

## ON A CLASS OF INTERMEDIATE LOCAL-NONLOCAL ELLIPTIC PROBLEMS

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ABSTRACT. This paper is concerned with the existence of solutions for a class of intermediate local-nonlocal boundary value problems of the following type:

$$(IP) \quad -\operatorname{div} \left[ a \left( \int_{\Omega(x,r)} u(y) dy \right) \nabla u \right] = f(x, u, \nabla u) \quad \text{in } \Omega, \quad u \in H_0^1(\Omega),$$

where  $\Omega$  is a bounded domain of  $\mathbb{R}^N$ ,  $a: \mathbb{R} \rightarrow \mathbb{R}$  is a continuous function,  $f: \Omega \times \mathbb{R} \times \mathbb{R}^N$  is a given function,  $r > 0$  is a fixed number,  $\Omega(x, r) = \Omega \cap B(x, r)$ , where  $B(x, r) = \{y \in \mathbb{R}^N : |y - x| < r\}$ . Here  $|\cdot|$  is the Euclidian norm,

$$\int_{\Omega(x,r)} u(y) dy = \frac{1}{\operatorname{meas}(\Omega(x,r))} \int_{\Omega(x,r)} u(y) dy$$

and  $\operatorname{meas}(X)$  denotes the Lebesgue measure of a measurable set  $X \subset \mathbb{R}^N$ .

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## 1. Introduction

In this work we will be concerned with the intermediate class of local-nonlocal elliptic problems

$$(IP) \quad -\operatorname{div}\left(a\left(\int_{\Omega(x,r)} u(y) dy\right)\nabla u\right) = f(x, u, \nabla u) \quad \text{in } \Omega, \quad u \in H_0^1(\Omega),$$

where  $\Omega \subset \mathbb{R}^N$ ,  $N \geq 1$ , is a bounded domain,  $a: \mathbb{R} \rightarrow \mathbb{R}$  is a continuous function,  $r > 0$  is a fixed real number,

$$\Omega(x, r) := \Omega \cap B(x, r), \quad \text{with } B(x, r) := \{y \in \mathbb{R}^N; |y - x| < r\}.$$

Here  $|\cdot|$  is the usual Euclidian norm of  $\mathbb{R}^N$  and

$$\int_{\Omega(x,r)} u(y) dy = \frac{1}{\operatorname{meas}(\Omega(x, r))} \int_{\Omega(x,r)} u(y) dy,$$

where  $\operatorname{meas}(\Omega(x, r))$  is the Lebesgue measure of the set  $\Omega(x, r)$ .

Note that (IP) is a class of interpolating problems between the purely local problems

$$(L) \quad -\operatorname{div}(a(u(x))\nabla u) = f(x, u, \nabla u) \quad \text{in } \Omega, \quad u \in H_0^1(\Omega),$$

and the nonlocal problems

$$(NL) \quad -\operatorname{div}\left(a\left(\int_{\Omega} u(x) dx\right)\nabla u\right) = f(x, u, \nabla u) \quad \text{in } \Omega, \quad u \in H_0^1(\Omega).$$

Note that in our case, we are considering a nonlocal quantity  $\int_{\Omega(x,r)} u(y) dy$  which is calculated locally in neighbourhoods of the form  $\Omega(x, r)$ .

REMARK 1.1. Although we are working in the space  $H_0^1(\Omega)$ , we may treat problem (IP) in the space  $H_0^1(\Omega; \Gamma_0)$ , where  $\Gamma_0 \subset \partial\Omega$  is a part of  $\partial\Omega$  of positive measure, that is,  $u = 0$  on  $\Gamma_0$  and  $\frac{\partial u}{\partial \nu} = 0$  on  $\partial\Omega \setminus \Gamma_0$ . See, for example [5].

The purely nonlocal counterpart of problem (IP) is problem (NL), it has been studied by several authors, see e.g. [9], [8] and [7] among others. Equations like (NL) appear in several phenomena. For instance,  $u = u(x)$  may represent a density of population (for instance of bacteria) subject to spreading and because we are considering homogeneous Dirichlet boundary condition ( $u \in H_0^1(\Omega)$ ) it means that the domain  $\Omega$  is surrounded by inhospitable environment. Contrary to the local model in which the crowding effect of the population  $u$  at  $x$  depends only on the value of the population in the same point, model (NL) considers the case in which the crowding effect depends on the total population in  $\Omega$ . In the present model (IP), the crowding effect depends also on the value of the population in neighbourhoods of  $x$ . According to [6], see also [1], such a model seems to be more realistic.

In the present paper, we use mainly Galerkin's method in order to approach problem (IP). For this, our approach relies on a variant of the Brouwer Fixed

Point Theorem which will be quoted below. Its proof may be found in Lions [12, p. 53].

PROPOSITION 1.2. *Suppose that  $F: \mathbb{R}^m \rightarrow \mathbb{R}^m$  is a continuous function such that  $(F(\xi), \xi) \geq 0$  on  $|\xi| = r$ , where  $(\cdot, \cdot)$  is the usual inner product in  $\mathbb{R}^m$  and  $|\cdot|$  its corresponding norm. Then there exists  $\xi_0 \in \overline{B_r(0)}$  such that  $F(\xi_0) = 0$ .*

This paper is organized as follows. In Section 2, we consider the existence of solution for a class of pseudo-linear problems, while in Section 3 we prove the existence of solution for a large class of nonlinearities involving a convective term.

### 2. A pseudo-linear problem

In order to illustrate the method, we first study a simpler case, namely, the pseudo-linear version of problem (IP). More precisely, for each  $f \in H^{-1}(\Omega)$ , we search weak solutions of the problem

$$(PL) \quad -\operatorname{div} \left( a \left( \int_{\Omega(x,r)} u(y) dy \right) \nabla u \right) = f(x) \quad \text{in } \Omega, \quad u \in H_0^1(\Omega).$$

Here  $H_0^1(\Omega)$  is understood as the closure of  $\mathcal{D}(\Omega)$  in  $H^1(\Omega)$  and is supposed to be equipped with the Dirichlet norm  $\|u\| = (\int_{\Omega} |\nabla u|^2)^{1/2}$ .  $H^{-1}(\Omega)$  denotes the dual space of  $H_0^1(\Omega)$  and  $\langle \cdot, \cdot \rangle$  will denote the duality bracket between these spaces.

We will suppose

(H<sub>1</sub>)  $a$  is continuous and there exists  $\lambda > 0$  such that  $a(s) \geq \lambda > 0$  for all  $s \in \mathbb{R}$ .

Moreover, we will say that  $\Omega$  is regular, if there is  $\tau > 0$  such that

$$(2.1) \quad \operatorname{meas}(\Omega(x,r)) \geq \tau = \tau(r) > 0, \quad \text{for all } x \in \overline{\Omega}.$$

Note that this is the case for a smooth domain.

Our main result in this section is the following:

THEOREM 2.1. *If  $a$  satisfies (H<sub>1</sub>) and if*

- (a)  $a$  is bounded, or
- (b)  $\Omega$  is regular,

*then for each  $f \in H^{-1}(\Omega)$ , problem (PL) possesses a weak solution  $u \in H_0^1(\Omega)$ .*

PROOF. Since the operator

$$Lu = -\operatorname{div} \left( a \left( \int_{\Omega(x,r)} u(y) dy \right) \nabla u \right)$$

has no variational structure, we will attack problem (PL) by using a Galerkin method. For that, let  $\mathbb{B} = \{e_1, e_2, \dots\}$  be a Hilbertian basis of  $H_0^1(\Omega)$  satisfying

$((e_i, e_j)) = \delta_{ij}$ , where  $((\cdot, \cdot))$  is the usual inner product in  $H_0^1(\Omega)$  and  $\delta_{ij}$  is the Kroenecker symbol. Setting  $\mathbb{V}_m := [e_1, \dots, e_m]$ , the span of the set  $\{e_1, \dots, e_m\}$ , for each  $u \in \mathbb{V}_m$  there is  $\xi = (\xi_1, \dots, \xi_m) \in \mathbb{R}^m$  such that  $u = \sum_{j=1}^m \xi_j e_j$ . Thus  $\|u\| = |\xi|$ , where

$$\|u\| = \left( \int_{\Omega} |\nabla u|^2 \right)^{1/2} \quad \text{and} \quad |\xi| = \left( \sum_{j=1}^m \xi_j^2 \right)^{1/2}.$$

Consequently,  $\mathbb{V}_m$  and  $\mathbb{R}^m$  are isometrically isomorphic finite dimensional vector spaces. Unless stated explicitly otherwise, we identify  $u \leftrightarrow \xi$ ,  $u \in \mathbb{V}_m$ ,  $\xi \in \mathbb{R}^m$ .

Let  $F: \mathbb{R}^m \rightarrow \mathbb{R}^m$ ,  $F = (F_1, \dots, F_m)$  be given by

$$F_i(\xi) = \int_{\Omega} a \left( \int_{\Omega(x,r)} u(y) dy \right) \nabla u \cdot \nabla e_i - \langle f, e_i \rangle, \quad i = 1, \dots, m,$$

so that

$$F_i(\xi)\xi_i = \int_{\Omega} a \left( \int_{\Omega(x,r)} u(y) dy \right) \nabla u \cdot \nabla (\xi_i e_i) - \langle f, (\xi_i e_i) \rangle, \quad i = 1, \dots, m.$$

Consequently,

$$((F(\xi), \xi)) = \int_{\Omega} a \left( \int_{\Omega(x,r)} u(y) dy \right) |\nabla u|^2 - \langle f, u \rangle, \quad \text{for all } u \in \mathbb{V}_m.$$

In view of assumption  $(H_1)$ ,  $((F(\xi), \xi)) \geq \lambda \|u\|^2 - \|f\|^* \|u\|$ , for all  $u$  in  $\mathbb{V}_m$ , where  $\|f\|^*$  denotes the strong dual norm of  $f$ . Then  $((F(\xi), \xi)) > 0$ , if  $\|u\| > \|f\|^* / \lambda$ . Therefore, there is  $u_m \in \mathbb{V}_m$  with  $\|u_m\| \leq \|f\|^* / \lambda$  such that  $F(u_m) = 0$ , i.e.

$$0 = F_i(u_m) = \int_{\Omega} a \left( \int_{\Omega(x,r)} u_m(y) dy \right) \nabla u_m \nabla e_i - \int_{\Omega} f e_i, \quad \text{for all } i = 1, \dots, m.$$

Hence,

$$(2.2) \quad \int_{\Omega} a \left( \int_{\Omega(x,r)} u_m(y) dy \right) \nabla u_m \nabla \varphi = \int_{\Omega} f \varphi, \quad \text{for all } \varphi \in \mathbb{V}_k, \quad k \leq m.$$

In what follows we fix  $k$ . From the boundedness of the real sequence  $(\|u_m\|)$ , it follows that there is a subsequence of  $(u_m)$ , still labelled by  $m$ , such that  $u_m \rightharpoonup u$  in  $H_0^1(\Omega)$  and  $u_m \rightarrow u$  in  $L^2(\Omega)$ . As  $u_m \rightarrow u$  also in  $L^1(\Omega)$  and  $\Omega$  is bounded, we have

$$\left| \int_{\Omega(x,r)} u_m dy - \int_{\Omega(x,r)} u dy \right| \leq \int_{\Omega(x,r)} |u_m - u| dy \leq \int_{\Omega} |u_m - u| dy \rightarrow 0,$$

uniformly for  $x \in \Omega$ . In view of continuity of  $a$  it follows that

$$a \left( \int_{\Omega(x,r)} u_m dy \right) \rightarrow a \left( \int_{\Omega(x,r)} u dy \right), \quad \text{for each } x \in \Omega.$$

It is easy to see that in both cases (a) or (b),  $a(\int_{\Omega(x,r)} u_m dy)$  is bounded independently of  $m$ . Thus by the Lebesgue theorem,

$$a\left(\int_{\Omega(x,r)} u_m(y) dy\right) \nabla \varphi \rightarrow a\left(\int_{\Omega(x,r)} u(y) dy\right) \nabla \varphi \quad \text{in } L^2(\Omega).$$

As  $\nabla u_m \rightharpoonup \nabla u$  in  $L^2(\Omega)$ , taking the limit as  $m \rightarrow +\infty$  in (2.2), we get

$$\int_{\Omega} a\left(\int_{\Omega(x,r)} u(y) dy\right) \nabla u \nabla \varphi = \int_{\Omega} f \varphi, \quad \text{for all } \varphi \in \mathbb{V}_k.$$

Since  $k$  is arbitrary, we obtain

$$\int_{\Omega} a\left(\int_{\Omega(x,r)} u(y) dy\right) \nabla u \nabla \varphi = \int_{\Omega} f \varphi, \quad \text{for all } \varphi \in H_0^1(\Omega),$$

showing that  $u$  is a weak solution of problem (PL). □

Here, we would like to point out that one could use also the Schauder Fixed Point Theorem in the spirit of [5] in order to get the existence result above. However, as we have said before, the technique we developed here will be useful in the second part of the paper.

### 3. A sublinear singular problem with a convective term

In this section, our main goal is to study a problem involving sublinear, singular and convective terms. More precisely, we will be concerned with the existence of positive solutions to the problem

$$(3.1) \quad \begin{cases} -\operatorname{div}\left(a\left(\int_{\Omega(x,r)} u(y) dy\right) \nabla u\right) = H(x)u^\alpha + \frac{K(x)}{u^\gamma} + L(x)|\nabla u|^\theta & \text{in } \Omega, \\ u \in H_0^1(\Omega), \end{cases}$$

where  $H(x), K(x), L(x) \geq 0$ , for all  $x \in \Omega$ , are given functions whose properties will be timely introduced and  $\alpha, \gamma$  and  $\theta$  are positive numbers suitably chosen.

REMARK 3.1. We should remark that it would be more natural, before studying problem (3.1), to attack problems like

$$\begin{cases} -\operatorname{div}\left(a\left(\int_{\Omega(x,r)} u(y) dy\right) \nabla u\right) = a(x)u^\alpha + b(x)u^\beta & \text{in } \Omega, \\ u > 0 & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

where  $a$  and  $b$  are given functions and  $\alpha, \beta > 0$  are real numbers. Note that if  $0 < \alpha < 1$  and  $b \equiv 0$  we have a typical sublinear problem. If  $a \equiv 0$  and  $1 < \beta \leq 2^*$  we are in the presence of a superlinear problem. If both  $a, b$  are not simultaneously vanishing and  $0 < \alpha < 1 < \beta \leq 2^*$  we have a concave-convex problem which was studied, for example, by Ambrosetti, Brezis and Cerami [2]. Due to some technical difficulties we were not able yet to deal with it.

In order to approach problem (3.1), let us begin by considering the auxiliary problem

$$(3.2) \quad \begin{cases} -\operatorname{div} \left( a \left( \int_{\Omega(x,r)} u(y) dy \right) \nabla u \right) \\ \qquad \qquad \qquad = H(x)(u^+)^{\alpha} + \frac{K(x)}{(|u| + \varepsilon)^{\gamma}} + L(x)|\nabla u|^{\theta} \quad \text{in } \Omega, \\ u \in H_0^1(\Omega), \end{cases}$$

where  $0 < \varepsilon < 1$  is a fixed number. We will pose the following assumptions:

$$(H_2) \quad 0 < \alpha, \gamma < 1,$$

$$(H_3) \quad H, K, L \in L^{\infty}(\Omega) \text{ and, for } h_0 > 0, H(x), K(x), L(x) \geq h_0 \text{ for almost every } x \in \Omega,$$

$$(H_4) \quad 0 < \theta < 1.$$

**THEOREM 3.2.** *Under the assumptions of Theorem 2.1 and (H<sub>1</sub>)–(H<sub>4</sub>), problem (3.2) possesses a positive solution.*

**PROOF.** As in the previous section, we introduce functions  $F_i(\xi)$ , given now by

$$F_i(\xi) = \int_{\Omega} a \left( \int_{\Omega(x,r)} u(y) dy \right) \nabla u \nabla e_i \\ - \int_{\Omega} H(u^+)^{\alpha} e_i - \int_{\Omega} \frac{K}{(|u| + \varepsilon)^{\gamma}} e_i - \int_{\Omega} L |\nabla u|^{\theta} e_i,$$

for all  $i = 1, \dots, m$ . Hence

$$((F(\xi), \xi)) = \int_{\Omega} a \left( \int_{\Omega(x,r)} u(y) dy \right) |\nabla u|^2 \\ - \int_{\Omega} H(u^+)^{\alpha} u - \int_{\Omega} K \frac{u}{(|u| + \varepsilon)^{\gamma}} - \int_{\Omega} L |\nabla u|^{\theta} u.$$

We recall that, as before, we are identifying  $u \in \mathbb{V}_m$  with  $\xi \in \mathbb{R}^m$ . As  $a(s) \geq \lambda > 0$  for all  $s \in \mathbb{R}$ , we have

$$\int_{\Omega} a \left( \int_{\Omega(x,r)} u(y) dy \right) |\nabla u|^2 \geq \lambda \int_{\Omega} |\nabla u|^2.$$

On the other hand, by the Sobolev continuous embedding and Poincaré inequality

$$\int_{\Omega} H(u^+)^{\alpha} u \leq C \|H\|_{\infty} (\|\nabla u\|^2)^{(\alpha+1)/2} = C \|H\|_{\infty} \|u\|^{\alpha+1}$$

and

$$\int_{\Omega} K \frac{u}{(|u| + \varepsilon)^{\gamma}} \leq \int_{\Omega} K |u|^{1-\gamma} \leq C \|K\|_{\infty} \|u\|^{1-\gamma},$$

for some positive constant  $C$ , which is independent of  $\varepsilon$ . Here, we point out that, at this stage,  $0 < \varepsilon < 1$  is fixed.

In view of  $(H_4)$ , one has  $0 < \theta < 1 < (N + 2)/N \leq 2$  if  $N \geq 2$ , that is, in particular  $\theta < 2$ . Thus

$$\left| \int_{\Omega} L |\nabla u|^{\theta} u \right| \leq \|L\|_{\infty} \left[ \int_{\Omega} (|\nabla u|^{\theta})^{2/\theta} \right]^{\theta/2} \left( \int_{\Omega} |u|^{2/(2-\theta)} \right)^{(2-\theta)/2}.$$

Since  $0 < \theta < (N + 2)/N$ ,  $N \geq 2$ , we also have  $2/(2 - \theta) < 2^* = 2N/(N - 2)$  and so  $H_0^1(\Omega) \hookrightarrow L^{2/(2-\theta)}(\Omega)$ . So,

$$\left| \int_{\Omega} L |\nabla u|^{\theta} u \right| \leq \|L\|_{\infty} \|u\|^{\theta} \|u\|_{2/(2-\theta)} \leq C \|u\|^{\theta+1}.$$

The last inequalities imply that

$$((F(\xi), \xi)) \geq \lambda \|u\|^2 - C \|H\|_{\infty} \|u\|^{\alpha+1} - C \|K\|_{\infty} \|u\|^{1-\gamma} - C \|u\|^{\theta+1}.$$

In view of assumptions  $(H_2)$ – $(H_4)$ , we may find a real constant  $R > 0$  such that  $((F(\xi), \xi)) > 0$  if  $\|u\| = |\xi| = R$ . Here it is important to observe that  $R$  does not depend on  $m$  or  $\varepsilon$ . By the Brouwer Fixed Point Theorem, there is  $u_{\varepsilon,m} \in \mathbb{V}_m$  such that  $F(u_{\varepsilon,m}) = 0$ ,  $\|u_{\varepsilon,m}\| \leq R$ ,  $m = 1, 2, \dots$ , that is, for all  $\varphi \in \mathbb{V}_m$ ,

$$\begin{aligned} \int_{\Omega} a \left( \int_{\Omega} u_{\varepsilon,m}(y) dy \right) \nabla u_m \nabla \varphi \\ = \int_{\Omega} H(u_{\varepsilon,m}^+)^{\alpha} \varphi + \int_{\Omega} \frac{K}{(|u_{\varepsilon,m}| + \varepsilon)^{\gamma}} \varphi + \int_{\Omega} L |\nabla u_{\varepsilon,m}|^{\theta} \varphi. \end{aligned}$$

Hereafter, we will denote by  $u_m$  the function  $u_{\varepsilon,m}$ . Since  $\|u_m\| \leq R$  for all  $m \in \mathbb{N}$ , there is  $u_{\varepsilon} \in H_0^1(\Omega)$  such that, perhaps for some subsequence,

$$\begin{aligned} u_m &\rightharpoonup u_{\varepsilon} && \text{in } H_0^1(\Omega), \\ u_m &\rightarrow u_{\varepsilon} && \text{in } L^q(\Omega), \quad 1 \leq q < 2^*, \\ u_m(x) &\rightarrow u_{\varepsilon}(x) && \text{a.e. in } \Omega. \end{aligned}$$

(We have a conflict of notation between  $u_m$  and  $u_{\varepsilon}$  but it should be no trouble). We now fix  $1 \leq k < m$  and  $\varphi \in \mathbb{V}_k$ . As in the previous section,

$$\int_{\Omega} a \left( \int_{\Omega((x,r))} u_m(y) dy \right) \nabla u_m \nabla \varphi \rightarrow \int_{\Omega} a \left( \int_{\Omega((x,r))} u_{\varepsilon}(y) dy \right) \nabla u_{\varepsilon} \nabla \varphi,$$

for all  $\varphi \in \mathbb{V}_k$ . At the expense of extracting a subsequence we can assume that  $u_m \rightarrow u_{\varepsilon}$  in  $L^q(\Omega)$  and  $|u_m| \leq h$  almost everywhere for some  $h \in L^q(\Omega)$ . Since for  $q > 2$ ,  $h^{\alpha} \varphi \in L^1(\Omega)$ , by the Lebesgue Dominated Convergence Theorem, for each  $\varphi \in \mathbb{V}_k$  we have

$$\int_{\Omega} H(u_m^+)^{\alpha} \varphi \rightarrow \int_{\Omega} H(u_{\varepsilon}^+)^{\alpha} \varphi \quad \text{and} \quad \int_{\Omega} \frac{K}{(|u_m| + \varepsilon)^{\gamma}} \varphi \rightarrow \int_{\Omega} \frac{K}{(|u_{\varepsilon}| + \varepsilon)^{\gamma}} \varphi.$$

Our next step is to pass to the limit in the gradient term. Since  $(u_m)$  is bounded in  $H_0^1(\Omega)$ , it is easy to prove that  $(|\nabla u_m|^{\theta})$  is bounded in  $L^{2/\theta}(\Omega)$ .

Then, there is  $g \in L^{2/\theta}(\Omega)$  such that

$$(3.3) \quad L|\nabla u_m|^\theta \rightharpoonup g \quad \text{in } L^{2/\theta}(\Omega),$$

or, equivalently,

$$\int_{\Omega} L|\nabla u_m|^\theta \varphi \rightarrow \int_{\Omega} g \varphi, \quad \text{for all } \varphi \in L^{(2/\theta)'}(\Omega),$$

where  $(2/\theta)' = 2/(2-\theta)$  is the conjugate exponent of  $2/\theta$ . Furthermore,

$$\int_{\Omega} L|\nabla u_m|^\theta u_m = \int_{\Omega} L|\nabla u_m|^\theta u_\varepsilon + \int_{\Omega} L|\nabla u_m|^\theta (u_m - u_\varepsilon).$$

In view of  $u_m \rightarrow u_\varepsilon$  in  $L^{2/(2-\theta)}(\Omega)$  (note that  $2/(2-\theta) < 2 < 2^*$ ), we obtain

$$\begin{aligned} & \left| \int_{\Omega} L|\nabla u_m|^\theta (u_m - u_\varepsilon) \right| \\ & \leq \|L\|_\infty \left( \int_{\Omega} (|\nabla u_m|^\theta)^{2/\theta} \right)^{\theta/2} \left( \int_{\Omega} |u_m - u_\varepsilon|^{2/(2-\theta)} \right)^{(2-\theta)/2} \\ & \leq C \|u_m - u_\varepsilon\|_{L^{2/(2-\theta)}} \rightarrow 0. \end{aligned}$$

Consequently,

$$\int_{\Omega} L|\nabla u_m|^\theta u_m \rightarrow \int_{\Omega} g u_\varepsilon.$$

Fixing  $e_j$ , we obtain, for  $1 \leq j \leq k$ ,

$$\begin{aligned} & \int_{\Omega} a \left( \int_{\Omega(x,r)} u_m(y) dy \right) \nabla u_m \nabla e_j \\ & = \int_{\Omega} H(u_m^+)^\alpha e_j + \int_{\Omega} \frac{K}{(|u_m| + \varepsilon)^\gamma} e_j + \int_{\Omega} L|\nabla u_m|^\theta e_j. \end{aligned}$$

Taking limits as  $m \rightarrow +\infty$ , we get

$$\int_{\Omega} a \left( \int_{\Omega(x,r)} u_\varepsilon(y) dy \right) \nabla u_\varepsilon \nabla e_j = \int_{\Omega} H(u_\varepsilon^+)^\alpha e_j + \int_{\Omega} \frac{K}{(|u_\varepsilon| + \varepsilon)^\gamma} e_j + \int_{\Omega} L g e_j.$$

Since  $k$  is arbitrary, the last equality becomes

$$\int_{\Omega} a \left( \int_{\Omega(x,r)} u_\varepsilon(y) dy \right) \nabla u_\varepsilon \nabla \varphi = \int_{\Omega} H(u_\varepsilon^+)^\alpha \varphi + \int_{\Omega} \frac{K}{(|u_\varepsilon| + \varepsilon)^\gamma} \varphi + \int_{\Omega} L g \varphi$$

for all  $\varphi \in H_0^1(\Omega)$ . Hence,  $u_\varepsilon$  is a weak solution of the problem

$$-\operatorname{div} \left( a \left( \int_{\Omega(x,r)} u_\varepsilon(y) dy \right) \nabla u_\varepsilon \right) = H(u_\varepsilon^+)^\alpha + \frac{K}{(|u_\varepsilon| + \varepsilon)^\gamma} + Lg$$

in  $\Omega$ ,  $u_\varepsilon \in H_0^1(\Omega)$ . Since  $a, H, K$  and  $g$  are nonnegative functions, the maximum principle (see [11, Theorem 8.1, p. 179] or [10, Theorem 1.14, p. 47]) ensures that  $u_\varepsilon \geq 0$ , and so,  $u_\varepsilon$  is a solution to

$$-\operatorname{div} \left( a \left( \int_{\Omega(x,r)} u_\varepsilon(y) dy \right) \nabla u_\varepsilon \right) = H(x) u_\varepsilon^\alpha + \frac{K(x)}{(u_\varepsilon + \varepsilon)^\gamma} + L(x)g$$

in  $\Omega$ ,  $u_\varepsilon \in H_0^1(\Omega)$ . Therefore,

$$(3.4) \quad \int_{\Omega} a\left(\int_{\Omega(x,r)} u_\varepsilon(y) dy\right) |\nabla u_\varepsilon|^2 = \int_{\Omega} H u_\varepsilon^{\alpha+1} + \int_{\Omega} \frac{K}{(u_\varepsilon + \varepsilon)^\gamma} u_\varepsilon + \int_{\Omega} L g u_\varepsilon.$$

On the other hand, we know that

$$\begin{aligned} \int_{\Omega} a\left(\int_{\Omega(x,r)} u_m(y) dy\right) |\nabla u_m|^2 \\ = \int_{\Omega} H(u_m^+)^{\alpha+1} + \int_{\Omega} \frac{K}{(|u_m| + \varepsilon)^\gamma} u_m + \int_{\Omega} L |\nabla u_m|^\theta u_m. \end{aligned}$$

Hence

$$(3.5) \quad \int_{\Omega} a\left(\int_{\Omega(x,r)} u_m(y) dy\right) |\nabla u_m|^2 \rightarrow \int_{\Omega} H u_\varepsilon^{\alpha+1} + \int_{\Omega} \frac{K}{(u_\varepsilon + \varepsilon)^\gamma} + \int_{\Omega} L g u_\varepsilon.$$

From (3.4) and (3.5),

$$\int_{\Omega} a\left(\int_{\Omega(x,r)} u_m(y) dy\right) |\nabla u_m|^2 \rightarrow \int_{\Omega} a\left(\int_{\Omega(x,r)} u_\varepsilon\right) |\nabla u_\varepsilon|^2.$$

Arguing as in Section 1, one has that  $a(\int_{\Omega(x,r)} u_m dy)$  is bounded independently of  $m$  and

$$(3.6) \quad a\left(\int_{\Omega(x,r)} u_m(y) dy\right) \rightarrow a\left(\int_{\Omega(x,r)} u_\varepsilon(y) dy\right), \quad \text{for all } x \in \Omega.$$

Hence

$$\begin{aligned} \int_{\Omega} |\nabla(u_m - u_\varepsilon)|^2 &\leq \frac{1}{\lambda} \int_{\Omega} a\left(\int_{\Omega(x,r)} u_m(y) dy\right) |\nabla u_m - u_\varepsilon|^2 \\ &= \frac{1}{\lambda} \int_{\Omega} a\left(\int_{\Omega(x,r)} u_m(y) dy\right) \{|\nabla u_m|^2 - 2\nabla u_m \cdot \nabla u_\varepsilon + |\nabla u_\varepsilon|^2\} \rightarrow 0, \end{aligned}$$

i.e.

$$(3.7) \quad u_m \rightarrow u_\varepsilon \quad \text{in } H_0^1(\Omega).$$

The above limit implies that up to a subsequence

$$(3.8) \quad L|\nabla u_m|^\theta \rightharpoonup L|\nabla u_\varepsilon|^\theta \quad \text{in } L^{2/\theta}(\Omega).$$

To see that, note first that from

$$\int_{\Omega} (|\nabla u_m| - |\nabla u_\varepsilon|)^2 \leq \int_{\Omega} |\nabla u_m - \nabla u_\varepsilon|^2$$

one derives that  $|\nabla u_m| \rightarrow |\nabla u_\varepsilon|$  in  $L^2(\Omega)$ . Thus, up to a subsequence one has  $|\nabla u_m| \rightarrow |\nabla u_\varepsilon|$  almost everywhere in  $\Omega$ ,  $|\nabla u_m| \leq h$  for some  $h \in L^2(\Omega)$ . This implies that, for any  $\varphi \in L^{(2/\theta)'}(\Omega)$ ,  $L|\nabla u_m|^\theta \varphi \leq Lh^\theta \varphi$  with  $Lh^\theta \varphi \in L^1(\Omega)$ . Then (3.8) follows from the Lebesgue Dominated Convergence Theorem. Now, we recall that for all  $j = 1, 2, \dots$ ,

$$\int_{\Omega} a\left(\int_{\Omega} u_m(y) dy\right) \nabla u_m \nabla e_j = \int_{\Omega} H(u_m^+)^{\alpha} e_j + \int_{\Omega} \frac{K}{(|u_m| + \varepsilon)^\gamma} e_j + \int_{\Omega} L |\nabla u_m|^\theta e_j.$$

Gathering (3.6), (3.7), (3.8) and taking limits as  $m \rightarrow +\infty$  on both sides of the last equality, we obtain

$$\int_{\Omega} a\left(\int_{\Omega} u_{\varepsilon}(y) dy\right) \nabla u_{\varepsilon} \nabla e_j = \int_{\Omega} H(u_{\varepsilon}^+)^{\alpha} e_j + \int_{\Omega} \frac{K}{(u_{\varepsilon} + \varepsilon)^{\gamma}} e_j + \int_{\Omega} L |\nabla u_{\varepsilon}|^{\theta} e_j.$$

So,  $u_{\varepsilon} \in H_0^1(\Omega)$  is a positive weak solution of auxiliary problem (3.2). □

Now, we are ready to prove the main result of this section

**THEOREM 3.3.** *Under the same assumptions as in Theorem 3.2, problem (3.1) possesses a weak positive solution.*

**PROOF.** First of all we note that that we will use the notation introduced in the previous sections. Thus, we recall that  $\|u_m\| \leq R$  for all  $m = 1, 2, \dots$ , and  $R$  does not depend on  $\varepsilon$ . Hence  $\|u_{\varepsilon}\| \leq \liminf \|u_m\| \leq R$ . Consequently, fixing  $\varepsilon_n = 1/n$  and  $v_n := u_{\varepsilon_n}$ , for some subsequence still denoted by  $n$ , there exists  $v \in H_0^1(\Omega)$  satisfying

$$\begin{aligned} v_n &\rightharpoonup v && \text{in } H_0^1(\Omega), \\ v_n &\rightarrow v && \text{in } L^q(\Omega), \quad 1 \leq q < 2^*, \\ v_n(x) &\rightarrow v(x) && \text{a.e. in } \Omega. \end{aligned}$$

Let us consider the function

$$M(t) = h_0 t^{\alpha} + \frac{h_0}{(t + 1)^{\gamma}}, \quad \text{for } t \geq 0,$$

where  $h_0$  is defined in assumption (H<sub>3</sub>). Thus, there is  $m_0 > 0$  such that  $M(t) \geq m_0 > 0$  for all  $t \geq 0$ . Noticing that

$$H(x)v_n^{\alpha} + \frac{K(x)}{(v_n + \varepsilon_n)^{\gamma}} + L|\nabla v_n|^{\theta} \geq h_0 v_n^{\alpha} + \frac{h_0}{(v_n + \varepsilon_n)^{\gamma}} \geq m_0,$$

for all  $n \in \mathbb{N}$ , we obtain

$$-\operatorname{div}\left(a\left(\int_{\Omega(x,r)} v_n(y) dy\right) \nabla v_n\right) \geq m_0 \quad \text{in } \Omega, \text{ for all } n \in \mathbb{N}.$$

Let  $\omega_n > 0$  be the unique solution of the problem

$$-\operatorname{div}\left(a\left(\int_{\Omega(x,r)} v_n(y) dy\right) \nabla w_n\right) = m_0 \quad \text{in } \Omega, \quad w_n \in H_0^1(\Omega).$$

Note that, for each  $n \in \mathbb{N}$ ,  $a(\int_{\Omega(x,r)} v_n(y) dy)$  is a positive function, which belongs to  $C(\overline{\Omega})$ , this implies positivity of  $w_n$ . Consequently,

$$-\operatorname{div}\left(a\left(\int_{\Omega(x,r)} v_n(y) dy\right) \nabla v_n\right) \geq -\operatorname{div}\left(a\left(\int_{\Omega(x,r)} v_n(y) dy\right) \nabla w_n\right),$$

i.e.

$$\int_{\Omega} a\left(\int_{\Omega(x,r)} v_n(y) dy\right) \nabla v_n \nabla \varphi \geq \int_{\Omega} a\left(\int_{\Omega(x,r)} v_n\right) \nabla w_n \nabla \varphi,$$

for all  $\varphi \in H_0^1(\Omega)$ ,  $\varphi \geq 0$ . This implies, by the aforementioned maximum principle, that

$$(3.9) \quad v_n \geq w_n \quad \text{in } \Omega.$$

Since

$$\int_{\Omega} a\left(\int_{\Omega(x,r)} v_n(y) dy\right) \nabla w_n \nabla \varphi = \int_{\Omega} m_0 \varphi, \quad \text{for all } \varphi \in H_0^1(\Omega),$$

we have  $\lambda \|w_n\|^2 \leq C \|w_n\|$  and so,  $\|w_n\| \leq C$  for all  $n \in \mathbb{N}$ . As before, there is  $w \in H_0^1(\Omega)$  such that  $w_n \rightharpoonup w$  in  $H_0^1(\Omega)$  and

$$-\text{div}\left(a\left(\int_{\Omega(x,r)} v(y) dy\right) \nabla w\right) = m_0 \quad \text{in } \Omega, \quad w \in H_0^1(\Omega).$$

Consequently,  $w > 0$  in  $\Omega$  and, thanks to the elliptic regularity,  $w \in C(\bar{\Omega})$ . In view of (3.9), if  $n \rightarrow \infty$ , we obtain

$$(3.10) \quad v(x) \geq w(x) > 0 \quad \text{a.e. in } \Omega.$$

We now claim that up to a subsequence  $\nabla v_n(x) \rightarrow \nabla v(x)$  almost everywhere in  $\Omega$ . Indeed, given  $\Omega' \Subset \Omega$ , there is  $\phi \in C_0^\infty(\Omega)$  such that  $\phi(x) = 1$  for all  $x \in \Omega'$ .

Repeating the arguments of the proof of the previous theorem and using (3.10) to control the singular term, we deduce also that for some  $g \in L^{2/\theta}(\Omega)$  (see (3.3))

$$\int_{\Omega} a\left(\int_{\Omega} v(y) dy\right) \nabla v \nabla \psi = \int_{\Omega} H v^\alpha \psi + \int_{\Omega} \frac{K}{v^\gamma} \psi + \int_{\Omega} g \psi,$$

for all  $\psi \in H_0^1(\Omega)$  with compact support. Taking  $\psi = v\phi$  leads to

$$\begin{aligned} \int_{\Omega} a\left(\int_{\Omega} v(y) dy\right) |\nabla v|^2 \phi + a\left(\int_{\Omega} v(y) dy\right) \nabla v \nabla \phi v \\ = \int_{\Omega} H v^\alpha v \phi + \int_{\Omega} \frac{K}{v^\gamma} v \phi + \int_{\Omega} g v \phi, \end{aligned}$$

Now, taking  $v_n \phi$  as a test function in the equation satisfied by  $v_n$ , one gets

$$\begin{aligned} \int_{\Omega} a\left(\int_{\Omega} v_n(y) dy\right) |\nabla v_n|^2 \phi + a\left(\int_{\Omega} v_n(y) dy\right) \nabla v_n \nabla \phi v \\ = \int_{\Omega} H v_n^\alpha v_n \phi + \int_{\Omega} \frac{K}{v_n^\gamma} v_n \phi + \int_{\Omega} |\nabla v_n|^\theta v_n \phi. \end{aligned}$$

Taking the limit in  $n$ , we deduce easily arguing as in the proof of Theorem 3.2 that

$$(3.11) \quad \int_{\Omega} a\left(\int_{\Omega} v_n(y) dy\right) |\nabla v_n|^2 \phi \rightarrow \int_{\Omega} a\left(\int_{\Omega} v(y) dy\right) |\nabla v|^2 \phi.$$

We have also

$$(3.12) \quad \int_{\Omega'} |\nabla(v_n - v)|^2 \leq \frac{1}{\lambda} \int_{\Omega} a\left(\int_{\Omega(x,r)} v_n(y) dy\right) |\nabla v_n - v|^2 \phi$$

and

$$\begin{aligned} & \int_{\Omega} a\left(\int_{\Omega(x,r)} v_n(y) dy\right) |\nabla v_n - v|^2 \phi \\ &= \int_{\Omega} a\left(\int_{\Omega(x,r)} v_n(y) dy\right) (|\nabla v_n|^2 - 2\nabla v_n \cdot \nabla v + |\nabla v|^2) \phi. \end{aligned}$$

From (3.11) and (3.12), taking the limit in  $n$ , we deduce  $|\nabla v_n - \nabla v| \rightarrow 0$  in  $L^2(\Omega')$ . Hence, for some subsequence,  $\nabla v_n(x) \rightarrow \nabla v(x)$  almost everywhere in  $\Omega'$ .

Since  $\Omega = \bigcup_{j=1}^{\infty} \Omega_j$  with  $\Omega_j = \{x \in \Omega : d(x, \partial\Omega) \geq 1/j\}$ , the above study implies that  $\nabla v_n(x) \rightarrow \nabla v(x)$  almost everywhere in  $\Omega_j$ , consequently, for some subsequence,  $\nabla v_n(x) \rightarrow \nabla v(x)$  almost everywhere in  $\Omega$ .

Now, gathering this with the boundedness of  $(|\nabla v_n|^\theta)$  in  $L^{2/\theta}(\Omega)$  we can conclude as below (3.8) that the weak limit of  $(|\nabla v_n|^\theta)$  in  $L^{2/\theta}(\Omega)$  is  $|\nabla v|^\theta$ , that is,

$$\int_{\Omega} |\nabla v_n|^\theta \psi \rightarrow \int_{\Omega} |\nabla v|^\theta \psi, \quad \text{for all } \psi \in L^{2/\theta}(\Omega).$$

Using this, we derive easily that  $v$  verifies

$$(3.13) \quad \begin{aligned} & \int_{\Omega} a\left(\int_{\Omega} v(y) dy\right) \nabla v \nabla \psi \\ &= \int_{\Omega} H v^\alpha \psi + \int_{\Omega} \frac{K}{v^\gamma} \psi + \int_{\Omega} L |\nabla v|^\theta \psi, \quad \text{for all } \psi \in C_0^\infty(\Omega). \end{aligned}$$

From the above equality, there is  $C > 0$  such that

$$\left| \int_{\Omega} \frac{K\psi}{v^\gamma} \right| \leq C \|\psi\|, \quad \text{for all } \psi \in C_0^\infty(\Omega).$$

Combining the density of  $C_0^\infty(\Omega)$  in  $H_0^1(\Omega)$  with the last inequality, we derive that

$$\left| \int_{\Omega} \frac{Kw}{v^\gamma} \right| \leq C \|w\|, \quad \text{for all } w \in H_0^1(\Omega).$$

Then, if  $w \in H_0^1(\Omega)$  and  $(\psi_n) \subset C_0^\infty(\Omega)$  verify  $\psi_n \rightarrow w$  in  $H_0^1(\Omega)$ , we can infer that

$$\lim_{n \rightarrow \infty} \int_{\Omega} \frac{K\psi_n}{v^\gamma} = \int_{\Omega} \frac{Kw}{v^\gamma}.$$

The last limit combined with equality (3.13) and the Sobolev embedding gives

$$\begin{aligned} & \int_{\Omega} a\left(\int_{\Omega} v(y) dy\right) \nabla v \nabla \psi \\ &= \int_{\Omega} H v^{\alpha} \psi + \int_{\Omega} \frac{K}{v^{\gamma}} \psi + \int_{\Omega} L |\nabla v|^{\theta} \psi, \quad \text{for all } \psi \in H_0^1(\Omega), \end{aligned}$$

showing that  $v$  is a solution of problem (3.1).  $\square$

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