# THE $K$-OPERATOR AND THE QUALOCATION METHOD FOR STRONGLY ELLIPTIC EQUATIONS ON SMOOTH CURVES 

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#### Abstract

Superconvergence in $L^{2}$-norm and max-norm is considered for the approximation of the equation $L u=f$ where $L$ is a strongly elliptic pseudo-differential operator. Let $u_{h}$ be the qualocation approximation to the solution $u$. The $K$-operator applied to $u_{h}$, by averaging the values of $u_{h}$, achieves a better approximation than $u_{h}$ itself. In this way, we have exploited the highest order of convergence (in negative norm) available for $u_{h}$ to get high order convergence in $L^{2}$ and maximum estimates. The same result is obtained for the approximation of the derivatives of $u$.


1. Introduction. In this paper we shall discuss a way of increasing the order of convergence (in $L^{2}$-norm and in max-norm) for the qualocation method, when used to approximate the solution of the integral equation

$$
\begin{equation*}
L u=f \tag{1.1}
\end{equation*}
$$

in which the operator $L$ is a pseudo-differential operator of any order on a smooth closed curve $\Gamma$ in $\mathbf{R}^{2}$. A common example of such operators is the integral operator with logarithmic kernel which occurs when a boundary-value problem for the Laplacian on a two-dimensional domain is reformulated as an integral equation on the boundary (see e.g. $[\mathbf{9}, \mathbf{1 0}, \mathbf{1 7}]$ ).

The qualocation method (see $[\mathbf{8}, \mathbf{1 4}-\mathbf{1 8}]$ ), which can be explained in short terms as a quadrature-based modification of the collocation method with unusual quadrature rules, aims to increase the order of convergence given by the collocation method while reducing the difficulty in implementation of the Galerkin method. Formally, the qualocation method is obtained by replacing the 'outer' integral in the approximate equation arising from the Galerkin method by a wellchosen quadrature rule. In some particular cases, it even gives higher

[^0][^1]order convergence than the Galerkin method itself. To illustrate, consider for example the case where $L$ is the logarithmic-kernel operator on a smooth curve $\Gamma$ in the plane and where the trial and test spaces are spaces of piecewise constant functions on a uniform mesh. The Galerkin and the collocation methods yield an $O\left(h^{3}\right)$ order of convergence in suitable negative norms (see e.g. $[\mathbf{1}, \mathbf{2}, \mathbf{1 3}, \mathbf{2 1}]$ ). Yet, it is shown in $[\mathbf{8}]$ that the quadrature rule for the qualocation method can be chosen so that the qualocation method yields an order $O\left(h^{5}\right)$ (in suitable negative norm). More precisely, a Simpson-type rule that achieves order $O\left(h^{5}\right)$ has just two points per interval, one at the break-point where the weight is $3 / 7$, and the other at the midpoint where the weight is $4 / 7$. For a systematic review of the qualocation method, see $[\mathbf{1 6}, \mathbf{1 7}]$.

The aim of this paper is to improve in an $L^{2}$ or pointwise sense the order of convergence of the approximation given by the qualocation method. More precisely, we will exploit the highest order convergence in negative norm of the qualocation method to obtain a higher order of convergence in the $L^{2}$-norm and the max-norm. Instead of using the qualocation approximation $u_{h}$ itself as the approximation to $u$, we shall consider $K_{h} * u_{h}$, where $K_{h}$ is a fixed function to be defined and * denotes convolution.

The function $K_{h}$ appeared in 1974 in the PDE literature in [4], and its theory was worked out in detail in [6]. It is defined as a linear combination of B-splines such that its support is small and that it reproduces certain polynomials under convolution. For some elliptic boundary value problems, Bramble and Schatz [6] approximate the solution $u$ by $K_{h} * u_{h}$, where $u_{h}$ is given by the Galerkin method, and get a local error of order $O\left(h^{2 r-2}\right)$ for both the $L^{2}$-norm and max-norm, compared to $O\left(h^{r}\right)$ for the Galerkin method itself.

An alternative construction of the function $K_{h}$ and hence an alternative proof was given by Thomée [19]. That author considered the error estimates not only for the approximate solution but also for the derivatives.

We will follow Bramble and Schatz in constructing the function $K_{h}$, and will prove error estimates in the $L^{2}$-norm and the max-norm for the solution and its derivatives. The keypoint of the proof is the invariance with respect to translation of a simplified form of the problem, the method and the test space. In the BIE literature, the application
of the $K$-operator to the Galerkin approximation of the logarithmickernel equation on a smooth curve has been discussed in unpublished work of Schatz, Sloan and Wahlbin. (See also [20] for a discussion of this application of the $K$-operator to obtain both the global and local estimates.) It is worth noting that superconvergence in max-norm for the Galerkin approximation to second kind integral equations has been proved by Chandler [7]. That author gave two methods to achieve superconvergence from the Galerkin approximate solutions: one is the natural iteration (the idea of which is due to Sloan, see [7]), the other is 'superinterpolation' (see [7, Section 5]). The latter alternative is an analog to the method of Bramble and Schatz [6] and Thomée [19].

This paper contains 5 sections. Section 2 gives some notations to be used and a brief review of the qualocation method. One can find the definition and properties of the $K$-operator in Section 3. The main result of the paper is in Section 4. Section 5 is devoted to a numerical experiment.
2. Notations and some preliminaries. We will consider in this paper complex valued functions which are periodic with period 1. Each periodic function $u$ has a Fourier expansion

$$
u(x) \sim \sum_{n \in \mathbf{Z}} \hat{u}(n) e^{2 \pi i n x}
$$

where the Fourier coefficients are given by the formula

$$
\hat{u}(n)=\int_{0}^{1} u(x) e^{-2 \pi i n x} d x
$$

provided $u$ is in $L^{1}(0,1)$. For $s \in \mathbf{R}$ we define the norm

$$
\|u\|_{s}^{2}=|\hat{u}(0)|^{2}+\sum_{n \neq 0}|n|^{2 s}|\hat{u}(n)|^{2}
$$

The Sobolev space $H^{s}$ consists of all periodic distributions $u$ for which the norm $\|u\|_{s}$ is finite. When $s=0, H^{0}$ is the usual $L^{2}$ space with norm denoted by $\|\cdot\|$. We will also use the following notations:

$$
\begin{aligned}
|v|_{0} & =\max _{0 \leq x \leq 1}|v(x)| \\
|v|_{s} & =\sum_{j=0}^{s}\left|D^{j} v\right|_{0}
\end{aligned}
$$

Throughout this paper $c$ denotes a generic constant which can take different values at different occurrences.

As in [8] we are concerned with pseudo-differential operators of the form

$$
L=L_{0}+L_{1}
$$

The principal part $L_{0}$ of the operator $L$ is defined by

$$
\begin{equation*}
L_{0} u(x):=\sum_{n \in \mathbf{Z}}[n]_{\beta} \hat{u}(n) e^{2 \pi i n x} \tag{2.1}
\end{equation*}
$$

where $\beta \in \mathbf{R}$ and $[n]_{\beta}$ is defined either by

$$
[n]_{\beta}:= \begin{cases}1 & \text { for } n=0 \\ |n|^{\beta} & \text { for } n \neq 0,\end{cases}
$$

(in which case $L_{0}$ is an even operator of order $\beta$ ) or by

$$
[n]_{\beta}:= \begin{cases}1 & \text { for } n=0 \\ (\operatorname{sign} n)|n|^{\beta} & \text { for } n \neq 0\end{cases}
$$

(in which case $L_{0}$ is an odd operator of order $\beta$ ). In either case $L_{0}$ is a pseudo-differential operator of order $\beta$, and is an isometry from $H^{s}$ to $H^{s-\beta}$ for all $s \in \mathbf{R}$.

In [8], the operator $L_{1}$ is required to be a continuous mapping

$$
L_{1}: H^{s} \longrightarrow H^{t} \quad \forall s, t \in \mathbf{R} .
$$

In fact, if we follow the perturbation argument used in [11] we need assume only that $L_{1}$ is a bounded operator

$$
\begin{equation*}
L_{1}: H^{s} \longrightarrow H^{s-\beta+\eta} \quad \forall s \in \mathbf{R} \tag{2.2}
\end{equation*}
$$

where $\eta$ is some positive number to be specified later. We then have $L_{0}^{-1} L_{1}$ bounded from $H^{s}$ to $H^{s+\eta}$ and compact on $H^{s}$ for all $s \in \mathbf{R}$. We also assume that $L$ is $1-1$, and thus by the Fredholm alternative

$$
\left(I+L_{0}^{-1} L_{1}\right)^{-1}: H^{s} \longrightarrow H^{s}
$$

is bounded for all $s \in \mathbf{R}$. For the convenience of the readers we recall here some main results obtained by Chandler and Sloan $[\mathbf{8}]$.

Let $x_{i}=i h, i \in \mathbf{Z}, h=1 / N$ be a uniform mesh with $N$ subintervals of the interval $[0,1]$. (Note that $x_{i}$ and $x_{i+N}$ denote the same points in this 1-periodic setting). Let $S_{h}$ be the trial space consisting of periodic splines of order $r$ (i.e. of degree $\leq r-1$ ) with knots $\left\{x_{i}\right\}$, which are $r-2$ times continuously differentiable. Similarly, let the test space $S_{h}^{\prime}$ be the set of periodic splines of order $r^{\prime}$ (i.e. of degree $\leq r^{\prime}-1$ ) with knots $\left\{x_{i}\right\}$ and $r^{\prime}-2$ continuous derivatives.

The qualocation method is a discrete Petrov-Galerkin method which approximates the outer integral by a composite quadrature rule determined by points $\left\{\xi_{j}: 1 \leq j \leq J\right\}$, where

$$
\begin{equation*}
0 \leq \xi_{1}<\xi_{2}<\cdots<\xi_{J}<1 \tag{2.3}
\end{equation*}
$$

and weights $\left\{\omega_{j}: 1 \leq j \leq J\right\}$ such that

$$
\omega_{j}>0, \quad \sum_{j=1}^{J} \omega_{j}=1
$$

The qualocation rule is defined by

$$
\begin{equation*}
Q_{N}(g):=h \sum_{i=1}^{N} \sum_{j=1}^{J} \omega_{j} g\left(x_{i}+h \xi_{j}\right) \tag{2.4}
\end{equation*}
$$

With this rule a discrete inner product is defined by

$$
\begin{equation*}
\langle u, v\rangle=Q_{N}(u \bar{v}), \tag{2.5}
\end{equation*}
$$

where $\bar{v}$ denotes the complex conjugate of $v$. The qualocation solution to the equation (1.1) is now defined by

$$
\begin{equation*}
u_{h} \in S_{h} \quad \text { and } \quad\left\langle L u_{h}, \psi^{\prime}\right\rangle=\left\langle f, \psi^{\prime}\right\rangle \quad \forall \psi^{\prime} \in S_{h}^{\prime} \tag{2.6}
\end{equation*}
$$

After choosing bases for $S_{h}$ and $S_{h}^{\prime}$, we deduce from (2.6) a system of $N$ linear equations in $N$ unknowns, which is referred to as the qualocation equation. The qualocation method is well defined if either

$$
r>\beta+1
$$

or

$$
r>\beta+1 / 2 \quad \text { and } \quad \xi_{1}>0 .
$$

The condition $\xi_{1}>0$ in the latter alternative is necessary because of the fact that if

$$
\beta+1 / 2<r \leq \beta+1,
$$

then $L \psi$ for $\psi \in S_{h}$ is not continuous at the knot points, so that in this case the knot points are not allowed as quadrature points. The condition $r>\beta+1 / 2$ ensures the continuity of $L \psi$ at points other than knot points for $\psi \in S_{h}$. (See $[\mathbf{2}, \mathbf{8}]$ for more details.)
Let

$$
D(y):=\sum_{j=1}^{J} w_{j}\left(1+\Omega\left(\xi_{j}, y\right)\right)\left(1+\overline{\Delta^{\prime}\left(\xi_{j}, y\right)}\right), \quad y \in\left[-\frac{1}{2}, \frac{1}{2}\right],
$$

and let

$$
E(y):=\sum_{j=1}^{J} w_{j} \Omega\left(\xi_{j}, y\right)\left(1+\overline{\Delta^{\prime}\left(\xi_{j}, y\right)}\right), \quad y \in\left[-\frac{1}{2}, \frac{1}{2}\right],
$$

where

$$
\Delta^{\prime}(\xi, y)=y^{r^{\prime}} \sum_{l \neq 0} \frac{1}{(l+y)^{r^{\prime}}} e^{2 \pi i l \xi},
$$

and where

$$
\Omega(\xi, y)=|y|^{r-\beta} \sum_{l \neq 0} \frac{1}{|l+y|^{r-\beta}} e^{2 \pi i l \xi}
$$

if $r$ and $L_{0}$ are both even or both odd, or

$$
\Omega(\xi, y)=(\operatorname{sign} y)|y|^{r-\beta} \sum_{l \neq 0} \frac{\operatorname{sign} l}{|l+y|^{r-\beta}} e^{2 \pi i l \xi}
$$

if $r$ and $L_{0}$ are of opposite parity. The qualocation method is stable if

$$
\inf \{|D(y)|: y \in[-1 / 2,1 / 2]\}>0
$$

It is said to be of order $r-\beta+b$ if

$$
E(y)=O\left(|y|^{r-\beta+b}\right), \quad y \in[-1 / 2,1 / 2] .
$$

We have the following theorem (cf. [8]):

Theorem A. Let (1.1) be solved by a well defined qualocation method which is stable and of order $r-\beta+b, b \geq 0$. Let the condition (2.2) hold for some $\eta>b+1 / 2$. Then for all $N$ sufficiently large $u_{h}$ is uniquely defined. Moreover, for all $s, t$ satisfying

$$
\begin{equation*}
s<r-\frac{1}{2}, \quad \beta+\frac{1}{2}<t, \quad \beta-b \leq s \leq t \leq r \tag{2.7}
\end{equation*}
$$

we have

$$
\begin{equation*}
\left\|u_{h}-u\right\|_{s} \leq c h^{t-s}\|u\|_{t+\max \{\beta-s, 0\}} . \tag{2.8}
\end{equation*}
$$

The case $L=L_{0}$ was proved in [8].
Proof for the case $L=L_{0}+L_{1}$. We give here a slightly different argument from that in [8] by using the reasoning used in [11]. Assume for the moment that (2.6) has a solution $u_{h} \in S_{h}$. Since we can write the defining equation as

$$
\left\langle\left(L_{0}+L_{1}\right) u_{h}, \psi^{\prime}\right\rangle=\left\langle\left(L_{0}+L_{1}\right) u, \psi^{\prime}\right\rangle \quad \forall \psi^{\prime} \in S_{h}^{\prime},
$$

or

$$
\begin{equation*}
\left\langle L_{0} u_{h}, \psi^{\prime}\right\rangle=\left\langle L_{0}\left(u+L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right), \psi^{\prime}\right\rangle \quad \forall \psi^{\prime} \in S_{h}^{\prime} \tag{2.9}
\end{equation*}
$$

we have from Theorem 2 in $[\mathbf{8}]$ for the special case $L=L_{0}$

$$
\begin{aligned}
\left\|u_{h}-u-L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right\|_{s} & \leq c h^{t-s}\left\|u+L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right\|_{t_{s}} \\
& \leq c h^{t-s}\left(\|u\|_{t_{s}}+\left\|L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right\|_{t_{s}}\right)
\end{aligned}
$$

where $t_{s}=t+\max \{\beta-s, 0\}$. Using (2.2) we then deduce

$$
\begin{equation*}
\left\|u_{h}-u-L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right\|_{s} \leq c h^{t-s}\left(\|u\|_{t_{s}}+\left\|u-u_{h}\right\|_{t_{s}-\eta}\right) \tag{2.10}
\end{equation*}
$$

On the other hand, since $\left(I+L_{0}^{-1} L_{1}\right)$ is an isomorphism on $H^{s}$ for all $s \in \mathbf{R}$ we have

$$
\begin{equation*}
\left\|u_{h}-u\right\|_{s} \leq c\left\|\left(I+L_{0}^{-1} L_{1}\right)\left(u_{h}-u\right)\right\|_{s} \tag{2.11}
\end{equation*}
$$

Inequalities (2.10) and (2.11) now give

$$
\begin{equation*}
\left\|u_{h}-u\right\|_{s} \leq c h^{t-s}\left(\|u\|_{t_{s}}+\left\|u_{h}-u\right\|_{t_{s}-\eta}\right) \tag{2.12}
\end{equation*}
$$

Note that (2.12) holds for all $s$ and $t$ satisfying (2.7). Also note that $\beta+1 / 2<t_{s} \leq r+b$. Since $\eta>b+1 / 2$ and $r>\beta+1 / 2$, we can choose $\eta^{\prime}$ such that

$$
1 / 2 \leq \eta^{\prime} \leq \eta \quad \text { and } \quad \beta \leq t_{s}-\eta^{\prime}<r-1 / 2
$$

Therefore we can write (2.12) with $s$ replaced by $t_{s}-\eta^{\prime}$ and $t$ by $\bar{t}=\min \left\{r, t_{s}\right\}$ to obtain

$$
\left\|u_{h}-u\right\|_{t_{s}-\eta^{\prime}} \leq c h^{\bar{t}-t_{s}+\eta^{\prime}}\left(\|u\|_{t^{*}}+\left\|u_{h}-u\right\|_{t^{*}-\eta}\right)
$$

where $t^{*}=\bar{t}+\max \left\{\beta-t_{s}+\eta^{\prime}, 0\right\}=\bar{t} \leq t_{s}$. Since $\left\|u_{h}-u\right\|_{t^{*}-\eta} \leq$ $\left\|u_{h}-u\right\|_{t_{s}-\eta^{\prime}}$ and $\bar{t}-t_{s}+\eta^{\prime} \geq 1 / 2$, we have, for sufficiently large $N$,

$$
\begin{equation*}
\left\|u_{h}-u\right\|_{t_{s}-\eta} \leq c h^{1 / 2}\|u\|_{t_{s}} \tag{2.13}
\end{equation*}
$$

Inequalities (2.12) and (2.13) now give the desired estimate (2.8). It remains to establish the existence and uniqueness of the solution $u_{h}$ of (2.6). Assume that there are two solutions $u_{h}^{(1)}$ and $u_{h}^{(2)}$ of (2.6). Then $u_{h}=u_{h}^{(1)}-u_{h}^{(2)}$ is the solution to (2.6) with $f=0$ on the right hand side. Since $L$ is $1-1$, we have the exact solution $u=0$ in that case; therefore we obtain from (2.8) $u_{h}=0$ for large $N$. Uniqueness (for large $N$ ) for equation (2.6) is proved. The existence of $u_{h}$ for large $N$ then follows because (2.6) is a system of $N$ equations in $N$ unknowns.

Results on max-norm estimates have been proved for the case in which the trial space is a space of smoothest splines of odd degree, the test space is space of trigonometric polynomials and $L=L_{0}$ is an even operator (see [15]). Actually the same argument can be used to prove the following theorem:

Theorem B. Let the conditions of Theorem A hold, and let $\delta>0$. If $u \in H^{t}$ with

$$
t \geq r+\max \{\beta, \delta\}+1 / 2
$$

then

$$
\begin{equation*}
\left|u_{h}-u\right|_{0} \leq c h^{\min \{r, r+b-\beta\}}\|u\|_{t} . \tag{2.14}
\end{equation*}
$$

We will in this paper exploit the highest order in negative norm given by the qualocation method to further develop the order of the $L^{2}$ and max-norm estimates. To do so we will approximate $u$ by, instead of $u_{h}$, an average of $u_{h}$ values defined by a convolution operator. If $\beta-b<0$, the order will be $O\left(h^{r-\beta+b}\right)$ compared to $O\left(h^{r}\right)$ given by the qualocation method. The case $\beta-b \geq 0$ is not interesting in our analysis since for both the $L^{2}$ - and max-norms the qualocation method itself gives optimal estimates of order $O\left(h^{r+b-\beta}\right.$ ) (see Theorems A and $B)$. In this case the averaging method gives the same results.

In the following section we will give the definition and some properties of the $K$-operator.
3. The $K$-operator and its properties. The $K$-operator acting on $u_{h}$ is defined by the convolution of $u_{h}$ with a function $K_{h}$ defined as a linear combination of B-splines such that it reproduces polynomials (up to some degree) under convolution. For the application to our problem we will give here its definition in the 1-dimensional case only.
Let

$$
\chi(x)= \begin{cases}1 & \text { if }-1 / 2<x<1 / 2 \\ \frac{1}{2} & \text { if } x=1 / 2 \text { or } x=-1 / 2 \\ 0 & \text { otherwise }\end{cases}
$$

and let

$$
\psi^{(l)}=\chi * \chi * \cdots * \chi, \quad \text { with }(l-1) \text { times of convolution, } \quad l \geq 1
$$

It is well known that $\psi^{(l)}$ is the B-spline of order $l$ symmetric about 0 with support $[-l / 2, l / 2]$. Let $q, l$ be arbitrary but fixed positive integers. We define $K_{q}^{l}$ by

$$
\begin{equation*}
K_{q}^{l}(x)=\sum_{j=-(q-1)}^{q-1} k_{j} \psi^{(l)}(x-j) \tag{3.1}
\end{equation*}
$$

where $k_{j}, j=-(q-1), \ldots, q-1$ are chosen such that

$$
\int_{-\infty}^{\infty} K_{q}^{l}(x) x^{i} d x= \begin{cases}1 & \text { if } i=0  \tag{3.2}\\ 0 & \text { if } i=1, \ldots, 2 q-1\end{cases}
$$

Since $\psi^{(l)}$ is an even function and since we want $K_{q}^{l}$ to have the same property, we impose a symmetry condition on $k_{j}$ :

$$
\begin{equation*}
k_{-j}=k_{j}, \quad j=1, \ldots, q-1 \tag{3.3}
\end{equation*}
$$

Then the condition (3.2) is equivalent to

$$
\int_{-\infty}^{\infty} K_{q}^{l}(x) x^{2 m} d x= \begin{cases}1 & \text { if } m=0  \tag{3.4}\\ 0 & \text { if } m=1, \ldots, q-1\end{cases}
$$

In fact (3.4) can be written as

$$
\sum_{j=0}^{q-1} k_{j}^{\prime} \int_{-\infty}^{\infty} \psi^{(l)}(x)(x+j)^{2 m} d x= \begin{cases}1 & \text { if } m=0  \tag{3.5}\\ 0 & \text { if } m=1, \ldots, q-1\end{cases}
$$

where $k_{0}^{\prime}=k_{0}, k_{j}^{\prime}=2 k_{j}, j=1, \ldots, q-1$. The system (3.5) is a system of $q$ equations with $q$ unknowns $k_{0}^{\prime}, \ldots, k_{q-1}^{\prime}$. It was proved in [ $\mathbf{5}$, Lemma 8.1] that the solutions exist uniquely.
Now for $0<h<1$, we define

$$
\begin{equation*}
K_{h}(x)=K_{h, q}^{l}(x)=\frac{1}{h} K_{q}^{l}\left(\frac{x}{h}\right) \tag{3.6}
\end{equation*}
$$

Then we have supp $K_{h, q}^{l}=[-(q-1+l / 2) h,(q-1+l / 2) h]$ and

$$
\int_{-\infty}^{\infty} K_{h}(x) x^{i} d x= \begin{cases}1 & \text { if } i=0  \tag{3.7}\\ 0 & \text { if } i=1, \ldots, 2 q-1\end{cases}
$$

As an example, we give here the graph of $K_{3}^{4}$ (Figure 1), a cubic spline. The coefficients $k_{j}$ in that case are $k_{0}=181 / 120, k_{1}=k_{-1}=$ $-17 / 60, k_{2}=k_{-2}=7 / 240$.

Representation of $K_{h} * u_{h}$. Let $\psi_{h, p}^{\left(l^{\prime}\right)}$ be 1-periodic functions defined by

$$
\psi_{h, p}^{\left(l^{\prime}\right)}(x)=\psi^{\left(l^{\prime}\right)}\left(x / h-l^{\prime} / 2\right) \quad \text { for } x \in[0,1) ; \quad l^{\prime}=2, \ldots, N
$$



FIGURE 1. Graph of $K_{3}^{4}$.
If $u_{h}$ is a solution to the equation (2.6), then since $u_{h} \in S_{h}$ we can write $u_{h}$ in the form

$$
u_{h}(x)=\sum_{i=0}^{N-1} c_{i} \psi_{h, p}^{(r)}(x-i h) .
$$

Hence $K_{h} * u_{h}$ can be represented as

$$
K_{h} * u_{h}(x)=\sum_{i=0}^{N-1} c_{i} \phi_{h, p}^{(r+l)}(x-(i+r / 2) h),
$$

where $\phi_{h, p}^{\left(l^{\prime}\right)}$ is a 1-periodic, even function defined by

$$
\phi_{h, p}^{\left(l^{\prime}\right)}(x)=\sum_{j=-(q-1)}^{q-1} k_{j} \psi_{h, p}^{\left(l^{\prime}\right)}\left(x-j h+\frac{l^{\prime} h}{2}\right) .
$$

To ensure that $\phi_{h, p}^{(r+l)}$, and hence $K_{h} * u_{h}$, is a periodic spline of order $r+l$, we require $q-1+(r+l) / 2 \leq N / 2$. If the inequality is strict, then


Figure 2. Graph of $\phi$.
the support of $\phi_{h, p}^{(r+l)}$ in $[-1 / 2,1 / 2]$ is $[-(q-1+(r+l) / 2) h,(q-1+$ $(r+l) / 2) h]$. As an example, the graph of $\phi(t)=\phi_{h, p}^{(5)}(t h)$ for the case $l=4, q=3$, and $r=1$ is given in Figure 2. Its support in $[-N / 2, N / 2]$ is $[-9 / 2,9 / 2]$.

Stability discussion. Assume that

$$
\bar{u}_{h}(x)=\sum_{i=0}^{N-1} \bar{c}_{i} \psi_{h, p}^{(r)}(x-i h)
$$

and that

$$
\left|c_{i}-\bar{c}_{i}\right| \leq \varepsilon \quad \text { for } i=0, \ldots, N-1
$$

Then

$$
\left|K_{h} * u_{h}(x)-K_{h} * \bar{u}_{h}(x)\right| \leq \varepsilon \sum_{i=0}^{N-1}\left|\phi_{h, p}^{(r+l)}\left(x-\left(i+\frac{r}{2}\right) h\right)\right|
$$

In the case of the example to be discussed in Section 5, we have $l=4$, $q=3$ and $r=1$. Elementary but lengthy calculation gives us

$$
\begin{aligned}
\max _{0 \leq x \leq 1} \sum_{i=0}^{N-1}\left|\phi_{h, p}^{(5)}\left(x-\left(i+\frac{1}{2}\right) h\right)\right| & =\sum_{i=0}^{N-1}\left|\phi_{h, p}^{(5)}\left(\left(i+\frac{1}{2}\right) h\right)\right| \\
& =1.2146
\end{aligned}
$$

Hence

$$
\left|K_{h} * u_{h}(x)-K_{h} * \bar{u}_{h}(x)\right| \leq 1.2146 \varepsilon,
$$

i.e. the $K$-operator method is quite stable in this case.

We will give here some properties of the $K$-operator. From (3.7) it is easy to see that $K_{h}$ reproduces polynomials of order $\leq 2 q$ (i.e. of degree $\leq 2 q-1$ ) under convolution, i.e.,

$$
K_{h} * v=v \quad \text { if } v \in \mathbf{P}_{2 q} .
$$

The following lemma is a consequence of the above property and the Bramble- Hilbert lemma [3]:

Lemma 3.1 (See [6]).

$$
\begin{array}{ll}
\left\|K_{h} * u-u\right\| \leq c h^{s}\|u\|_{s}, & \text { for } 0 \leq s \leq 2 q, \\
\left|K_{h} * u-u\right|_{0} \leq c h^{s}|u|_{s}, & \text { for } 0 \leq s \leq 2 q . \tag{3.9}
\end{array}
$$

Another interesting property of $K_{h}$ is that the differential operator when applied to $K_{h}$ is changed to the central differential operator applied to a somewhat similar function. More precisely, letting

$$
\begin{aligned}
\partial_{h} v(x) & =\frac{1}{h}\left\{v\left(x+\frac{h}{2}\right)-v\left(x-\frac{h}{2}\right)\right\}, \\
\partial_{h}^{\alpha} v & =\partial_{h}^{\alpha-1}\left(\partial_{h} v\right), \quad \alpha=2,3, \ldots
\end{aligned}
$$

we have the following easily proved lemma:

Lemma 3.2 (See [6]). For any $\alpha=0,1, \ldots$, l, we have

$$
\begin{equation*}
D^{\alpha} K_{h}=\partial_{h}^{\alpha} V_{h, q}^{l-\alpha} \tag{3.10}
\end{equation*}
$$

where

$$
V_{h, q}^{\beta}(x)=\frac{1}{h} \sum_{j=-(q-1)}^{q-1} k_{j} \psi^{(\beta)}\left(\frac{x}{h}-j\right)
$$

Note that in this notation $K_{h}=V_{h, q}^{l}$. From (3.10) we have

## Lemma 3.3.

$$
D^{\alpha}\left(K_{h} * v\right)=V_{h, q}^{l-\alpha} * \partial_{h}^{\alpha} v \quad \text { for } \alpha=1, \ldots, l
$$

Before going to the main results of this paper we need the following lemmas:

Lemma 3.4 (cf. [6]). Let $\tau>0$, and let $\tau^{*}=\lceil\tau\rceil$, the least integer greater than or equal to $\tau$. Then

$$
\|v\| \leq c \sum_{\gamma=0}^{\tau^{*}}\left\|D^{\gamma} v\right\|_{-\tau}
$$

Proof. The result comes directly from the definition of the Sobolev norms.

Let $T_{h}$ be the translation operator defined by

$$
T_{h} v(x)=v(x+h)
$$

Then the following property for the discrete inner product defined by (2.3), (2.4) and (2.5) is easily proved.

Lemma 3.5. For any $u, v$,

$$
\left\langle T_{h} u, v\right\rangle=\left\langle u, T_{-h} v\right\rangle
$$

4. Application to the qualocation method. For the reason given in the comment following Theorem B, we consider only the case $\beta-b<0$.

Theorem 4.1. Assume that the conditions of Theorem A hold. Let $\tau=b-\beta>0$ and let $\tau^{*}$ be defined as in Lemma 3.4. Let $\alpha$, $l$, and $q$ be integers satisfying

$$
\begin{equation*}
\alpha \geq 0, \quad l \geq \tau^{*}+\alpha, \quad 2 q \geq r+\tau \tag{4.1}
\end{equation*}
$$

Assume further that the condition (2.2) holds for some $\eta>b+1 / 2+$ $\tau^{*}+\alpha$. Then

$$
\begin{equation*}
\left\|D^{\alpha} u-D^{\alpha} K_{h} * u_{h}\right\| \leq c h^{r+\tau}\|u\|_{R} \tag{4.2}
\end{equation*}
$$

where $R=r+b+\tau^{*}+\alpha$.

Proof. By the triangle inequality we have

$$
\begin{aligned}
\left\|D^{\alpha} u-D^{\alpha} K_{h} * u_{h}\right\| & \leq\left\|D^{\alpha} u-D^{\alpha} K_{h} * u\right\|+\left\|D^{\alpha} K_{h} *\left(u-u_{h}\right)\right\| \\
& =I+I I
\end{aligned}
$$

We will prove separately that $I$ and $I I$ satisfy (4.2). By the property of the convolution operator and by Lemma 3.1 we have

$$
\begin{equation*}
I=\left\|D^{\alpha} u-K_{h} * D^{\alpha} u\right\| \leq c h^{s}\left\|D^{\alpha} u\right\|_{s} \leq c h^{s}\|u\|_{s+\alpha} \quad \text { for } 0 \leq s \leq 2 q \tag{4.3}
\end{equation*}
$$

To estimate $I I$, we assume first that $L=L_{0}$, i.e. $L_{1}=0$. Then by Lemmas 3.4 and 3.3

$$
\begin{aligned}
I I & \leq c \sum_{\gamma=0}^{\tau^{*}}\left\|\left(D^{\alpha+\gamma} K_{h}\right) *\left(u_{h}-u\right)\right\|_{-\tau} \\
& =c \sum_{\gamma=0}^{\tau^{*}}\left\|V_{h, q}^{l-\alpha-\gamma} * \partial_{h}^{\alpha+\gamma}\left(u_{h}-u\right)\right\|_{-\tau} .
\end{aligned}
$$

Hence, from the definitions of the convolution operator and Sobolev norms, we have

$$
\begin{equation*}
I I \leq c \sum_{\gamma=0}^{\tau^{*}}\left\|\partial_{h}^{\alpha+\gamma}\left(u_{h}-u\right)\right\|_{-\tau}=c \sum_{\gamma=0}^{\tau *}\left\|\tilde{\partial}_{h}^{\alpha+\gamma}\left(u_{h}-u\right)\right\|_{-\tau} \tag{4.4}
\end{equation*}
$$

where $\tilde{\partial}_{h}$ is the forward difference operator defined by

$$
\tilde{\partial}_{h} v(x)=\frac{1}{h}\{v(x+h)-v(x)\}=T_{h / 2} \partial_{h} v(x)
$$

and

$$
\tilde{\partial}_{h}^{j} v=\tilde{\partial}_{h}^{j-1}\left(\tilde{\partial}_{h} v\right), \quad j=2,3, \ldots
$$

We will prove that for any $j \in \mathbf{N}, \tilde{\partial}_{h}^{j} u_{h}$ is the qualocation approximant to $\tilde{\partial}_{h}^{j} u$, i.e., $\tilde{\partial}_{h}^{j} u_{h} \in S_{h}$ and

$$
\begin{equation*}
\left\langle L_{0} \tilde{\partial}_{h}^{j}\left(u_{h}-u\right), \psi^{\prime}\right\rangle=0 \quad \text { for } \psi^{\prime} \in S_{h}^{\prime} \tag{4.5}
\end{equation*}
$$

The proof is carried out for $j=1$; the general case is then obtained by induction. That $\tilde{\partial}_{h} u_{h}$ belongs to $S_{h}$ follows from the definition of $\tilde{\partial}_{h}$ and the fact that the space $S_{h}$ is invariant under translation by $h$. By the definition of the forward difference operator and the fact that $u_{h}$ satisfies (2.6) with $L=L_{0}$, we have

$$
\begin{aligned}
\left\langle L_{0} \tilde{\partial}_{h}\left(u_{h}-u\right), \psi^{\prime}\right\rangle & =\frac{1}{h}\left\{\left\langle L_{0} T_{h}\left(u_{h}-u\right), \psi^{\prime}\right\rangle-\left\langle L_{0}\left(u_{h}-u\right), \psi^{\prime}\right\rangle\right\} \\
& =\frac{1}{h}\left\langle L_{0} T_{h}\left(u_{h}-u\right), \psi^{\prime}\right\rangle \quad \text { for any } \psi^{\prime} \in S_{h}^{\prime}
\end{aligned}
$$

Since $L_{0}$ commutes with $T_{h}$ (which can be proved directly from the definition (2.1) of $L_{0}$ or by using the fact that $L_{0}$ is a multiplier operator, see e.g. [12]) we obtain by using Lemma 3.5

$$
\left\langle L_{0} \tilde{\partial}_{h}\left(u_{h}-u\right), \psi^{\prime}\right\rangle=\frac{1}{h}\left\langle L_{0}\left(u_{h}-u\right), T_{-h} \psi^{\prime}\right\rangle \quad \text { for any } \psi^{\prime} \in S_{h}^{\prime}
$$

Since $S_{h}^{\prime}$ is invariant under translation by $h$, we conclude that

$$
\left\langle L_{0} \tilde{\partial}_{h}\left(u_{h}-u\right), \psi^{\prime}\right\rangle=0 \quad \text { for any } \psi^{\prime} \in S_{h}^{\prime}
$$

Hence (4.5) is proved.
We can now use the estimate $(2.8)$ for $\tilde{\partial}_{h}^{\alpha+\gamma}\left(u_{h}-u\right)$ to obtain

$$
\begin{equation*}
\left\|\tilde{\partial}_{h}^{\alpha+\gamma}\left(u_{h}-u\right)\right\|_{-\tau} \leq c h^{r+\tau}\left\|\tilde{\partial}_{h}^{\alpha+\gamma} u\right\|_{r+b} \leq c h^{r+\tau}\|u\|_{r+b+\alpha+\gamma} . \tag{4.6}
\end{equation*}
$$

Now (4.4) and (4.6) give the required estimate for $I I$ and hence the theorem is proved in case $L=L_{0}$. For the general case, a familiar argument is used. From the equation (2.9), we see that $u_{h}$ is the qualocation approximant to $u+L_{0}^{-1} L_{1}\left(u-u_{h}\right)$ in the case $L=L_{0}$ and hence by the first part of the proof we have

$$
\left\|D^{\alpha}\left(u+L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right)-D^{\alpha} K_{h} * u_{h}\right\| \leq c h^{r+\tau}\left\|u+L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right\|_{R}
$$

By the triangle inequality and (2.2) we have

$$
\begin{align*}
\left\|D^{\alpha} u-D^{\alpha} K_{h} * u_{h}\right\| \leq & \left\|D^{\alpha}\left(u+L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right)-D^{\alpha} K_{h} * u_{h}\right\| \\
& +\left\|D^{\alpha} L_{0}^{-1} L_{1}\left(u-u_{h}\right)\right\| \\
\leq & c h^{r+\tau}\|u\|_{R}+c h^{r+\tau}\left\|u-u_{h}\right\|_{R-\eta}+\left\|u-u_{h}\right\|_{\alpha-\eta} . \tag{4.7}
\end{align*}
$$

Since $\eta>b+1 / 2+\tau^{*}+\alpha$, it follows that $R-\eta<r-1 / 2$; hence Theorem A gives

$$
\begin{equation*}
\left\|u-u_{h}\right\|_{R-\eta} \leq c h^{r-R+\eta}\|u\|_{r} \leq c h^{r-R+\eta}\|u\|_{R} \tag{4.8}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|u-u_{h}\right\|_{\alpha-\eta} \leq\left\|u-u_{h}\right\|_{-\tau} \leq c h^{r+\tau}\|u\|_{r+b} \leq c h^{r+\tau}\|u\|_{R} \tag{4.9}
\end{equation*}
$$

Inequalities (4.7)-(4.9) now give the desired result.

Theorem 4.2. Let the conditions of Theorem 4.1 hold. For $\delta>0$,

$$
\begin{equation*}
\left|D^{\alpha} u-D^{\alpha} K_{h} * u_{h}\right|_{0} \leq c h^{r+\tau}\|u\|_{R^{\prime}} \tag{4.10}
\end{equation*}
$$

where $R^{\prime}=r+\tau^{*}+\alpha+\max \{b+1, \max \{\beta, \delta\}+1 / 2\}$.

Proof. By the triangle inequality we have

$$
\begin{aligned}
\left|D^{\alpha} u-D^{\alpha} K_{h} * u_{h}\right|_{0} & \leq\left|D^{\alpha} u-D^{\alpha} K_{h} * u\right|_{0}+\left|D^{\alpha} K_{h} *\left(u-u_{h}\right)\right|_{0} \\
& =I+I I
\end{aligned}
$$

As in the proof of Theorem 4.1 we have

$$
\begin{equation*}
I \leq c h^{s}|u|_{s+\alpha} \quad \text { for } 0 \leq s \leq 2 q \tag{4.11}
\end{equation*}
$$

To estimate $I I$ we use Bramble and Schatz's trick [6]. Let $k_{h}(x)=$ $K_{h, q}^{1}(x)$. Then we have

$$
\begin{align*}
I I \leq & \left|k_{h} * D^{\alpha} K_{h} *\left(u_{h}-u\right)\right|_{0} \\
& +\left|k_{h} * D^{\alpha} K_{h} *\left(u_{h}-u\right)-D^{\alpha} K_{h} *\left(u_{h}-u\right)\right|_{0} \\
= & I I I+I V . \tag{4.12}
\end{align*}
$$

We will prove separately that $I I I$ and $I V$ satisfy (4.10). Since

$$
I I I \leq c\left\|k_{h} * D^{\alpha} K_{h} *\left(u_{h}-u\right)\right\|_{1}=c \sum_{\gamma=0}^{1}\left\|D^{\gamma} k_{h} * D^{\alpha} K_{h} *\left(u_{h}-u\right)\right\|
$$

from Lemmas 3.3 and 3.4 we infer

$$
\begin{equation*}
I I I \leq c \sum_{\gamma=0}^{\tau^{*}+1}\left\|\partial_{h}^{\alpha+\gamma}\left(u_{h}-u\right)\right\|_{-\tau} \tag{4.13}
\end{equation*}
$$

Again consider first the case $L=L_{0}$. By (4.13), (4.5) and (2.8) we have

$$
\begin{equation*}
I I I \leq c h^{r+\tau}\|u\|_{r+b+\tau^{*}+1+\alpha} \tag{4.14}
\end{equation*}
$$

To estimate $I V$, again we use Lemmas 3.1 and 3.3 to obtain

$$
\begin{align*}
I V & \leq c h^{\tau^{*}}\left|D^{\alpha} K_{h} *\left(u_{h}-u\right)\right|_{\tau^{*}} \\
& =c h^{\tau^{*}} \sum_{\gamma=0}^{\tau^{*}}\left|D^{\alpha+\gamma} K_{h} *\left(u_{h}-u\right)\right|_{0} \\
& \leq c h^{\tau^{*}} \sum_{\gamma=0}^{\tau^{*}}\left|\partial_{h}^{\alpha+\gamma}\left(u_{h}-u\right)\right|_{0} \tag{4.15}
\end{align*}
$$

Using (4.5) and (2.14) we have

$$
\begin{align*}
\left|\partial_{h}^{\alpha+\gamma}\left(u_{h}-u\right)\right|_{0} & \leq c h^{r}\left\|\partial_{h}^{\alpha+\gamma} u\right\|_{r+\max \{\beta, \delta\}+1 / 2} \\
& \leq c h^{r}\|u\|_{r+\max \{\beta, \delta\}+\alpha+\gamma+1 / 2} \tag{4.16}
\end{align*}
$$

From (4.15) and (4.16) we infer

$$
I V \leq c h^{r+\tau}\|u\|_{r+\max \{\beta, \delta\}+\alpha+\tau^{*}+1 / 2}
$$

Hence the result is proved in case $L_{1}=0$. The case $L_{1} \neq 0$ is treated by the familiar argument used in the proof of Theorem 4.1.
5. An example. In this section we test the averaging method when $L$ is the logarithmic-kernel integral operator, for which the principal part $L_{0}$ is an even operator of order -1 . This operator arises in the boundary integral formulation of the Dirichlet problem for Laplace's equation. Consider the boundary value problem

$$
\begin{equation*}
\Delta U=0 \quad \text { in } \Omega, \quad U=g \quad \text { on } \Gamma \tag{5.1}
\end{equation*}
$$

where $\Omega$ is a bounded domain in $\mathbf{R}^{2}$ whose boundary $\Gamma$ is a simple smooth closed curve, parametrized by $\gamma:[0,1] \longrightarrow \mathbf{R}^{2}$ with $\left|\gamma^{\prime}\right|>0$. To avoid the problem of ' $\Gamma$-contours' (see e.g. $[\mathbf{9}, \mathbf{1 7}]$ ), we assume that the transfinite diameter of $\Gamma$ is different from 1.

By Green's theorem we can express $U$ in the form

$$
U(t)=\frac{1}{2 \pi} \int_{\Gamma}\left\{\left(\frac{\partial}{\partial n_{s}} \log |t-s|\right) U(s)-\log |t-s| \frac{\partial U(s)}{\partial n_{s}}\right\} d l_{s}, \quad t \in \Omega
$$

where $d l_{s}$ is the element of arc length and $\partial / \partial n_{s}$ denotes the directional derivative operator in the direction of the outward normal at $s$. By letting $t$ approach the boundary $\Gamma$ and using the continuity properties of the single and double layer potentials (see e.g. $[\mathbf{1 0}, \mathbf{1 7}]$ ) we obtain

$$
\begin{equation*}
U(t)=\frac{1}{\pi} \int_{\Gamma}\left\{\left(\frac{\partial}{\partial n_{s}} \log |t-s|\right) U(s)-\log |t-s| \frac{\partial U(s)}{\partial n_{s}}\right\} d l_{s}, \quad t \in \Gamma \tag{5.2}
\end{equation*}
$$

Letting $z=\partial U / \partial n$ and using the boundary condition for $U$ we infer from (5.2) an integral equation for $z$ :
$-\frac{1}{\pi} \int_{\Gamma} \log |t-s| z(s) d l_{s}=g(t)-\frac{1}{\pi} \int_{\Gamma}\left(\frac{\partial}{\partial n_{s}} \log |t-s|\right) g(s) d l_{s}, \quad t \in \Gamma$.
Using the parametrization for $\Gamma$ we can rewrite (5.3) in the form

$$
\begin{equation*}
L u(x)=f(x) \quad \text { for } x \in[0,1] \tag{5.4}
\end{equation*}
$$

where

$$
u(x)=(2 \pi)^{-1} z[\gamma(x)]\left|\gamma^{\prime}(x)\right|
$$

and

$$
\begin{align*}
L u(x)= & -2 \int_{0}^{1} \log (|\gamma(x)-\gamma(y)|) u(y) d y \\
= & -2 \int_{0}^{1} \log \left|2 e^{-1 / 2} \sin \pi(x-y)\right| u(y) d y  \tag{5.5}\\
& +2 \int_{0}^{1} \log \left(\frac{\left|2 e^{-1 / 2} \sin \pi(x-y)\right|}{|\gamma(x)-\gamma(y)|}\right) u(y) d y \\
= & L_{0} u(x)+L_{1} u(x) \quad \text { for } 0 \leq x \leq 1 .
\end{align*}
$$

It is known that $L_{0}$ is expressible as (see e.g. $[\mathbf{8}, \mathbf{1 7}]$ )

$$
L_{0} u(x)=\hat{u}(0)+\sum_{n \neq 0} \frac{1}{|n|} \hat{u}(n) e^{2 \pi i n x}
$$

We have therefore a special case of (1.1) with $L_{0}$ an even operator of order $\beta=-1$.

We solve (5.4) using piecewise constant splines as trial and test functions. Let $u_{h}$ be given by the qualocation method and let $U_{h}$ be the approximate potential given by

$$
\begin{align*}
U_{h}(t)= & \frac{1}{2 \pi} \int_{\Gamma}\left(\frac{\partial}{\partial n_{s}} \log |t-s|\right) g(s) d l_{s} \\
& -\int_{0}^{1} \log |t-\gamma(x)| u_{h}(x) d x, \quad t \in \Omega \tag{5.6}
\end{align*}
$$

As proved in [8], if we use the Simpson-type quadrature rule with just two points per interval, one at the break-point where the weight is $3 / 7$ and the other at the mid-point where the weight is $4 / 7$, then the additional order of convergence is $b=3$, i.e. the highest order achieved is

$$
\left\|u-u_{h}\right\|_{-4} \leq c h^{5}\|u\|_{4}
$$

Therefore we can investigate $U$ inside the boundary $\Gamma$ by writing

$$
\begin{aligned}
U(t)-U_{h}(t) & =-\int_{0}^{1} \log |t-\gamma(x)|\left(u(x)-u_{h}(x)\right) d x \\
& =\left(u-u_{h}, G(t-\gamma(\cdot))\right) \quad \text { for } t \in \Omega
\end{aligned}
$$

(where $G(t)=-\log |t|$ ) and then using the Cauchy-Schwarz inequality to obtain

$$
\begin{gathered}
\left|U_{h}(t)-U(t)\right| \leq\left\|u_{h}-u\right\|_{-4}\|G(t-\gamma(\cdot))\|_{4} \leq c h^{5}\|u\|_{4}\|G(t-\gamma(\cdot))\|_{4} \\
\text { for } t \in \Omega .
\end{gathered}
$$

However, for $t \in \Gamma$ the use of Cauchy-Schwarz inequality is not possible because of the nonsmoothness of the logarithmic-kernel on the boundary. If we approximate $U$ by $U_{h}^{*}$ defined by (5.6) with $u_{h}$ replaced by $K_{h} * u_{h}$, where $K_{h}=K_{h, 3}^{4}$ as given by Theorem 4.1, we can now make use of (4.2) (with $\alpha=0$ ) to obtain

$$
\begin{aligned}
\left|U_{h}^{*}(t)-U(t)\right| & =\left|\left(K_{h} * u_{h}-u, G(t-\gamma(\cdot))\right)\right| \\
& \leq\left\|K_{h} * u_{h}-u\right\|\|G(t-\gamma(\cdot))\| \\
& \leq c h^{5}\|u\|_{8}\|G(t-\gamma(\cdot))\| \quad \text { for } t \in \Omega \cup \Gamma .
\end{aligned}
$$

Hence the averaging method gives an order of convergence in maxnorm in $\bar{\Omega}$ for the approximation of the potential $U$. However, high smoothness is required for the exact solution $u$ of (5.4).

Order of convergence. Consider now the case $\Gamma$ is the ellipse $\left(t_{1} / 2\right)^{2}+\left(t_{2} / 3\right)^{2}=1$ and $g\left(t_{1}, t_{2}\right)=\sin \left(t_{1}-0.1\right) \cosh \left(t_{2}-0.2\right)$. We use the qualocation package written by B. Burn and D. Dowsett (University of New South Wales, Australia) to carry out the numerical experiment. Note that the exact solution of (5.4) is

$$
\begin{aligned}
u(x)= & 3 \cos 2 \pi x \cos (2 \cos 2 \pi x-0.1) \cosh (3 \sin 2 \pi x-0.2) \\
& +2 \sin 2 \pi x \sin (2 \cos 2 \pi x-0.1) \sinh (3 \sin 2 \pi x-0.2)
\end{aligned}
$$

The numerical results shown in Table 1 are :
(1) The max-errors and the estimated orders of convergence for the qualocation solution,
(2) The errors and estimated orders of convergence at midpoints for the qualocation solution,
(3) The max-errors and the estimated orders of convergence given by the $K$-operator.
The results are as expected. Superconvergence at midpoints given by the qualocation method was proved in [15]. Slow asymptotic
achievement for the $K$-operator is due to the requirement that $N \geq 16$ (see Section 3).

TABLE 1. Errors of the Approximations of Solution.

| $N$ | $\left\|u_{h}-u\right\|_{0}$ | $\max \left\|u_{h}\left(x_{i+1 / 2}\right)-u\left(x_{i+1 / 2}\right)\right\|$ | $\left\|K_{h} * u_{h}-u\right\|_{0}$ |  |  |  |
| ---: | :---: | :---: | :---: | :--- | :--- | :--- |
| 16 | 8.17 | $0.59 \mathrm{E}-00$ | $0.92 \mathrm{E}-00$ |  |  |  |
| 32 | 4.22 | 0.95 | $0.24 \mathrm{E}-00$ | 1.28 | $4.26 \mathrm{E}-02$ | 4.43 |
| 64 | 2.08 | 1.02 | $6.10 \mathrm{E}-02$ | 2.00 | $9.57 \mathrm{E}-04$ | 5.48 |
| 128 | 1.05 | 0.99 | $1.54 \mathrm{E}-02$ | 1.99 | $1.97 \mathrm{E}-05$ | 5.60 |
| 256 | 0.52 | 1.00 | $3.85 \mathrm{E}-03$ | 2.00 | $4.35 \mathrm{E}-07$ | 5.50 |
| 512 | 0.26 | 1.00 | $9.62 \mathrm{E}-04$ | 2.00 | $1.07 \mathrm{E}-08$ | 5.34 |

Approximation of the first derivative. To approximate $u^{\prime}(x)$, by Theorem 4.1 we take $l=5$ and $p=3$. Hence

$$
K_{h}(x)=\frac{1}{h} \sum_{j=-2}^{2} k_{j} \psi^{(5)}\left(\frac{x}{h}-j\right)
$$

where

$$
k_{0}=\frac{319}{192}, \quad k_{1}=k_{-1}=-\frac{107}{288}, \quad k_{2}=k_{-2}=\frac{47}{1152}
$$

The numerical results yield the expected $O\left(h^{5}\right)$ convergence (see Table $2)$.

TABLE 2. Errors of the Approximation of Derivative.

| $N$ | Maximum Errors | Orders of Convergence |
| ---: | :---: | :---: |
| 16 | $39.9 \mathrm{E}-00$ |  |
| 32 | $2.29 \mathrm{E}-00$ | 4.12 |
| 64 | $5.70 \mathrm{E}-02$ | 5.33 |
| 128 | $1.16 \mathrm{E}-03$ | 5.62 |
| 256 | $2.37 \mathrm{E}-05$ | 5.61 |
| 512 | $5.38 \mathrm{E}-07$ | 5.46 |

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