The existence of the global solutions to semilinear wave equations with a class of cubic nonlinearities in 2-dimensional space

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Abstract. This paper deals with the Cauchy problem of the semilinear wave equation with a small initial data in 2-dimensional space. When the nonlinearity is cubic, we can not expect the global existence of smooth solutions, in general. However, Godin [1] showed that if the nonlinearity has the null-form, the solution exists globally. In this paper, we will show the global solvability for the other type of nonlinearities which does not have null-form.

Key words: semilinear wave equation, null-condition, global solvability.

1. Introduction

Let us consider the following Cauchy problem;

$$\Box u = \partial_t^2 u - \Delta u = F(\partial u) \quad (x, t) \in \mathbf{R}^2 \times (0, \infty), \tag{1.1}$$

$$u(x, 0) = \varepsilon f(x), \ \partial_t u(x, 0) = \varepsilon g(x) \quad x \in \mathbf{R}^2.$$
 (1.2)

Here, $\partial = (\partial_0, \partial_1, \partial_2)$, $\partial_0 = \partial_t = \partial/\partial t$, $\partial_j = \partial/\partial x_j$ (j = 1, 2) and ε is a positive small parameter. We assume $f, g \in C_0^{\infty}(\mathbf{R}^2; \mathbf{R})$, $|f| + |g| \not\equiv 0$ and $\sup\{f, g\} \subset \{x \in \mathbf{R}^2; |x| \leq M\}$. We also assume that

$$F \in C^{\infty}(\mathbf{R}^3; \mathbf{R}),$$

 $F(\partial u) = O(|\partial u|^3)$ near $\partial u = 0.$

More presicely, we assume

$$F(\partial u) = \sum_{\alpha,\beta,\gamma=0}^{2} A^{\alpha\beta\gamma} \partial_{\alpha} u \partial_{\beta} u \partial_{\gamma} u + O(|\partial u|^{4}) \quad \text{near} \quad \partial u = 0, \quad (1.3)$$

where $A^{\alpha\beta\gamma}$ are real constants.

The aim of this paper is to estimate the lifespan T_{ε} of the smooth solution to the Cauchy problem (1.1) and (1.2), which is defined for each

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 $\varepsilon > 0$ as follows;

$$T_{\varepsilon} = \sup\{T > 0; \text{there exists a smooth solution}\}$$

to (1.1) and (1.2) in
$$C^{\infty}(\mathbf{R}^2 \times [0, T); \mathbf{R})$$
.

In order to state the results which we have already known about the lifespan, we introduce some notations.

A. Hoshiqa

For vectors $X = (X_0, X_1, X_2) \in \mathbf{R}^3$, we define

$$C(X) = \sum_{\alpha, \beta, \gamma = 0}^{2} A^{\alpha\beta\gamma} X_{\alpha} X_{\beta} X_{\gamma}.$$

This function characterizes the essential cubic terms of $F(\partial u)$ along the light cone.

On the other hand, let $u^0 = u^0(x, t)$ be the solution to the Cauchy problem;

$$\Box u^{0} = 0 \quad (x, t) \in \mathbf{R}^{2} \times (0, \infty),$$

$$u^{0}(x, 0) = f(x), \ \partial_{t} u^{0}(x, 0) = q(x)$$

and set $r = |x| \ge 0$, $\omega = x/r \in S^1$ and $\rho = r - t \in \mathbf{R}$. Then we define

$$\mathcal{F}(\omega, \rho) = \lim_{r \to \infty} r^{1/2} u^0(r\omega, r - \rho),$$

which is called the Friedlander radiation field. Hörmander showed in [2] the following properties of \mathcal{F} .

$$|\partial_{\rho}^{k} \mathcal{F}(\omega, \rho)| \le C(1+|\rho|)^{-1/2-k} \quad \rho \in \mathbf{R},$$
 (1.4)

$$\mathcal{F}(\omega, \rho) = 0 \quad \text{for} \quad \rho \ge M.$$
 (1.5)

By (1.4) and (1.5), we find that the constant

$$H = \max_{\rho \in \mathbf{R}, \ \omega \in S^1} \left\{ -C(-1, \ \omega)(\partial_{\rho} \mathcal{F}(\omega, \ \rho))^2 \right\}$$
 (1.6)

is well-defined and nonnegative.

Then Godin proved the following (a) and (b) in [1].

(a) If H > 0, then

$$\liminf_{\varepsilon \to +0} \varepsilon^2 \log T_{\varepsilon} \ge \frac{1}{H}$$

holds.

(b) If $C(-1, \omega) \equiv 0$ holds for $\omega \in S^1$, then $T_{\varepsilon} = \infty$ holds for sufficiently small $\varepsilon > 0$.

The condition $C(-1, \omega) \equiv 0$ is called the *null-condition*. If the null-condition is satisfied, then we can write the cubic part of the nonlinear term $F(\partial u)$ as a linear combination of terms $\partial_{\alpha} u((\partial_0 u)^2 - |\nabla u|^2)$, i.e.,

$$F(\partial u) = \sum_{\alpha=0}^{2} C_{\alpha} \partial_{\alpha} u((\partial_{0} u)^{2} - |\nabla u|^{2}) + O(|\partial u|^{4}),$$

where C_{α} are real constants. It follows from (1.6) that the null-condition is a sufficient condition of H=0 and not a necessary condition. In the case where $|f|+|g| \not\equiv 0$, we find that H=0 is equivalent to the condition

$$C(-1,\omega) \ge 0$$
 for any $\omega \in S^1$. (1.7)

This means that there is a nonlinearity $F(\partial u)$ which does not satisfy the assumptions of both (a) and (b). For example,

$$F(\partial u) = -(\partial_0 u)^3, \quad F(\partial u) = -\partial_0 u |\nabla u|^2,$$

$$F(\partial u) = -(\partial_0 u)^3 + \partial_0 u \partial_1 u \partial_2 u$$

and so on. For the above $F(\partial u)$, we can easily show that

$$C(X) = -X_0^3$$
, $C(X) = -X_0(X_1^2 + X_2^2)$,
 $C(X) = -X_0^3 - X_0X_1X_2$,

respectively. Hence we find that these $F(\partial u)$ satisfy (1.7). From the results (a) and (b), we can expect the solution exists longer than the time $\exp(C/\varepsilon^2)$ for any constant C > 0 in such cases. Our purpose of this paper is to show the global existence of the smooth solution to (1.1) and (1.2), when the condition (1.7) holds.

2. Statement of the main theorem

We introduce generalized differential operators;

$$\Omega = x_1 \partial_2 - x_2 \partial_1, \quad L_i = t \partial_i + x_i \partial_0, \quad S = t \partial_0 + x_1 \partial_1 + x_2 \partial_2$$

and denote

$$\Gamma_0 = S, \ \Gamma_1 = \Omega, \ \Gamma_2 = L_1, \ \Gamma_3 = L_2, \ \Gamma_4 = \partial_0, \ \Gamma_5 = \partial_1, \ \Gamma_6 = \partial_2.$$

We can verify the following commutator relations;

$$[\partial_{\alpha}, \, \partial_{\beta}] = 0, \ [S, \, \partial_{\alpha}] = -\partial_{\alpha}, \ [\Omega, \, \partial_{1}] = -\partial_{2}, \ [\Omega, \, \partial_{2}] = \partial_{1},$$

$$[\Omega, \, \partial_{0}] = 0, \ [L_{i}, \, \partial_{j}] = -\delta_{ij}\partial_{0}, \ [L_{i}, \, \partial_{0}] = -\partial_{i}, \ [S, \, \Omega] = 0,$$

$$[\Omega, \, L_{1}] = -L_{2}, \ [\Omega, \, L_{2}] = L_{1}, \ [S, \, L_{i}] = 0, \quad [L_{1}, \, L_{2}] = \Omega,$$

$$[\Gamma_{\sigma}, \, \Box] = -2\delta_{0\sigma}\Box$$

$$(2.1)$$

for α , $\beta = 0, 1, 2$, i, j = 1, 2 and $\sigma = 0, 1, \ldots, 6$. Here [,] denotes the usual commutator of linear operators and $\delta_{\alpha\beta}$ is the Kronecker delta. We also write $\Gamma^a u = \Gamma_0^{a_0} \Gamma_1^{a_1} \cdots \Gamma_6^{a_6} u$ for a multi-index $a = (a_0, a_1, \ldots, a_6)$.

Next we define some generalized Sobolev norms as

$$|v(x,t)|_{k} = \sum_{|a| \le k} |\Gamma^{a}v(x,t)|$$

$$|v(t)|_{k} = \sum_{|a| \le k} ||\Gamma^{a}v(\cdot,t)||_{L_{x}^{\infty}(\mathbf{R}^{2};\mathbf{R})}$$

$$[v(t)]_{k} = \sum_{|a| \le k} ||(1+|\cdot|+t)^{1/2}(1+||\cdot|-t|)^{1/2}\Gamma^{a}v(\cdot,t)||_{L_{x}^{\infty}(\mathbf{R}^{2};\mathbf{R})}$$

$$||v(t)||_{k} = \sum_{|a| \le k} ||\Gamma^{a}v(\cdot,t)||_{L_{x}^{2}(\mathbf{R}^{2};\mathbf{R})}.$$
(2.2)

Note that by (2.1) and the definition of norms, we have

$$|v(x,t)| \le \frac{[v(t)]_0}{(1+t)^{1/2}} \tag{2.3}$$

and

$$\alpha |\partial v(t)|_{k} \leq \sum_{|a| \leq k} |\partial \Gamma^{a} v(t)|_{0} \leq \beta |\partial v(t)|_{k}$$

$$\alpha [\partial v(t)]_{k} \leq \sum_{|a| \leq k} [\partial \Gamma^{a} v(t)]_{0} \leq \beta [\partial v(t)]_{k}$$

$$\alpha ||\partial v(t)||_{k} \leq \sum_{|a| \leq k} ||\partial \Gamma^{a} v(t)||_{0} \leq \beta ||\partial v(t)||_{k}$$

$$(2.4)$$

for some positive constants α and β . Now we state the main theorem.

Theorem 2.1 Assume that (1.7) holds. Then there exists a constant $\varepsilon_* > 0$ such that $T_{\varepsilon} = \infty$ holds for $\varepsilon \in (0, \varepsilon_*)$. Moreover, for any integer $k \geq 5$, there exists constants $c_l > 0$ and $m_l \geq 0$ (l = 0, 1, ..., k) such that the

solution satisfies

$$\|\partial u(t)\|_{l} \le c_{l} \varepsilon (1 + \varepsilon^{2} \log(1 + t))^{m_{l}} \quad 0 \le t < \infty, \ l = 0, 1, \dots, k$$
 (2.5)

for $\varepsilon \in (0, \varepsilon_*)$. Here $c_l > c_{l-1}$ and $c_0 = c_0(f, g)$, $m_l > m_{l-1}$ and $m_0 = 0$.

Remark Kubo proved the same result in [8], in which he also showed an asymptotic behavior of solutions.

3. Proof of the main theorem

First of all, the following local existence result has been shown in Theorem 6.5.3 in [3].

Proposition 3.1 There exist constants D > 0 and $\varepsilon' > 0$ such that for $0 < \varepsilon < \varepsilon'$ the Cauchy problem (1.1) and (1.2) has a local solution $u \in C^{\infty}(\mathbf{R}^2 \times [0, \varepsilon^{-1}]; \mathbf{R})$ which satisfies

$$\|\partial u(t)\|_k \le D\varepsilon \quad 0 \le t \le \frac{1}{\varepsilon}$$
 (3.1)

for $k = 0, 1, 2, \dots$

Combining Proposition 3.1 with the following lemma, we can show the Theorem 2.1.

Lemma 3.1 Assume that (1.7) holds, choose an integer $k \geq 5$. Let $u \in C^{\infty}(\mathbf{R}^2 \times [0, T); \mathbf{R})$ be a solution to (1.1) and (1.2). Then, there exist constants $\varepsilon_0 > 0$, $c_l > 0$ and $m_l \geq 0$ ($c_l > c_{l-1}$, $c_0 > D$, $m_l > m_{l-1}$, $m_0 = 0$, $l = 0, 1, \ldots, k$) independent of T such that if

$$\|\partial u(t)\|_{l} \le 2c_{l}\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{l}}$$

 $0 \le t < T, \ l = 0, 1, \dots, k \quad (3.2)$

holds for an $\varepsilon \in (0, \varepsilon_0)$, then

$$\|\partial u(t)\|_{l} \le c_{l} \varepsilon (1 + \varepsilon^{2} \log(1 + t))^{m_{l}}$$

 $0 \le t < T, \ l = 0, 1, \dots, k \quad (3.3)$

holds for the same ε .

Proof of Theorem 2.1. Choose an integer $k \geq 5$ and define a set

$$U_{\varepsilon} = \{t; t \leq T_{\varepsilon}, \|\partial u(t)\|_{l} \leq 2c_{l}\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{l}},$$

$$l = 0, 1, ..., k$$
 $(\subset [0, \infty))$

for each $\varepsilon \in (0, \varepsilon_*)$, where $\varepsilon_* = \min\{\varepsilon_0, \varepsilon'\}$. Proposition 3.1 implies that U_{ε} is not empty. Furthermore, by Lemma 3.1, we can easily show that U_{ε} is open and closed in $[0, \infty)$ from the usual argument. Namely we have $U_{\varepsilon} = [0, \infty)$ and therefore we find that $T_{\varepsilon} = \infty$ and (2.5) hold for any $\varepsilon \in (0, \varepsilon_*)$. This completes the proof of Theorem 2.1.

In what follows we concentrate on showing (3.3) under the assumption (3.2). For this purpose, we will use the following propositions.

Proposition 3.2 Let $v \in C^2(\mathbf{R}^2 \times [0, T); \mathbf{R})$ be a function satisfying $\sup_{0 \le t \le T} \|v(t)\|_2 < \infty$. Then there exists a constant K > 0 such that

$$[\partial v(t)]_0 \le K \|\partial v(t)\|_2 \quad 0 \le t < T \tag{3.4}$$

holds.

Proposition 3.3 Let $v \in C^1(\mathbf{R}^2 \times [0, T); \mathbf{R})$ be a function satisfying v(x, t) = 0 when $|x| \ge t + R$ for some constant R > 0. Then there exists a constant L = L(R) > 0 such that

$$\left\| \frac{v(t)}{1 + |t - |x||} \right\|_{0} \le L \|\partial v(t)\|_{0} \quad 0 \le t < T$$
(3.5)

holds.

Proposition 3.4 Let $u \in C^2(\mathbf{R}^2 \times [0, T); \mathbf{R})$ be a solution to (1.1) and (1.2). Then

$$u(x, t) = 0 \quad for \quad |x| \ge t + M \tag{3.6}$$

holds.

See Corollary 1 in [7], Lemma 3.2 in [6] and Theorem 4a in [5] for the proof of Propositions 3.2, 3.3 and 3.4, respectively.

Proof of Lemma 3.1. Now we show Lemma 3.1 by 3 steps.

Step 1. There exist constants C' > 0 and $\varepsilon_1 > 0$ such that

$$|\partial u(x,t)|_j \le \frac{C'\varepsilon(1+\varepsilon^2\log(1+t))^j}{(1+|x|+t)^{1/2}} \quad (x,t) \in \mathbf{R}^2 \times [0,T)$$
 (3.7)

holds for $\varepsilon \in (0, \varepsilon_1)$ and j = 0, 1, 2. Here C' is independent of the constants c_l and m_l (l = 0, 1, 2, ..., k).

Firstly, by (3.1), we know that

$$|\partial u(x, t)|_2 \le \frac{KD\varepsilon}{(1+t)^{1/2}} \quad 0 \le t \le \frac{1}{\varepsilon}$$
 (3.8)

holds for $\varepsilon \in (0, \varepsilon')$. This implies that (3.7) is true when $0 \le t \le 1/\varepsilon$. Secondly, by (3.2) and Proposition 3.2, we find that

$$(1 + |x| + t)^{1/2} (1 + ||x| - t|)^{1/2} |\partial u(t, x)|_{l-2}$$

$$\leq [\partial u(t)]_{l-2}$$

$$\leq K ||\partial u(t)||_{l}$$

$$\leq 2Kc_{l}\varepsilon (1 + \varepsilon^{2} \log(1 + t))^{m_{l}}$$
(3.9)

holds for $(x, t) \in \mathbf{R}^2 \times [0, T)$ and $l = 2, 3, \ldots, k$. Then, setting

$$\Lambda_0 = \left\{ (y, s) \mid y \in \mathbf{R}^2, \ \frac{1}{\varepsilon} \le s < T, \ ||y| - s| \ge \frac{s}{2} \right\},$$
$$\Lambda_1 = \left\{ (y, s) \mid y \in \mathbf{R}^2, \ \frac{1}{\varepsilon} \le s < T, \ ||y| - s| \le \frac{s}{2} \right\},$$

we find that (3.7) is true for $(x, t) \in \Lambda_0$. In fact, by (3.9) and the fact that

$$1 + ||x| - t| \ge 1 + \frac{t}{2} \ge \frac{1}{2\varepsilon} \quad (x, t) \in \Lambda_0,$$

we have

$$(1+|x|+t)^{1/2}|\partial u(t,x)|_{2} \leq \frac{2Kc_{4}\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{4}}}{(1+t/2)^{1/2}}$$

$$\leq \frac{2^{5/4}Kc_{4}\varepsilon^{5/4}(1+\log(1+t))^{m_{4}}}{(1+t/2)^{1/4}} \quad (3.10)$$

$$\leq \varepsilon$$

for $\varepsilon \in (0, \varepsilon'')$, if we take ε'' as

$$\varepsilon'' < \min \Big\{ 1, \, \Big(2^{5/4} K c_4 \sup_{0 \le s} \frac{(1 + \log(1 + s))^{m_4}}{(1 + s/2)^{1/4}} \Big)^{-4} \Big\}.$$

In order to show (3.7) for $(x, t) \in \Lambda_1$, we prepare an estimate of u. It

676 A. Hoshiga

follows from (3.9) and Proposition 3.4 that

$$|\Gamma^{a}u(x,t)| = \left| -\int_{|x|}^{t+M} \partial_{r}\Gamma^{a}u(\lambda\omega,t)d\lambda \right|$$

$$\leq [\partial u(t)]_{4} \int_{|x|}^{t+M} \frac{1}{(1+t+\lambda)^{1/2}(1+|\lambda-t|)^{1/2}}d\lambda$$

$$\leq 2Kc_{4}\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{6}}$$

$$\times \int_{|x|-t}^{M} \frac{1}{(1+2t+\mu)^{1/2}(1+|\mu|)^{1/2}}d\mu$$

$$\leq 2Kc_{4}\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{6}} \left(\int_{0}^{M} d\mu + B\left(\frac{1}{2}, \frac{1}{2}\right)\right)$$

$$\leq 2Kc_{4}(M+\pi)\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{6}}$$

for $|a| \leq 4$. Here, B(p, q) stands for the beta function. Hence we have

$$|u(x, t)|_4 \le \tilde{C}_* \varepsilon (1 + \log(1 + t))^{m_6} \quad (x, t) \in \Lambda_1,$$
 (3.11)

if $\varepsilon \in (0, 1)$. Here \tilde{C}_* is a constant depending on c_4 . Since the operator ∂_j (j = 1, 2) can be written as

$$\partial_j = -\omega_j \partial_0 + \frac{1}{t} L_j + \frac{\omega_j}{t+r} S - \sum_{i=1}^2 \frac{r \omega_i \omega_j}{t(t+r)} L_i,$$

we find

$$\partial_{\alpha}v = -\omega_{\alpha}\partial_{0}v + O\left(\sum_{|\alpha|=1} \frac{|\Gamma^{a}v|_{0}}{t}\right) \quad \text{with} \quad \omega_{0} = -1, \tag{3.12}$$

$$(\partial_0 + \partial_r)v = O\left(\sum_{|a|=1} \frac{|\Gamma^a v|_0}{t}\right),\tag{3.13}$$

$$(\partial_0 + \partial_r)^2 v = O\left(\sum_{|a|=2} \frac{|\Gamma^a v|_0}{t^2}\right)$$
(3.14)

and therefore we obtain

$$\Box u = -C(-1, \omega)(\partial_0 u)^3 + O\left(\sum_{|b|=1} \left\{ \frac{|\partial u|_0^2 |\Gamma^b u|_0}{t} + \frac{|\partial u|_0 |\Gamma^b u|_0^2}{t^2} + \frac{|\Gamma^b u|_0^3}{t^3} \right\} + |\partial u|_0^4 \right). \quad (3.15)$$

We also introduce characteristic lines of (1.1). By (1.1), we can show that

$$r^{-1/2}(\partial_0 + \partial_r)(r^{1/2}\partial_0 v) = \frac{1}{2}\Box v + \frac{1}{2}(\partial_0 + \partial_r)^2 v + \frac{1}{2r}(\partial_0 + \partial_r)v + \frac{1}{2r^2}\Omega^2 v \quad (3.16)$$

holds for a function v(x, t). Then, setting $r_{\lambda}(s) = s + \lambda$ for each $\lambda \in \mathbf{R}$ and denoting $v(s) = v(r_{\lambda}(s)\omega, s)$, we obtain

$$2\frac{d}{ds}(r_{\lambda}(s)^{1/2}\partial_{0}v(s)) = r_{\lambda}(s)^{1/2}\Box v + r_{\lambda}(s)^{1/2}(\partial_{0} + \partial_{r})^{2}v + \frac{1}{r_{\lambda}(s)^{1/2}}(\partial_{0} + \partial_{r})v + \frac{1}{r_{\lambda}(s)^{3/2}}\Omega^{2}v \quad (3.17)$$

for each $\omega \in S^1$ and $\lambda \in \mathbf{R}$. For each $(x, t) \in \Lambda_1$, settting $x = r\omega$ and $\lambda = r - t$, we find that $r_{\lambda}(t) = r$. Then we call the line $(r_{\lambda}(s)\omega, s)$ $(0 \le s < T)$ the characteristic line of (1.1) passing through the point (x, t). Moreover, we denote

$$t_{\lambda} = \inf\{s \mid (r_{\lambda}(s)\omega, s) \in \Lambda_1\}$$

which is the time when $(r_{\lambda}(s), s) \in \partial \Lambda_1$. Note that $t_{\lambda} \geq 1/\varepsilon$ and $(r_{\lambda}(s)\omega, s) \in \Lambda_1$ holds for $t_{\lambda} \leq s < T$.

Hence, for any $(x, t) \in \Lambda_1$, by multiplying

$$P(s) = \exp\left(\frac{1}{2} \int_{t_0}^{s} C(-1, \omega) (\partial_0 u(\tau))^2 d\tau\right)$$

to the both sides of (3.17) with v = u and integrating it from t_{λ} to t, we have

$$P(t)|x|^{1/2}|\partial_{0}u(x,t)|$$

$$\leq r_{\lambda}(t_{\lambda})^{1/2}|\partial_{0}u(r_{\lambda}(t_{\lambda})\omega,t_{\lambda})|$$

$$+ C_{1}\int_{t_{\lambda}}^{t}P(s)\left(\frac{|\partial u|_{0}^{2}|u|_{1}}{(1+s)^{1/2}}+r_{\lambda}(s)^{1/2}|\partial u|_{0}^{4}+\frac{|u|_{2}}{(1+s)^{3/2}}\right)ds,$$

where we have used (3.13), (3.14), (3.15) and the fact

$$\frac{1}{As} \le \frac{1}{r_{\lambda}(s)} \le \frac{A}{1+s} \quad (r_{\lambda}(s)\omega, s) \in \Lambda_1 \tag{3.18}$$

for some constnat A > 0. Since P(t) is monotonously increasing and $P(t) \ge$

1, we have by (3.8), (3.9), (3.10) and (3.11),

$$|x|^{1/2} |\partial_0 u(x, t)| \le \frac{1}{P(t)} \Big\{ r_{\lambda}(t_{\lambda})^{1/2} |\partial_0 u(r_{\lambda}(t_{\lambda})\omega, t_{\lambda})|$$

$$+ C_1 \int_{t_{\lambda}}^t P(s) \Big(\frac{|\partial u|_0^2 |u|_1}{r_{\lambda}(s)^{1/2}} + r_{\lambda}(s)^{1/2} |\partial u|_0^4 + \frac{|u|_2}{r_{\lambda}(s)^{3/2}} \Big) ds \Big\}$$

$$\le (1 + KD)\varepsilon + \tilde{C}_1 \int_{t_{\lambda}}^t \frac{\varepsilon (1 + \log(1+s))^{4m_6}}{(1+s)^{3/2}} ds$$

$$\le (1 + KD)\varepsilon + \frac{\tilde{C}_2 \varepsilon}{(1+t_{\lambda})^{1/4}}$$

$$\le (1 + KD)\varepsilon + \tilde{C}_2 \varepsilon^{1+1/4} \quad (x, t) \in \Lambda_1.$$

Hereafter, C_j stands for constnats independent of c_l and m_l , while \tilde{C}_j stands for constants depending on c_l or m_l (l = 0, 1, 2, ..., k). Therefore, taking

$$\varepsilon_0''' \le \min \Big\{ \varepsilon', \, \varepsilon'', \, \frac{1}{\tilde{C}_2^4} \Big\},$$

we have

$$|x|^{1/2}|\partial_0 u(x,t)| \le (2+KD)\varepsilon \quad (x,t) \in \Lambda_1$$
 (3.19)

for $\varepsilon \in (0, \varepsilon_0''')$. Moreover, by (3.11), (3.12) and (3.19), we have

$$|x|^{1/2}|\partial_j u(x,t)| = |x|^{1/2}|\partial_0 u(x,t)| + O\left(\frac{|u|_1}{t}\right)$$

$$\leq (3 + KD)\varepsilon \quad (x,t) \in \Lambda_1 \tag{3.20}$$

for $\varepsilon \in (0, \varepsilon_0''')$, taking ε_0''' smaller if necessary. Therefore, by (3.18), (3.19) and (3.20), we find that

$$(1+t)^{1/2} |\partial u(x,t)|_0 \le C_0' \varepsilon \quad (x,t) \in \Lambda_1$$
 (3.21)

holds for $\varepsilon \in (0, \varepsilon_0''')$, if we take $C_0' > 3(3 + KD)A^{1/2}$.

Next we take $v = \Gamma u$ in (3.16). Here Γ stands for any one of Γ_{α} ($\alpha = 0, 1, \ldots, 6$). By (1.3), (2.1) and (3.15), we have

$$\Box \Gamma u = \Gamma \Box u + C \Box u$$

= $-3C(-1, \omega)(\partial_0 u)^2(\partial_0 \Gamma u) +$ (3.22)

$$+ O \bigg(|\partial u|_{0}^{3} + |\partial u|_{0}^{3} |\partial u|_{1} \\ + \sum_{\stackrel{|b|=1}{|c|=1}} \bigg\{ \frac{|\partial u|_{0}^{2} |\Gamma^{b}u|_{1} + |\partial u|_{0} |\Gamma^{b}u|_{0} |\partial u|_{1}}{t} + \\ + \frac{|\partial u|_{1} |\Gamma^{b}u|_{0}^{2} + |\partial u|_{0} |\Gamma^{b}u|_{0} |\Gamma^{c}u|_{1}}{t^{2}} + \frac{|\Gamma^{b}u|_{0}^{2} |\Gamma^{c}u|_{1}}{t^{3}} \bigg\} \bigg).$$

Hence, for any $(x, t) \in \Lambda_1$, by multiplying

$$P(s) = \exp\left(\frac{3}{2} \int_{t_{\lambda}}^{s} C(-1, \omega) (\partial_{0} u(\tau))^{2} d\tau\right)$$

to the both sides of (3.17)) with $v = \Gamma u$ and integrating it from t_{λ} to t, we have

$$\begin{split} &P(t)|x|^{1/2}|\partial_{0}\Gamma u(x,\,t)|\\ &\leq r_{\lambda}(t_{\lambda})^{1/2}|\partial_{0}\Gamma u(r_{\lambda}(t_{\lambda})\omega,\,t_{\lambda})|\\ &+C_{2}\int_{t_{\lambda}}^{t}P(s)\Big(r_{\lambda}(s)^{1/2}(|\partial u|_{0}^{3}+|\partial u|_{0}^{3}|\partial u|_{1})\\ &+\frac{|\partial u|_{0}^{2}|u|_{2}+|\partial u|_{0}|u|_{1}|\partial u|_{1}}{(1+s)^{1/2}}+\frac{|u|_{3}}{(1+s)^{3/2}}\Big)ds, \end{split}$$

where we have used (3.13), (3.14), (3.15) and (3.18). Since P(t) is monotonously increasing and $P(t) \geq 1$, we have by (3.8), (3.9), (3.10), (3.11), (3.21) and (3.22),

$$|x|^{1/2} |\partial_0 \Gamma u(x, t)| \le \frac{1}{P(t)} \left\{ r_\lambda(t_\lambda)^{1/2} |\partial_0 \Gamma u(r_\lambda(t_\lambda)\omega, t_\lambda)| + C_2 \int_{t_\lambda}^t P(s) \left(r_\lambda(s)^{1/2} (|\partial u|_0^3 + |\partial u|_0^3 |\partial u|_1) + \frac{|\partial u|_0^2 |u|_2 + |\partial u|_0 |u|_1 |\partial u|_1}{(1+s)^{1/2}} + \frac{|u|_3}{(1+s)^{3/2}} \right) ds \right\}$$

$$\le (1+KD)\varepsilon + \int_{t_\lambda}^t \left(\frac{C_3\varepsilon^3}{1+s} + \frac{\tilde{C}_3\varepsilon(1+\log(1+s))^{4m_6}}{(1+s)^{3/2}} \right) ds$$

$$\le (1+KD)\varepsilon + C_3\varepsilon^3 \log(1+t) + \frac{\tilde{C}_4\varepsilon}{(1+t_\lambda)^{1/4}}$$

$$\leq (1 + KD + C_3)\varepsilon(1 + \varepsilon^2 \log(1 + t)) + \tilde{C}_4 \varepsilon^{1 + 1/4}$$

$$\leq (2 + KD + C_3)\varepsilon(1 + \varepsilon^2 \log(1 + t)) \quad (x, t) \in \Lambda_1$$

for $\varepsilon \in (0, \varepsilon_1''')$, if we take

$$\varepsilon_1^{\prime\prime\prime} < \min \left\{ \varepsilon_0^{\prime\prime\prime}, \, \frac{1}{\tilde{C}_4^4} \right\}.$$

Namely we have

$$|x|^{1/2} |\partial_0 u(x, t)|_1$$

 $\leq 7(2 + KD + C_3)\varepsilon(1 + \varepsilon^2 \log(1 + t)) \quad (x, t) \in \Lambda_1 \quad (3.23)$

and therefore by using (3.11), (3.12) and (3.23), we have

$$|x|^{1/2} |\partial_j u(x, t)|_1 \le (7(2 + KD + C_3) + 1)$$

 $\times \varepsilon (1 + \varepsilon^2 \log(1 + t)) \quad (x, t) \in \Lambda_1 \quad (3.24)$

for $\varepsilon \in (0, \varepsilon_1''')$, taking ε_1''' smaller if necessary. Hence, by (3.18) and (3.24), we have

$$(1+t)^{1/2} |\partial u(x,t)|_1 \le C_1' \varepsilon (1+\varepsilon^2 \log(1+t)) \quad (x,t) \in \Lambda_1$$
 (3.25)

for $\varepsilon \in (0, \varepsilon_1''')$, if we take $C_1' > (21(2 + KD + C_3) + 2)A^{1/2}$.

Repeating the same argument, we find that there exist positive constants C_2' and ε_2''' ($<\varepsilon_1'''$) such that

$$(1+t)^{1/2} |\partial u(x, t)|_2 \le C_2' \varepsilon (1+\varepsilon^2 \log(1+t))^2 \quad (x, t) \in \Lambda_1$$

for $\varepsilon \in (0, \varepsilon_2''')$. Therefore, taking

$$C' = \max\{KD, 1, C'_0, C'_1, C'_2\}$$
 and $\varepsilon_1 = \varepsilon'''_2$,

we obtain (3.7).

Next we show the following

Step 2. Let ν be a small positive number. Then, there exist constants $\tilde{C}'' > 0$ and $\varepsilon_2 > 0$ such that

$$\|\partial u(t)\|_{k+1} \le \tilde{C}'' \varepsilon (1+t)^{\nu} \quad 0 \le t < T \tag{3.26}$$

holds for $\varepsilon \in (0, \varepsilon_2)$. Here \tilde{C}'' depends on the constants m_l and c_l (l = 0, 1, 2, ..., k).

By (1.1) and (2.1), we have

$$\Box \Gamma^a u = \Gamma^a F(\partial u) + \sum_{|b| < |a|}' \Gamma^b F(\partial u). \tag{3.27}$$

Here $\sum_{|b|<|a|}' A_b = \sum_{b<|a|} \gamma_b A_b$ with certain constants γ_b . Furthermore, by (1.3), we have

$$\Gamma^{a}F(\partial u) + \sum_{|b|<|a|}^{\prime} \Gamma^{b}F(\partial u)
= \sum_{\alpha,\beta,\gamma=0,1,2}^{\prime} \partial_{\alpha}u\partial_{\beta}u\Gamma^{a}\partial_{\gamma}u + \sum_{\substack{|c|+|d|+|e|\leq|a|\\|c|,|d|,|e|\leq|a|-1}}^{\prime} \Gamma^{c}\partial u\Gamma^{d}\partial u\Gamma^{e}\partial u
+ O\left(\sum_{|c|+|d|+|e|+|f|\leq|a|}^{\prime} \Gamma^{c}\partial u\Gamma^{d}\partial u\Gamma^{e}\partial u\Gamma^{f}\partial u\right)$$

$$= O\left(\left(|\partial u(t)|_{0}^{2} + |\partial u(t)|_{[|a|/2]}^{3}\right)|\partial u(x,t)|_{|a|} + |\partial u(t)|_{[|a|/2]}^{2}|\partial u(x,t)|_{|a|-1}\right).$$
(3.28)

Thus, multiplying $\partial_0 \Gamma^a u$ to both sides of (3.27), integrating it over \mathbf{R}^2 and summing up with respect to a over $|a| \leq k+1$, we have

$$\frac{d}{dt} \|\partial u(t)\|_{k+1} \le C_4 \Big(\big(|\partial u(t)|_0^2 + |\partial u(t)|_{[(k+1)/2]}^3 \big) \|\partial u(t)\|_{k+1} + |\partial u(t)|_{[(k+1)/2]}^2 \|\partial u(t)\|_k \Big).$$
(3.29)

Therefore, by (2.3), (2.4), (3.2), (3.7), (3.9), (3.29), the Gronwall inequality and the fact that $[(k+1)/2] \le k-2$ holds for $k \ge 4$, we have

$$\|\partial u(t)\|_{k+1} \le C_5 \Big(\|\partial u(0)\|_{k+1} + \int_0^t \frac{[\partial u(s)]_{[(k+1)/2]}^2}{1+s} \|\partial u(s)\|_k ds \Big)$$

$$\times \exp\left(C_4 \int_0^t \Big(|\partial u(s)|_0^2 + \frac{[\partial u(s)]_{[(k+1)/2]}^3}{(1+s)^{3/2}} \Big) ds \Big)$$

$$\le \Big(C_6 \varepsilon + \tilde{C}_5 \varepsilon^3 \int_0^t \frac{(1+\varepsilon^2 \log(1+s))^{3m_k}}{1+s} ds \Big)$$

$$\times \exp\left(\int_0^t \Big(\frac{C_7 \varepsilon^2}{1+s} + \frac{\tilde{C}_6 \varepsilon^3 (1+\varepsilon^2 \log(1+s))^{3m_k}}{(1+s)^{3/2}} \Big) ds \Big)$$

$$\le (C_6 \varepsilon + \tilde{C}_7 \varepsilon (1+\varepsilon^2 \log(1+t))^{3m_k+1})$$

$$\times \exp(C_7 \varepsilon^2 \log(1+t) + \tilde{C}_8 \varepsilon^3)$$

$$\leq \tilde{C}_9 \varepsilon (1+\varepsilon^2 \log(1+t))^{3m_k+1} \times (1+t)^{C_7 \varepsilon^2}$$

for $\varepsilon \in (0, \varepsilon_1)$. Hence, setting

$$\tilde{C}_{10} = \sup_{0 \le t} \left(\frac{(1 + \log(1+t))^{3m_k + 1}}{(1+t)^{\nu/2}} \right), \quad \tilde{C}'' = \tilde{C}_9 \tilde{C}_{10},$$

$$\varepsilon_2 = \min \left\{ \varepsilon_1, \frac{\sqrt{\nu}}{\sqrt{2C_7}} \right\},$$

we have

$$\|\partial u(t)\|_{k+1} \le \tilde{C}_9 \tilde{C}_{10} \varepsilon (1+t)^{\nu/2 + C_7 \varepsilon^2}$$

$$\le \tilde{C}'' \varepsilon (1+t)^{\nu}$$

for $\varepsilon \in (0, \varepsilon_2)$. This implies (3.26).

Finally, we show (3.3).

Step 3. We can determine constants c_l , m_l (l = 0, 1, 2, ..., k) and ε_0 so that (3.3) holds under the assumption (3.2).

By (3.1), we know that

$$\|\partial u(t)\|_k \le D\varepsilon \quad 0 \le t < \frac{1}{\varepsilon}$$
 (3.30)

for $\varepsilon \in (0, \varepsilon')$. Hence we have only to consider the case $1/\varepsilon \le t < T$.

In order to estimate $\|\partial u(t)\|_l$, we make use of (3.27) again, but estimate the right hand side more precisely this time.

When $(x, t) \in \Lambda_0$, we find that $P(1 + ||x| - t|) \ge (1 + |x| + t)$ holds for a certain constant P > 0. Hence we have

$$|\partial u(x,t)|_{l} \le \frac{P}{1+|x|+t}[\partial u(t)]_{l} \quad (x,t) \in \Lambda_{0}. \tag{3.31}$$

On the other hand, when $(x, t) \in \Lambda_1$, we find that

$$\frac{1}{O}(1+|x|+t) \le r \le Qt \quad (x,t) \in \Lambda_1$$
 (3.32)

for a certain constant Q > 0. Thus we have by (3.12) and (3.32),

$$\sum_{\alpha,\beta,\gamma=0}^{2} A^{\alpha\beta\gamma} (\partial_{\alpha} \Gamma^{a} u \partial_{\beta} u \partial_{\gamma} u + \partial_{\alpha} u \partial_{\beta} \Gamma^{a} u \partial_{\gamma} u + \partial_{\alpha} u \partial_{\beta} u \partial_{\gamma} \Gamma^{a} u)$$

$$= -3C(-1, \omega)(\partial_{0}u)^{2}\partial_{0}\Gamma^{a}u$$

$$+ O\left(\sum_{\stackrel{|b|=1}{|c|=1}} \left\{ \frac{|\partial u|_{0}^{2}|\Gamma^{b}u|_{|a|} + |\partial u|_{0}|\Gamma^{b}u|_{0}|\partial u|_{|a|}}{t} + \frac{|\partial u|_{0}|\Gamma^{b}u|_{0}|\Gamma^{c}u|_{|a|} + |\Gamma^{b}u|_{0}^{2}|\partial u|_{|a|}}{t^{2}} + \frac{|\Gamma^{b}u|_{0}^{2}|\Gamma^{c}u|_{|a|}}{t^{3}} \right) \right)$$

$$= -3C(-1, \omega)(\partial_{0}u)^{2}\partial_{0}\Gamma^{a}u$$

$$+ O\left(\frac{|\partial u|_{0}^{2}|u|_{|a|+1} + |\partial u|_{0}|u|_{1}|\partial u|_{|a|}}{1 + |x| + t}\right) (x, t) \in \Lambda_{1}.$$
(3.33)

Hence, combining (3.28) and (3.33), we obtain

$$\Gamma^{a}F(\partial u) + \sum_{|b|<|a|}^{\prime} \Gamma^{b}F(\partial u)
= -3C(-1, \omega)(\partial_{0}u)^{2}\partial_{0}\Gamma^{a}u
+ O\left(\frac{|\partial u|_{0}^{2}|u|_{|a|+1} + |\partial u|_{0}|u|_{1}|\partial u|_{|a|}}{1 + |x| + t}
+ |\partial u|_{||a|/2|}^{3}|\partial u|_{|a|} + (1 - \delta_{0|a|})|\partial u|_{||a|/2|}^{2}|\partial u|_{|a|-1}\right) \quad (x, t) \in \Lambda_{1}.$$

Firstly, we consider the case l=0. By (1.7), (3.27), (3.30), (3.31), (3.34) with a=0 and Propositions 3.3 and 3.4, we have

$$\begin{split} &\frac{d}{dt} \|\partial u(t)\|_{0}^{2} \\ &\leq \int_{\mathbf{R}^{2}} F(\partial u) \partial_{0} u dx \\ &\leq C_{8} \int_{||x|-t| \geq t/2} |\partial u(x,t)|_{0}^{4} dx - 3 \int_{||x|-t| \leq t/2} C(-1,\omega) |\partial u(x,t)|_{0}^{4} dx \\ &+ C_{8} \int_{||x|-t| \leq t/2} \frac{|\partial u(x,t)|_{0}^{3} |u(x,t)|_{1}}{1+|x|+t} dx + C_{8} \int_{\mathbf{R}^{2}} |\partial u(x,t)|_{0}^{5} dx \\ &\leq C_{9} \int_{||x|-t| \geq t/2} \frac{[\partial u(t)]_{0}^{2} |\partial u(x,t)|_{0}^{2}}{(1+t)^{2}} dx \\ &+ C_{8} \int_{||x|-t| \leq t/2} \frac{[\partial u(t)]_{0}^{2} |u(x,t)|_{1} |\partial u(x,t)|_{0}}{(1+t)^{2}(1+||x|-t|)} dx \\ &+ C_{10} \int_{\mathbf{R}^{2}} \frac{(C')^{3} \varepsilon^{3} |\partial u(x,t)|_{0}^{2}}{(1+t)^{3/2}} dx \end{split}$$

$$\leq C_{11} \left(\frac{\varepsilon^2}{(1+t)^{3/2}} \|\partial u(t)\|_0^2 + \sum_{|c|=1} \frac{[\partial u(t)]_0^2}{(1+t)^2} \left\| \frac{\Gamma^c u(t)}{1+||x|-t|} \right\|_0 \|\partial u(t)\|_0 \right)$$

$$\leq C_{12} \left(\frac{\varepsilon^2}{(1+t)^{3/2}} \|\partial u(t)\|_0^2 + \frac{[\partial u(t)]_0^2}{(1+t)^2} \|\partial u(t)\|_1 \|\partial u(t)\|_0 \right),$$

which implis

$$\frac{d}{dt} \|\partial u(t)\|_{0} \le C_{13} \left(\frac{\varepsilon^{2}}{(1+t)^{3/2}} \|\partial u(t)\|_{0} + \frac{[\partial u(t)]_{0}^{2}}{(1+t)^{2}} \|\partial u(t)\|_{1} \right).$$
(3.35)

Therefore, it follows from (3.2), (3.9), (3.30), (3.35) and the Gronwall inequality that

$$\|\partial u(t)\|_{0} \leq \left(\|\partial u(\varepsilon^{-1})\|_{0} + C_{13} \int_{1/\varepsilon}^{t} \frac{[\partial u(s)]_{0}^{2}}{(1+s)^{2}} \|\partial u(s)\|_{1} ds\right)$$

$$\times \exp\left(\int_{1/\varepsilon}^{t} \frac{C_{13}\varepsilon^{2}}{(1+s)^{3/2}} ds\right)$$

$$\leq \left(C_{14}\varepsilon + \tilde{C}_{11} \int_{1/\varepsilon}^{t} \frac{\varepsilon^{3} (1+\varepsilon^{2} \log(1+s))^{3m_{2}}}{(1+s)^{2}} ds\right)$$

$$\times \exp(C_{15}\varepsilon^{5/2})$$

$$\leq \left(C_{14}\varepsilon + \int_{1/\varepsilon}^{t} \frac{\tilde{C}_{12}\varepsilon^{3}}{(1+s)^{3/2}} ds\right) \exp(C_{15}\varepsilon^{5/2})$$

$$\leq (C_{14}\varepsilon + \tilde{C}_{13}\varepsilon^{7/2}) \exp(C_{15}\varepsilon^{5/2})$$

$$\leq (C_{14}+1)e\varepsilon$$

for $\varepsilon \in (0, \varepsilon_{00})$, if we take $\varepsilon_{00} = \min\{\varepsilon_2, 1/\tilde{C}_{13}^{2/5}, 1/C_{15}^{2/5}\}$. Hence we find that (3.3) holds for l = 0, if we take $c_0 = (C_{14} + 1)e$ and $m_0 = 0$.

Next, we estimate $\|\partial u(t)\|_l$ assuming c_i and m_i (i = 0, 1, 2, ..., l-1) are determined. By (1.7), (2.4), (3.27), (3.31), (3.34) and Proposition 3.3, we have

$$\begin{split} &\frac{d}{dt} \|\partial u(t)\|_{l}^{2} \\ &\leq \sum_{|a|\leq l}^{\prime} \int_{\mathbf{R}^{2}} (\Gamma^{a} F(\partial u) + \sum_{|b|<|a|}^{\prime} \Gamma^{b} F(\partial u)) \partial_{0} \Gamma^{a} u dx \\ &\leq \sum_{|a|\leq l}^{\prime} \left(\int_{||x|-t|\geq t/2} |\partial u(x,\,t)|_{[|a|/2]}^{2} |\partial u(x,\,t)|_{|a|}^{2} dx \right) \end{split}$$

$$\begin{split} &-3\int_{||x|-t|\leq t/2}C(-1,\omega)|\partial u(x,t)|_0^2|\partial u(x,t)|_{|a|}^2dx\\ &+\int_{||x|-t|\leq t/2}\frac{1}{1+|x|+t}(|\partial u(x,t)|_0^2|u(x,t)|_{|a|+1}\\ &+|\partial u(x,t)|_0|u(x,t)|_1|\partial u(x,t)|_{|a|})|\partial u(x,t)|_{|a|}dx\\ &+\int_{||x|-t|\leq t/2}|\partial u(x,t)|_{||a|/2|}^2|\partial u(x,t)|_{|a|-1}|\partial u(x,t)|_{|a|}dx\\ &+\int_{\mathbf{R}^2}|\partial u(x,t)|_{||a|/2|}^3|\partial u(x,t)|_{|a|}^2dx\\ &+\int_{\mathbf{R}^2}|\partial u(x,t)|_{||a|/2|}^3|\partial u(x,t)|_{|a|}^2dx\\ &+\int_{||x|-t|\leq t/2}\frac{[\partial u(t)]_{||a|/2|}^2|\partial u(x,t)|_{|a|}}{(1+t)^2}dx\\ &+\int_{||x|-t|\leq t/2}\frac{[\partial u(t)]_{||a|/2|}^2|\partial u(x,t)|_{|a|+1}|\partial u(x,t)|_{|a|}}{(1+t)^{3/2}}dx\\ &+\int_{||x|-t|\leq t/2}\frac{[\partial u(t)]_{||a|/2|}^2|\partial u(x,t)|_{|a|}}{(1+t)^{3/2}}dx\\ &+\int_{\mathbf{R}^2}\frac{[\partial u(t)]_{||a|/2|}^3|\partial u(x,t)|_{|a|}}{(1+t)^{3/2}}dx\\ &+\int_{\mathbf{R}^2}\frac{[\partial u(t)]_{||a|/2|}^3|\partial u(x,t)|_{|a|}^2}{(1+t)^{3/2}}dx\\ &\leq C_{16}\Big(\frac{[\partial u(t)]_{||t/2|}^2}{(1+t)^2}\|\partial u(t)\|_t^2+\sum_{|c|\leq t+1}\frac{[\partial u(t)]_0^2}{(1+t)^2}\Big\|\frac{\Gamma^c u(x,t)}{1+||x|-t|}\Big\|_0\|\partial u(t)\|_t\\ &+\frac{[\partial u(t)]_{||t/2|}^3}{(1+t)^{3/2}}\|\partial u(t)\|_t^2+|\partial u(t)|_{||t/2|}^2\|\partial u(t)\|_{l-1}\|\partial u(t)\|_t\\ &+\frac{[\partial u(t)]_{||t/2|}^3}{(1+t)^{3/2}}\|\partial u(t)\|_t^2\Big)\\ &\leq C_{17}\Big(\frac{[\partial u(t)]_{||t/2|}^2}{(1+t)^2}\|\partial u(t)\|_{l-1}\|\partial u(t)\|_t\\ &+\frac{[\partial u(t)]_0^2}{(1+t)^2}\|\partial u(t)\|_{l+1}\|\partial u(t)\|_t\\ &+\frac{[\partial u(t)]_0^2}{(1+t)^2}\|\partial u(t)\|_{l+1}\|\partial u(t)\|_t\\ &+\frac{[\partial u(t)]_0^2}{(1+t)^2}\|\partial u(t)\|_{l-1}\|\partial u(t)\|_t\Big), \end{split}$$

which implies

686 A. Hoshiga

$$\frac{d}{dt} \|\partial u(t)\|_{l} \leq C_{18} \left(\frac{[\partial u(t)]_{[l/2]}^{2} + [\partial u(t)]_{[l/2]}^{3} + [\partial u(t)]_{0} |u(t)|_{1}}{(1+t)^{3/2}} \|\partial u(t)\|_{l} + \frac{[\partial u(t)]_{0}^{2}}{(1+t)^{2}} \|\partial u(t)\|_{l+1} + |\partial u(t)|_{[l/2]}^{2} \|\partial u(t)\|_{l-1} \right). (3.36)$$

Therefore, when l = 1, 2, 3, 4, it follows from (3.2), (3.7), (3.11), (3.26) with $\nu < 1/4$, (3.30), (3.36), Proposition 3.4 and the Gronwall inequality that

$$\begin{split} \|\partial u(t)\|_{l} &\leq \left\{ \|\partial u(\varepsilon^{-1})\|_{l} + \int_{1/\varepsilon}^{t} \left(\frac{C_{18}[\partial u(s)]_{0}^{2}}{(1+s)^{2}} \|\partial u(s)\|_{l+1} \right. \right. \\ &\quad + \frac{C_{19}[\partial u(s)]_{[l/2]}^{2}}{1+s} \|\partial u(s)\|_{l-1} ds \right\} \\ &\quad \times \exp\left(C_{18} \int_{1/\varepsilon}^{t} \frac{[\partial u(s)]_{2}^{2} + [\partial u(s)]_{3}^{3} + [\partial u(s)]_{0}|u(s)|_{1}}{(1+s)^{3/2}} ds \right) \\ &\leq \left\{ C_{20\varepsilon} + \int_{1/\varepsilon}^{t} \left(\frac{\tilde{C}_{14}\varepsilon^{3}(1+\varepsilon^{2}\log(1+s))^{2m_{2}}}{(1+s)^{2-\nu}} \right. \\ &\quad + \frac{C_{21}c_{l-1}\varepsilon^{3}(1+\varepsilon^{2}\log(1+s))^{m_{l-1}}}{1+s} \right) ds \right\} \\ &\quad \times \exp\left(\int_{1/\varepsilon}^{t} \frac{\tilde{C}_{15}(1+\log(1+s))^{3m_{6}}\varepsilon^{2}}{(1+s)^{3/2}} ds \right) \\ &\leq \left\{ C_{20\varepsilon} + \tilde{C}_{16}\varepsilon^{3} + \frac{C_{21}c_{l-1}\varepsilon}{m_{l-1}+1} (1+\varepsilon^{2}\log(1+t))^{m_{l-1}+1} \right\} \\ &\quad \times \exp(\tilde{C}_{17}\varepsilon^{2}) \\ &\leq c_{l}\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{l}} \end{split}$$

for $\varepsilon \in (0, \varepsilon_{0l})$, if we take

$$\varepsilon_{0l} = \min \left\{ \varepsilon_2, \, \frac{1}{\tilde{C}_{16}^{1/2}}, \, \frac{1}{\tilde{C}_{17}^{1/2}} \right\}, \quad c_l = \left(C_{20} + 1 + \frac{C_{21}c_{l-1}}{m_{l-1} + 1} \right) e,$$

$$m_l = m_{l-1} + 1.$$

On the other hand, when $l \geq 5$, it follows from (3.2), (3.9), (3.26) with $\nu < 1/4$, (3.30), (3.36), Proposition 3.4, the Gronwall inequality and the fact $\lfloor l/2 \rfloor + 2 \leq l - 1$ that

$$\|\partial u(t)\|_{l} \leq \left\{ \|\partial u(\varepsilon^{-1})\|_{l} + \int_{1/\varepsilon}^{t} C_{22} \left(\frac{[\partial u(s)]_{0}^{2}}{(1+s)^{2}} \|\partial u(s)\|_{l+1} \right) \right\}$$

$$+ \frac{[\partial u(s)]_{[l/2]}^{2}}{1+s} \|\partial u(s)\|_{l-1} ds$$

$$\times \exp\left(C_{23} \int_{1/\varepsilon}^{t} \frac{[\partial u(s)]_{[l/2]}^{2} + [\partial u(s)]_{[l/2]}^{3} + [\partial u(s)]_{0} |u(s)|_{1}}{(1+s)^{3/2}} ds\right)$$

$$\leq \left\{C_{24}\varepsilon + \int_{1/\varepsilon}^{t} \left(\frac{\tilde{C}_{18}\varepsilon^{3}(1+\varepsilon^{2}\log(1+s))^{2m_{2}}}{(1+s)^{2-\nu}} + \frac{C_{25}c_{l-1}^{3}\varepsilon^{3}(1+\varepsilon^{2}\log(1+s))^{3m_{l-1}}}{1+s}\right) ds\right\}$$

$$\times \exp\left(\int_{1/\varepsilon}^{t} \frac{\tilde{C}_{19}(1+\log(1+s))^{3m_{l-1}+2m_{6}}\varepsilon^{2}}{(1+s)^{3/2}} ds\right)$$

$$\leq \left\{C_{24}\varepsilon + \tilde{C}_{20}\varepsilon^{3} + \frac{C_{25}c_{l-1}^{3}\varepsilon}{3m_{l-1}+1}(1+\varepsilon^{2}\log(1+t))^{3m_{l-1}+1}\right\}$$

$$\times \exp(\tilde{C}_{21}\varepsilon^{2})$$

$$\leq c_{l}\varepsilon(1+\varepsilon^{2}\log(1+t))^{m_{l}}$$

holds for $\varepsilon \in (0, \varepsilon_{0l})$, if we take

$$\varepsilon_{0l} = \min \left\{ \varepsilon_2, \, \frac{1}{\tilde{C}_{20}^{1/2}}, \, \frac{1}{\tilde{C}_{21}^{1/2}} \right\}, \quad c_l = \left(C_{24} + 1 + \frac{C_{25} c_{l-1}^3}{3m_{l-1} + 1} \right) e,$$

$$m_l = 3m_{l-1} + 1.$$

Hence, we obtain (3.3) if we take $\varepsilon_0 = \min\{\varepsilon_{00}, \varepsilon_{01}, \dots, \varepsilon_{0k}\}$. This completes the proof of Lemma 3.1.

References

[1] Godin P., Lifespan of solutions of semilinear wave equations in two space dimensions. Comm. in P.D.E. (5 and 6) 18 (1993), 895–916.

Hörmander L., The lifespan of classical solutions of nonlinear hyperbolic equations.
 Lecture Note in Math. vol. 1256, pp. 214–280, (1987).

- [3] Hörmander L., Lectures on Nonlinear Hyperbolic Differential Equations. Springer, 1997.
- [4] Hoshiga A., Existence and blowing up of solutions to systems of quasilinear wave equations in two space dimensions. Advances in Math. Sci. and Appli. (1) 15 (2005), 69–110.

- [5] John F., Blow-up for quasi-linear wave equations in three space dimensions. Comm. Pure Appli. Math. **34** (1981), 29–51.
- [6] Katayama S., Global existence for systems of nonlinear wave equations in two space dimensions, II. Publ. RIMS Kyoto Univ. 31 (1995), 645–665.
- [7] Klainerman S., Remarks on the global Sobolev inequalities in the Minkowski space \mathbf{R}^{n+1} . Comm. Pure Appli. Math. **40** (1987), 111–117.
- [8] Kubo H., Large time behavior of solutions to semilinear wave equations with dissipative structure. to appear in Expanded Volume of DCDS.

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