# ON A BARGMANN-TYPE TRANSFORM AND A HILBERT SPACE OF HOLOMORPHIC FUNCTIONS 

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0. Introduction. In [1], Bargmann studied an integral transform of $L^{2}\left(\boldsymbol{R}^{n}\right)$ onto a Hilbert space consisting of entire holomorphic functions on $\boldsymbol{C}^{n}$. His transform may be regarded as a half-form pairing between real and complex polarizations of $\boldsymbol{R}^{2 n} \cong \boldsymbol{C}^{n}$ (see [8, §2]). In [7], Rawnsley showed that $\stackrel{\circ}{T}^{*} S^{n-1}$ (the cotangent bundle of the ( $n-1$ )-sphere minus its zero section) has a Kaehler structure with the Kaehler form equal to the natural symplectic form. Furthermore, he studied in [8] the halfform pairing between real and complex polarizations of $\stackrel{\circ}{T}^{*} S^{n-1}$, but it is not unitary. Also, we know that there does not exist a distinguished kernel, the definition of which is given in [2, IV. 5], for these polarizations. More precisely, there does exist a "distinguished kernel" defined in a neighborhood of the diagonal of $\stackrel{\circ}{T}^{*} S^{n-1} \times \stackrel{\circ}{T}^{*} S^{n-1}$, but it does not extend globally. This "kernel", however, suggests us to consider an integral transform:

$$
\mathscr{F}: f \mapsto \hat{f}(z)=\int_{S^{n-1}} e^{x \cdot z} f(x) d S(x)
$$

where $z \in \boldsymbol{C}^{n}, z^{2}=0$ (for the notations, see Section 1). Incidentally, transformations of the same form as $\mathscr{F}$ have been studied by several authors (see, for example, [3, §4], [6, §7], [4, Theorem 2.10], [8, p. 175] and $[5, \S 4])$. In the present note, motivated by these works, we consider the integral transform $\mathscr{F}$ of $L^{2}\left(S^{n-1}\right)$ into a space consisting of holomorphic functions on the Kaehler manifold $\stackrel{\circ}{T}^{*} S^{n-1}=\left\{z \in C^{n} \mid z^{2}=0, z \neq 0\right\}$. $\mathscr{F}$ is injective. In Section 2, we construct, in the case of even-dimensional spheres, a "Plancherel measure" on $\stackrel{\circ}{T}^{*} S^{n-1}$ to describe the image of $L^{2}\left(S^{n-1}\right)$ under this transform. The "inversion formula" is also obtained. As an application, we give in Section 3 an integral representation of a oneparameter group of unitary transformations on $L^{2}\left(S^{n-1}\right)$ generated by a pseudo-differential operator $-i\left\{\Delta+(n-2)^{2} / 4\right\}^{1 / 2}$, where $\Delta$ is the LaplaceBeltrami operator on $S^{n-1}$ (cf. [8, p. 177]).

[^0]For the sake of simplicity, we shall assume $n \geqq 3$ throughout this paper. In Sections 2 and 3, we furthermore assume that $n$ is odd. The reason why we exclude the case of even $n$ is that, in this case, we cannot identify the function $\rho_{n}$ which satisfies the equation in Lemma 2.1. The Lie differentiation with respect to a vector field $X$ is denoted by $\mathscr{L}_{x}$. Volume forms and measures are used interchangeably.

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Added September 6, 1985. R. Wada has informed us that she is able to identify the function $\rho_{n}$ also for even $n$, and to remove the assumption in Section 2 that $n$ is odd.

1. Preliminaries. Let $\boldsymbol{R}^{n}=\left\{x=\left(x_{1}, \cdots, x_{n}\right)\right\}$ be the $n$-space with the inner product $x \cdot y=\sum x_{i} y_{i}$ and the norm $|x|=(x \cdot x)^{1 / 2}$, and $S^{n-1}=$ $\left\{x \in \boldsymbol{R}^{n}| | x \mid=1\right\}$ be the unit sphere. The volume element on $S^{n-1}$ is denoted by $d S$. The volume of $S^{n-1}$ is given by $\operatorname{vol}\left(S^{n-1}\right)=2 \pi^{n / 2} / \Gamma(n / 2)$. Let $L^{2}\left(S^{n-1}\right)$ be the Hilbert space of square-integrable functions on $S^{n-1}$ with the following inner product and norm:

$$
\langle f, g\rangle_{S}=\int_{S^{n-1}} \bar{f} g d S, \quad\|f\|_{S}=\langle f, f\rangle_{S}^{1 / 2}
$$

The subspace of $L^{2}\left(S^{n-1}\right)$ consisting of spherical harmonics of degree $m$ is denoted by $H_{m}\left(S^{n-1}\right), m=0,1,2, \cdots$. The following is well-known (see, for example, $[4, \S 3]):$

Lemma 1.1. (i) $\operatorname{dim} H_{m}\left(S^{n-1}\right)=(2 m+n-2) \Gamma(m+n-2) / \Gamma(n-1) \times$ $\Gamma(m+1)$.
(ii) The subspaces $H_{m}\left(S^{n-1}\right), m=0,1,2, \cdots$, are mutually orthogonal with respect to the inner product $\langle,\rangle_{s}$.
(iii) Let $f_{m} \in H_{m}\left(S^{n-1}\right), m=0,1,2, \cdots$. Then $f=\sum f_{m}$ belongs to $L^{2}\left(S^{n-1}\right)$ if and only if $\sum\left\|f_{m}\right\|_{s}^{2}<\infty$, and in that case, $\|f\|_{s}^{2}=\sum\left\|f_{m}\right\|_{s}^{2}$.
(iv) For any $z \in \boldsymbol{C}^{n}$ with $z^{2}=0$, the function on $S^{n-1}, x \mapsto(x \cdot z)^{m}$, belongs to $H_{m}\left(S^{n-1}\right)$, where $z^{2}=\sum z_{i}^{2}$ and $x \cdot z=\sum x_{i} z_{i}$. Furthermore, $H_{m}\left(S^{n-1}\right)$ is spanned by these functions.

For any $1 \leqq i_{1}, \cdots, i_{m} \leqq n$, let us define an element of $H_{m}\left(S^{n-1}\right)$ by

$$
h_{i_{i} \cdots i_{m}}=\left(\prod_{k=0}^{m-1}(2-n-2 k)\right)^{-1}\left[\partial^{m}|x|^{2-n} / \partial x_{i_{1}} \cdots \partial x_{i_{m}}\right] \mid S^{n-1},
$$

where we assume $n \geqq 3$. Note that $H_{m}\left(S^{n-1}\right)$ is spanned by these functions.

LEMMA 1.2. (i) $h_{i_{1} \cdots i_{m}}(x)=x_{i_{1}} \cdots x_{i_{m}}-1 /(2(2 m+n-4)) \sum_{a \neq b} \delta_{i_{a} i_{b}} x_{i_{1}} \cdots$ $\hat{x}_{i_{a}} \cdots \hat{x}_{i_{b}} \cdots x_{i_{m}}+h^{\prime}(x)$, where $h^{\prime} \in \bigoplus_{k=0}^{m-4} H_{k}\left(S^{n-1}\right)$.
(ii) $x_{j} h_{i_{1} \cdots i_{m}}=h_{i_{1} \cdots i_{m} j}+1 /(2 m+n-2) \sum_{a=1}^{m} \delta_{i_{a} j} h_{i_{1} \cdots \hat{i}_{a} \cdots i_{m}}-1 /((2 m+n-$ 2)( $2 m+n-4$ )) $\sum_{a \neq b} \delta_{i_{a} i_{b}} h_{i_{1} \cdots \hat{i}_{a} \cdots \hat{i}_{b} \cdots i_{m} j}$, where $x_{j}$ denotes the function $x \mapsto x_{j}$ on $S^{n-1}$.

For the proof of (ii), recall that the multiplication of $x_{j}$ is a symmetric operator on $L^{2}\left(S^{n-1}\right)$. Then, from (i) and the orthogonality of the subspaces $H_{m}\left(S^{n-1}\right), m=0,1,2, \cdots$, we have $x_{j} h_{i_{1} \cdots i_{m}} \in H_{m+1}\left(S^{n-1}\right) \oplus H_{m-1}\left(S^{n-1}\right)$.

Lemma 1.3. For any $m=1,2, \cdots$, we have

$$
\begin{aligned}
\left\langle h_{i_{1} \cdots i_{m}},\right. & \left.h_{j_{1} \cdots j_{m}}\right\rangle_{S} \\
= & \frac{1}{2 m+n-2} \sum_{a=1}^{m} \delta_{i_{a} j_{m}}\left\langle h_{i_{1} \cdots \hat{i}_{a} \cdots i_{m}}, h_{j_{1} \cdots j_{m-1}}\right\rangle_{S} \\
& -\frac{1}{(2 m+n-2)(2 m+n-4)} \sum_{a \neq b} \delta_{i_{a} i_{b}}\left\langle h_{i_{1} \cdots \hat{i}_{a} \cdots \hat{i}_{b} \cdots i_{m} i_{m}}, h_{j_{1} \cdots j_{m-1}}\right\rangle_{S}
\end{aligned}
$$

Proof. Let $\xi_{j}$ denote the restriction to $S^{n-1}$ of the vector field $\sum_{i=1}^{n}\left(\delta_{j i}-x_{j} x_{i}\right) \partial / \partial x_{i}$ on $\boldsymbol{R}^{n}$. Then $\xi_{j}$ is tangent to $S^{n-1}$. Since $\mathscr{L}_{\xi j} d S=$ $-(n-1) x_{j} d S$, we have from

$$
\int_{S^{n-1}} \mathscr{L}_{\xi_{j}}(\bar{f} g d S)=0
$$

that $\xi_{j}-(n-1) x_{j} / 2$ is a skew-symmetric operator on $\left(C^{\infty}\left(S^{n-1}\right),\langle,\rangle_{s}\right)$. Then, by Lemma 1.2, we have

$$
\begin{aligned}
\left(\xi_{j}-\right. & \left.\frac{n-1}{2} x_{j}\right) h_{i_{1} \cdots i_{m}} \\
& =-\left(m+\frac{n-1}{2}\right) h_{i_{1} \cdots i_{m} j}+\frac{2 m+n-3}{2(2 m+n-2)} \sum_{a=1}^{m} \delta_{i_{a} j} h_{i_{1} \cdots \hat{i}_{a} \cdots i_{m}} \\
& \quad-\frac{2 m+n-3}{2(2 m+n-2)(2 m+n-4)} \sum_{a \neq b} \delta_{i_{a} i_{b}} h_{i_{1} \cdots \hat{i}_{a} \cdots \hat{i}_{b} \cdots i_{m} j}
\end{aligned}
$$

Using this formula on both sides of the equation:

$$
\begin{aligned}
\left\langle\left(\xi_{j_{m}}-\frac{n-1}{2} x_{j_{m}}\right)\right. & \left.h_{i_{1} \cdots i_{m}}, h_{j_{1} \cdots j_{m-1}}\right\rangle_{s} \\
& =-\left\langle h_{i_{1} \cdots i_{m}},\left(\xi_{j_{m}}-\frac{n-1}{2} x_{j_{m}}\right) h_{j_{1} \cdots j_{m-1}}\right\rangle_{s},
\end{aligned}
$$

we obtain our lemma.
Now, we shall consider an integral transform (cf. [3, §4] and [6, § 7]): For any $f \in L^{2}\left(S^{n-1}\right)$ and $z \in \boldsymbol{C}^{n}$, let us define

$$
\hat{f}(z)=\int_{s^{n}-1} e^{x^{x \cdot z} f(x) d S(x) .}
$$

Then we have:
Lemma 1.4. (i) $\hat{f}$ is an entire function on $\boldsymbol{C}^{n}$.
(ii) $|\hat{f}(z)| \leqq\left(\operatorname{vol}\left(S^{n-1}\right)\right)^{1 / 2}\|f\|_{S} e^{|\mathrm{Re} z|}$.
(iii) If $f \in L^{2}\left(S^{n-1}\right), f=\sum f_{m}$ with $f_{m} \in H_{m}\left(S^{n-1}\right)$, then $\sum \hat{f}_{m}$ converges to $\hat{f}$ uniformly on any bounded set in $\boldsymbol{C}^{n}$.

Proposition 1.5. If $z^{2}=0$, then

$$
\hat{h}_{i_{1} \cdots i_{m}}(z)=\frac{\operatorname{vol}\left(S^{n-1}\right) \Gamma(n / 2)}{2^{m} \Gamma(m+n / 2)} z_{i_{1}} \cdots z_{i_{m}}
$$

Proof. We shall prove this by induction on $m$. If $m=0$, then both sides of the equation are equal to $\operatorname{vol}\left(S^{n-1}\right)$. Now, let $m>0$ and assume that the proposition holds for $m-1$. Since, by (i) of Lemma 1.2,

$$
\int_{S^{n-1}} x_{j_{1}} \cdots x_{j_{m}} h_{i_{1} \cdots i_{m}}(x) d S(x)=\left\langle h_{i_{1} \cdots i_{m}}, h_{j_{1} \cdots j_{m}}\right\rangle_{S},
$$

we have

$$
\hat{h}_{i_{1} \cdots i_{m}}(z)=\frac{1}{m!} \sum_{j_{1}, \cdots, j_{m}}\left\langle h_{i_{1} \cdots i_{m}}, h_{j_{1} \cdots j_{m}}\right\rangle z_{j_{1}} \cdots z_{j_{m}} .
$$

Then, since $z^{2}=0$, by Lemma 1.3 and the induction assumption, we see that the proposition holds also for $m$.

Let $\pi: \stackrel{\circ}{T^{*}} S^{n-1} \rightarrow S^{n-1}$ be the bundle consisting of non-zero cotangent vectors to $S^{n-1}$. The canonical one-form $\theta$ on $T^{*} S^{n-1}$ is defined by $\theta_{\alpha}(X)=$ $\alpha\left(\pi_{*} X\right)$ for any $\alpha \in \stackrel{\circ}{T}^{*} S^{n-1}$ and $X \in T_{\alpha}\left(\stackrel{\circ}{T}^{*} S^{n-1}\right)$. The symplectic form and the Liouville volume form on $\stackrel{\circ}{T}^{*} S^{n-1}$ are given by $\Omega=-d \theta$ and $d M=$ $(-1)^{(n-1)(n-2) / 2}((n-1)!)^{-1} \Omega^{n-1}$, respectively. For any real-valued function $h \in C^{\infty}\left(\stackrel{\circ}{T}^{*} S^{n-1}\right)$, the unique vector field $X_{h}$ on $\stackrel{\circ}{T}^{*} S^{n-1}$ for which $\left.X_{h}\right\lrcorner \Omega=d h$ is called the Hamiltonian vector field of $h$. By means of the metric, we may identify $\stackrel{\circ}{T}^{*} S^{n-1}$ with the space $\stackrel{\circ}{T} S^{n-1}$ consisting of non-zero tangent vectors to $S^{n-1}$, which is identified with

$$
M=\left\{(x, y) \in \boldsymbol{R}^{n} \times \boldsymbol{R}^{n}| | x \mid=1, x \cdot y=0, y \neq 0\right\}
$$

Furthermore, by an injection $(x, y) \mapsto z=|y| x+\sqrt{-1} y$ of $M$ into $\boldsymbol{C}^{n}, M$ is identified with a complex cone $\left\{z \in \boldsymbol{C}^{n} \mid z^{2}=0, z \neq 0\right\}$. This identification gives $M$ a complex structure $J$. It is known (see [7] and [8, p. 173]) that $J$ is compatible with the symplectic structure, i.e., $(X, Y) \mapsto-\Omega(J(X), Y)$ is a Kaehler metric on $M$.

Let $\operatorname{Holo}(M)$ and $P_{m}(M)$ denote, respectively, the restrictions to $M$ of entire holomorphic functions and homogeneous polynomials of degree $m$ on $C^{n}$. For any $\varphi \in \operatorname{Holo}(M)$, there exists a unique $\varphi_{m} \in P_{m}(M)$ such that $\varphi=\sum_{m=0}^{\infty} \varphi_{m}$; uniformly convergent on any bounded set in $M$. If we define $\psi_{i_{1} \cdots i_{m}} \in P_{m}(M)$ by $\psi_{i_{1} \cdots i_{m}}(z)=z_{i_{1}} \cdots z_{i_{m}}$, then $P_{m}(M)$ is spanned by these functions. Since $z_{1}^{2}+\cdots+z_{n}^{2}=0$, we have $\operatorname{dim} P_{m}(M)=$ $\operatorname{dim} H_{m}\left(S^{n-1}\right)$. (Cf. [5, § 3].)

The unit cotangent bundle $T_{1}^{*} S^{n-1}$ to $S^{n-1}$ is identified with $N=$ $\{(x, y) \in M||y|=1\}$. The canonical volume element on $N$ is denoted by $d N$. If we define a function $r \in C^{\infty}(M)$ and a projection $p: M \rightarrow N$ by $r(x, y)=|y|$ and $p(x, y)=\left(x,|y|^{-1} y\right)$, respectively, then we have $d M=$ $p^{*} d N \wedge r^{n-2} d r$. An inner product in $C^{\infty}(N)$ is defined by

$$
\langle\varphi, \psi\rangle_{N}=\int_{N} \bar{\varphi} \psi d N
$$

The restriction of $\psi_{i_{1} \cdots i_{m}}$ onto $N$ will also be denoted by the same letter.
Lemma 1.6. (i) If $l \neq m$, then $\left\langle\psi_{i_{1} \cdots i_{l}}, \psi_{j_{1} \cdots j_{m}}\right\rangle_{N}=0$.
(ii) For any $m=1,2, \cdots$, we have

$$
\begin{aligned}
& \frac{(m+n-3)(2 m+n-2)}{2}\left\langle\dot{\psi}_{i_{1} \cdots i_{m}}, \dot{\psi}_{j_{1} \cdots j_{m}}\right\rangle_{N} \\
&=(2 m+n-4) \sum_{a=1}^{m} \delta_{i_{a} j_{m}}\left\langle\psi_{i_{1} \cdots \hat{i}_{a} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N} \\
&-\sum_{a \neq b} \delta_{i_{a} i_{b}}\left\langle\psi_{i_{1} \cdots \hat{i}_{a} \cdots \hat{i}_{b} \cdots i_{m} j_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}
\end{aligned}
$$

Proof. (i) The Hamiltonian vector field of the function $r$ is given by

$$
\left.X_{r}\right|_{(x, y)}=\sum_{i=1}^{n}\left(\frac{1}{|y|} y_{i} \frac{\partial}{\partial x_{i}}-|y| x_{i} \frac{\partial}{\partial y_{i}}\right)
$$

$X_{r}$ is tangent to $N$. Denoting the restriction of $X_{r}$ to $N$ also by $X_{r}$, we have $\mathscr{L}_{X_{r}} d N=0$. It follows that $\left\langle X_{r} \varphi, \psi\right\rangle_{N}=-\left\langle\varphi, X_{r} \psi\right\rangle_{N}$ for any $\varphi, \psi \in C^{\infty}(N)$. On the other hand, from $X_{r} z_{i}=-\sqrt{-1} z_{i}$ we have $X_{r} \psi_{i_{1} \cdots i_{m}}=-\sqrt{-1} m \psi_{i_{1} \cdots i_{m}}$. Then (i) follows immediately.
(ii) The Hamiltonian vector field $X_{j}$ of the function $(x, y) \mapsto y_{j}$ on $M$ is given by

$$
\left.X_{j}\right|_{(x, y)}=\sum_{i=1}^{n}\left\{\left(\delta_{i i}-x_{j} x_{i}\right) \frac{\partial}{\partial x_{i}}+\left(x_{j} y_{i}-y_{j} x_{i}\right) \frac{\partial}{\partial y_{i}}\right\} .
$$

Since $\left[X_{j}, \sum_{k=1}^{n} y_{k}\left(\partial / \partial y_{k}\right)\right]=0, X_{j}$ induces a tangent vector field $\eta_{j}$ to $N$. $\eta_{j}$ is given by

$$
\left.\eta_{j}\right|_{(x, y)}=\sum_{i=1}^{n}\left\{\left(\delta_{j i}-x_{j} x_{i}\right) \frac{\partial}{\partial x_{i}}-y_{j} x_{i} \frac{\partial}{\partial y_{i}}\right\}
$$

Since $\mathscr{L}_{\eta_{j}} d N=-2^{-1}(n-1)\left(z_{j}+\bar{z}_{j}\right) d N$, we have from

$$
\begin{gathered}
\int_{N} \mathscr{L}_{\eta_{j}}(\bar{\varphi} \psi d N)=0 \\
\left\langle\eta_{j} \varphi, \psi\right\rangle_{N}+\left\langle\varphi, \eta_{j} \psi\right\rangle_{N}=\frac{n-1}{2}\left(\left\langle z_{j} \varphi, \psi\right\rangle_{N}+\left\langle\varphi, z_{j} \psi\right\rangle_{N}\right)
\end{gathered}
$$

for any $\varphi, \psi \in C^{\infty}(N)$, where $z_{j}$ denotes the function $(x, y) \mapsto z_{j}=x_{j}+$ $\sqrt{-1} y_{j}$ on $N$. If we put $j=j_{m}, \varphi=\psi_{i_{1} \cdots i_{m}}$ and $\psi=\psi_{j_{1} \cdots j_{m-1}}$, then using (i) we have

$$
\left\langle\eta_{j_{m}} \psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}+\left\langle\psi_{i_{1} \cdots i_{m}}, \eta_{j_{m}} \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}=\frac{n-1}{2}\left\langle\psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m}}\right\rangle_{N}
$$

Now, since $\eta_{j}\left(z_{k}\right)=\delta_{j k}-(1 / 2) z_{j} z_{k}-(1 / 2) z_{j} \bar{z}_{k}$, we have

$$
\begin{aligned}
& \left\langle\eta_{j_{m}} \psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N} \\
& \quad=\sum_{a=1}^{m} \delta_{i_{a} j_{m}}\left\langle\psi_{i_{1} \cdots \hat{i}_{a} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}-\frac{1}{2} \sum_{a=1}^{m}\left\langle\psi_{i_{1} \cdots i_{a} \cdots i_{m} j_{m}}, \psi_{j_{1} \cdots j_{m-1} i_{a}}\right\rangle_{N}
\end{aligned}
$$

and

$$
\left\langle\psi_{i_{1} \cdots i_{m}}, \eta_{j_{m}} \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}=-\frac{m-1}{2}\left\langle\psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m}}\right\rangle_{N} .
$$

It follows that

$$
\begin{aligned}
& \frac{m+n-2}{2}\left\langle\psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m}}\right\rangle_{N} \\
& \quad=\sum_{a=1}^{m} \delta_{i_{a} j_{m}}\left\langle\psi_{i_{1} \cdots \hat{i}_{a} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}-\frac{1}{2} \sum_{a=1}^{m}\left\langle\psi_{i_{1} \cdots \hat{i}_{a} \cdots i_{m} j_{m}}, \psi_{j_{1} \cdots j_{m-1} i_{a}}\right\rangle_{N}
\end{aligned}
$$

from which we obtain (ii).
Lemma 1.7. For any $m=0,1,2, \cdots$, we have

$$
\left\langle h_{i_{1} \cdots i_{m}}, h_{j_{1} \cdots j_{m}}\right\rangle_{s}=c_{m}\left\langle\hat{h}_{i_{1} \cdots i_{m}}, \hat{h}_{j_{1} \cdots j_{m}}\right\rangle_{N},
$$

where

$$
c_{m}=\frac{\Gamma(m+n / 2) \Gamma(m+1) \operatorname{dim} H_{m}\left(S^{n-1}\right)}{\left(\operatorname{vol}\left(S^{n-1}\right)\right)^{2} \operatorname{vol}\left(S^{n-2}\right) \Gamma(n / 2)}
$$

Proof. By Proposition 1.5,

$$
\hat{h}_{i_{1} \cdots i_{m}}=\frac{\operatorname{vol}\left(S^{n-1}\right) \Gamma(n / 2)}{2^{m} \Gamma(m+n / 2)} \psi_{i_{1} \cdots i_{m}}
$$

Hence it suffices to show that

$$
\left\langle h_{i_{1} \cdots i_{m}}, h_{j_{1} \cdots j_{m}}\right\rangle_{S}=c_{m}^{\prime}\left\langle\psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m}}\right\rangle_{N},
$$

where

$$
c_{m}^{\prime}=\frac{\Gamma(n / 2) \Gamma(m+1) \operatorname{dim} H_{m}\left(S^{n-1}\right)}{2^{2 m} \operatorname{vol}\left(S^{n-2}\right) \Gamma(m+n / 2)}
$$

We shall show this by induction on $m$. If $m=0$, then both sides of the equation are equal to $\operatorname{vol}\left(S^{n-1}\right)$. Now, let $m>0$ and assume that the equality holds for $m-1$. Then by Lemma 1.3 and (ii) of Lemma 1.6, we have

$$
\begin{aligned}
&\left\langle h_{i_{1} \cdots i_{m}}, h_{j_{1} \cdots j_{m}}\right\rangle_{S} \\
&= \frac{c_{m-1}^{\prime}}{(2 m+n-2)(2 m+n-4)}\left\{(2 m+n-4) \sum_{a=1}^{m} \delta_{i_{a} j_{m}}\left\langle\psi_{i_{1} \cdots \hat{i}_{a} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}\right. \\
&\left.-\sum_{a \neq b} \delta_{i_{a} i_{b}}\left\langle\psi_{i_{1} \cdots \hat{i}_{a} \cdots \hat{i}_{b} \cdots i_{m} j_{m}}, \psi_{j_{1} \cdots j_{m-1}}\right\rangle_{N}\right\} \\
&= \frac{(m+n-3) c_{m-1}^{\prime}}{2(2 m+n-4)}\left\langle\psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m}}\right\rangle_{N}=c_{m}^{\prime}\left\langle\psi_{i_{1} \cdots i_{m}}, \psi_{j_{1} \cdots j_{m}}\right\rangle_{N} .
\end{aligned}
$$

Proposition 1.8 (cf. [5, §4]). For any $x \in S^{n-1}$, we have

$$
\int_{N}(x \cdot \bar{z})^{m} \psi_{i_{1} \cdots i_{m}}(z) d N(z)=\frac{\operatorname{vol}\left(S^{n-1}\right) \operatorname{vol}\left(S^{n-2}\right) 2^{m}}{\operatorname{dim} H_{m}\left(S^{n-1}\right)} h_{i_{1} \cdots i_{m}}(x) .
$$

This proposition is proved by induction on $m$, where we use (ii) of Lemma 1.2 and (ii) of Lemma 1.6.
2. Hilbert space $P(M)$ and integral transform $\mathscr{F}$. From now on, we shall assume that $n=3,5,7, \cdots$.

Lemma 2.1. There exists a unique polynomial $\rho_{n}$ which satisfies

$$
\int_{0}^{\infty} r^{2 m+n-2} e^{-2 r} \rho_{n}(r) d r=c_{m}
$$

for all $m=0,1,2, \cdots$.
Proof. If there exists a polynomial $\rho_{n}(r)=\sum_{k} a_{n, k} r^{k}$ which satisfies the condition in our lemma, then the coefficients must satisfy $\sum_{k} a_{n, k} 2^{-(k+2 m+n-1)} \Gamma(k+2 m+n-1)=c_{m}$ for all $m$. This condition is rewritten as

$$
\begin{aligned}
& \sum_{k} a_{n, k} \frac{\Gamma(k+2 m+n-1)}{2^{k+1} \Gamma(2 m+n-1)} \\
& \quad \quad=\frac{\pi^{1 / 2}(2 m+n-2) \Gamma(m+n-2)}{\left(\operatorname{vol}\left(S^{n-1}\right)\right)^{2} \operatorname{vol}\left(S^{n-2}\right) \Gamma(n-1) \Gamma(n / 2) \Gamma(m+(n-1) / 2)}
\end{aligned}
$$

for all $m=0,1,2, \cdots$. Since $n$ is odd, both sides of the equation above are polynomials of $m$. Hence, $a_{n, k}$ are determined uniquely. The existence of $\rho_{n}$ also follows from the above equation.

Note that the degree of the polynomial $\rho_{n}$ is $(n-1) / 2$, and the coefficient of the highest degree is positive. For example, we have $\rho_{3}(r)=a_{3,1}(r-1 / 2), \quad \rho_{5}(r)=a_{5,2}\left(r^{2}-r\right) \quad$ and $\quad \rho_{7}(r)=a_{7,3}\left(r^{3}-r^{2}-r / 2\right)$. Unfortunately, since $a_{n, 0}=0$ and $a_{n, 1}<0$ for $n \geqq 5, \rho_{n} \mid(0, \infty)$ is not a positive function. It is to be desired that there exists a positive function on ( $0, \infty$ ) which satisfies the equation in Lemma 2.1. We also remark that for even $n$, there does not exist any polynomial which satisfies the condition in Lemma 2.1. This is the reason why we restrict our attention to the case of odd $n$.

Now, for any $\varphi, \psi \in \operatorname{Holo}(M)$, let us define

$$
\langle\varphi, \psi\rangle_{M}=\int_{M} \overline{\varphi(z)} \psi(z) d \mu_{n}(z),
$$

where $d \mu_{n}(z)=e^{-2|y|} \rho_{n}(|y|) d M(z), z=|y| x+\sqrt{-1} y \in M$ (cf. [8, p. 174]). Although the measure $d \mu_{n}$ is not positive, we have:

Theorem 2.2. For any $\varphi \in P_{l}(M)$ and $\psi \in P_{m}(M)$,

$$
\langle\varphi, \psi\rangle_{M}=c_{m}\langle\varphi, \psi\rangle_{N},
$$

where $\varphi$ and $\psi$ on the right hand side stand for the restrictions of $\varphi$ and $\psi$ onto $N$, respectively. In particular, $\langle,\rangle_{M}$ is positive definite on $P_{m}(M)$, and $f \mapsto \hat{f}$ is a unitary isomorphism of $\left(H_{m}\left(S^{n-1}\right),\langle,\rangle_{s}\right)$ onto $\left.P_{m}(M),\langle,\rangle_{M}\right)$.

Proof. Since $d M=p^{*} d N \wedge r^{n-2} d r$, we have, by (i) of Lemma 1.6 and Lemma 2.1,

$$
\langle\varphi, \psi\rangle_{M}=\int_{0}^{\infty} r^{l+m+n-2} e^{-2 r} \rho_{n}(r) d r \int_{N} \bar{\varphi} \psi d N=c_{m}\langle\varphi, \psi\rangle_{N} .
$$

Then, the unitarity of $f \mapsto \hat{f}$ follows from Lemma 1.7.
The following lemma is due to Bargmann [1, p. 190].
Lemma 2.3. Let $S=\sum_{k=1}^{\infty} b_{k}$ be a series with non-negative real terms, let $\gamma_{k}(t), t>0$, be so chosen that (1) $0 \leqq \gamma_{k}(t) \leqq 1$, (2) $\lim _{t \rightarrow \infty} \gamma_{k}(t)=1$, and set $S(t)=\sum \gamma_{k}(t) b_{k}$. S converges if and only if $S(t)$ are uniformly bounded, and in that case $S=\lim S(t)$.

Proposition 2.4. Let $\varphi \in \operatorname{Holo}(M), \varphi=\sum \varphi_{m}$ with $\varphi_{m} \in P_{m}(M)$. Then

$$
\langle\varphi, \varphi\rangle_{M}=\sum\left\langle\varphi_{m}, \varphi_{m}\right\rangle_{M},
$$

i.e., either both sides are infinite, or both sides are finite and equal.

Proof. For any $\sigma>0$, let

$$
I(\sigma)=\int_{M(\sigma)}|\varphi|^{2} d \mu_{n},
$$

where $M(\sigma)=\{z=|y| x+\sqrt{-1} y \in M| | y \mid \leqq \sigma\}$. Then $\sigma \mapsto I(\sigma)$ is, for large $\sigma$, monotone increasing and $\langle\varphi, \varphi\rangle_{\mu}=\lim _{\sigma \rightarrow \infty} I(\sigma)$. Since $\sum \varphi_{m}$ converges uniformly to $\varphi$ on $M(\sigma)$, we have by (i) of Lemma 1.6 and Theorem 2.2,

$$
\begin{aligned}
I(\sigma) & =\sum_{l, m=0}^{\infty} \int_{M(\sigma)} \overline{\varphi_{l}(z)} \varphi_{m}(z) d \mu_{n}(z)=\sum_{l, m=0}^{\infty} \int_{0}^{\sigma} r^{l+m+n-2} e^{-2 r} \rho_{n}(r) d r \int_{N} \overline{\varphi_{l}} \varphi_{m} d N \\
& =\sum_{m=0}^{\infty} \int_{0}^{\sigma} r^{2 m+n-2} e^{-2 r} \rho_{n}(r) d r\left\langle\varphi_{m}, \varphi_{m}\right\rangle_{N}=\sum_{m=0}^{\infty} \frac{c_{m}(\sigma)}{c_{m}}\left\langle\varphi_{m}, \varphi_{m}\right\rangle_{M},
\end{aligned}
$$

where

$$
c_{m}(\sigma)=\int_{0}^{\sigma} r^{2 m+n-2} e^{-2 r} \rho_{n}(r) d r
$$

Since there exists $\sigma_{n}>0$ such that $c_{m}(\sigma)>0$ for all $\sigma>\sigma_{n}$ and $m=$ $0,1,2, \cdots$, applying Lemma 2.3, we have the desired result.

Now, let us define

$$
P(M)=\left\{\varphi \in \operatorname{Holo}(M) \mid\langle\varphi, \varphi\rangle_{M}<\infty\right\} .
$$

Then it follows from Theorem 2.2 and Proposition 2.4 that $\langle,\rangle_{M}$ is a Hermitian inner product in $P(M)$. The corresponding norm is denoted by $\left\|\|_{M}\right.$.

THEOREM 2.5. $\mathscr{F}: f \mapsto \hat{f}$ is a unitary isomorphism of $\left(L^{2}\left(S^{n-1}\right),\langle,\rangle_{s}\right)$ onto $\left(P(M),\langle,\rangle_{M}\right)$.

Proof. Let $f \in L^{2}\left(S^{n-1}\right), f=\sum f_{m}$ with $f_{m} \in H_{m}\left(S^{n-1}\right)$. Then, by (iii) of Lemma 1.4, Proposition 2.4, Theorem 2.2 and (iii) of Lemma 1.1, we have

$$
\|\widehat{f}\|_{M}^{2}=\sum\left\|\widehat{f}_{m}\right\|_{M}^{2}=\sum\left\|f_{m}\right\|_{S}^{2}=\|f\|_{S}^{2}<\infty
$$

It follows that $\hat{f} \in P(M)$ and that $\mathscr{F}$ is unitary. The surjectivity of $\mathscr{F}$ is also shown easily.

We have from Theorem 2.5 and (ii) of Lemma 1.4 the following:
Corollary 2.6. (i) $\left(P(M),\langle,\rangle_{s}\right)$ is a Hilbert space. (ii) For any $\varphi \in P(M)$ and $z=|y| x+\sqrt{-1} y \in M$,

$$
|\varphi(z)| \leqq\left(\operatorname{vol}\left(S^{n-1}\right)\right)^{1 / 2}\|\varphi\|_{M} e^{|y|} .
$$

From (ii) of Corollary 2.6, it follows that, for a fixed $w \in M$, the map $\varphi \mapsto \varphi(w)$ defines a bounded linear functional on $P(M)$. It is necessarily of the form

$$
\varphi(w)=\left\langle e_{w}, \varphi\right\rangle_{M}
$$

with a uniquely defined $e_{w} \in P(M)$. If we define function on $M \times M$ by

$$
K(w, z)=\int_{S^{n-1}} e^{x \cdot w} e^{x \cdot \bar{z}} d S(x)
$$

then $\overline{K(w, z)}=K(z, w)$ and $\overline{K(w, \cdot)} \in P(M)$ immediately from the definition.
Lemma 2.7 (cf. [1, § 1 c$]$ ).

$$
e_{w}(z)=\overline{K(w, z)} .
$$

Proof. It is sufficient to show that

$$
\left\langle\overline{K(w, \cdot)}, \psi_{i_{1} \cdots i_{m}}\right\rangle_{M}=\psi_{i_{1} \cdots i_{m}}(w) .
$$

Making use of Theorem 2.2, Lemma 1.6 and Propositions 1.8 and 1.5, we have

$$
\begin{aligned}
\left\langle\overline{K(w, \cdot)}, \dot{\psi}_{i_{1} \cdots i_{m}}\right\rangle_{M} & =\int_{M}\left(\int_{S^{n-1}} e^{x \cdot w} e^{x \cdot \bar{z}} d S(x)\right) \psi_{i_{1} \cdots i_{m}}(z) d \mu_{n}(z) \\
& =\int_{S^{n-1}} e^{x \cdot w}\left(\int_{M} e^{x \cdot \bar{z}} \psi_{i_{1} \cdots i_{m}}(z) d \mu_{n}(z)\right) d S(x) \\
& =\frac{1}{m!} \int_{S^{n-1}} e^{x \cdot w}\left(\int_{M}(x \cdot \bar{z})^{m} \psi_{i_{1} \cdots i_{m}}(z) d \mu_{n}(z)\right) d S(x) \\
& =\frac{c_{m}}{m!} \int_{S^{n-1}} e^{x \cdot w}\left(\int_{N}(x \cdot \bar{z})^{m} \psi_{i_{1} \cdots i_{m}}(z) d N(z)\right) d S(x) \\
& =\frac{c_{m} \operatorname{vol}\left(S^{n-1}\right) \operatorname{vol}\left(S^{n-2}\right) 2^{m}}{m!\operatorname{dim} H_{m}\left(S^{n-1}\right)} \int_{S^{n-1}} e^{x \cdot w} h_{i_{1} \cdots i_{m}}(x) d S(x) \\
& =\psi_{i_{1} \cdots i_{m}}(w) .
\end{aligned}
$$

$K$ is the reproducing kernel for $P(M)$, i.e.,

$$
\varphi(w)=\int_{M} K(w, z) \varphi(z) d \mu_{n}(z)
$$

Now, we shall consider the inverse operator $\mathscr{F}^{-1}$. Let $P^{(\alpha)}(M)=\{\rho \in$ $\operatorname{Holo}(M) \mid$ for a suitable $c>0,|\varphi(z)| \leqq c e^{\lambda|y|}$ for all $\left.z=|y| x+\sqrt{-1} y \in M\right\}$ $(0<\lambda<1)$. Then $P^{(\lambda)}(M)$ is a subspace of $P(M)$. If, for each $\varphi \in P(M)$, we define $\varphi^{(\lambda)}$ by $\varphi^{(\lambda)}(z)=\varphi(\lambda z)$, then $\varphi^{(\lambda)} \in P^{(\lambda)}(M)$.

Lemma 2.8 (cf. [1, p. 197]). (i) $\varphi \in P(M)$ if and only if all $\varphi^{(\lambda)} \in$ $P(M), 0<\lambda<1$, and their norms $\left\|\varphi^{(\lambda)}\right\|_{M}$ are uniformly bounded.
(ii) If $\varphi \in P(M)$, then $\left\|\varphi-\varphi^{(\lambda)}\right\|_{M} \rightarrow 0$ as $\lambda \rightarrow 1$.

Proof. Let $\varphi \in \operatorname{Holo}(M), \varphi=\sum \varphi_{m}$ with $\varphi_{m} \in P_{m}(M)$. Then we have $\varphi^{(\lambda)}(z)=\varphi(\lambda z)=\sum \lambda^{m} \varphi_{m}(z)$. It follows from Proposition 2.4 that $\left\|\varphi^{(\lambda)}\right\|_{\mathbb{M}}^{2}=$ $\sum \lambda^{2 m}\left\|\varphi_{m}\right\|_{M}^{2}$. Then by Lemma 2.3 we have (i). (ii) follows immediately from $\left\|\varphi-\varphi^{(\lambda)}\right\|_{M}^{2}=\sum\left(1-\lambda^{m}\right)^{2}\left\|\varphi_{m}\right\|_{M}^{2}$.

Theorem 2.9 (cf. [1, p. 202]). If $\varphi \in P^{(\lambda)}(M)$ for some $\lambda, 0<\lambda<1$, then

$$
\left(\mathscr{F}^{-1} \varphi\right)(x)=\int_{M} e^{x \cdot \bar{z}} \varphi(\boldsymbol{z}) d \mu_{n}(\boldsymbol{z}),
$$

for any $x \in S^{n-1}$.
Proof. Since $\varphi \in P^{(\lambda)}(M)$, the integration converges absolutely. It suffices to prove that

$$
\int_{S^{n-1}} e^{x \cdot w}\left(\int_{M} e^{x \cdot \bar{z}} \mathcal{P}(z) d \mu_{n}(z)\right) d S(x)=\varphi(w),
$$

which we show easily by interchanging integrations and using the reproducing property of $K$.

Corollary 2.10 (cf. [1, (2.14)]). For any $\varphi \in P(M)$,

$$
\left(\mathscr{F}^{-1} \varphi\right)(x)=\operatorname{Lim}_{\lambda \rightarrow 1} \int_{M} e^{x \cdot \bar{z}} \varphi(\lambda z) d \mu_{n}(z),
$$

where Lim means the strong convergence in $L^{2}\left(S^{n-1}\right)$.
We also have another explicit expression for $\mathscr{F}^{-1}$.
Theorem 2.11 (cf. [1, (2.15)]). For any $\varphi \in P(M)$,

$$
\left(\mathscr{F}^{-1} \varphi\right)(x)=\operatorname{Lim}_{\sigma \rightarrow \infty} \int_{M(\sigma)} e^{x \cdot \bar{z}} \varphi(z) d \mu_{n}(z)
$$

Proof. Let $\varphi=\sum \varphi_{m}$ with $\varphi_{m} \in P_{m}(M)$. Define, for $x \in S^{n-1}$,

$$
f^{(\sigma)}(x)=\int_{M(o)} e^{x \cdot \bar{z}} \varphi(z) d \mu_{n}(z)
$$

and

$$
f_{m}^{(o)}(x)=\int_{\boldsymbol{M}(\sigma)} e^{x \cdot \bar{z}} \varphi_{m}(z) d \mu_{n}(z) .
$$

Then, by Propositions 1.5 and 1.8 , we have for any $w \in M$,

$$
\begin{aligned}
\left(\mathscr{F} f_{m}^{(\sigma)}\right)(w) & =\int_{S^{n-1}} e^{x \cdot w}\left(\int_{M(\sigma)} e^{x \cdot \bar{z}} \varphi_{m}(z) d \mu_{n}(z)\right) d S(x) \\
& =\frac{c_{m}(\sigma)}{m!} \int_{S^{n-1}} e^{x \cdot w}\left(\int_{N}(x \cdot \bar{z})^{m} \varphi_{m}(z) d N(z)\right) d S(x)=\frac{c_{m}(\sigma)}{c_{m}} \varphi_{m}(w)
\end{aligned}
$$

By the uniform convergence of $\varphi=\sum \varphi_{m}$ on $M(\sigma)$, we have

$$
\begin{aligned}
\left(\mathscr{F} f^{(\sigma)}\right)(w) & =\int_{S^{n-1 \times M(\sigma)}} e^{x \cdot w} e^{x \cdot \bar{z}} \varphi(z) d \mu_{n}(z) d S(x) \\
& =\sum \int_{S^{n-1 \times M(\sigma)}} e^{x \cdot w} e^{x \cdot \bar{z}} \varphi_{m}(z) d \mu_{n}(z) d S(x) \\
& =\sum\left(\mathscr{F} f_{m}^{(\sigma)}\right)(w)=\sum \frac{c_{m}(\sigma)}{c_{m}} \varphi_{m}(w)
\end{aligned}
$$

It follows from Proposition 2.4 that

$$
\left\|\varphi-\mathscr{F} f^{(\sigma)}\right\|_{M}^{2}=\sum\left(1-\frac{c_{m}(\sigma)}{c_{m}}\right)^{2}\left\|\varphi_{m}\right\|_{M}^{2} \rightarrow 0
$$

as $\sigma \rightarrow \infty$. Here recall that there exists a constant $\sigma_{n}>0$ such that $c_{m}(\sigma)>0$ for any $\sigma>\sigma_{n}$ and $m=0,1,2, \cdots$. Since $\mathscr{F}$ is a unitary isomorphism, we have $\mathscr{F}^{-1} \varphi=\operatorname{Lim}_{\sigma \rightarrow \infty} f^{(o)}$.
3. An application. The mapping $\mathscr{F}$ establishes a unitary isomorphism between the linear operators on $P(M)$ and those on $L^{2}\left(S^{n-1}\right)$. In this section, we shall consider a one-parameter group of unitary transformations, which is easily analyzed on $P(M)$, and translate the results into the language of $L^{2}\left(S^{n-1}\right)$ (see [1, §3] and [8, p. 177]).

The one-parameter group of canonical transformations on $M$ generated by the Hamiltonian vector field $X_{r}$ is given by $\phi_{t}: z \mapsto e^{i t} z$. Since $X_{r} r=0$ and $\mathscr{L}_{x_{r}} d M=0, \phi_{t}$ preserves the measure $d \mu_{n}$ as well as the complex structure $J$ on $M$. Hence $\phi_{t}$ induces a unitary transformation $\varphi \mapsto \varphi \circ \phi_{-t}$ on $P(M)$. Let us define a one-parameter group $\left\{V_{t} \mid t \in \boldsymbol{R}\right\}$ of unitary transformations on $P(M)$ by

$$
\left(V_{t} \varphi\right)(z)=e^{-i(n-2) t / 2} \varphi\left(e^{-i t} z\right)
$$

(see [8, p. 177]). Then

$$
V_{t} \varphi_{m}=e^{-i \mid m+(n-2) / 2 t t} \varphi_{m}
$$

for any $\varphi_{m} \in P_{m}(M)$, and $\left\{V_{t}\right\}$ is strongly continuous in $t$. The infinitesimal generator of $\left\{V_{t}\right\}$ is given by $X_{r}-i(n-2) / 2$. Now, let $U_{t}=\mathscr{F}^{-1} \circ V_{t} \circ \mathscr{F}$ be the operator corresponding to $V_{t}$ under the unitary isomorphism $\mathscr{F}$. Then, for any $f \in L^{2}\left(S^{n-1}\right)$ and $x^{\prime} \in S^{n-1}$, we have from Theorem 2.11

$$
\begin{aligned}
\left(U_{t} f\right)\left(x^{\prime}\right) & =\operatorname{Lim}_{\sigma \rightarrow \infty} \int_{M(\sigma)} e^{x^{\prime} \cdot \bar{z}} e^{-i(n-2) t / 2} \int_{S^{n-1}} e^{x \cdot \exp (-i t) z} f(x) d S(x) d \mu_{n}(z) \\
& =\operatorname{Lim}_{\sigma \rightarrow \infty} \int_{S^{n-1}} U^{(\sigma)}\left(t, x^{\prime}, x\right) f(x) d S(x),
\end{aligned}
$$

where

$$
U^{(o)}\left(t, x^{\prime}, x\right)=e^{-i(n-2) t / 2} \int_{M(\sigma)} e^{x^{\prime} \cdot \bar{z}+\exp (-i t) x \cdot z} d \mu_{n}(z)
$$

(cf. [1, (3.10a)]). Since $U_{t} f_{m}=e^{-i(m+\langle n-2) / 2\rangle t} f_{m}$ for any $f_{m} \in H_{m}\left(S^{n-1}\right)$, we have $U_{t}=\exp \left[-i\left\{\Delta+(n-2)^{2} / 4\right\}^{1 / 2} t\right]$, where $\Delta$ is the Laplace-Beltrami operator on $S^{n-1}$ (see [8, p. 177]). Thus, we have the following:

Theorem 3.1. The one-parameter group of unitary transformations, $U_{t}=\exp \left[-i\left\{\Delta+(n-2)^{2} / 4\right\}^{1 / 2} t\right]$, on $L^{2}\left(S^{n-1}\right)$ generated by the operator $-i\left\{\Delta+(n-2)^{2} / 4\right\}^{1 / 2}$ is represented by

$$
\left(U_{t} f\right)\left(x^{\prime}\right)=\operatorname{Lim}_{\sigma \rightarrow \infty} \int_{S^{n-1}} U^{(\sigma)}\left(t, x^{\prime}, x\right) f(x) d S(x)
$$

## References

[1] V. Bargmann, On a Hilbert space of analytic functions and an associated integral transform. Part I, Comm. Pure Appl. Math. 14 (1961), 187-214.
[2] K. Gawedzki, Fourier-like kernels in geometric quantization, Dissertation Math. 128, Warszawa, 1976.
[3] M. Hashizume, A. Kowata, K. Minemura and K. Okamoto, An integral representation of an eigenfunction of the Laplacian on the Euclidean space, Hiroshima Math. J. 2 (1972), 535-545.
[4] S. Helgason, Topics in harmonic analysis on homogeneous spaces. Progress in Math. 13, Birkhäuser, Boston-Basel-Stuttgart, 1981.
[5] A. Kowata and K. Okamoto, Harmonic functions and the Borel-Weil theorem, Hiroshima Math. J. 4 (1974), 89-97.
[6] M. Morimoto, Analytic functions on the sphere and their Fourier-Borel transformations, Banach Center Publications, 11, PWN-Polish Scientific Publishers, Warsaw, 1983, 223-250.
[7] J. H. Rawnsley, Coherent states and Kaehler manifolds, Quart. J. Math. Oxford Ser. 28 (1977), 403-415.
[8] J. H. Rawnsley, A nonunitary pairing of polarizations for the Kepler problem, Trans. Amer. Math. Soc. 250 (1979), 167-180.
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