# NEWTON FLOWS FOR REAL EQUATIONS 

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1. Introduction. Let $G: \mathbf{R}^{n} \rightarrow \mathbf{R}^{n}$ be a smooth mapping with Jacobian matrix $D G(x)$. In this paper we shall discuss the dynamical system

$$
\begin{equation*}
N(x)=x-D G(x)^{-1} G(x) \tag{1}
\end{equation*}
$$

provided by Newton's method for the system of equations

$$
\begin{equation*}
G(x)=0 . \tag{2}
\end{equation*}
$$

If $n=2$ and $G$ is a rational mapping $R$ of the complex plane $\mathbf{C}$, then the dynamics of (1), though possibly very complicated and delicate, is understood in terms of the classical and recent theory of Julia sets [ $\mathbf{3}$, 4, 1]. In particular, since $\infty$ is typically a repelling fixed point of $N$ one has that

$$
\begin{equation*}
J_{N}=\text { closure }\left\{x \in \overline{\mathbf{C}}: N^{k}(x)=\infty, \text { for some } k \in \mathbf{N}\right\} \tag{3}
\end{equation*}
$$

is the Julia set of $N(x)=x-R(x) / R^{\prime}(x)$ (here $\overline{\mathbf{C}}=\mathbf{C} \cup\{\infty\}$ and $N^{k}=N \circ \cdots \circ N k$-times). Moreover, if $\bar{x} \in \mathbf{C}$ is a simple zero of $R$, i.e., $R^{\prime}(\bar{x}) \neq 0$, then $\bar{x}$ is an attractive fixed point of $N$; if

$$
\begin{equation*}
A(\bar{x})=\left\{x \in \mathbf{C}: N^{k}(x) \rightarrow \bar{x} \text { as } k \rightarrow \infty\right\}, \tag{4}
\end{equation*}
$$

is its basin of attraction, then

$$
\begin{equation*}
\partial A(\bar{x})=J_{N} . \tag{5}
\end{equation*}
$$

Since (5) is true for any attractive fixed point of $N$ (or even cycles), $J_{N}$ is typically a fractal set which in addition has the interesting property that Newton's method clearly will diverge for initial values in $J_{N}$. On the other hand, if $n$ is not restricted to be 1 or 2 and $G$ is simply

[^0]smooth, the dynamics of (2) is much more delicate and far from being understood. For example:
(a) $N$ may allow strange attractors (see [5]) which is not possible in the complex case.
(b) What is the appropriate analogue to a Julia set? Is there a result similar to (5)?

Associated with the dynamical system (1) there is the system of ordinary differential equations

$$
\begin{equation*}
\dot{x}(t)=-D G(x(t))^{-1} G(x(t)) \tag{6}
\end{equation*}
$$

Knowledge of the flow defined by this system contributes much to the understanding of the orbit structure of (1). We observe that (1) is simply a particular case $(h=1)$ of an Euler method

$$
\begin{equation*}
N_{h}(x)=x-h D G(x)^{-1} G(x) \tag{7}
\end{equation*}
$$

for (6).
The boundary of the domain of definition of (6) is the singular set

$$
\begin{equation*}
S=\left\{x \in \mathbf{R}^{n}: \operatorname{det} D G(x)=0\right\} \tag{8}
\end{equation*}
$$

(typically (i.e., if 0 is a regular value of $\operatorname{det} D G: \mathbf{R}^{n} \rightarrow \mathbf{R}$ ) a collection of smooth $n-1$ manifolds); this set plays an important role in relating the systems (6) and (7) (see [5] for details).

Our objective here is to give some evidence for an interesting conjecture (which is true for Newton's method for rational mappings of C) for the mapping $N$.

Define the Julia-like set of $N$ by

$$
\begin{equation*}
J_{N}=\text { closure }\left\{x \in \mathbf{R}^{n}: N^{k}(x) \in S, \text { some } k \in \mathbf{N} \cup\{0\}\right\} \tag{9}
\end{equation*}
$$

generated by the preimages of $S$. Define the exploding set of $N$ by

$$
\begin{equation*}
E_{N}=\text { closure }\left\{x \in \mathbf{R}^{n}:\left\|N^{k}(x)\right\| \rightarrow \infty \text { as } k \rightarrow \infty\right\} \tag{10}
\end{equation*}
$$

(where $\|\cdot\|$ is some norm for $\mathbf{R}^{n}$ ). While it is apparent that $J_{N} \neq \emptyset$, it is by no means clear or obvious that $E_{N} \neq \emptyset$. We then have the following conjecture.


Figure 1. Bifurcation diagrams for (11) and (12).

Conjecture. $J_{N}=E_{N}$.
(Observe that in the complex case $\infty$ typically has a dense inverse orbit in $J_{N}($ see (3)).)
2. A special case. In this section we shall discuss the above conjecture for a particular model problem in $\mathbf{R}^{2}$.


Figure 2(a). Phase Portrait of (6) with $G$ as in (11), $\mu=2.1$, two sinks.

Let

$$
\begin{equation*}
G(x)=A x-\mu F(x) \tag{11}
\end{equation*}
$$

where

$$
\begin{aligned}
A & =\left(\begin{array}{cc}
2 & -1 \\
-1 & 2
\end{array}\right) \\
F(x) & =\binom{f\left(x_{1}\right)}{f\left(x_{2}\right)}
\end{aligned}
$$

and

$$
f(s)=s-s^{2}
$$



Figure 2(b). Phase Portrait of (6) with $G$ as in (11), $\mu=3.2$, four sinks.

We note that (11) is a standard two point difference approximation for the boundary value problem

$$
\left\{\begin{array}{l}
u^{\prime \prime}+\lambda f(u)=0  \tag{12}\\
u(0)=0=u(\pi)
\end{array}\right.
$$

where $\mu=\lambda \delta^{2}$ and $\delta=\pi / 3$. The bifurcation diagrams for (11) and (12) are given by Figure 1, and Figure 2 shows the continuous time flow of (6) for two choices of $\mu$ and $G$ as in (11).
In this example the singular set $S$ is given by a pair of hyperbolas $S^{+}$ and $S^{-}$. One can easily show that $S^{+}$behaves like a global repeller in both cases and $S^{-}$like a global attractor for $\mu<3$. For $\mu>3$, however, $S^{-}$has passed through a bifurcation state (at $\mu=3$ ) and as a result decomposes into repelling and attracting components.
Figure 3 and 4 show plots of delicate computer experiments displaying $J_{N}$ for various choices of $h$ and $\mu=2.1$ and $\mu=3.2$.
Apparent from these experiments is the crucial role of the singular set $S$ which generates Cantor sets of curves. In addition, Figure 3 demonstrates the importance of the straight line

$$
\begin{equation*}
\mathbf{G}_{\mu}=\left\{x=\left(x_{1}, x_{2}\right): x_{1}+x_{2}+(3-\mu) / \mu=0\right\} . \tag{13}
\end{equation*}
$$

One of the results from [5] is the following theorem.

Theorem. Let $0<\mu<3$ and $0<h<2$.
(a) $\mathbf{G}_{\mu} \subset J_{N}$ and

$$
J_{N}=\text { closure }\left\{x \in \mathbf{R}^{2}: N^{k}(x)=P_{\mu}, \text { some } k \in \mathbf{N}\right\}
$$

where

$$
\left\{P_{\mu}\right\}=S^{-} \cap \mathbf{G}_{\mu} .
$$

(b) $\left.N\right|_{\mathbf{G}_{\mu}}$ is equivalent to a Newton method on the real line

$$
r(s)=s-h k(s) / k^{\prime}(s),
$$

where $k(s)=\mu s^{2}-(\mu+1)(\mu-3) / 4 \mu$.

$h=0.3$

$h=1.4$

$h=1.0$

$h=1.6$

$h=1.7$


Figure 3. The Julia like set $J_{N}$ for (7), $G$ as in (11) and $\mu=2.1$.
(c) $\left.N\right|_{\mathbf{G}_{\mu}}$ is chaotic, i.e., $N$ restricted to $\mathbf{G}_{\mu}$ is equivalent to $z \rightarrow z^{2}$ on the unit circle.

With these observations we are now in a position to discuss the main point of this paper.
3. Theorem and conjecture. Let $G$ be as in (11) and $0<\mu<$ $3,0<h<2$.
(a) There is a dense set $\mathbf{H}_{\mu} \subset \mathbf{G}_{\mu}$ such that each $Q \in \mathbf{H}_{\mu}$ is a periodic repeller of $N$ (see (7)).
(b) Each $Q \in \mathbf{H}_{\mu}$ distinguishes a smooth 1 -manifold $M_{Q}$ which is

- diffeomorphic to $[0, \infty)$
- invariant under $N^{p}$, where $Q$ has period $p$.
(c) For each $x \in M_{Q}-\{Q\},\left\|N^{k p}(x)\right\| \rightarrow \infty$ as $k \rightarrow \infty$.


## REMARK.

(i) Note that (c) means that $E_{N}$ is not empty. Computer experiments based on (c) have provided strong evidence that, in the above case, indeed

$$
J_{N}=E_{N}
$$


$h=0.3$

$h=1.4$

$h=1.7$

$h=1.0$

$h=1.6$

$h=1.8$


Figure 4. The Julia like set $J_{N}$ for (7), $G$ as in (11), $\mu=3.2$.
(ii) The 1-manifolds above are similar to the "hairs" as discussed in [2] on Julia sets for the exponential family

$$
E_{\lambda}(x)=\lambda \exp x
$$

in $\mathbf{C}$.
(iii) It is not difficult to show that the continuous time flow (6) remains bounded for all time. In that regard Euler's method is surprisingly different for all $h>0$ (see [5])!

We shall now give a sketch of a proof of our main result:
Step 1. On shows that the dynamics of $N$ restricted to $\mathbf{G}_{\mu}$ is equivalent to the dynamics

$$
\begin{aligned}
& \alpha \rightarrow 2 \alpha(\bmod 1) \\
& \alpha \in[0,1] .
\end{aligned}
$$



Figure 5.
In this equivalence the point $\left\{P_{\mu}\right\}=S^{-} \cap \mathbf{G}_{\mu}$ corresponds to $1 / 2$. As a consequence one obtains the dense subset $\mathbf{H}_{\mu} \subset \mathbf{G}_{\mu}$ as asserted in (a).

Step 2. Let $Q \in \mathbf{H}_{\mu}$ be a point of period $p$. Using binary operations from step 1 , one can find sequences $\left\{Q_{n}\right\}_{n=1}^{\infty}$ and $\left\{Q^{n}\right\}_{n=1}^{\infty}$ in $\mathbf{G}_{\mu}$ such that
$-\lim Q_{n}=Q=\lim Q^{n}$
$-N^{p}\left(Q_{n}\right)=Q_{n-1}, N^{p}\left(Q^{n}\right)=Q^{n-1}$

$$
\bullet\left\{Q_{n}\right\},\left\{Q^{n}\right\} \subset\left\{x: N^{k}(x)=P_{\mu}, \text { some } k \in N\right\}
$$

Step 3. Each of the points $Q_{n}$ and $Q^{n}$ is a touch point of a component of the iterated inverse images of $S^{-}$. The latter is the set

$$
\cup_{k \geq 0} N^{-k}\left(S^{-}\right)
$$

where $N^{-k}(X)=\left\{x: N^{k}(x) \in X\right\}$. Using these components one may construct sets $M_{n}$ as given in Figure 5, and show that $N^{p}\left(M_{n}\right) \subset$ $M_{n-1}$. Now one defines

$$
M_{Q}=\cap_{n \geq 1} M_{n}
$$

so that $N^{p}\left(M_{Q}\right) \subset M_{Q}$ will follow by construction.

Step 4. On finally must show that

$$
\left\|N^{k p}(x)\right\| \rightarrow \infty \text { as } k \rightarrow \infty
$$

whenever $x \in M_{Q}$. This part of the proof is supported by experimental evidence, except for special choices of $h$ and $\mu$, e.g., this step is not too difficult to establish in case $\mu=2, h=1$ and $p=2$.

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