{\begin{array}{l}
\lambda_{b}^{a} \lambda_{c}^{b}=\varepsilon \delta_{c}^{a} \\
\mu_{q}^{p} \mu_{r}^{a}=\varepsilon^{\prime} \delta_{r}^{p}
\end{array}\right.
$$

Consequently, if we change our notations and write $\xi_{\alpha}^{i}, \eta_{j}^{\alpha}$ instead of $\overline{\xi_{\alpha}^{i}}, \overline{\eta_{j}^{\alpha}}$, we see that our theorem is true.
Q.E.D.

We call the frame such that the tensors $\phi_{j}^{i}, \psi_{j}^{i}$ take the form (2.7) satisfying (2.8) an adapted frame of the first order.

REMARK 1. The above demonstration shows that the conditions (2.2) and (2.6) follow from the conditions (2.1), (2.3), (2.4) and (2.5).

REMARK 2. From (2.8) we see that if $\varepsilon=-1$, then $l$ is even and if $\varepsilon^{\prime}=-1, m$ is even.
$3^{\circ}$. Suppose that $M^{n}$ be a differentiable manifold with a $(\phi, \psi)$-structure. Then, at evary point $P$ of $M^{n}$, the set of vectors such that

$$
\begin{equation*}
\phi_{j}^{i} \xi^{j}=0 \tag{2.20}
\end{equation*}
$$

is an $m$-dimensional vector subspace $V_{m}$ spanned by $\xi_{p}^{l}$ at $P$. In the same way we can see from (2.2), that the set of vectors at $P$ such that

$$
\begin{equation*}
\psi_{j}^{\prime} \xi_{a}^{j}=0 \tag{2.21}
\end{equation*}
$$

is an $l$-dimensional vector subspace $V_{l}$ spanned by $\xi_{a}^{i}$. Hence, $V_{l}$ and $V_{m}$ are disjoint and complementary. In other words, if we denote the tangent space at $P$ by $T_{P}$, then

$$
\begin{equation*}
T_{P}=V_{l} \oplus V_{m} \tag{2.22}
\end{equation*}
$$

The correspondence $P \in M^{n}$ to $V_{\iota}$ at $P$ and the correspondence $P \in M^{n}$ to $V_{m}$ at $P$ define the so-called $l$ - and $m$-dimensional distributions over $M^{n}$. We call them $D_{l}$ and $D_{m}$. Then we get the following

THEOREM 2. Suppose $M^{n}$ be a differentiable manifold with a $(\phi, \psi)$-structure. Then, the two distributions $D_{l}$ and $D_{m}$ are disjoint and complementary.

Adapted frames of the first order in a coordinate neighborhood $U$ are nothing but frames whose first $l$ vectors span the vector space $V_{l}$ of $D_{l}$ and whose last $m$ vectors span the vector space $V_{m}$ of $D_{m}$ at every point of $U$.
$4^{\circ}$. Now, we consider a transformation of adapted frames of the first order

$$
\left\{\begin{array}{l}
\overline{\xi_{b}^{i}}=\alpha_{b}^{i} \xi_{a}^{i},  \tag{2.23}\\
\bar{\xi}_{q}^{i}=\beta_{q}^{p \xi_{p}^{i}}
\end{array}\right.
$$

where $\alpha$ and $\beta$ are non-singular matrices. Then, it induces a transformation of $\eta$ of the form

$$
\left\{\begin{array}{l}
\overline{\eta_{j}^{\bar{z}}}={ }^{\prime} \alpha_{b}^{\tau} \eta_{j}^{b},  \tag{2.24}\\
\overline{\eta_{j}^{p}}={ }^{\prime} \beta_{q}^{p} \eta_{j}^{\eta},
\end{array}\right.
$$

where ' $\alpha$ and ' $\beta$ are inverse matrices of $\alpha$ and $\beta$ respectively. Putting (2.23) and (2.24) into

$$
\begin{aligned}
\phi_{j}^{b} & =\bar{\lambda}_{b}^{\bar{a}} \bar{\xi}_{a} \bar{\eta}_{j}^{b} \\
& =\lambda_{b}^{a} \xi_{a}^{i} \eta_{j}^{b}
\end{aligned}
$$

we see that

$$
\begin{equation*}
\boldsymbol{\alpha}_{a}^{c} \bar{\lambda}_{b}^{a}=\lambda_{a}^{c} \alpha_{b}^{a}, \tag{2.25}
\end{equation*}
$$

which shows that $\lambda_{3}^{a}$, transform like components of a mixed tensor under transformations (2.23). Hence, if $\varepsilon=-1$, we can take $\alpha$ so that the matrix $\lambda$ takes the form

$$
\lambda=\left(\begin{array}{cc}
0 & -E_{l^{\prime}}  \tag{2.26}\\
E_{l^{\prime}} & 0
\end{array}\right)
$$

where $E_{l^{\prime}}$ is a unit matrix of dimension $l^{\prime}=l / 2$. In this case $\phi^{i}$, reduces to the form

$$
\begin{equation*}
\phi_{j}^{i}=-\bar{\xi}_{A}^{i} \bar{\eta}_{j}^{A *}+\bar{\xi}_{A+}^{\eta_{j}^{A}}, \tag{2.27}
\end{equation*}
$$

where $A$ runs over $1,2, \ldots \ldots, l^{\prime}$ and $A^{*}=A+l^{\prime}$. If we take $\beta$ arbitrary, then with respect to the frame $\left(\overline{\xi_{A}^{i}}, \overline{\xi_{A}}, \overline{\xi_{p}^{i}}\right)$ thus deternined, $\phi$ has the following components:

$$
\phi=\left(\begin{array}{cc:c}
0 & -E_{l^{\prime}} & 0  \tag{2.28}\\
E_{l^{\prime}} & 0 & 0 \\
\hdashline 0 & 0 & 0
\end{array}\right)
$$

Especially, we can see that the characteristic roots of the matrix $\phi$ are equal to $0, i$ and $-i$ with multiplicities $m, l^{\prime}$ and $l^{\prime}$ respectively.

On the other hand, if $\varepsilon=+1$, then the characteristic roots of the matix $\lambda$ are equal to -1 or +1 . If multiplicities of the roots -1 and +1 are $l_{1}$ and $l_{2}\left(l_{1}+l_{2}=l\right)$ respectively, then we can take $\alpha$ so that the matrix $\lambda$ takes the form

$$
\lambda=\left(\begin{array}{cc}
-E_{l_{1}} & 0  \tag{2.29}\\
0 & E_{l_{2}}
\end{array}\right)
$$

where $E_{l_{1}}$ and $E_{l_{2}}$ are unit matrices of dimensions $l_{1}$ and $l_{2}$. In this case $\phi$ takes the form

$$
\begin{equation*}
\phi_{l}^{i}=-\overline{\xi_{E}} \overline{\eta_{j}^{k}}+\overline{\xi_{H}^{i}} \overline{\eta_{j}^{H}} . \tag{2.30}
\end{equation*}
$$

Even if we take $\beta$ arbitrary, with respect to the frame ( $\left.\overline{\xi_{H}^{l}}, \overline{\xi_{K}}, \overline{\xi_{p}}\right)$ thus determined, $\phi$ has the following components:

$$
\phi=\left(\begin{array}{cc:c}
-E_{l_{1}} & 0 & 0  \tag{2.31}\\
0 & E_{l_{2}} & \\
\hdashline 0 & 0 & 0
\end{array}\right)
$$

Similar facts hold good for $\psi$ too. Hence, summarizing the above results, we get the following

THEOREM 3. Suppose $M^{n}$ be a differentiable manifold with a $(\phi, \psi)$-structure of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$. Then we can take frames over every coordinate neighborhood $U$ so that $\phi$ and $\psi$ take the following forms:

$$
\left\{\begin{align*}
\phi_{j}^{i} & =-\xi_{A}^{i} \eta_{j}^{A^{*}}+\xi_{A}^{i} \eta_{l}^{A}, & & \text { for } \varepsilon=-1  \tag{2.32}\\
& =-\xi_{H}^{i} \eta_{j}^{H}+\xi_{K}^{i} \eta_{j}^{K}, & & \text { for } \varepsilon=+1 \\
\psi_{j}^{i} & =-\xi_{L}^{i} \eta_{j}^{L^{*}}+\xi_{L^{*}}^{i} \eta_{j}^{L}, & & \text { for } \varepsilon^{\prime}=-1 \\
& =-\xi_{M}^{i} \eta_{j}^{H L}+\xi_{s}^{i} \eta_{J .}^{S} . & & \text { for } \varepsilon^{\prime}=+1
\end{align*}\right.
$$

REMARK. If $M^{n}$, is a differentiable manifold with a ( $\phi, \psi$ ) -structure of type $(-,+)$ such that rank of $\left|\boldsymbol{\psi}_{j}^{i}\right|$ is equal to 1 , then with respect to adapted frames of the first order in a coordinate neighborhood $U$ of $M^{n}, \psi_{j}^{i}$ may take the form $\xi^{i} \eta_{j}$ as $\left(\phi, \psi_{)}\right)$-structure is identical with $(\phi,-\psi)$-structure. Hence, the conditions (2.1) to (2.6) reduce in this case to

$$
\left\{\begin{align*}
\text { rank }\left|\phi_{j}^{i}\right| & =2 l^{\prime}, & n & =2 l^{\prime}+1  \tag{2.33}\\
\phi_{j}^{i} \xi^{j} & =0, & \phi_{j}^{i} \eta_{i} & =0,
\end{align*}\right.
$$

$$
\phi_{j}^{i} \phi_{k}^{\prime}=-\delta_{k}^{i}+\xi^{i} \eta_{j} \xi^{j} \eta_{k} .
$$

and trivial equations. These combined with $\xi^{i} \eta_{i}=1$ are nothing but the defining equations of the $(\phi, \xi, \eta)$-structure for $M^{2 l^{\prime}+1}$. However, contrary to the case of $(\phi, \xi, \eta)$-structure, our vector fields $\xi^{i}$ and $\eta_{j}$ are defined locally. They do not in general constitute a vector field over $M^{2 l^{\prime}+1}$. Hence, the set of differentiable manifolds with $(\phi, \psi)$-structures of type $(-,+)$ such that the rank of $\left|\psi^{i}\right|$ is equal to 1 is somewhat wider than the set of differentiable manifolds with $(\phi, \xi, \eta)$ structures.

Formulas in (2.32) are canonical forms of the tensors $\phi$ and $\psi$. We call any frame with respect to which $\phi$ and $\psi$ take such canonical forms an adapted frame of the second order of the given $(\phi, \psi)$-structure.
$5^{\circ}$. Now, the tensor fields $\phi$ and $\psi$ define linear maps of tangent vectors at every point of $M^{n}$ by $v \rightarrow \phi v$ and $v \rightarrow \psi v$.

THEOREM 4. Suppose $M^{n}$ be a differentiable manifold with a $(\phi, \psi)$-structure and $V_{l}, V_{m}$ are associated vector spaces at any point $P$ of $M^{n}$. Then

$$
\begin{equation*}
\phi \phi v=\varepsilon v, \quad \text { for } v \in V_{l} \tag{i}
\end{equation*}
$$

$$
\begin{equation*}
\psi \psi v=\varepsilon^{\prime} v, \quad \text { for } v \in V_{m} \tag{ii}
\end{equation*}
$$

$$
\begin{equation*}
\phi \psi v=0, \psi \phi v=0, \quad \text { for } v \in T_{P} \tag{iii}
\end{equation*}
$$

PROOF. If $v \in V_{l}$, then $\psi v=0$ and the converse is also true. In this case we see easily that

$$
\phi \phi v=\varepsilon\left(\delta-\varepsilon^{\prime} \psi \psi\right) v=\varepsilon v .
$$

Hence (i) is proved. In the same way we can prove (ii). (iii) follows immediately from (2.3) and (2.4).

THEOREM 5. The linear maps $v \rightarrow(\phi+\psi) v$ and $v \rightarrow(\phi-\psi) v$ of tangent spaces are non-singular.

PROOF. By virtue of (2.3), (2.4) and (2.5), we can verify that

$$
(\phi+\psi)\left(\varepsilon \phi+\varepsilon^{\prime} \psi\right)=\varepsilon \phi \phi+\varepsilon^{\prime} \psi \psi=\delta .
$$

Hence, $\phi+\psi$ and $\varepsilon \phi+\varepsilon^{\prime} \psi$ are non-singular and inverse to each other. Similarly, $\phi-\psi$ is non-singular.

## 3. Associated Riemannian metric $g$.

$6^{\circ}$. In this section we study if we can associate a positive definite Riemannian metric $g$ to any differentiable manifold $M^{n}$ with a $(\phi, \psi)$-structure or not. We begin with a lemma.

LEMMA 1. Suppose $M^{n}$ be a differentiable manifold such that there exist two distributions $D_{l}$ and $D_{m}$ of dimensions $l$ and $m$ which are disjoint and complementary. Then there exists a positive definite Riemannian metric $h$ with respect to which the vector spaces $V_{l}$ and $V_{m}$ of the distributions at every point of $M^{\prime n}$ are orthogonal to each other.

PROOF. First we introduce an arbitrary positive definite Riemannian metric $f$ over $M^{n}$. Suppose $\left\{U_{\alpha}\right\}$ be a sufficiently fine open covering of $M^{n}$ by coordinate neighborgoods.

Now we take $l$ (resp. $m$ ) orthonormal vector fields $\boldsymbol{\xi}_{a}^{i}$ (resp. $\boldsymbol{\xi}_{p}^{i}$ ) over $U_{\alpha}$ with respect to $f$ so that they span the vector space $V_{l}$ (resp. $V_{m}$ ) of the distribution $D_{l}$ (resp. $D_{m}$ ) at every point of $U_{\alpha}$. Of course, $\xi_{a}^{i}$ and $\xi_{p}^{i}$ are not orthogonal to each other in general. We define

$$
\begin{equation*}
h^{i j}\left(U_{\alpha}\right)=\sum_{a} \xi_{a}^{\prime} \xi_{a}^{j}+\sum_{p} \xi_{p}^{\prime} \xi_{p}^{j} \tag{3.1}
\end{equation*}
$$

On the other hand, let $U_{\beta}$ be another coordinate neighborhood which belongs to $\left\{U_{\alpha}\right\}$ such that $U_{a} \cap U_{\beta}$ is not empty and $\overline{\xi_{a}^{i}}, \overline{\xi_{p}^{i}}$ are vector fields over $U_{\beta}$ defined in the same way as above. Then, it is evident that

$$
\begin{aligned}
\bar{\xi}_{a}^{i} & =\sum_{b} u_{a b} \xi_{b}^{i} \\
\bar{\xi}_{p}^{i} & =\sum_{q} u_{p q} \xi_{q}^{i}
\end{aligned}
$$

hold good over $U_{a} \cap U_{\beta}$, where ( $u_{a b}$ ) and ( $u_{p q}$ ) are orthogonal matrices. We can easily verify that

$$
h^{i j}\left(U_{\alpha}\right)=h^{i j}\left(U_{\beta}\right)
$$

holds good over $U_{\alpha} \cap U_{\beta}$. This shows that the set of tensor fields $h^{i j}\left(U_{\alpha}\right), U_{\alpha}$ $\in\left\{U_{\alpha}\right\}$, constitutes a global tensor field over $M^{n}$. The inverse $h_{i j}$ of the tensor field $h^{i j}$, then determine s a positive definite Riemannian metric over $M^{n}$. We can easily verify that, with respect to $h=\left(h_{i j}\right)$, the two vector spaces $V_{l}$ and $V_{m}$ are orthogonal at every point of $M^{n}$.
Q. E. D.

LEMMA 2. With respect to the metric $h$ over $M^{n}$ defined in the proof of Lemma 1, the relation

$$
\begin{equation*}
h_{i j} \phi_{n}^{i} \boldsymbol{\psi}_{k}^{j}=0 \tag{3.2}
\end{equation*}
$$

holds good.
PROOF. Let $U$ be an arbitrary coordinate neighborhood of $M^{n}$. We take frames over $U$ so that orthonormal vectors $\xi_{a}^{i}$ (resp. $\xi_{p}^{i}$ ) with respect to the metric $h$ span the vector space $V_{l}$ (resp. $V_{m}$ ).

On the other hand, as ( $\xi_{a}^{i}, \xi_{p}^{i}$ ) are adapted frames of the first order, we have

$$
\left\{\begin{aligned}
\phi_{j}^{i} & =\lambda_{b}^{a} \xi_{a}^{i} \eta_{j}^{b}, \\
\psi_{j}^{i} & =\mu_{q}^{p} \xi_{p}^{i} \eta_{j}^{q} .
\end{aligned}\right.
$$

Hence, making use of the fact

$$
h_{i j} \xi_{a}^{\xi} \xi_{p}^{j}=0,
$$

we can easily verify that our Lemma is true.
$7^{\circ}$. Now let us prove one of our main theorems.
THEOREM 6. Suppose $M^{n}$ be a differentiable manifold with a $(\phi, \psi)$-structure of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$. Then, there exists a positive definite Riemannian metric $g$ over $M^{n}$ such that the relations

$$
\left\{\begin{array}{l}
g_{i j} \phi_{h}^{i} \psi_{k}^{j}=0,  \tag{3.3}\\
g_{i h} \phi_{j}^{h}=\varepsilon g_{j h} \phi_{l}^{h}, \\
g_{i h} \psi_{j}^{h}=\varepsilon^{\prime} g_{j h} \psi_{i}^{h}, \\
g_{i j} \phi_{h}^{i} \phi_{k}^{j}+g_{i j} \psi_{h}^{i} \psi_{k}^{j}=g_{h k}
\end{array}\right.
$$

hold good.
Proof. We put

$$
\begin{equation*}
g_{i j}=\frac{1}{2}\left(h_{i j}+h_{\alpha \beta} \phi_{i}^{\alpha} \phi_{j}^{\beta}+h_{\alpha \beta} \psi_{i}^{\alpha} \psi_{j}^{\beta}\right) . \tag{3.4}
\end{equation*}
$$

Then, first by virtue of (2.3), (2.4) and (3.2), (3.3) ${ }_{1}$ is easily seen to be true.
Secondly, we see by virtue of (2.4) and (2.5) that

$$
\begin{aligned}
g_{i l} \phi_{j}^{h} & =\frac{1}{2} h_{i l} \phi_{j}^{h}+\frac{1}{2} h_{\alpha \beta} \phi_{i}^{\alpha}\left(\varepsilon \delta_{j}^{\beta}-\varepsilon \varepsilon^{\prime} \boldsymbol{\psi}_{l}^{\beta} \psi_{j}^{h}\right) \\
& =\frac{1}{2} h_{i l} \phi_{j}^{h}+\frac{\varepsilon}{2} h_{j h} \phi_{i}^{h} \\
& =\varepsilon g_{j h} \phi_{i}^{h} .
\end{aligned}
$$

In the same way, $(3.3)_{3}$ can be proved.
Thirdly, we see by virtue of (3.3) $)_{2,3}$ and (2.5) that

$$
\begin{aligned}
g_{i j} \phi_{h}^{i} \phi_{k}^{j} & +g_{i j} \psi_{h}^{i} \psi_{k}^{j} \\
& =\varepsilon g_{i h} \phi_{j}^{i} \phi_{k}^{j}+\varepsilon^{\prime} g_{i n} \psi_{j}^{i} \psi_{k}^{j} \\
& =g_{h k} .
\end{aligned}
$$

Hence, (3.3) ${ }_{4}$ is proved.
Q. E. D.

We shall call the Riemannian metric $g$ whose existence is insured by Theorem 6 the associated Riemannian metric of the $(\boldsymbol{\phi}, \boldsymbol{\psi})$-structure in consideration. And the differentiable manifold with the $(\phi, \psi)$-structure and its associated metric $g$ is called a manifold with $a(\phi, \psi, g)$-structure. It is an analogue of the almost Hermitian manifold for almost complex structure.

TheOrem 7. Suppose $M^{n}$ be a differentiable manifold with a $(\phi, \psi, g)$ structure. Then, tensor equations

$$
\left\{\begin{array}{l}
g_{i j}\left(\phi_{l l}^{i}+\psi_{h}^{i}\right)\left(\phi_{k}^{j}+\psi_{k}^{j}\right)=g_{l k},  \tag{3.5}\\
g_{i j}\left(\phi_{l h}^{i}-\psi_{h}^{i}\right)\left(\phi_{k}^{j}-\psi_{k}^{j}\right)=g_{h k}
\end{array}\right.
$$

hold good.
PROOF. By virtue of (2.3), (2.4) and (3.3) $)_{4}$, we can easily see that

$$
\begin{aligned}
& g_{i j}\left(\phi_{h}^{i}+\psi_{h}^{i}\right)\left(\phi_{k}^{j}+\psi_{k}^{j}\right) \\
& \quad=g_{i j} \phi_{h}^{i} \phi_{k}^{j}+g_{i j} \psi_{h}^{i} \psi_{k}^{j}=g_{h k} .
\end{aligned}
$$

In the same way $(3.5)_{2}$ can be proved.
THEOREM 8. Suppose $M^{n}$ be a differentiable manifold with a $(\phi, \psi, g)$ structure, then the two distributions $D_{l}$ and $D_{m}$ are orthogonal with respect to the metric $g$ at every point of $M^{n}$.

Proof. Take a point $P$ of $M^{n}$ and $\xi_{a}^{i}$ are vectors which span the vector space $V_{l}$ of $D_{l}$ at $P$ and $\xi_{p}^{i}$ are vectors which span the vector space $V_{m}$ of $D_{m}$ at $P$. Then, by virtue of (3.3) ${ }_{4}$,

$$
\begin{aligned}
g_{i j} \xi_{a} \xi_{p}^{j} & =\left(g_{\alpha \beta} \phi_{i}^{\alpha} \psi_{j}^{\beta}+g_{\alpha \beta} \psi_{i}^{\alpha} \psi_{j}^{\beta}\right) \xi_{a}^{i} \xi_{p}^{j} \\
& =0,
\end{aligned}
$$

which shows that $V_{l}$ and $V_{m}$ at $P$ are orthogonal to each other.
Q. E. D.

## 4. Associated tensor fields.

$8^{\circ}$. Suppose $M^{n}$ be a differentiable manifold with a $(\phi, \psi, g)$-structure. If we put

$$
\begin{equation*}
\phi_{i j}=g_{i n} \phi_{j}^{n}, \tag{4.1}
\end{equation*}
$$

then, by virtue of $(3.3)_{1}$, we get

$$
\begin{equation*}
\phi_{i j}=\varepsilon \phi_{j i} . \tag{4.2}
\end{equation*}
$$

Such a tensor will be called as $\varepsilon$-symmetric with respect to its indices. Of course, $\varepsilon$-symmetry means symmetry if $\varepsilon=+1$ and skew-symmetry if $\varepsilon=-1$. In the
same way, if we put

$$
\begin{equation*}
\psi_{i j}=g_{i h} \psi_{j}^{h}, \tag{4.3}
\end{equation*}
$$

then, by virtue of (3.3) $)_{2}, \boldsymbol{\psi}_{i j}$ is $\boldsymbol{\varepsilon}^{\prime}$-symmetric, i. e.

$$
\begin{equation*}
\psi_{i j}=\varepsilon^{\prime} \psi_{j i} . \tag{4.4}
\end{equation*}
$$

We can solve (4.1) and (4.3) with respect to $\phi_{j}^{i}$ and $\psi_{j}^{i}$ getting

$$
\left\{\begin{array}{l}
\phi_{j}^{i}=g^{i n} \phi_{h j}  \tag{45}\\
\psi_{j}^{i}=g^{i n} \psi_{h j}
\end{array}\right.
$$

Now, we put

$$
\begin{equation*}
\phi^{i j}=\phi_{h}^{i} g^{n j}, \phi_{j}^{i}=\phi^{i h} g_{l j,}, \tag{4.6}
\end{equation*}
$$

then

$$
\begin{align*}
\phi^{i j} & =g^{i k} \phi_{k h} g^{n_{j}} \\
& =\varepsilon g^{i k} \phi_{h k} g^{j h} \\
& =\varepsilon \phi_{k}^{j} g^{k i} \\
\therefore \quad \phi^{i j} & =\varepsilon \phi^{j i}, \tag{4.7}
\end{align*}
$$

So, $\phi^{i j}$ is an $\varepsilon$-symmetric contravariant tensor fields. In the same way, if we put

$$
\begin{equation*}
\psi^{i j}=\psi_{n}^{i} g^{h j}, \psi_{j}^{i}=\psi^{i h} g_{n j} \tag{4.8}
\end{equation*}
$$

then we get

$$
\begin{equation*}
\psi^{i j}=\varepsilon^{\prime} \psi^{j i} \tag{4.9}
\end{equation*}
$$

We call four tensor fields $\phi_{i j}, \psi_{i j}, \phi^{i j}, \psi^{i j}$ the associated tensor fields of the ( $\phi, \psi, g$ )-structure in consideration.

THEOREM 9. Let $\phi_{i j}, \psi_{i j}, \phi^{i j}, \psi^{i j}$ be associated tensor fields of a differentiable manifold with $a(\phi, \psi, g)$-structure. Then, the relations

$$
\left\{\begin{array}{l}
\phi^{i h} \psi_{h j}=0,  \tag{4.10}\\
\psi^{i h} \phi_{h j}=0
\end{array}\right.
$$

hold good.
PROOF. By virtue of (4.6) 2 , (4.3) and (2.3), we see that

$$
\begin{aligned}
\phi^{i h} \psi_{h j} & =\phi_{\alpha}^{i} g^{\alpha h} g_{h \beta} \psi_{j}^{\beta} \\
& =\phi_{a}^{i} \psi_{j}^{\alpha}=0 .
\end{aligned}
$$

Hence (4.10) ${ }_{1}$ is proved. In the same way, we can prove (4.10) $)_{2}$.
$9^{\circ}$. In the next place, we shall study the converse problem. We assume that $M^{n}$ be a differentiable manifold with a $\varepsilon$-symmetric tensor field $\phi_{i j}$ and $\varepsilon^{\prime}$-symmetric tensor field $\psi^{i j}$ such that

$$
\left\{\begin{align*}
\text { rank }\left|\phi_{i j}\right| & =l  \tag{4.11}\\
\text { rank }\left|\psi^{i j}\right| & =m,(l+m=n) \\
\psi^{i n} \phi_{h j} & =0
\end{align*}\right.
$$

hold good. And we shall study if we can find a positive definite Riemannian metric $g$ such that the tensor field

$$
\phi_{j}^{i}=g^{t h} \phi_{h j}, \psi_{j}^{i}=\psi^{i h} g_{h j}
$$

and $g$ define a $(\phi, \psi, g)$-structure over $M^{n}$ of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$ or not.
First, we introduce an arbitrary positive definite Riemannian metric $h$ over $M^{n}$. Then, $\phi_{i h}{ }^{h k} \phi_{k j}$ is a symmetric tensor field over $M^{n}$. Hence all roots of the characteristic equation

$$
\begin{equation*}
\left|\phi_{i h} h^{h k} \phi_{k j}+\rho h_{i j}\right|=0 \tag{4.12}
\end{equation*}
$$

are real. As $\left|\phi_{i j}\right|$ is of rank $l, 0$ is a root with multiplicity $m$. It is easily seen that all other roots have the opposite sign to $\varepsilon$.

Now, we denote all distinct non-zero roots by $\rho_{1}, \rho_{2}, \ldots \ldots, \rho_{r}$, their multiplicities by $\mu_{1}, \mu_{2}, \ldots \ldots, \mu_{r}$ and the characteristic spaces corresponding to the roots 0 and $\rho_{1}, \ldots \ldots, \rho_{r}$ by $V_{0}, V_{1}, \ldots \ldots, V_{r}$ respectively. Then, we see that

$$
\begin{gathered}
\operatorname{dim} V_{0}=m, \quad \operatorname{dim} V_{\lambda}=\mu_{\lambda}, \quad(\lambda=1, \ldots \ldots, r) \\
\mu_{1}+\mu_{2}+\cdots \cdots+\mu_{r}=l
\end{gathered}
$$

As is well known, all different characteristic spaces are orthogonal to each other with respect to the metric $h$.

In the same way, we take up the symmetric tensor field $\psi^{i k} h_{h k} \psi^{k j}$ and consider the characteristic equation

$$
\begin{equation*}
\left|\psi^{i h} h_{h k} \psi^{k j}+\sigma h^{i j}\right|=0 \tag{4.13}
\end{equation*}
$$

Then, it has 0 as a root of multiplicity $l$. Other roots are all real and they have the opposite $\operatorname{sign}$ to $\varepsilon^{\prime}$. We shall denote all distinct non-zero roots by $\sigma_{1}, \sigma_{2}, \ldots$ $\ldots, \sigma_{s}$, their multiplicities by $\nu_{1}, \nu_{2}, \ldots \ldots, \nu_{s}$ and the characteristic spaces of characteristic covectors corresponding to the roots 0 and $\sigma_{1}, \sigma_{2}, \ldots ., \sigma_{s}$ by $W_{0}^{*}$, $W_{1}^{*}, \ldots \ldots, W_{s}^{*}$. Then we see that

$$
\begin{gathered}
\operatorname{dim} W_{0}^{*}=l, \quad \operatorname{dim} W_{\lambda}^{*}=\nu_{\lambda}, \quad(\lambda=1,2, \ldots \ldots, s) \\
\nu_{1}+\nu_{2}+\cdots \cdots+\nu_{s}=m
\end{gathered}
$$

and $W_{u}^{*}, \ldots \ldots, W_{s}^{*}$ are orthogonal to each other.
$10^{\circ}$. Now, we consider a linear map of the tangent space $T_{P}$ at a point $P$ of $M^{n}$ to its dual space $T_{P}^{*}$ defined by

$$
\begin{equation*}
h: X^{i} \longrightarrow h_{i j} X^{j} . \tag{4.14}
\end{equation*}
$$

Then, we get the following
LEMMA 3. Let $h V_{\lambda}(\lambda=1, \ldots \ldots, r)$ be the image of the vector space $V_{\lambda}$ under the map $h$, then

$$
\begin{equation*}
W_{v}^{*}=h V_{1} \oplus \cdots \cdots \oplus h V_{r} . \tag{4.15}
\end{equation*}
$$

PROOF. Suppose that $X \in V_{\lambda}(\lambda=1,2, \ldots \ldots, r$ fixed $)$, then

$$
\phi_{i h} h^{h k} \phi_{k j} X^{j}=-\rho_{\lambda} h_{i j} X^{j} .
$$

By virtue of the last equation and (2.4) we see that

$$
\begin{aligned}
& \psi^{i h} h_{h k} \psi^{k j}\left(-\rho_{\lambda} h_{j m} X^{m}\right) \\
= & \psi^{i h} h_{h k} \psi^{k j} \phi_{j x} h^{\alpha \beta} \phi_{\beta \gamma} X^{\gamma}=0 .
\end{aligned}
$$

As $\rho_{\lambda} \neq 0$, this shows that $h X \in W_{0}^{*}$, hence $h V_{\lambda} \subset W_{0}^{*}$. Therefore, $h V_{1} \oplus \cdots$ $\cdots \oplus h V_{r} \cong W_{0}^{*}$. However, taking account of the dimensions of $W_{0}^{*}$ and $h V_{\mathbf{1}} \oplus \cdots$ $\cdots \oplus h V_{r}$, we see that the equality sign holds good.
Q. E. D.

LEMMA 4. The tangent space $T_{P}$ at any point $P$ of $M^{n}$ decomposes into the form

$$
\begin{equation*}
T_{P}=V_{1} \oplus \cdots \cdots \oplus V_{r} \oplus h^{-1} W_{1}^{*} \oplus \cdots \cdots \oplus h^{-1} W_{s}^{*} \tag{4.16}
\end{equation*}
$$

and any two of these component spaces are orthogonal to each other.
PROOF. Quite analogously to Lemma 3, we can prove that

$$
V_{0}=h^{-1} W_{1}^{*} \oplus \cdots \cdots \oplus h^{-1} W_{s}^{*}
$$

As $W_{1}^{*}, \ldots \ldots, W_{s}^{*}$ are orthogonal to each other, $h^{-1} W_{1}^{*}, \ldots \ldots, h^{-1} W_{s}^{*}$ are orthogonal to each other too. On the other hand, $V_{0} \oplus V_{1} \oplus \cdots \cdots \oplus V_{r}$ is a decomposition of the tangent space at any point of $M^{n}$. So our assertion is true. Q. E. D.

Now, suppose $\left\{U_{\alpha}\right\}$ be an open covering of $M^{n}$. We take frames $\boldsymbol{\xi}_{1} \ldots \ldots, \boldsymbol{\xi}_{n}$ over $U_{\alpha}$ such that

$$
\begin{array}{cc}
\xi_{1}, \ldots \ldots, \xi_{\mu_{1}} & \text { span } V_{1}, \\
\xi_{\mu_{1}+1}, \ldots \ldots, \xi_{\mu_{1}+\mu_{2}} & \text { span } V_{2}, \\
\ldots \ldots \ldots \ldots \ldots \ldots \ldots & \\
\xi_{\mu_{1}+\ldots+\mu_{r-1+1}, \ldots \ldots, \xi_{l}} & \text { span } V_{r},
\end{array}
$$

| $\xi_{l+1}, \ldots \ldots, \xi_{l+\nu_{1}}$ | span $h^{-1} W_{1}^{*}$, |
| :---: | :---: |
| $\ldots \ldots \ldots \ldots \ldots \ldots \ldots .$. |  |
| $\xi_{l+\nu_{1}+\ldots+\nu_{l-1}+1}, \ldots \ldots, \xi_{n}$ | span $h^{-1} W_{s}^{*}$. |

As the terminology which is used only in this section, we call such frame an adapted frame. Then, with respect to adapted frames, the components of the tensor field $h$ have the form :

where $h_{1}, \ldots \ldots, h_{r}, h_{r+1}, \ldots \ldots, h_{r+s}$ are matrices of order $\mu_{1}, \ldots \ldots, \mu_{r}$ and $\nu_{1}, \ldots \ldots, \nu_{s}$ respectively.
$11^{\circ}$. To find the form of components of $\phi_{i j}$ with respect to adapted frames, we consider the linear map $X \rightarrow \widetilde{X}$ of the tangent space defined by

$$
\begin{equation*}
\widetilde{X}^{i}=h^{i h} \phi_{h j} X^{j} \tag{4.18}
\end{equation*}
$$

Lemma 5. If $X \in V_{\lambda}(\lambda=1, \ldots \ldots, r)$, then $\widetilde{X} \in V_{\lambda}$.
PROOF. By virtue of

$$
\begin{equation*}
\phi_{i h} h^{n k} \phi_{k j} X^{j}=-\rho h_{i j} X^{j}, \tag{4.19}
\end{equation*}
$$

we can easily verify that

$$
\begin{aligned}
& \left(\phi_{i h} h^{h k} \phi_{k j}+\rho_{\lambda} h_{i j} h^{j l} \phi_{l m} X^{m}\right. \\
& \quad=\phi_{i h} h^{h k}\left(-\rho_{\lambda} h_{k m} X^{m}\right)+\rho_{\lambda} \phi_{i m} X^{m} \\
& \quad=0 .
\end{aligned}
$$

Q. E. D.

LEMMA 6. If $X \in V_{\lambda}(\lambda=1, \ldots \ldots, r)$, then

$$
\begin{equation*}
{\widetilde{X^{i}}}^{i}=-\rho_{\lambda} X^{i} \tag{4.20}
\end{equation*}
$$

PROOF. By virtue of (4.19), we see that

$$
\begin{aligned}
\widetilde{\widetilde{X}}^{i} & =h^{i n} \phi_{h j}\left(h^{j l} \phi_{l m} X^{m}\right) \\
& =h^{i n}\left(-\rho_{\lambda} h_{h m} X^{m}\right) \\
& =-\rho_{\lambda} X^{i} .
\end{aligned}
$$

Q. E. D.

Now, we put components of the tensor $\phi$ with respect to adapted frames in the form

where $\phi_{u v}(u, v=1, \ldots \ldots, r)$ is a $\left(\mu_{u}, \mu_{v}\right)$-matrix, $\phi_{r+u, v}(u=1, \ldots \ldots, s, v=1, \ldots$ $\ldots, r$ ) is a $\left(\nu_{u}, \mu_{v}\right)$-matrix, $\phi_{u, r+v}(u=1, \ldots \ldots, r, v=1, \ldots \ldots, s)$ is a ( $\mu_{u}, \nu_{v}$ )-matrix, and $\phi_{r+u, r+v}(u, v=1, \ldots \ldots, s)$ is a ( $\left.\nu_{u}, \boldsymbol{\nu}_{v}\right)$-matrix. Then, we see that

$$
h^{-1} \phi=\left(\begin{array}{cc}
h_{1}^{-1} \phi_{11} \ldots \ldots \ldots \ldots h_{1}^{-1} \phi_{1 r+s} \\
\vdots & \vdots \\
\vdots & \\
h_{r+s}^{-1} \phi_{r+s 1} \ldots \ldots \ldots h_{r+s}^{-1} \phi_{r+s}
\end{array}\right) .
$$

If we assume that $X \in V_{1}$, then its components with respect to adapted frames are of the form ( $X_{1}, 0, \ldots \ldots, 0$ ), where $X_{1}$ is a vector with $\mu_{1}$ components. So, in this case the vector $\widetilde{X}=h^{-1} \phi X$ has components $\left(h_{1}^{-1} \phi_{11} X_{1}, h_{2}^{-1} \phi_{21} X_{1}, \ldots\right.$ $\ldots, h_{r+s}^{-1} \phi_{r+s_{1}} X_{1}$ ). However, as $\widetilde{X} \in V_{1}$, we see that

$$
h^{-1} \phi_{\lambda 1} X_{1}=0, \quad(\lambda=2, \ldots \ldots, r+s) .
$$

Since $X_{1}$ is an arbitrary vector of $V_{1}$, we get

$$
\phi_{21}=\phi_{31}=\cdots \cdots=\phi_{r+s 1}=0
$$

In the same way, by considering vectors of $V_{2}, \ldots \ldots, V_{r}$, we get

$$
\begin{cases}\phi_{u v}=0 & (u \neq v ; u, v=1, \ldots \ldots, r) \\ \phi_{r+u, v}=0 & (u=1, \ldots \ldots, s, v=1, \ldots \ldots, r) .\end{cases}
$$

Lemma 7. If $Y \in h^{-1} W_{\lambda}^{*}(\lambda=1, \ldots \ldots, s)$, then

$$
\begin{equation*}
\phi_{i j} Y^{j}=0 . \tag{4.21}
\end{equation*}
$$

Proof. By assumption $Y^{*}=h Y \in W_{j}^{*}$, so

$$
\psi^{i n} h_{h k} \psi^{k j} Y_{j}^{*}=-\sigma_{\lambda} h^{i j} Y_{j}^{*}=-\sigma_{1} Y^{i} .
$$

By virtue of this and $(4.11)_{3}$, we see that

$$
-\sigma_{\lambda} \phi_{i j} Y^{j}=\phi_{i j} \psi^{j \alpha} h_{\alpha \beta} \psi^{\beta \gamma} Y_{\gamma}^{*}=0,
$$

which is to be proved.
Q. E. D.

Now, suppose $Y \in h^{-1} W_{1}^{*} \oplus \cdots \cdots \oplus h^{-1} W_{s}^{*}$, then the components of $Y$ with respect to an adapted frame have the form ( $0, \ldots \ldots, 0, Y_{r+1}, \ldots \ldots, Y_{r+s}$ ), where $Y_{r+u}(u=1, \ldots \ldots, s)$ are vectors with $\nu_{u}$ components. Hence, (4.21) shows us that

$$
\phi_{u r+v}=0, \quad \phi_{r+u r+v}=0,
$$

where $u=1, \ldots \ldots, r$, and $v=1, \ldots \ldots, s$. Consequently, we see that $\phi$ has the form

$$
\phi=\left(\begin{array}{ccc:c}
\phi_{1} & & & 0  \tag{4.22}\\
& \phi_{2} & & \\
& \ddots & & 0 \\
0 & & \phi_{r} & \\
\hdashline-\cdots & - & & \\
\hdashline & 0 & & 0
\end{array}\right)
$$

with respect to adapted frames, where $\phi_{\lambda}(\lambda=1, \ldots \ldots, r)$ is a $\varepsilon$-symmetric $\mu_{\lambda^{-}}$ matrix.

To find components of $\psi$ with respect to adapted frames, we consider the linear map

$$
\begin{equation*}
\widetilde{Y}_{j}^{*}=\psi^{i n} g_{n j} Y_{i}^{*} . \tag{4.23}
\end{equation*}
$$

Then, we can easily prove lemmas analogous to Lemmas 5, 6 and 7. Making use of these facts, we can similarly prove that the components of $\psi$ with respect to adapted frames have the following form :

$$
\psi=\left(\begin{array}{c:ccc}
0 & & 0 &  \tag{4.24}\\
\hdashline & \psi^{1} & \cdots & - \\
0 & \psi^{2} & & \\
& 0 & \ddots & 0 \\
& 0 & & \psi^{s}
\end{array}\right),
$$

where $\psi^{\lambda}(\lambda=1, \ldots \ldots, s)$ is a $\varepsilon^{\prime}$-symmetric $\nu_{\lambda}$-matrix.
$12^{\circ}$. Now, making use of the decomposition of the tangent spaces stated in Lemma 4, let us introduce a new Riemannian metric $g$ over $M^{n}$ by

$$
g=\left(\begin{array}{cccccc}
g_{1} & & & & &  \tag{4.25}\\
& \ddots & & & & \\
& & \ddots & & & \\
\\
& & & g_{r} & & \\
\\
& & & & g_{r+1} & \\
\\
& & & & & \\
& & & & \\
& & & & g_{r+s}
\end{array}\right)
$$

where we have put

$$
\begin{cases}g_{u}=\sqrt{-\varepsilon \rho_{u}} h_{u} & (u=1, \ldots \ldots, r),  \tag{4.26}\\ g_{r+v}=h_{r+v} / \sqrt{-\varepsilon^{\prime} \sigma_{v}} & (v=1, \ldots \ldots, s) .\end{cases}
$$

As characteristic roots and characteristic spaces are independent upon the choice of coordinate neighborhoods, the Riemannian metrics defined thus for every coordinate neighborhood of the covering $\left\{U_{a}\right\}$ constitute a single globally defined Riemannian metric over $M^{n}$.

We put
(4.27)

$$
\phi_{j}^{i}=g^{i n} \phi_{h j}
$$

Then, $\phi_{j}^{i}$ is a globally defined tensor field over $M^{n}$ too. By virtue of (4.22) and (4.25), we see that $\phi_{j}^{i}$ has components of the form

$$
\left(\phi_{j}^{i}\right)=\left(\begin{array}{ccc:c}
g_{1}^{-1} \phi_{1} & & 0 & 0  \tag{4.28}\\
& \cdot & & \\
0 & & \dot{g}_{r}^{-1} \phi_{r} & \\
\hdashline 0 & & 0
\end{array}\right)
$$

with respect to adapted frames. We now define a modified linear map of (4.18) by

$$
\begin{equation*}
\hat{X}^{i}=g^{i n} \phi_{h j} X^{j} \tag{4.29}
\end{equation*}
$$

If we assume $X \in V_{1}$, then with respect to adapted frames, we see that

$$
\begin{aligned}
\hat{X}^{u} & =g^{u v} \phi_{v w} X^{w} \quad\left(u, v, w=1, \ldots \ldots, \mu_{1}\right) \\
& =\frac{1}{\sqrt{-\varepsilon \rho_{1}}} \widetilde{X}^{u} .
\end{aligned}
$$

Hence, we have

$$
\hat{X}^{u}=\frac{1}{-\varepsilon \rho_{1}} \widetilde{\widetilde{X}}^{u}=\varepsilon X^{u} .
$$

This shows that

$$
\left(g_{1}^{-1} \phi_{1}\right)^{2}=\varepsilon
$$

Similar formulas hold good for matrices $g_{\lambda}^{-1} \phi_{\lambda}(\lambda=1, \ldots \ldots, r)$ too. Hence, we see that

$$
\left(\phi_{j}^{i} \phi_{k}^{j}\right)=\left(\begin{array}{ccc:c}
\varepsilon & & 0 & 0  \tag{4.30}\\
0 & \cdot & \varepsilon & 0 \\
\hdashline 0 & \cdots & \varepsilon & \\
\hdashline & 0 & & 0
\end{array}\right)
$$

Finally, we put

$$
\begin{equation*}
\psi_{j}^{i}=\psi^{i n} g_{n j}, \tag{4.31}
\end{equation*}
$$

then, $\boldsymbol{\psi}_{j}^{i}$ is a globally defined tensor field over $M^{n}$. And, in the same way as above, we can prove that ( $\psi_{j}^{i}$ ) and ( $\psi_{j}^{i} \psi_{k}^{j}$ ) have the forms
and

$$
\left(\psi_{j}^{j} \psi_{k}^{j}\right)=\left(\begin{array}{c:ccc}
0 & & 0 &  \tag{4.33}\\
\hdashline & \varepsilon^{\prime} & & 0 \\
0 & \cdot & \cdot & \\
& 0 & & \varepsilon^{\prime}
\end{array}\right)
$$

with respect to adapted frames.
Combining (4.30) and (4.33), we see that

$$
\begin{equation*}
\varepsilon \phi_{j}^{i} \phi_{k}^{j}+\varepsilon^{\prime} \psi_{j}^{i} \psi_{k}^{j}=\delta_{k}^{i} \tag{4.34}
\end{equation*}
$$

holds good with respect to adapted frames. However, as (4.34) is a tensor equation, it does hold for any frame, especially for natural frames. We can easily prove, by virtue of (4.28) and (4.32) that

$$
\begin{equation*}
\phi_{j}^{i} \psi_{k}^{j}=0, \quad \psi_{j}^{i} \phi_{k}^{j}=0 . \tag{4.35}
\end{equation*}
$$

As the ranks of $\left|\phi_{j}^{i}\right|,\left|\psi_{j}^{i}\right|$ are $l$ and $m$ respectively, we see that our $\phi_{j}^{i}$ and $\psi_{j}^{i}$ give a $(\phi, \psi)$-structure of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$ to $M^{n}$.

We can easily prove that

$$
\begin{align*}
\phi_{i j} & =g_{i l} \phi_{j}^{h}=\varepsilon g_{j h} \phi_{i}^{h}  \tag{4.36}\\
\psi^{i j} & =\psi_{h}^{i} g^{n j}=\varepsilon^{\prime} \psi_{h}^{j} g^{h i}
\end{align*}
$$

and

$$
g_{i j} \phi_{h}^{i} \psi_{k}^{j}=0 .
$$

Finally, as we have shown in the proof of Theorem 6, the relation

$$
g_{i j} \phi_{h}^{i} \phi_{k}^{j}+g_{i j} \psi_{h}^{i} \psi_{k}^{j}=g_{h k}
$$

follows from (4.35) and (4.36). Consequently, we get the following
THEOREM 10. Suppose $M^{n}$ be a differentiable manifold and there exist $\varepsilon$-symmetric tensor field $\phi_{i j}$ and $\varepsilon^{\prime}$-symmetric tensor field $\psi^{i j}$ such that

$$
\operatorname{rank}\left|\phi_{i j}\right|=l, \quad \text { rank }\left|\psi^{i j}\right|=m, \quad l+m=n
$$

and

$$
\psi^{i n} \phi_{h g}=0
$$

hold good. Then, we can find a positive definite Riemannian metric $g$ such that the tensor fields

$$
\phi_{j}^{i}=g^{i n} \phi_{h j}, \quad \psi_{j}^{i}=\psi^{i n} g_{h j}
$$

and $g$ define $a(\phi, \psi, g)$-structure over $M^{n}$ of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$.
COROLLARY. Suppose $M^{n}$ be a differentiable manifold and there exists $\varepsilon$-symmetric tensor field $\phi_{i j}$ over $M^{n}$ such that

$$
\operatorname{rank}\left|\phi_{i j}\right|=l, \quad l \leqq n
$$

Then we can find a symmetric tensor field $\psi^{i j}$ such that

$$
\operatorname{rank}\left|\psi^{i j}\right|=m, \quad m=n-l
$$

and a positive definite Riemannian metric $g$ so that

$$
\phi_{j}^{2}=g^{i n} \phi_{h j}, \quad \psi_{j}^{i}=\psi^{i h} g_{h j}
$$

and $g$ define $a(\phi, \psi, g)$-structure of type $(\operatorname{sgn} \varepsilon,+)$.
Proof. By assumption, at any point $P$ of $M^{n}$, the set of vectors such that

$$
\phi_{i j} \xi^{j}=0
$$

span an $m$-dimensional vector subspace $V_{m}(P)$ of the tangent space $T_{P}$ at $P$. Now, suppose $h$ be a positive definite Riemannian metric $h$ over $M^{n}$ and take $m$ orthonormal vectors $\xi_{p}^{i}$ with respect to the metric $h$ and put

$$
\psi^{i j}(P)=\sum_{p} \xi_{p}^{\prime} \xi_{p}^{\prime}
$$

When $P$ moves over $M^{n}$, it is evident that $\psi^{i j}$,s constitute a globally defined symmetric tensor field over $M^{n}$. We can easily verify that

$$
\begin{aligned}
\operatorname{rank}\left|\psi^{i j}\right| & =m, \\
\psi^{i h} \phi_{h j} & =0 .
\end{aligned}
$$

Hence, by virtue of Theorem 10, we can conclude that our assertion is true.
Q. E. D.

## 5. The structure groups of the tangent bundles of manifolds with ( $\phi, \psi, g$ )-structures.

$13^{\circ}$. The structure group of the tangent bundle $T\left(M^{n}\right)$ of a differentiable
manifold $M^{n}$ is in general the general linear group $G L(n)$. However, we can prove the following

THEOREM 11. Suppose $M^{n}$ be a differentiable manifold with ( $\phi, \psi, g$ )structure of type $\left(\operatorname{sgn} \varepsilon, \operatorname{sgn} \varepsilon^{\prime}\right)$. Then, the structure group of the tangent bundle $T\left(M^{n}\right)$ is reducible to the following one:

$$
\begin{equation*}
U(l / 2) \times U(m / 2) \tag{i}
\end{equation*}
$$

$$
\text { if } \varepsilon=\varepsilon^{\prime}=-1
$$

$$
\begin{equation*}
U(l / 2) \times O\left(m_{1}\right) \times O\left(m_{2}\right) \quad \text { if } \varepsilon=-1, \varepsilon^{\prime}=+1 \tag{ii}
\end{equation*}
$$

$$
\begin{equation*}
O\left(l_{1}\right) \times O\left(l_{2}\right) \times O\left(m_{1}\right) \times O\left(m_{2}\right) \quad \text { if } \varepsilon=\varepsilon^{\prime}=+1 \tag{iii}
\end{equation*}
$$

where $l_{1}\left(l_{2}\right)$ is the number of negative (positive) roots of the characteristic equation of $\phi_{j}^{i}$ and $m_{1}\left(m_{2}\right)$ is that of $\psi_{j}^{i}$.

PROOF. We shall prove only the case (ii). The other cases can be proved quite analogously.

We take sufficiently fine open covering $\left\{U_{\alpha}\right\}$ of $M^{n}$ and determine in every $U_{\alpha}$ suitable frames. To do so, we take first a unit vector field $\xi_{1}^{i}$ over $U_{\alpha}$ contained in $D_{l}$ and put

$$
\begin{equation*}
\xi_{1 *}^{l}=\phi_{j}^{\prime} \xi_{1}^{\prime}, \quad 1^{*}=l^{\prime}+1, \quad l^{\prime}=l / 2 . \tag{5.1}
\end{equation*}
$$

Then, we can easily see that $\xi_{1 *}^{i}$ is a unit vector field orthogonal to $\xi_{1}^{i}$ and contained in $D_{l}$. Secondly, if we take a unit vector field $\xi_{2}^{i}$ orthogonal to $\xi_{1}^{i}$ and $\xi_{1 *}^{i}$ and contained in $D_{l}$, then

$$
\begin{equation*}
\xi_{2^{*}}^{i}=\phi_{j}^{\prime} \xi_{2}^{\prime}, \quad 2^{*}=l^{\prime}+2 \tag{5.2}
\end{equation*}
$$

is a unit vector field orthogonal to $\xi_{1}^{i}, \xi_{2}^{i}, \xi_{1^{*}}^{i}$ and contained in $D_{l}$. Continuing this process, we can find orthonormal vector fields $\xi_{\lambda}^{\prime}\left(\lambda=1, \ldots \ldots, l^{\prime}\right)$ and

$$
\begin{equation*}
\xi_{\lambda^{*}}^{i}=\phi_{j}^{i} \xi_{\lambda}^{j}, \quad \lambda^{*}=l^{\prime}+\lambda \tag{5.3}
\end{equation*}
$$

so that they span $D_{l}$ in $U_{\alpha}$. By virtue of (2.5) we can solve (5.3) as follows :

$$
\begin{equation*}
\xi_{\lambda}^{i}=-\phi_{j}^{i} \xi_{\lambda^{*}}^{i} . \tag{5.4}
\end{equation*}
$$

Next, we consider the characteristic equation

$$
\begin{equation*}
\left|\psi_{j}^{i}+\sigma \delta_{i}^{j}\right|=0 . \tag{5.5}
\end{equation*}
$$

Then, the characteristic roots are $0,-1$ or +1 . The characteristic space corresponding to 0 is the vector space $V_{l}$ of the distribution $D_{l}$. We denote the multiplicities of the roots -1 and +1 by $m_{1}$ and $m_{2}$ and the characteristic spaces corresponding to -1 and +1 by $W_{-}$and $W_{+}$. Then, as is easily seen, $V_{l}, W_{-}$and $W_{+}$are orthogonal to each other.

Now, we take orthonormal frames $\left(\xi_{\lambda}^{i}, \xi_{\lambda^{*}}^{i}, \xi_{\boldsymbol{\mu}}^{i}, \xi_{s}^{i}\right)$ over $U_{\alpha}$ so that $\xi_{\lambda}^{i}, \xi_{\lambda^{*}}^{i}$
are related by (5.3) and span $V_{l}, \xi_{m}^{\prime}\left(M=l+1, \ldots \ldots, l+m_{1}\right)$ span $W_{-}$and $\xi_{s}^{l}$ $\left(S=l+m_{1}+1, \ldots \ldots, n\right)$ span $W_{+}$. Then, we can easily see that $g, \phi$ and $\psi$ have the following forms with respect to such frames:

$$
\begin{align*}
& g=\left(\begin{array}{cccc}
E_{l^{\prime}} & & & 0 \\
& E_{l^{\prime}} & & \\
0 & & E_{m_{1}} & E_{m_{2}}
\end{array}\right),  \tag{5.6}\\
& \phi=\left(\begin{array}{ccc}
0 & -E_{l^{\prime}} & 0 \\
E_{l^{\prime}} & 0 & 0 \\
\hdashline-- & 0 & 0
\end{array}\right),  \tag{5.7}\\
& \psi=\left(\begin{array}{c:cc}
0 & 0 & \\
\hdashline 0 & -E_{m_{1}} & 0 \\
0 & 0 & E_{m_{2}}
\end{array}\right) . \tag{5.8}
\end{align*}
$$

Suppose $U_{\alpha} \cap U_{\beta}$ is not empty and $\overline{\xi_{\lambda}^{i}}, \overline{\xi_{\lambda *}}, \overline{\xi_{m f}^{i}}, \overline{\xi_{s}^{i}}$ be vector fields over $U_{\beta}$ defined in the same way as above, then over $U_{\alpha} \cap U_{\beta}$ we get

$$
\begin{aligned}
& \left\{\begin{array}{l}
\xi_{\mu}^{i}=a_{\mu}^{\lambda} \bar{\xi}_{\lambda}^{i}+b_{\mu}^{\lambda} \overline{\xi_{\lambda *}^{i}} \\
\xi_{\mu^{*}}^{i}=c_{\mu}^{\lambda} \bar{\xi}_{\lambda}^{i}+d_{\mu}^{\lambda} \bar{\xi}_{\lambda^{*}},
\end{array}\right. \\
& \left\{\begin{array}{l}
\xi_{N}^{i}=u_{N}^{n \overline{\xi_{i}^{i}}} \\
\xi_{T}^{i}=u_{T}^{S} \xi_{S}^{i},
\end{array}\right.
\end{aligned}
$$

where $\left(\begin{array}{ll}a & b \\ b & d\end{array}\right)$ and $\left(U_{N}^{Z}\right)$ and $\left(U_{T}^{S}\right)$ are orthogonal matrices. However, by virtue of (5.3), we can easily see that $d=a, c=-b$. Hence, the structure group of $T\left(M^{n}\right)$ in consideration is reducible to $U\left(l^{\prime}\right) \times O\left(m_{1}\right) \times O\left(m_{2}\right)$.
14. Converse to Theorem 11, we can prove the following

THEOREM 12. Suppose $M^{n}$ be a differentiable manifold such that the structure group of the tangent bundle $T\left(M^{n}\right)$ is reducible to

$$
\begin{equation*}
U\left(l^{\prime}\right) \times U\left(m^{\prime}\right), \quad 2\left(l^{\prime}+m^{\prime}\right)=n, \quad \text { or } \tag{i}
\end{equation*}
$$

(ii) $\quad U\left(l^{\prime}\right) \times O\left(m_{1}\right) \times O\left(m_{2}\right), \quad 2 l^{\prime}+m_{1}+m_{2}=n$, or
(iii) $\quad O\left(l_{1}\right) \times O\left(l_{2}\right) \times O\left(m_{1}\right) \times O\left(m_{2}\right), \quad l_{1}+l_{2}+m_{1}+m_{2}=n$.

Then, we can introduce $(\phi, \psi, g)$-structure over $M^{n}$ of type $(-,-),(-,+)$ or $(+,+)$ according as the structure group is of type (i), (ii) or (iii).

PROOF. We shall prove only the case (ii). The other cases can be proved
quite analogously. Let $\left\{U_{\alpha}\right\}$ be sufficiently fine open covering of $M^{n}$ by coordinate neighborhoods. Then, in every $U_{\alpha}$, we can take frames $\xi_{\lambda}^{i}, \xi_{\lambda^{*}}^{i}(\lambda=1, \ldots$ $\left.\ldots, l^{\prime}, \lambda^{*}=l^{\prime}+\lambda\right), \quad \xi_{m}^{i}\left(M=l+1, \ldots \ldots, l+m_{1}\right), \quad \xi_{S}^{i}\left(S=l+m_{1}+1, \ldots \ldots, n\right)$ so that the transformation of frames of $U_{\alpha}$ and of $U_{\beta}$ over non-empty $U_{\alpha} \cap U_{\beta}$ is given by an orthogonal transformation of the form

$$
\begin{aligned}
& \left\{\begin{array}{l}
\xi_{\mu}^{i}=a_{\mu}^{\lambda} \overline{\xi_{\lambda}^{i}}+b_{\mu}^{\lambda} \bar{\xi}_{\lambda^{*}}, \\
\xi_{\mu^{*}}^{i}=-b_{\mu}^{\lambda} \bar{\xi}_{\lambda}^{i}+a_{\mu}^{\lambda} \bar{\xi}_{\lambda^{*}},
\end{array}\right. \\
& \left\{\begin{array}{l}
\xi_{N}^{i}=u_{N}^{\pi} \bar{\xi}_{i}^{i}, \\
\xi_{T}^{i}=u_{T}^{s} \bar{\xi}_{s .} .
\end{array}\right.
\end{aligned}
$$

We denote the inverse matrix of $\left(\xi_{\alpha}^{i}\right)$ by $\left(\eta_{j}^{\alpha}\right)$ and define over every $U_{\alpha}$ tensor fields by

$$
\begin{aligned}
g_{i j} & =\sum_{\alpha=1}^{n} \eta_{i}^{\alpha} \eta_{j}^{\alpha}, \\
\phi_{j}^{i} & =-\xi_{\lambda}^{i} \eta_{j}^{\lambda^{*}}+\xi_{\lambda^{*}}^{i} \eta_{j}^{\lambda}, \\
\psi_{j}^{i} & =-\xi_{M}^{i} \eta_{j}^{M}+\xi_{s}^{i} \eta_{j}^{s},
\end{aligned}
$$

then all $g_{i j}$ 's, all $\phi_{j}^{\prime \prime}$ s and all $\psi_{j}^{i}$,s corresponding to $U_{a}$ 's constitute global tensor fields $g, \phi$ and $\psi$ respectively. This can be easily proved by virtue of the relations

$$
\begin{aligned}
& \left\{\begin{array}{l}
\overline{\bar{\eta}_{j}^{\lambda}}=a_{\mu}^{\lambda} \eta_{j}^{\mu}-b_{\mu}^{\lambda} u_{j}^{u^{*}}, \\
\boldsymbol{\eta}_{j}^{\lambda_{i}^{*}}
\end{array}=b_{\mu}^{\lambda} \eta_{i}^{\mu}+a_{\mu}^{\lambda} \eta_{i}^{u^{*},},\right.
\end{aligned}\left\{\begin{array}{l}
\overline{\boldsymbol{\eta}_{i}^{\mu}}=u_{N}^{M} \eta_{j}^{N}, \\
\overline{\boldsymbol{\eta}_{j}^{s}}=u_{\eta}^{S} \eta_{j}^{T} .
\end{array}\right.
$$

As $g, \phi, \psi$ have components of the form (5.5), (5.6) and (5.7) with respect to our frames in consideration, we can easily verify that (2.3), (2.5), (3.3) $(\varepsilon=-1$, $\varepsilon^{\prime}=1$ ) hold good. However, these equations are all tensor equations. Hence, they all hold good for any frames, especially for natural frames. Consequently, our Theorem is proved.
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