

## Seiberg-Witten theory and the geometric structure $\mathbf{R} \times H^2$

Mitsuhiro ITOH and Takahisa YAMASE

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**Abstract.** The moduli space of the solutions to the monopole equations over an oriented closed 3-manifold  $M$  carrying the geometric structure  $\mathbf{R} \times H^2$  is studied. Solving the parallel spinor equation, we obtain an explicit solution to the monopole equations. The moduli space consists of a single point with the Seiberg-Witten invariant  $\pm 1$ . Further, the (anti-)canonical line bundle  $K_M^{\pm 1}$  gives a monopole class of  $M$ .

*Key words:* Seiberg-Witten theory, geometric structure, monopole class, parallel spinor.

### 1. Introduction

Similar to the four-dimensional Seiberg-Witten theory, the study of solutions to the three-dimensional monopole equations

$$\begin{cases} c(*F_A) = \Phi \otimes \Phi^* - \frac{1}{2}|\Phi|^2 \text{Id}_W \\ D_A \Phi = 0 \end{cases}$$

over an oriented closed 3-manifold provides a new invariant of topology, the so-called Seiberg-Witten invariant. A class  $\alpha = c_1(L) \in H^2(M; \mathbf{R})$  is called a basic class if the Seiberg-Witten invariant is non-trivial. Furthermore, as a larger class,  $\alpha$  is called a monopole class if the monopole equations associated with  $\alpha$  have a solution for any metric  $h$  on  $M$ .

A generalization of Lichnerowicz's theorem holds also in the three-dimensional monopole equations as

$$0 = D_A D_A \Phi = \nabla_A^* \nabla_A \Phi + \frac{1}{4} s_h \Phi + \frac{1}{2} c(*F_A) \Phi$$

which leads to the well-known strong maximum principle that  $M$  with a metric of positive scalar curvature does not admit an irreducible solution. Another implication of this formula is the  $L^2$ -inequality

$$4 \int_M |F_A|^2 dv_h \leq \int_M s_h^2 dv_h \quad \text{so that} \quad \|\alpha_h\|_{(L^2, h)} \leq \frac{1}{4\pi} \|s_h\|_{(L^2, h)}$$

for the  $h$ -harmonic part  $\alpha_h$  of the 2-form representing  $\alpha$ . In [3] we obtained that if we assume the extremal situation above, namely, a solution satisfying

$$\|\alpha_h\|_{(L^2, h)} = \frac{1}{4\pi} \|s_h\|_{(L^2, h)}, \quad (1.1)$$

then  $\Phi$  and  $F_A$  are parallel and the scalar curvature  $s_h$  of  $h$  is negative constant so that the 3-manifold  $(M, h)$  must carry the geometric structure  $\mathbf{R} \times H^2$ . In this article, we call (1.1) the monopole extremal condition.

The main aims of this article are to determine the monopole class  $\alpha$  satisfying the monopole extremal condition above and to exhibit that under this condition the moduli space of solutions to the monopole equations consists of a single point, cut out transversely so that the Seiberg-Witten invariant is  $\pm 1$ .

**Main Theorem** *Let  $M$  be an oriented closed 3-manifold carrying the geometric structure  $\mathbf{R} \times H^2$  with the (anti-)canonical line bundle  $K_M^{\pm 1}$ . Here,  $K_M^{\pm 1} \rightarrow M$  is a complex line bundle naturally induced from the (anti-)canonical line bundle  $K_{H^2}^{\pm 1}$  over  $H^2$  by the quotient map:  $\mathbf{R} \times H^2 \rightarrow M$ . Suppose  $b_1(M) > 1$ . It follows then that (1) the moduli space of solutions to the monopole equations associated with the class  $\alpha = c_1(K_M^{\pm 1})$  and the metric  $h$  such that  $\pi^*h = dt^2 \oplus a^2 g_H$  consists of a single point and is transversal at this point and that (2)  $\alpha$  is a monopole class.*

**Remark** Proposition 5.1 in [4] is similar to the above theorem, although its proof is quite different from ours.

In Section 2, we review the three-dimensional Seiberg-Witten theory with the result of [3] and determine the monopole class  $\alpha = c_1(L)$  under the monopole extremal condition as  $L = K_M^{\pm 1}$ . An explicit form of spinor fields  $\Phi \in \Gamma(M; W)$ , which are parallel with respect to the canonical metric  $h$  is given in Section 3. Making use of these parallel spinor fields which turn out to be solutions to the monopole equations, we examine in Section 4 the moduli space  $\mathcal{M}(M; \alpha, h)$  of solutions associated with the metric  $h$  stated in Main Theorem (1). We can furthermore exhibit by applying the perturbation argument which is a typical device in the Seiberg-Witten theory

that the moduli space  $\mathcal{M}(M; \alpha, h')$  of solutions associated with an arbitrary metric  $h'$  cut out transversely so that the invariant  $SW(M, K_M^{\pm 1}) = \pm 1$  and as a byproduct that  $\alpha = c_1(K_M^{\pm 1})$  becomes a monopole class of  $M$ . Here, we need the topological restriction  $b_1(M) > 1$  for the perturbation trick being valid.

## 2. The monopole class and the (anti-)canonical line bundle

First, we will outline the three-dimensional Seiberg-Witten theory.

Let  $M$  be an oriented closed 3-manifold. Then there exists a  $Spin(3)^c$  structure on  $M$  defining the principal  $Spin(3)^c$ -bundle  $P$  associated with the orthonormal frame bundle  $SO(TM)$ . Let  $W$  be the spinor bundle associated with  $P$  and  $L = \det(W)$  be the determinant line bundle of  $W$ . The monopole equations are for a unitary connection  $A$  on  $L$  and a section  $\Phi$  of  $W$  as follows.

$$\begin{cases} c(*F_A) = \Phi \otimes \Phi^* - \frac{1}{2}|\Phi|^2 \text{Id}_W \\ D_A \Phi = 0 \end{cases}$$

Here,  $c : T^*M \rightarrow \text{End}(W)$  denotes the Clifford multiplication and  $*$  is the Hodge star operation. Further  $F_A$  is the curvature form of  $A$  and  $D_A$  is the Dirac operator twisted with  $A$ :

$$D_A : \Gamma(M; W) \xrightarrow{\nabla_A} \Gamma(M; T^*M \otimes W) \xrightarrow{c} \Gamma(M; W),$$

where  $\nabla_A$  is the spin connection on  $W$ .

As is well known, the monopole equations are invariant under the gauge action

$$(A, \Phi) \mapsto (A + g^{-1}dg, g^{-1}\Phi), \quad g \in \mathcal{G} = \Gamma(M; U(1))$$

so that we can define the moduli space of solutions to the monopole equations by the gauge action, namely,  $\mathcal{M} = \mathcal{S}/\mathcal{G}$ . Here,  $\mathcal{S}$  is the set of the solutions. It is known that  $\mathcal{M}$  has 0-dimensional compact oriented manifold structure ([1]).  $b_1(M) > 0$  guarantees that every solution  $(A, \Phi)$  is irreducible, that is,  $\Phi \neq 0$ . We usually define the Seiberg-Witten invariant  $SW(M, L)$  as the number of irreducible points, counted with sign in  $\mathcal{M}$ . Notice that  $\mathcal{M}$  has

irreducible points, provided  $SW(M, L) \neq 0$ .

In [3], for an oriented closed 3-manifold  $M$  and a monopole class  $\alpha$  of  $M$  we obtained that if  $M$  has a smooth metric  $h$  satisfying the monopole extremal condition, then  $\Phi$  and  $F_A$  are parallel and the scalar curvature  $s_h$  is negative constant. Moreover, in this article, we will get the explicit form of the monopole class  $\alpha$  when the monopole extremal condition is fulfilled. For this, denote by  $\pi : \mathbf{R} \times H^2 \rightarrow M$  the universal covering projection of  $M$ .

**Proposition 2.1** *Let  $\alpha$  be a monopole class of an oriented closed 3-manifold  $M$ . Suppose that there exists a smooth metric  $h$  on  $M$  which satisfies the monopole extremal condition. Then, the harmonic part  $\alpha_h$  is*

$$\alpha_h = \pm \frac{1}{2\pi} d\sigma_H,$$

where  $d\sigma_H$  denotes the 2-form on  $M$  whose lift  $\pi^*d\sigma_H$  is the area form of  $(H^2, g_H)$ .

*Proof.* From the argument in [3], the  $h$ -harmonic part  $\alpha_h$  of the monopole class  $\alpha$  is  $\alpha_h = \frac{i}{2\pi} F_A$  which is parallel. Then the lift  $\pi^*\alpha_h$  is parallel and symplectic over  $H^2$  so that it is proportional to the area form of  $(H^2, g_H)$ :

$$\alpha_h = \frac{i}{2\pi} F_A = \frac{c}{2\pi} d\sigma_H$$

for some real constant  $c$ . To determine  $c$ , we take the pull-back metric  $\pi^*h$  described as

$$\pi^*h = dt^2 \oplus a^2 g_H,$$

where  $a > 0$  and  $g_H = (dx^2 + dy^2)/y^2$  is the hyperbolic metric. Regarding  $H^2$  as the upper half plane  $\{z = x + iy \mid y > 0\}$ , we see that the scalar curvature  $s_h$  of  $h$  is  $-2/a^2$ . Since  $s_h$  is constant, we obtain

$$\|s_h\|_{(L^2, h)} = \sqrt{\int_M s_h^2 dv_h} = |s_h| \sqrt{\text{Vol}(M, h)} = \frac{2}{a^2} \sqrt{\text{Vol}(M, h)}.$$

Therefore, we get from the monopole extremal condition

$$\frac{|c|}{2\pi a^2} = \frac{1}{4\pi} \cdot \frac{2}{a^2}$$

so that  $c = \pm 1$  and hence  $\alpha_h = \pm \frac{1}{2\pi} d\sigma_H$ .  $\square$

Conversely, if the monopole class  $\alpha = c_1(L)$  satisfies  $\alpha_h = \pm \frac{1}{2\pi} d\sigma_H$ , then, as is easily seen, the monopole extremal condition holds.

Assume that  $M$  admits the geometric structure  $\mathbf{R} \times H^2$ . The (anti-) canonical line bundle  $K_{H^2}^{\pm 1}$  over  $H^2$  induce a complex line bundle denoted by  $K_M^{\pm 1}$ . This is because  $M$  is a  $\Gamma$ -quotient of  $\mathbf{R} \times H^2$ , where  $\Gamma$  is the discrete subgroup of  $Isom^+(\mathbf{R} \times H^2)$ , and the frame field  $\frac{1}{y} dz$  of  $K_{H^2}$  or  $\frac{1}{y} d\bar{z}$  of  $K_{H^2}^{-1}$  is invariant respectively under the action

$$z \mapsto \frac{az + b}{cz + d}, \quad a, b, c, d \in \mathbf{R}, \quad ad - bc = 1$$

so that  $K_{H^2}^{\pm 1}$  over  $\mathbf{R} \times H^2$  well descends to the bundle  $K_M^{\pm 1}$  over  $M$ , which we call the (anti-)canonical line bundle over  $M$ .

As a corollary of Main Theorem in [3], we can determine the complex line bundle  $L$  under the monopole extremal condition.

**Corollary 2.2** *Let  $L$  be a complex line bundle over an oriented closed 3-manifold  $M$ . Assume that the first Chern class of  $L$  is a monopole class  $\alpha$  of  $M$  and satisfies the monopole extremal condition. Then,  $L$  must be bundle-isomorphic to  $F \otimes K_M^{\pm 1}$ , where  $F$  is a complex line bundle with a flat connection and  $K_M^{\pm 1}$  is the (anti-)canonical line bundle over  $M$ .*

*Proof.* For simplicity, we write  $K_M^{\pm 1} = K$ . It suffices to show that

$$c_1(L \otimes K^{-1}) = c_1(L) - c_1(K) = 0 (= c_1(F)),$$

since over  $M$  the multiplicative group  $H^1(M; \mathcal{D}^\times)$ , the space of all equivalence classes of complex line bundle over  $M$ , is isomorphic to  $H^2(M; \mathbf{Z})$  via the map assigning a complex line bundle to its first Chern class (see [2]).

For this purpose, let  $D$  be the Hermitian holomorphic connection on  $K_{H^2}$  induced from the Levi-Civita connection  $\nabla$ . Its connection form  $A$  is easily computed as  $A = -\frac{i}{y} dz$  and  $F_A = dA = -i(dx \wedge dy)/y^2$  so that  $c_1(K_{H^2})$  coincides with  $\frac{1}{2\pi} [d\sigma_H]$ . This completes the proof.  $\square$

### 3. Parallel spinor solutions to the monopole equations

From now on, we take  $L = K_M^{\pm 1}$  and investigate an explicit form of the solutions to the monopole equations. The spinor bundle  $W$  is described as  $W = W_0 \otimes L_1$ , where  $W_0$  is the product bundle  $W_0 = M \times \mathbf{C}^2$  and  $L_1$  is some complex line bundle. Taking care that  $L = \det(W)$ , we obtain  $K_M^{\pm 1} = L_1^2$  so that  $L_1 = K_M^{\pm 1/2}$ . Hence we can take spinor fields

$$\Phi_0 = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \otimes \sqrt{dz}, \quad \Phi_0^- = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \otimes \sqrt{d\bar{z}} \in \Gamma(M; W), \quad W = W_0 \otimes K_M^{\pm 1/2},$$

where  $dz$  and  $d\bar{z}$  are regarded as sections of  $K_M^{\pm 1}$ . Under these conditions, we can show the following proposition.

**Proposition 3.1** *If  $\nabla_{A_0} \Phi_0 = \nabla_{A_0^-} \Phi_0^- = 0$ , where  $A_0$  and  $A_0^-$  are the connections of  $K_M^{\pm 1}$  associated with the Levi-Civita connection of  $(M, h)$ , then*

$$\Phi_0 = \begin{pmatrix} C/\sqrt{y} \\ 0 \end{pmatrix} \otimes \sqrt{dz}, \quad \Phi_0^- = \begin{pmatrix} 0 \\ C/\sqrt{y} \end{pmatrix} \otimes \sqrt{d\bar{z}}, \quad C = \pm\sqrt{-s_h}.$$

*Proof.* We consider the case for  $(A_0, \Phi_0)$  with  $L = K_M$ . (The case for  $(A_0^-, \Phi_0^-)$  is similar.) First, we see

$$(\nabla_{A_0})_X \Phi_0 = \left( \nabla_X \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \right) \otimes \sqrt{dz} + \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \otimes ((\nabla_{A_0})_X \sqrt{dz}), \quad (3.1.1)$$

where  $X$  is any tangent vector to  $M$ . By the definition of the spin connection, the first term is computed as follows.

$$\begin{aligned} \nabla_X \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} &= \begin{pmatrix} X\phi_1 \\ X\phi_2 \end{pmatrix} - \frac{1}{2} \sum_{i < j}^3 \omega_{ij}(X) c(e_i) c(e_j) \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \\ &= \begin{pmatrix} X\phi_1 \\ X\phi_2 \end{pmatrix} - \frac{i}{2} \omega_{23}(X) \begin{pmatrix} \phi_1 \\ -\phi_2 \end{pmatrix}, \end{aligned} \quad (3.1.2)$$

where  $\omega_{ij}$  are the connection forms of  $(M, h)$  with respect to the orthonormal frame  $\{dt, \frac{1}{\alpha y} dx, \frac{1}{\alpha y} dy\}$ . Here, the lift of  $h$  is  $dt^2 \oplus a^2 g_H$ . On the other

hand, since  $dz$  is regarded as a section of  $K_M$ , the second term of (3.1.1) is computed for  $X = \frac{\partial}{\partial z}$  as

$$(\nabla_{A_0})_{\frac{\partial}{\partial z}} \sqrt{dz} = \left( \frac{\partial}{\partial z} \log y \right) \sqrt{dz}, \quad (3.1.3)$$

where  $z = x + iy$ . Moreover from the local product structure of  $M$ , we see

$$(\nabla_{A_0})_{\frac{\partial}{\partial t}} \sqrt{dz} = 0, \quad (\nabla_{A_0})_{\frac{\partial}{\partial \bar{z}}} \sqrt{dz} = 0.$$

Substituting (3.1.2) and (3.1.3) into (3.1.1), from  $(\nabla_{A_0})_X \Phi_0 = 0$  for  $X = \frac{\partial}{\partial z}$ , we obtain

$$\begin{aligned} X\phi_1 - \frac{i}{2}\omega_{23}(X)\phi_1 + \left( \frac{\partial}{\partial z} \log y \right) \phi_1 &= 0, \\ X\phi_2 + \frac{i}{2}\omega_{23}(X)\phi_2 + \left( \frac{\partial}{\partial z} \log y \right) \phi_2 &= 0. \end{aligned}$$

Using  $\omega_{23} = -\frac{1}{y}dx$ , we get

$$\frac{\partial\phi_1}{\partial z} - \frac{i}{4y}\phi_1 = 0, \quad \frac{\partial\phi_2}{\partial z} - \frac{3i}{4y}\phi_2 = 0.$$

Similarly, for  $X = \frac{\partial}{\partial \bar{z}}$  we get

$$\frac{\partial\phi_1}{\partial \bar{z}} + \frac{i}{4y}\phi_1 = 0, \quad \frac{\partial\phi_2}{\partial \bar{z}} - \frac{i}{4y}\phi_2 = 0.$$

Solving the simultaneous equations for  $\phi_1$  and  $\phi_2$ , we get  $\phi_1 = \frac{C}{\sqrt{y}}$  and  $\phi_2 = 0$ .

Now we obtain

$$\Phi_0 = \begin{pmatrix} C/\sqrt{y} \\ 0 \end{pmatrix} \otimes \sqrt{dz}.$$

On the other hand, using  $\nabla_{A_0}\Phi_0 = 0$  and Lichnerowicz's formula, we get

$$s_h = -|\Phi_0|^2 \quad \text{so that} \quad C = \pm\sqrt{-s_h}. \quad \square$$

Moreover, we can show the following corollary.

**Corollary 3.2** *For the monopole class  $\alpha$  whose  $h$ -harmonic part is*

$$\alpha_h = \frac{i}{2\pi}F_{A_0} = \frac{1}{2\pi}d\sigma_H \quad \text{or} \quad \alpha_h = \frac{i}{2\pi}F_{A_0^-} = -\frac{1}{2\pi}d\sigma_H,$$

$(A_0, \Phi_0)$  or  $(A_0^-, \Phi_0^-)$  with  $|\Phi_0| = |\Phi_0^-| = \sqrt{2}$  is a solution to the monopole equations for  $L = K_M$  or  $L = K_M^{-1}$ , respectively.

*Proof.* We consider the case for  $(A_0, \Phi_0)$  with  $L = K_M$ . (The case for  $(A_0^-, \Phi_0^-)$  is similar.) Since  $D_A = c \circ \nabla_A$ , it is clear that  $D_{A_0}\Phi_0 = 0$ . On the other hand, the curvature form of  $A_0$  is described as

$$F_{A_0} = \pm id\sigma_H = \pm ie^2 \wedge e^3,$$

where  $e^1, e^2, e^3$  are the dual orthonormal frame of  $(M, h)$ . Therefore we obtain

$$c(*F_{A_0}) = \pm \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}.$$

In general for  $\Phi = \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix}$ , we get

$$\Phi \otimes \Phi^* - \frac{1}{2}|\Phi|^2 \text{Id}_W = \begin{pmatrix} \frac{1}{2}(|\varphi_1|^2 - |\varphi_2|^2) & \varphi_1 \overline{\varphi_2} \\ \overline{\varphi_1} \varphi_2 & \frac{1}{2}(|\varphi_2|^2 - |\varphi_1|^2) \end{pmatrix}.$$

Therefore for  $\Phi_0 = \begin{pmatrix} \varphi \\ 0 \end{pmatrix}$ , we obtain

$$\begin{aligned} \Phi_0 \otimes \Phi_0^* - \frac{1}{2}|\Phi_0|^2 \text{Id}_W &= \begin{pmatrix} \frac{1}{2}|\varphi|^2 & 0 \\ 0 & -\frac{1}{2}|\varphi|^2 \end{pmatrix} \\ &= -\frac{1}{2}|\varphi|^2 \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} = -\frac{1}{2}|\Phi_0|^2 \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}, \end{aligned}$$

Hence in order for  $(A_0, \Phi_0)$  to satisfy the first monopole equation, we may take

$$F_{A_0} = -id\sigma_H, \quad |\Phi_0| = \sqrt{2}. \quad \square$$

#### 4. The moduli space

*Proof of Main Theorem (1).* We consider the case for  $\alpha = c_1(K_M)$ . (The case for  $\alpha = c_1(K_M^{-1})$  is similar.) In this case,  $\alpha_h = \frac{1}{2\pi}d\sigma_H$  so that by the proof of Proposition 2.1, we obtain the monopole extremal condition

$$\|\alpha_h\|_{(L^2, h)} = \frac{1}{4\pi} \|s_h\|_{(L^2, h)}$$

for the metric  $h$  whose lift  $\pi^*h$  has the form  $dt^2 \oplus a^2g_H$ . Let  $(A, \Phi)$  be an arbitrary solution associated with the class  $\alpha$  and the metric  $h$ . In this case, recall that  $\nabla_A\Phi = 0$  holds. Now we take  $(A_0, \Phi_0)$  which satisfies

$$\alpha = \frac{1}{2\pi}d\sigma_H = \frac{i}{2\pi}F_{A_0} \quad \text{and} \quad \Phi_0 = \begin{pmatrix} \sqrt{2/y} \\ 0 \end{pmatrix} \otimes \sqrt{dz}.$$

From Corollary 3.2,  $(A_0, \Phi_0)$  is a solution. We can show that any solution  $(A, \Phi)$  is gauge equivalent to  $(A_0, \Phi_0)$ .

For this, we take  $A = A_0 + ia$ ,  $a \in \Omega^1(M)$  so that  $F_A = F_{A_0} + ida$ . Since  $F_A$  and  $F_{A_0}$  are harmonic, we obtain  $da = 0$  and  $F_A = F_{A_0}$ . Moreover, by the first monopole equation, we get  $|\Phi_0|^2 = 2|F_{A_0}| = 2|F_A| = |\Phi|^2$  and

$$\Phi_0 \otimes \Phi_0^* - \frac{1}{2}|\Phi_0|^2\text{Id}_W = c(*F_{A_0}) = c(*F_A) = \Phi \otimes \Phi^* - \frac{1}{2}|\Phi|^2\text{Id}_W$$

so that  $\Phi_0 \otimes \Phi_0^* = \Phi \otimes \Phi^*$ . Taking  $\Phi_0 = \begin{pmatrix} \varphi \\ 0 \end{pmatrix}$  and  $\Phi = \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix}$ , we get  $|\varphi|^2 = |\varphi_1|^2$  and  $|\varphi_2|^2 = 0$  so that there exists  $g \in \mathcal{G}$  such that  $\Phi = g^{-1}\Phi_0$ . Therefore

$$\begin{aligned} 0 &= \nabla_A\Phi = dg^{-1} \otimes \Phi_0 + g^{-1}\nabla_A\Phi_0 \\ &= dg^{-1} \otimes \Phi_0 + g^{-1}\nabla_{A_0}\Phi_0 + g^{-1}ia \otimes \Phi_0 \\ &= dg^{-1} \otimes \Phi_0 + g^{-1}ia \otimes \Phi_0 \end{aligned}$$

and hence

$$ia \otimes \Phi_0 = -gdg^{-1} \otimes \Phi_0 = g^{-1}dg \otimes \Phi_0, \quad \text{namely, } (A, \Phi) = (A_0 + g^{-1}dg, g^{-1}\Phi_0),$$

which implies that  $(A, \Phi)$  is gauge equivalent to  $(A_0, \Phi_0)$ .

From now on, we will show the transversality of the moduli space  $\mathcal{M}$ . To show this, we consider the following complex which turns out to be elliptic by the subsequent lemma, Lemma 4.1.

$$\mathcal{C} : 0 \rightarrow \Omega^0(M) \xrightarrow{G} \Omega^1(M) \oplus \Gamma(W) \xrightarrow{T} \Omega^1(M) \oplus \Gamma(W) \xrightarrow{S} \Omega^0(M) \rightarrow 0,$$

where

$$G_{(A_0, \Phi_0)}(u) = (du, -iu\Phi_0), \quad S_{(A_0, \Phi_0)}(a, \varphi) = \delta a + i\text{Im}\langle \Phi_0, \varphi \rangle$$

and  $T_{(A_0, \Phi_0)}(a, \varphi) = (b, \psi)$ , where

$$b = c(i * da) - \Phi_0 \otimes \varphi^* - \varphi \otimes \Phi_0^* + \frac{1}{2}(\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle)\text{Id}_W,$$

$$\psi = D_{A_0}\varphi + ic(a)\Phi_0.$$

**Lemma 4.1** (1)  $T \circ G = 0$ , (2)  $S \circ T = 0$ , (3)  $\text{Index}(\mathcal{C}) = 0$ .

*Proof of (1).* By definition,

$$(b, \psi) = T \circ G_{(A_0, \Phi_0)}(u) = T_{(A_0, \Phi_0)}(du, -iu\Phi_0),$$

where

$$b = c(i * d(du)) - iu|\Phi_0|^2 + iu|\Phi_0|^2 + \frac{1}{2}(iu|\Phi_0|^2 - iu|\Phi_0|^2)\text{Id}_W,$$

$$\psi = D_{A_0}(-iu\Phi_0) + ic(du)\Phi_0 = -i(c(du)\Phi_0 + uD_{A_0}\Phi_0) + ic(du)\Phi_0.$$

Obviously  $b$  and  $\psi$  vanish. □

*Proof of (2).* Let  $S \circ T_{(A_0, \Phi_0)}(a, \varphi) = S_{(A_0, \Phi_0)}(b, \psi)$ . It is sufficient to show

$$\int_M \langle S(b, \psi), u \rangle dv_h = 0 \quad \text{for any } u \in \Omega^0(M).$$

By definition,

$$b = *da + ic^{-1} \left( \Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^* - \frac{1}{2}(\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle)\text{Id}_W \right),$$

$$\psi = D_{A_0}\varphi + ic(a)\Phi_0$$

so that

$$\begin{aligned} S(b, \psi) &= \delta \left( * da + ic^{-1} \left( \Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^* - \frac{1}{2} (\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle) \text{Id}_W \right) \right) \\ &\quad + i \text{Im} \langle \Phi_0, D_{A_0}\varphi + ic(a)\Phi_0 \rangle \\ &= i \delta \left( c^{-1} \left( \Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^* - \frac{1}{2} (\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle) \text{Id}_W \right) \right) \\ &\quad + i \text{Im} \langle \Phi_0, D_{A_0}\varphi \rangle. \end{aligned}$$

Therefore

$$\begin{aligned} &\int_M \langle S(b, \psi), u \rangle dv_h \\ &= \int_M \left\langle i \delta \left( c^{-1} \left( \Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^* - \frac{1}{2} (\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle) \text{Id}_W \right) \right) \right. \\ &\quad \left. + i \text{Im} \langle \Phi_0, D_{A_0}\varphi \rangle, u \right\rangle dv_h \\ &= i \int_M \langle \Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^*, c(du) \rangle dv_h \\ &\quad + i \int_M \frac{1}{2} u (\langle \Phi_0, D_{A_0}\varphi \rangle - \langle D_{A_0}\varphi, \Phi_0 \rangle) dv_h \\ &= i \int_M \langle \Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^*, c(du) \rangle dv_h \\ &\quad + \frac{i}{2} \int_M (\langle c(du)\Phi_0, \varphi \rangle - \langle \varphi, c(du)\Phi_0 \rangle) dv_h \\ &= i \int_M \frac{1}{2} \text{tr} \begin{pmatrix} \varphi_{01}\overline{\varphi_1} + \varphi_1\overline{\varphi_{01}} & \varphi_{01}\overline{\varphi_2} \\ \varphi_2\overline{\varphi_{01}} & 0 \end{pmatrix} \begin{pmatrix} -ia_1 & -a_2 - ia_3 \\ a_2 - ia_3 & ia_1 \end{pmatrix} dv_h \\ &\quad + \frac{i}{2} \int_M \left\langle \left\langle \begin{pmatrix} ia_1\varphi_{01} \\ -a_2\varphi_{01} + ia_3\varphi_{01} \end{pmatrix}, \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix} \right\rangle \right. \\ &\quad \left. - \left\langle \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix}, \begin{pmatrix} ia_1\varphi_{01} \\ -a_2\varphi_{01} + ia_3\varphi_{01} \end{pmatrix} \right\rangle \right\rangle dv_h = 0. \end{aligned}$$

□

*Proof of (3).* Split  $\mathcal{C}$  into the direct sum of the following complexes;

$$\begin{aligned} \mathcal{C}_1 &: 0 \rightarrow \Omega^0(M) \xrightarrow{d} \Omega^1(M) \xrightarrow{*d} \Omega^1(M) \xrightarrow{\delta} \Omega^0(M) \rightarrow 0 \\ \mathcal{C}_2 &: 0 \rightarrow \Gamma(W) \xrightarrow{D_{A_0}} \Gamma(W) \rightarrow 0. \end{aligned}$$

The first one is equivalent to the de Rham complex so that  $\text{Index}(\mathcal{C}_1) = \chi(M) = 0$ . The second one is the spin complex and so  $\text{Index}(\mathcal{C}_2) = \text{Index}D_{A_0} = 0$ . Therefore

$$\text{Index}(\mathcal{C}) = \text{Index}(\mathcal{C}_1) + \text{Index}(\mathcal{C}_2) = 0. \quad \square$$

Using Lemma 4.1 (3), by definition,

$$\text{Index}(\mathcal{C}) = \dim H^0(\mathcal{C}) - \dim H^1(\mathcal{C}) + \dim H^2(\mathcal{C}) - \dim H^3(\mathcal{C}) = 0.$$

Since the solution is irreducible, if  $u \in \Omega^0(M)$  satisfies

$$G_{(A_0, \Phi_0)}(u) = (du, -iu\Phi_0) = (0, 0),$$

then  $u = 0$  so that  $H^0(\mathcal{C}) = \text{Ker}G = \{0\}$ . Moreover since  $S_{(A_0, \Phi_0)}^*(u) = (du, iu\Phi_0)$ , we have  $S_{(A_0, \Phi_0)}^* = G_{(A_0, -\Phi_0)}$  so that  $\text{Ker}S^* = \{0\}$ . Therefore  $H^3(\mathcal{C}) = \Omega^0(M)/\text{Im}S$  is isomorphic to  $\text{Ker}S^* = \{0\}$  and hence  $H^3(\mathcal{C}) = \{0\}$ . Consequently,  $H^1(\mathcal{C}) \cong H^2(\mathcal{C})$ . Therefore the surjectivity of  $T$  is equivalent to  $\text{Ker}S/\text{Im}T = \{0\}$  which is equivalent to  $\text{Ker}T/\text{Im}G = \{0\}$ . This is also equivalent to

$$\{(a, \varphi) \mid T_{(A_0, \Phi_0)}(a, \varphi) = (0, 0), G_{(A_0, \Phi_0)}^*(a, \varphi) = 0\} = \{(0, 0)\},$$

where

$$G_{(A_0, \Phi_0)}^*(a, \varphi) = \delta a - i\text{Im}\langle \Phi_0, \varphi \rangle.$$

It is clear that  $T_{(A_0, \Phi_0)}(a, \varphi) = 0$  implies

$$\begin{aligned} c(da) &= -i\left(\Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^* - \frac{1}{2}(\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle)\text{Id}_W\right), \\ D_{A_0}\varphi &= -ic(a)\Phi_0 \end{aligned}$$

and that  $G_{(A_0, \Phi_0)}^*(a, \varphi) = 0$  implies

$$\delta a = i \operatorname{Im} \langle \Phi_0, \varphi \rangle = \frac{i}{2} (\langle \Phi_0, \varphi \rangle - \langle \varphi, \Phi_0 \rangle).$$

By the direct computation together with the fact that  $\nabla_{A_0} \Phi_0 = D_{A_0} \Phi_0 = 0$ , we get

$$\begin{aligned} D_{A_0} D_{A_0} \varphi &= -i D_{A_0} (c(a) \Phi_0) \\ &= -i ((\delta a) \Phi_0 - 2(\nabla_{A_0})_{a^\#} \Phi_0 + c(da) \Phi_0 - c(a) D_{A_0} \Phi_0) \\ &= \frac{1}{2} (\langle \Phi_0, \varphi \rangle - \langle \varphi, \Phi_0 \rangle) \Phi_0 - ic(da) \Phi_0. \end{aligned}$$

Here we made use of the formula:

$$D_A(c(a)\Phi) = (\delta a)\Phi - 2(\nabla_A)_X \Phi + c(da)\Phi,$$

$a \in \Omega^1(M)$  and  $X = a^\# \in \mathcal{X}(M)$ . Now we have

$$\begin{aligned} c(da)\Phi_0 &= -i \left( \Phi_0 \otimes \varphi^* + \varphi \otimes \Phi_0^* - \frac{1}{2} (\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle) \operatorname{Id}_W \right) \Phi_0 \\ &= -i \left( \langle \Phi_0, \varphi \rangle \Phi_0 + |\Phi_0|^2 \varphi - \frac{1}{2} (\langle \Phi_0, \varphi \rangle + \langle \varphi, \Phi_0 \rangle) \Phi_0 \right) \\ &= -i \left( |\Phi_0|^2 \varphi + \frac{1}{2} (\langle \Phi_0, \varphi \rangle - \langle \varphi, \Phi_0 \rangle) \Phi_0 \right), \end{aligned}$$

so that the term  $D_{A_0} D_{A_0} \varphi$  becomes

$$\begin{aligned} D_{A_0} D_{A_0} \varphi &= \frac{1}{2} (\langle \Phi_0, \varphi \rangle - \langle \varphi, \Phi_0 \rangle) \Phi_0 - ic(da) \Phi_0 \\ &= \frac{1}{2} (\langle \Phi_0, \varphi \rangle - \langle \varphi, \Phi_0 \rangle) \Phi_0 - |\Phi_0|^2 \varphi - \frac{1}{2} (\langle \Phi_0, \varphi \rangle - \langle \varphi, \Phi_0 \rangle) \Phi_0 \\ &= -|\Phi_0|^2 \varphi. \end{aligned}$$

Therefore

$$\int_M \langle D_{A_0} D_{A_0} \varphi, \varphi \rangle dv_h = - \int_M |\Phi_0|^2 |\varphi|^2 dv_h \quad \text{or}$$

$$\int_M |D_{A_0} \varphi|^2 dv_h = - \int_M |\Phi_0|^2 |\varphi|^2 dv_h.$$

$|\Phi_0|$  is positive constant because the solution  $(A_0, \Phi_0)$  is irreducible and  $\Phi_0$  is parallel. Hence we conclude  $\varphi = 0$  so that  $a = 0$  by  $-ic(a)\Phi_0 = D_{A_0}\varphi = 0$ . From the above arguments, the transversality of  $\mathcal{M}$  is completely derived.  $\square$

*Proof of Main Theorem (2).* In order to see that the class  $\alpha = c_1(K_M)$  is a monopole class, we show that the Seiberg-Witten invariant does not vanish with respect to an arbitrary metric on  $M$ . We consider the case where a given metric  $h$  is arbitrary. In this case, we cannot always make use of the condition  $\nabla_A \Phi = 0$ . We usually think of the perturbed monopole equations as follows.

$$\begin{cases} c(*F_A + i\rho) = \Phi \otimes \Phi^* - \frac{1}{2}|\Phi|^2 \text{Id}_W \\ D_A \Phi = 0 \end{cases}$$

Here,  $\rho$  is a co-closed 1-form. With respect to these perturbed equations, it is known that the Seiberg-Witten invariant is independent of metrics  $g$  and perturbations  $\rho$  ([1]). More precisely, given a generic path  $(g_t, \rho_t)$ ,  $t \in [0, 1]$  connecting  $(g_0, \rho_0)$  and  $(g_1, \rho_1)$ , it is known that

$$SW_{(g_0, \rho_0)}(M, L) = SW_{(g_1, \rho_1)}(M, L).$$

To apply the perturbed argument to our case, we take  $L = K_M$  and  $(g_0, \rho_0) = (h, 0)$ . Main Theorem (1) together with the definition of the Seiberg-Witten invariant implies  $SW_{(h, 0)}(M, K_M) = \pm 1$  so that

$$SW(M, K_M) = \pm 1 (\neq 0).$$

This implies that the monopole equations associated with  $\alpha = c_1(K_M)$  has solutions which are irreducible by  $b_1(M) > 1$ . Hence  $\alpha$  is a monopole class.  $\square$

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Mitsuhiro Itoh  
Institute of Mathematics  
University of Tsukuba  
305-8571, TSUKUBA, JAPAN  
e-mail: itohm@sakura.cc.tsukuba.ac.jp

Takahisa Yamase  
Graduate School of Pure and Applied Sciences  
University of Tsukuba  
305-8571, TSUKUBA, JAPAN  
e-mail: ks7gauge@math.tsukuba.ac.jp