# INEQUALITIES OF THE MARKOV TYPE FOR PARTIAL DERIVATIVES OF POLYNOMIALS IN SEVERAL VARIABLES

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1. Introduction and main results. Markov-type inequalities give upper bounds for the derivatives of an algebraic polynomial by the polynomial itself. To be more precise, they provide a constant C such that  $||D^{\nu}f|| \leq C||f||$  for all polynomials of degree at most n, where D is the operator of differentiation. The constant C depends on n, on the order  $\nu$  of the derivative, and on the norm  $||\cdot||$ . We here consider the case where  $||\cdot||$  is one of the classical  $L^2$  norms and study the problem of extending such inequalities to the situation when f is a polynomial of several variables and  $D^{\nu}$  is replaced by a partial differential operator.

Let  $\mathcal{P}_n$  be the linear space of all polynomials  $f(t) = \sum_{j=0}^n f_j t^j$  of degree at most n with complex coefficients  $f_j$ . We equip  $\mathcal{P}_n$  with one of the classical Hermite, Laguerre, or Gegenbauer norms. These are defined by

(1) 
$$||f||^2 = \int_{-\infty}^{\infty} |f(t)|^2 e^{-t^2} dt,$$

(2) 
$$||f||^2 = \int_0^\infty |f(t)|^2 t^\alpha e^{-t} dt,$$

(3) 
$$||f||^2 = \int_{-1}^1 |f(t)|^2 (1-t^2)^\alpha dt,$$

where  $\alpha > -1$  is a parameter. Given a polynomial

$$p(\xi) = \xi^m + p_{m-1}\xi^{m-1} + \dots + p_0,$$

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we can consider the differential operator p(D) on  $\mathcal{P}_n$ . The best constant C such that  $||p(D)f|| \leq C||f||$  for all  $f \in \mathcal{P}_n$  is clearly nothing but the norm of the operator  $p(D) : \mathcal{P}_n \to \mathcal{P}_n$ . This constant will be denoted by  $\eta_n(p(D))$ ,  $\lambda_n(p(D))$ ,  $\gamma_n(p(D))$  in dependence on whether the norm  $||\cdot||$  is (1), (2), (3). In [4], we showed that

$$\begin{split} &\lim_{n\to\infty}\frac{\eta_n(p(D))}{n^{m/2}}=2^{m/2},\\ &\lim_{n\to\infty}\frac{\lambda_n(p(D))}{n^m}=\|L_{m,\alpha}\|_{\infty},\\ &\lim_{n\to\infty}\frac{\gamma_n(p(D))}{n^{2m}}=\|G_{m,\alpha}\|_{\infty}, \end{split}$$

where  $L_{m,\alpha}$  and  $G_{m,\alpha}$  are the Volterra integral operators on  $L^2(0,1)$  that are given by

(4)  

$$(L_{m,\alpha}f)(x) = \frac{1}{\Gamma(m)} \int_{x}^{1} x^{\alpha/2} y^{-\alpha/2} (y-x)^{m-1} f(y) \, dy,$$
(5)  

$$(G_{m,\alpha}f)(x) = \frac{1}{2^{m-1}\Gamma(m)} \int_{x}^{1} x^{1/2+\alpha} y^{1/2-\alpha} (y^2 - x^2)^{m-1} f(y) \, dy$$

and  $\|\cdot\|_{\infty}$  denotes the operator norm. Note that these operators are just the iterates (= powers) of their m=1 versions, that is,  $L_{m,\alpha}=L_{1,\alpha}^m$  and  $G_{m,\alpha}=G_{1,\alpha}^m$ .

Let  $[0,1]^N$  be the N-dimensional cube. Given a closed subset E of  $[0,1]^N$ , we define  $\mathcal{P}_n(E)$  as the linear space of all polynomials f of the form

$$f(t_1, \dots, t_N) = \sum_{(i_1/n, \dots, i_N/n) \in E} f_{i_1, \dots, i_N} t_1^{i_1} \cdots t_N^{i_N}$$

with complex coefficients. We will always assume that E contains a point in the interior of  $[0,1]^N$  and that E contains together with each of its points  $(x_1,\ldots,x_N)$  also the cube  $[0,x_1]\times\cdots\times[0,x_N]$ . The most canonical choice of E is

$$\Omega_N := \{(x_1, \dots, x_N) \in [0, 1]^N : x_1 + \dots + x_N \le 1\}.$$

For  $\delta > 0$ , we define  $E^{\delta} = \{(x_1^{\delta}, \dots, x_N^{\delta}) : (x_1, \dots, x_N) \in E\}$ . For example,

$$\Omega_N^{1/2} = \{ (x_1, \dots, x_N) \in [0, 1]^N : x_1^2 + \dots + x_N^2 \le 1 \},$$
  
$$\Omega_N^2 = \{ (x_1, \dots, x_N) \in [0, 1]^N : \sqrt{x_1} + \dots + \sqrt{x_N} \le 1 \}.$$

We endow  $\mathcal{P}_n(E)$  with the N-dimensional versions of the norms (1), (2), (3):

(6)  

$$||f||^{2} = \int_{\mathbf{R}^{N}} |f(t_{1}, \dots, t_{N})|^{2} e^{-t_{1}^{2}} \cdots e^{-t_{N}^{2}} dt_{1} \cdots dt_{N},$$
(7)  

$$||f||^{2} = \int_{(0,\infty)^{N}} |f(t_{1}, \dots, t_{N})|^{2} t_{1}^{\alpha_{1}} \cdots t_{N}^{\alpha_{N}} e^{-t_{1}} \cdots e^{-t_{N}} dt_{1} \cdots dt_{N},$$
(8)  

$$||f||^{2} = \int_{(-1,1)^{N}} |f(t_{1}, \dots, t_{N})|^{2} (1 - t_{1}^{2})^{\alpha_{1}} \cdots (1 - t_{N}^{2})^{\alpha_{N}} dt_{1} \cdots dt_{N},$$

where  $\alpha_j > -1$  for all j.

Take a polynomial

(9) 
$$p(\xi_1, \dots, \xi_N) = \sum_{\nu_1 + \dots + \nu_N \le M} p_{\nu_1, \dots, \nu_N} \xi_1^{\nu_1} \dots \xi_N^{\nu_N}.$$

Here  $\nu_1 + \cdots + \nu_n \leq M$  means that the sum is taken over all N-tuples  $(\nu_1, \ldots, \nu_N)$  of nonnegative integers  $\nu_j$  whose sum does not exceed M. The differential operator on  $\mathcal{P}_n(E)$  given by

$$p(\partial_1, \dots, \partial_N) = \sum_{\nu_1 + \dots + \nu_N < M} p_{\nu_1, \dots, \nu_N} \partial_1^{\nu_1} \dots \partial_N^{\nu_N}$$

may be written in the form

$$p(\partial_1, \dots, \partial_N) = \sum_{\nu_1 + \dots + \nu_N < M} p_{\nu_1, \dots, \nu_N} D^{\nu_1} \otimes \dots \otimes D^{\nu_N} \mid \mathcal{P}_n(E),$$

where | denotes restriction to a subspace. In dependence of the choice of the norm  $||\cdot||$  from (6), (7), (8), we let

$$\eta(p(\partial_1,\ldots,\partial_N)\mid \mathcal{P}_n(E)), \ \lambda(p(\partial_1,\ldots,\partial_N)\mid \mathcal{P}_n(E)), 
\gamma(p(\partial_1,\ldots,\partial_N)\mid \mathcal{P}_n(E))$$

denote the best constant C for which

$$||p(\partial_1,\ldots,\partial_N)f|| \leq C||f||$$
 for all  $f \in \mathcal{P}_n(E)$ .

Of course, this constant is just the norm of  $p(\partial_1, \ldots, \partial_N)$  as an operator on  $\mathcal{P}_n(E)$ . Our standing assumption that E contains the cube  $[0, x_1] \times \cdots \times [0, x_N]$  together with each of its points  $(x_1, \ldots, x_N)$  guarantees that  $p(\partial_1, \ldots, \partial_N)$  maps  $\mathcal{P}_n(E)$  into itself.

If  $A_1, \ldots, A_k$  are operators on  $L^2(0,1)$ , their tensor product  $A := A_1 \otimes \cdots \otimes A_k$  on  $L^2((0,1)^k)$  is defined in the usual way. If  $L^2(E)$  is an invariant subspace for A, we denote by  $A \mid L^2(E)$  the restriction of A to  $L^2(E)$ . The notation  $a_n \sim b_n$  means that  $a_n/b_n \to 1$  as  $n \to \infty$ .

We need one more definition. Let  $\nu := (\nu_1, \ldots, \nu_N)$  be an N-tuple of nonnegative integers. We put  $|\nu| = \nu_1 + \cdots + \nu_N$  and always suppose that  $|\nu| \geq 1$ . Let  $\nu_{j_1}, \ldots, \nu_{j_k}$   $(j_1 < \cdots < j_k)$  denote the nonzero integers among  $\nu_1, \ldots, \nu_N$ . Given a point  $(x_1, \ldots, x_k) \in [0, 1]^k$ , we define  $(x_1, \ldots, x_k)_{\nu}$  as the point in  $[0, 1]^N$  whose  $\ell$ th coordinate is 0 if  $\nu_{\ell} = 0$  and  $x_m$  if  $\ell = j_m$ . For example, if  $\nu = (0, \nu_2, 0, 0, \nu_5)$  with nonzero  $\nu_2$  and  $\nu_5$ , then k = 2,  $j_1 = 2$ ,  $j_2 = 5$  and  $(x_1, x_2)_{\nu} = (0, x_1, 0, 0, x_2)$ . Finally, given  $E \subset [0, 1]^N$ , we put

$$E_{\nu} = \{(x_1, \dots, x_k) \in [0, 1]^k : (x_1, \dots, x_k)_{\nu} \in E\}.$$

Note that if E is  $[0,1]^N$ ,  $\Omega_N$ ,  $\Omega_N^{\delta}$ , then  $E_{\nu}$  is simply  $[0,1]^k$ ,  $\Omega_k$ ,  $\Omega_k^{\delta}$ , respectively.

**Theorem 1.1.** The best constants for  $\partial_1^{\nu_1} \cdots \partial_N^{\nu_N}$  have the asymptotic behavior

(10)  

$$\eta(\partial_{1}^{\nu_{1}} \cdots \partial_{N}^{\nu_{N}} \mid \mathcal{P}_{n}(E)) \sim n^{|\nu|/2} \max_{(x_{1}, \dots, x_{k}) \in E_{\nu}} (2x_{1})^{\nu_{j_{1}}/2} \cdots (2x_{k})^{\nu_{j_{k}}/2},$$
  
(11)  
 $\lambda(\partial_{1}^{\nu_{1}} \cdots \partial_{N}^{\nu_{N}} \mid \mathcal{P}_{n}(E)) \sim n^{|\nu|} \|L_{\nu_{j_{1}}, \alpha_{j_{1}}} \otimes \cdots \otimes L_{\nu_{j_{k}}, \alpha_{j_{k}}} \mid L^{2}(E_{\nu})\|_{\infty},$   
(12)  
 $\gamma(\partial_{1}^{\nu_{1}} \cdots \partial_{N}^{\nu_{N}} \mid \mathcal{P}_{n}(E)) \sim n^{2|\nu|} \|G_{\nu_{j_{1}}, \alpha_{j_{1}}} \otimes \cdots \otimes G_{\nu_{j_{k}}, \alpha_{j_{k}}} \mid L^{2}(E_{\nu})\|_{\infty}.$ 

If E is  $[0,1]^N$ , then the maximum in (10) is  $2^{|\nu|/2}$ , and if  $E = \Omega_N^{\delta}$ , then this maximum equals

$$\left(\frac{2^{1/\delta}\nu_{j_1}}{|\nu|}\right)^{\delta\,\nu_{j_1}/2}\cdots\left(\frac{2^{1/\delta}\nu_{j_k}}{|\nu|}\right)^{\delta\,\nu_{j_k}/2}.$$

Thus, in the Hermite case the coefficient in the asymptotic formula is explicitly available. Theorem 1.2 will show that the Gegenbauer case can be reduced to the Laguerre case.

When dealing with the coefficients in the asymptotics, we have k dimensions instead of N. To avoid double subscripts, we assume in this context that we are given a k-tuple  $\nu = (\nu_1, \ldots, \nu_k)$  of positive integers and a k-tuple  $\alpha = (\alpha_1, \ldots, \alpha_k)$  of real numbers such that  $\alpha_j > -1$  for all j. The set E is now a closed subset of  $[0,1]^k$  which contains a point of  $(0,1)^k$  and which contains  $[0,x_1]\times\cdots\times[0,x_k]$  together with each of its points  $(x_1,\ldots,x_k)$ . Recall that  $E^2 = \{(x_1^2,\ldots,x_k^2):(x_1,\ldots,x_k)\in E\}$ . Henceforth we also make use of the abbreviations

$$L_{\nu,\alpha} = L_{\nu_1,\alpha_1} \otimes \cdots \otimes L_{\nu_k,\alpha_k}, \quad G_{\nu,\alpha} = G_{\nu_1,\alpha_1} \otimes \cdots \otimes G_{\nu_k,\alpha_k}.$$

**Theorem 1.2.** The Gegenbauer and Laguerre cases are related by the equality

$$||G_{\nu,\alpha}||L^2(E)||_{\infty} = 2^{-|\nu|}||L_{\nu,\alpha}||L^2(E^2)||_{\infty}.$$

Bounds for  $||L_{\nu,\alpha}||L^2(E)||_{\infty}$  are delivered by the trivial inequalities

$$(13) \quad \frac{\|L_{\nu,\alpha}^* L_{\nu,\alpha} \mid L^2(E)\|_2}{\|L_{\nu,\alpha} \mid L^2(E)\|_2} \le \|L_{\nu,\alpha} \mid L^2(E)\|_{\infty} \le \|L_{\nu,\alpha} \mid L^2(E)\|_2,$$

where  $\|\cdot\|_2$  stands for the Hilbert-Schmidt norm. Here  $L_{\nu,\alpha}^*$  is the adjoint of  $L_{\nu,\alpha} \mid L^2(E)$ , that is, by  $L_{\nu,\alpha}^*$  we actually mean  $(L_{\nu,\alpha} \mid L^2(E))^*$ .

Theorem 1.3. We have

$$||L_{\nu,\alpha}| L^{2}(\Omega_{k}^{\delta})||_{2}^{2} = \frac{\delta^{k}}{\Gamma(2\delta|\nu|+1)} \prod_{j=1}^{k} \frac{\Gamma(\alpha_{j}+1)\Gamma(2\nu_{j}-1)\Gamma(2\delta\nu_{j})}{\Gamma(\alpha_{j}+2\nu_{j})\Gamma(\nu_{j})^{2}}$$

and

$$\begin{split} \|L_{\nu,\alpha}^* L_{\nu,\alpha} \mid L^2(\Omega_k^{\delta})\|_2^2 \\ & \geq \frac{(2\delta)^k}{\Gamma(4\delta|\nu|+1)} \prod_{j=1}^k \frac{\Gamma(\alpha_j+1)^2 \Gamma(2\nu_j-1)^2 \Gamma(4\delta\nu_j)}{\Gamma(\alpha_j+2\nu_j)^2 (\alpha_j+2\nu_j+1) \Gamma(\nu_j)^4}. \end{split}$$

Inserting the expressions from Theorem 1.3 in (13) we get bounds

$$b_1(\nu, \alpha, \Omega_k^{\delta}) \le ||L_{\nu, \alpha}| L^2(\Omega_k^{\delta})||_{\infty} \le b_2(\nu, \alpha, \Omega_k^{\delta}).$$

If  $k, \delta, \alpha_1, \ldots, \alpha_k$  remain fixed and  $\nu_j \to \infty$  for all j, then

(14) 
$$\frac{b_2(\nu,\alpha,\Omega_k^{\delta})}{b_1(\nu,\alpha,\Omega_k^{\delta})} \sim (2\pi\delta)^{(k-1)/4} \frac{\nu_1^{1/4} \cdots \nu_k^{1/4}}{(\nu_1 + \cdots + \nu_k)^{1/4}}.$$

For k=1, the right-hand side of (14) is 1, but for  $k\geq 2$  it increases (though moderately). The following result reveals that the upper bound  $b_2(\nu,\alpha,\Omega_k^{\delta})$  is asymptotically sharp as  $\nu$  goes to infinity along a straight line.

**Theorem 1.4.** Let  $\nu = (\varrho_1 \tau, \dots, \varrho_k \tau)$  where  $\varrho_1, \dots, \varrho_k$  are positive real numbers. Then as  $\tau \to \infty$ ,

$$||L_{\nu,\alpha}||L^2(\Omega_k^{\delta})||_{\infty} \sim ||L_{\nu,\alpha}||L^2(\Omega_k^{\delta})||_2.$$

If, for example,  $k=2,\ \delta=1,\ \alpha_1=\alpha_2=0,\ \nu_1=\nu_2=:\nu,$  then Theorems 1.1, 1.3, 1.4 yield

(15) 
$$\lambda(\partial_1^{\nu}\partial_2^{\nu} \mid \mathcal{P}_n(\Omega_2)) \sim \frac{(2\pi\nu)^{1/4}}{4} \frac{1}{2^{2\nu}\nu^2} \frac{1}{\Gamma(\nu)^2} n^{2\nu}.$$

In the case k = 1 we infer from the same theorems that

(16) 
$$\lambda(\partial_1^{\nu} \mid P_n(\Omega_2)) = \lambda(\partial_2^{\nu} \mid P_n(\Omega_2)) \sim \frac{1}{2\nu\Gamma(\nu)} n^{\nu}.$$

It is easily seen that always

(17) 
$$\lambda(\partial_1^{\nu}\partial_2^{\nu} \mid \mathcal{P}_n(\Omega_2)) \leq \lambda(\partial_1^{\nu} \mid P_n(\Omega_2))\lambda(\partial_2^{\nu} \mid P_n(\Omega_2)).$$

However, from (15) and (16) we obtain that

$$\frac{\lambda(\partial_1^{\nu}\partial_2^{\nu}\mid \mathcal{P}_n(\Omega_2))}{\lambda(\partial_1^{\nu}\mid P_n(\Omega_2))\lambda(\partial_2^{\nu}\mid P_n(\Omega_2))}\sim \frac{(2\pi\nu)^{1/4}}{2^{2\nu}},$$

which is much smaller than 1 if  $\nu$  is large and hence much sharper than (17).

We finally turn to linear combinations of partial derivatives. Let

$$p_0(\xi_1, \dots, \xi_N) = \sum_{\nu_1 + \dots + \nu_N = M} p_{\nu_1, \dots, \nu_N} \xi_1^{\nu_1} \dots \xi_N^{\nu_N}$$

be the principal part of polynomial (9). We write  $a_n \simeq b_n$  if there exist constants  $0 < c_1 < c_2 < \infty$  such that  $c_1b_n \leq a_n \leq c_2b_n$  for all sufficiently large n. The following theorem reveals that the asymptotic behavior of the best constants for linear partial differential operators with constant coefficients is completely determined by their principal parts.

**Theorem 1.5.** Let C stand for  $\eta$ ,  $\lambda$ , or  $\gamma$  and put  $\sigma=1/2$  in the Hermite case,  $\sigma=1$  in the Laguerre case, and  $\sigma=2$  in the Gegenbauer case. Then

(18) 
$$C(p_0(\partial_1, \dots, \partial_N) \mid \mathcal{P}_n(E)) \simeq n^{\sigma M}$$

and

(19) 
$$C(p(\partial_1, \dots, \partial_N) \mid \mathcal{P}_n(E)) \sim C(p_0(\partial_1, \dots, \partial_N) \mid \mathcal{P}_n(E)).$$

The paper is organized as follows. Section 2 is devoted to comments on previous work and the results and methods of this paper. Sections 3 to 7 contain the proofs of Theorems 1.1 to 1.5 and in Section 8 we list some problems we have to leave open.

2. Remarks on the history. The problem of finding upper bounds for the derivatives of functions in terms of the functions themselves has a long and rich history. Nowadays one speaks of Bernstein-type inequalities if the functions are trigonometric polynomials and of Markov-type inequalities in the case of algebraic polynomials. The Markov brothers [24, 25] found the best constant C in the inequality  $||D^m f|| \leq C||f||$  when  $||\cdot||$  is the  $L^{\infty}$  norm on some finite interval. We refer to [2, 15, 26, 27, 29, 30] for more on the subject and here confine ourselves to the asymptotics of the best constants in Markov-type inequalities with  $L^2$  norms.

An  $L^2$  version of a Markov-type inequality was first established by Erhard Schmidt [31] and subsequently for  $L^p$  norms by Hille, Szegő and Tamarkin [17]. In 1944, Schmidt [32] proved that

$$\eta_n(D) = \sqrt{2n}, \quad \lambda_n(D) \sim \frac{2}{\pi} n, \quad \gamma_n(D) \sim \frac{1}{\pi} n^2,$$

assuming  $\alpha=0$ , and even provided the next two terms in the asymptotics of  $\lambda_n(D)$  and  $\gamma_n(D)$ . Schmidt also observed that for the Hermite weight the problem is more or less trivial. Shampine [33] then showed that, again for  $\alpha=0$ ,  $\lambda_n(D^2)\sim n^2/\mu_0^2$  and  $\gamma_n(D^2)\sim n^4/(4\mu_0^2)$  where  $\mu_0$  is the smallest positive root of the equation  $1+\cos\mu\cosh\mu=0$ . For the exact values of  $\lambda_n(D)$  and  $\gamma_n(D)$  in the case  $\alpha=0$  see [17, 21, 36]. The idea that the best constants in question are the largest singular value (= operator norm) of certain matrices was developed in [7, 8, 9] and used to derive bounds for

$$\liminf_{n \to \infty} \frac{\lambda_n(D^m)}{n^m}, \quad \limsup_{n \to \infty} \frac{\lambda_n(D^m)}{n^m}$$

for general m and  $\alpha$ . In [3, 4] we proved that  $\lambda_n(D^m)/n^m$  and  $\gamma_n(D^m)/n^{2m}$  converge to a limit as  $n \to \infty$  and identified these limits as the operator norms of certain Volterra integral operators.

The appearance of Volterra operators in this context connects us with another field of research. Paul Halmos [16] was probably the first to state explicitly that the operator norm of the operator

$$(L_{1,0}f)(x) = \int_{x}^{1} f(y) dy$$

on  $L^2(0,1)$  is  $2/\pi$ . Combining this result with our formula  $\lambda_n(D) \sim \|L_{1,0}\|_{\infty}$  n reproduces Schmidt's formula  $\lambda_n(D) \sim (2/\pi) \, n$ . Halmos also raised the question of determining the operator norms of the iterates  $L_{1,0}^m = L_{m,0}$ . This problem was subsequently studied by many authors, including [1, 12, 18, 22, 23, 35]. The much earlier paper [14] was detected by Thorpe [35] to be also of relevancy in connection with the matter. The reader may consult [3, 4] for details. In the course of these investigations sharp bounds for  $\|L_{m,0}\|_{\infty}$  and the asymptotic formula  $\|L_{m,0}\|_{\infty} \sim 1/(2 \, m!)$  were established. In [4] we solved the corresponding problems for the norms  $\|L_{m,0}\|_{\infty}$  and  $\|G_{m,\alpha}\|_{\infty}$ . In particular, the N=1 versions of Theorems 1.2, 1.3, 1.4 are already in [4]. We also want to note that Halmos's  $\|L_{1,0}\|_{\infty} = 2/\pi$  was in [10] and [4] extended to the equality  $\|L_{1,\alpha}\|_{\infty} = 1/j_0(\alpha)$  where  $j_0(\alpha)$  is the smallest positive zero of the Bessel function  $J_{(\alpha-1)/2}$ .

The literature on multivariate Markov-type inequalities is immense, a main topic of research being inequalities for the  $L^{\infty}$  norm on multidimensional regions and manifolds. See, for example, [19, 20, 28]. However, we are not aware of publications dealing with best constants in multivariate Markov-type inequalities with the  $L^2$  norm and for partial derivatives of arbitrary order. Note that even the one-dimensional versions of the results of this paper were established only in [3, 4]. Clearly, for  $E = [0,1]^N$  our Theorems 1.1 to 1.4 simply follow from their one-dimensional counterparts by taking tensor products. However, passage to the simplex  $E = \Omega_N$  makes things nontrivial. Moreover, Theorem 1.2 even motivates consideration of the entire scale  $E = \Omega_N^{\delta}$ .

The approach employed in  $[\mathbf{3}, \mathbf{4}]$  and also in Section 4 of this paper is based on an idea by Harold Widom  $[\mathbf{37}, \mathbf{38}, \mathbf{39}]$ , which was independently also introduced by Lawrence Shampine  $[\mathbf{33}, \mathbf{34}]$ . In order to find the asymptotic behavior of spectral quantities of a sequence of  $n \times n$  matrices  $A_n$ , they associated an integral operator  $W_{A_n}$  on  $L^2(0,1)$  with each matrix  $A_n$  and then studied whether, after appropriate scaling, the operators  $W_{A_n}$  converge uniformly to some limiting operator. In this way Widom and Shampine were able to express asymptotic properties of  $A_n$  in terms of the limiting operator. In particular, Shampine  $[\mathbf{33}]$  considered  $W_{A_n}$  where  $A_n$  is the matrix representation of the operator  $(D^2)^*D^2$ . For  $(D^m)^*D^m$  with  $m \geq 2$ , the matrices become more complicated and hence we have the limitation to m = 2 in  $[\mathbf{33}]$ . What is new in our approach is that we exploit the

fact that the replacement  $A_n \mapsto W_{A_n}$  is an algebraic homomorphism that preserves also tensor products. Thus, we simply consider the operator  $W_{A_n}$  for  $A_n$  being the matrix representation of D and show that, after scaling,  $W_{A_n}$  converges to some limiting operator. Once this has been done, we can easily pass to adjoints, sums, products, and tensor products.

The reasoning in Section 5 is similar to the one of [3, 4]. The argument used in Section 6 is different from [3, 4] and based on a strategy that was in another context pursued in [6].

**3.** Hermite weights. In this section we prove the assertion of Theorem 1.1 that concerns the Hermite case. This case is particularly simple.

An orthonormal basis in  $\mathcal{P}_n$  with the norm (1) is  $\{h_0, h_1, \ldots, h_n\}$  where  $h_k$  is the kth normalized Hermite polynomial. We have

(20) 
$$D^{\nu}h_{i} = 2^{\nu/2} \sqrt{\frac{\Gamma(i+1)}{\Gamma(i-\nu+1)}} h_{i-\nu}$$

for  $\nu \leq i$ . As usual,  $h_{i_1} \otimes \cdots \otimes h_{i_N}$  is defined by

$$(h_{i_1} \otimes \cdots \otimes h_{i_N})(t_1, \ldots, t_N) = h_{i_1}(t_1) \cdots h_{i_N}(t_N).$$

Then  $\{h_{i_1} \otimes \cdots \otimes h_{i_N} : (i_1/n, \dots, i_N/n) \in E\}$  is an orthonormal basis in  $\mathcal{P}_n(E)$ . For

$$f = \sum f_{i_1,\dots,i_N} h_{i_1} \otimes \dots \otimes h_{i_N} \in \mathcal{P}_n(E),$$

we then get

(21) 
$$\|(D^{\nu_{1}} \otimes \cdots \otimes D^{\nu_{N}})f\|^{2}$$

$$= ((D^{\nu_{1}} \otimes \cdots \otimes D^{\nu_{N}})f, (D^{\nu_{1}} \otimes \cdots \otimes D^{\nu_{N}})f)$$

$$= \left(\sum f_{i_{1}, \dots, i_{N}} D^{\nu_{1}} h_{i_{1}} \otimes \cdots \otimes D^{\nu_{N}} h_{i_{N}}, \right.$$

$$\sum f_{k_{1}, \dots, k_{N}} D^{\nu_{1}} h_{k_{1}} \otimes \cdots \otimes D^{\nu_{N}} h_{k_{N}} \right)$$

$$= \sum \sum f_{i_{1}, \dots, i_{N}} \overline{f_{k_{1}, \dots, k_{N}}} (D^{\nu_{1}} h_{i_{1}}, D^{\nu_{1}} h_{k_{1}}) \cdots$$

$$\times (D^{\nu_{N}} h_{i_{N}}, D^{\nu_{N}} h_{k_{N}}).$$

Using (20) and the orthonormality of  $h_0, h_1, \ldots, h_n$ , we see that (21) equals

$$\sum |f_{i_1,\dots,i_N}|^2 2^{\nu_1} \frac{\Gamma(i_1+1)}{\Gamma(i_1-\nu_1+1)} \cdots 2^{\nu_N} \frac{\Gamma(i_N+1)}{\Gamma(i_N-\nu_N+1)},$$

the sum over  $i_1 \geq \nu_1, \ldots, i_N \geq \nu_N, (i_1/n, \ldots, i_N/n) \in E$ . It follows that the operator norm of  $D^{\nu_1} \otimes \cdots \otimes D^{\nu_N}$  on  $\mathcal{P}_n(E)$  is given by

$$\begin{split} \|D^{\nu_1} \otimes \cdots \otimes D^{\nu_N} &\mid \mathcal{P}_n(E) \|_{\infty}^2 \\ &= \max 2^{|\nu|} \frac{\Gamma(i_1 + 1)}{\Gamma(i_1 - \nu_1 + 1)} \cdots \frac{\Gamma(i_N + 1)}{\Gamma(i_N - \nu_N + 1)} \\ &= 2^{|\nu|} \max i_1(i_1 - 1) \cdots (i_1 - \nu_1 + 1) \cdots \\ &\qquad \times i_N(i_N - 1) \cdots (i_N - \nu_N + 1), \end{split}$$

the maximum over  $i_1 \geq \nu_1, \ldots, i_N \geq \nu_N, (i_1/n, \ldots, i_N/n) \in E$ . Consequently,

(22) 
$$\eta^{2}(\partial_{1}^{\nu_{1}}\cdots\partial_{N}^{\nu_{N}}\mid\mathcal{P}_{n}(E))/(2n)^{|\nu|}$$

$$= \max\frac{i_{1}}{n}\left(\frac{i_{1}}{n} - \frac{1}{n}\right)\cdots\left(\frac{i_{1}}{n} - \frac{\nu_{1} - 1}{n}\right)\cdots$$

$$\times \frac{i_{N}}{n}\left(\frac{i_{N}}{n} - \frac{1}{n}\right)\cdots\left(\frac{i_{N}}{n} - \frac{\nu_{N} - 1}{n}\right)$$

$$= \max\left(\frac{i_{1}}{n}\right)^{\nu_{1}}\cdots\left(\frac{i_{N}}{n}\right)^{\nu_{N}}\left(1 + O\left(\frac{1}{n}\right)\right).$$

The limit of (22) is

$$\max_{(x_1,\dots,x_N)\in E} x_1^{\nu_1}\cdots x_N^{\nu_N} = \max_{(x_1,\dots,x_k)\in E_{\nu}} x_1^{\nu_{j_1}}\cdots x_k^{\nu_{j_k}},$$

which proves (10).

4. Laguerre and Gegenbauer weights. This section is devoted to the proof of Theorem 1.1 in the Laguerre and Gegenbauer cases and to the proof of Theorem 1.2.

Suppose first that  $\nu_j \geq 1$  for all j. In that case  $E_{\nu} = E$ . We denote by  $p_{\alpha,i}$  the ith normalized Laguerre or Gegenbauer polynomial

in the norm (2) or (3), respectively. Then  $\{p_{\alpha,0}, p_{\alpha,1}, \ldots, p_{\alpha,n}\}$  is an orthonormal basis in  $\mathcal{P}_n$ . Let  $D_{\alpha,n}$  be the matrix representation of the operator  $D: \mathcal{P}_n \to \mathcal{P}_n$  in this basis. Given a matrix  $A_n = (a_{ik})_{i,k=0}^n$ , we denote by  $W_{A_n}$  the integral operator on  $L^2(0,1)$  with the piecewise constant kernel  $(n+1)a_{[(n+1)x],[(n+1)y]}$ , where  $[\xi]$  is the integral part of  $\xi$ . Recall that  $L_{m,\alpha}$  and  $G_{m,\alpha}$  are given by (4) and (5). Put  $\sigma=1$  and  $T_{m,\alpha}=L_{m,\alpha}$  in the Laguerre case and  $\sigma=2$  and  $T_{m,\alpha}=G_{m,\alpha}$  in the Gegenbauer case. In [4] we showed that

(23) 
$$||(n+1)^{-\sigma m} W_{D_{\alpha,n}^m} - T_{m,\alpha}||_{\infty} \longrightarrow 0 \text{ as } n \to \infty.$$

An orthonormal basis in  $\mathcal{P}_n([0,1]^N)$  is

$$\mathcal{F} := \{ p_{\alpha_1, i_1} \otimes \cdots \otimes p_{\alpha_N, i_N} : (i_1, \dots, i_N) \in S_n^N \}$$

where  $S_n := \{0, 1, ..., n\}$ . If A is a linear operator on  $\mathcal{P}_n([0, 1]^N)$ , we denote by  $A_n = (a_{i,k})_{i,k \in S_n^N}$  its matrix representation in this basis. Thus, if  $i = (i_1, ..., i_N)$  and  $k = (k_1, ..., k_N)$ , then

$$a_{i,k} = (A(p_{\alpha_1,k_1} \otimes \cdots \otimes p_{\alpha_N,k_N}), p_{\alpha_1,i_1} \otimes \cdots \otimes p_{\alpha_N,i_N}).$$

In the case where  $A = \partial_1^{\nu_1} \cdots \partial_N^{\nu_N} = D^{\nu_1} \otimes \cdots \otimes D^{\nu_N}$ , the matrix representation  $A_n$  is just the Kronecker product  $D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N}$ . We associate with  $A_n$  the integral operator  $W_{A_n}$  on  $L^2((0,1)^N)$  given by

$$(W_{A_n}f)(x_1,\ldots,x_N)$$

$$= (n+1)^N \int_{(0,1)^N} a_{[(n+1)x_1],\ldots,[(n+1)x_N],[(n+1)y_1],\ldots,[(n+1)y_N]} \times f(y_1,\ldots,y_N) dy.$$

Throughout what follows  $||A||_{\infty}$  denotes the operator norm if A is an operator and the spectral norm in case A is a matrix.

**Lemma 4.1** (Widom and Shampine). If  $C_1, \ldots, C_N$  are linear operators on  $\mathcal{P}_n$ , A, B are linear operators on  $\mathcal{P}_n([0,1]^N)$ , and  $\alpha, \beta \in \mathbf{C}$ , then

$$W_{\alpha A_n + \beta B_n} = \alpha W_{A_n} + \beta W_{B_n},$$

$$W_{A_n B_n} = W_{A_n} W_{B_n},$$

$$W_{(C_1)_n \otimes \cdots \otimes (C_N)_n} = W_{(C_1)_n} \otimes \cdots \otimes W_{(C_N)_n},$$

$$\|W_{A_n}\|_{\infty} = \|A_n\|_{\infty}.$$

*Proof.* Let  $I_k$  be the interval (k/(n+1), (k+1)/(n+1)), denote by  $\chi_k$  the characteristic function of  $I_k$ , and consider the operators

$$R: \mathbf{C}^{n+1} \longrightarrow L^2(0,1), \quad \{x_k\}_{k=0}^n \longmapsto \sqrt{n+1} \sum_{k=0}^n x_k \chi_k,$$

$$S: L^2(0,1) \longrightarrow \mathbf{C}^{n+1}, \quad f \longmapsto \left\{ \sqrt{n+1} \int_{I_k} f(x) \, dx \right\}_{k=0}^n.$$

It can be verified straightforwardly that  $\|R\|_{\infty} = \|S\|_{\infty} = 1$ , that SR = I, and that  $R(C_i)_n S = W_{(C_i)_n}$  and  $(R \otimes \cdots \otimes R) X_n (S \otimes \cdots \otimes S) = W_{X_n}$  for every linear operator X on  $\mathcal{P}_n([0,1]^N)$ . It follows that  $W_{\alpha A_n + \beta B_n} = \alpha W_{A_n} + \beta W_{B_n}$ ,  $W_{A_n B_n} = W_{A_n} W_{B_n}$ , and  $W_{(C_1)_n \otimes \cdots \otimes (C_N)_n} = W_{(C_1)_n} \otimes \cdots \otimes W_{(C_N)_n}$ . Since

$$||W_{A_n}||_{\infty} = ||(R \otimes \cdots \otimes R)A_n(S \otimes \cdots \otimes S)||_{\infty} \leq ||A_n||_{\infty}$$
$$= ||(S \otimes \cdots \otimes S)W_{A_n}(R \otimes \cdots \otimes R)||_{\infty} \leq ||W_{A_n}||_{\infty},$$

we arrive at the equality  $||W_{A_n}||_{\infty} = ||A_n||_{\infty}$ .

From (23) we infer that

$$\|(n+1)^{-\sigma|\nu|} W_{D_{\alpha_1,n}^{\nu_1}} \otimes \cdots \otimes W_{D_{\alpha_N,n}^{\nu_N}} - T_{\nu_1,\alpha_1} \otimes \cdots \otimes T_{\nu_N,\alpha_N}\|_{\infty} \longrightarrow 0,$$

and from Lemma 4.1 we therefore deduce that

$$(24) \|(n+1)^{-\sigma|\nu|} W_{D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N}} - T_{\nu_1,\alpha_1} \otimes \cdots \otimes T_{\nu_N,\alpha_N} \|_{\infty} \longrightarrow 0.$$

**Lemma 4.2.** Let X be a Banach space and suppose X is the direct sum of two closed subspaces U and V,  $X = U \oplus V$ . Let K be a bounded linear operator on X which has U as an invariant subspace. Then, with  $P_U$  denoting the projection of X onto U parallel to V,

$$||K \mid U||_{\infty} = ||KP_U||_{\infty}.$$

*Proof.* The decomposition  $X = U \oplus V$  allows us to represent K by a  $2 \times 2$  operator matrix. Since U is an invariant subspace, the 2,1 entry of this matrix is zero. Thus,

$$K = \begin{pmatrix} A & B \\ 0 & C \end{pmatrix}.$$

Clearly,  $K \mid U = A$  and

$$KP_U = \left( egin{array}{cc} A & B \\ 0 & C \end{array} \right) \left( egin{array}{cc} I & 0 \\ 0 & 0 \end{array} \right) = \left( egin{array}{cc} A & 0 \\ 0 & 0 \end{array} \right).$$

This shows that  $||K||U||_{\infty} = ||KP_U||_{\infty}$ .

We may think of  $\mathcal{P}_n(E)$  as a subspace of  $\mathcal{P}_n([0,1]^N)$ . Moreover, in the orthonormal basis  $\mathcal{F}$  we may identify  $\mathcal{P}_n([0,1]^N)$  with  $\ell^2(S_n^N)$  and  $\mathcal{P}_n(E)$  with  $\ell^2(\Pi_n)$  where  $S_n = \{0,1,\ldots,n\}$  and

$$\Pi_n = \{(i_1, \dots, i_N) \in S_n^N : (i_1/n, \dots, i_N/n) \in E\}.$$

We are interested in the norm  $||D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N}|| \ell^2(\Pi_n)||_{\infty}$ . Let  $P_{\Pi_n}$  be the orthogonal projection of  $\ell^2(S_n^N)$  onto  $\ell^2(\Pi_n)$ . By Lemma 4.2,

$$||D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N} | \ell^2(\Pi_n)||_{\infty}$$

$$= ||(D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N}) P_{\Pi_n}||_{\infty}$$

where the norm on the right is taken over  $\ell^2(S_n^N)$ . We denote the matrix representation of  $P_{\Pi_n}$  in the orthonormal basis  $\mathcal{F}$  also by  $P_{\Pi_n}$ , although strict use of notation would require to denote it by  $(P_{\Pi_n})_n$ . A little thought reveals that

(25) 
$$(P_{\Pi_n})_{i,k} = \begin{cases} 1 & \text{if } i = k \in \Pi_n, \\ 0 & \text{otherwise.} \end{cases}$$

**Lemma 4.3.** The operators  $W_{P_{\Pi_n}}$  converge strongly (= pointwise) on  $L^2((0,1)^N)$  to the operator

$$P_E: L^2((0,1)^N) \longrightarrow L^2((0,1)^N), \quad (P_E f)(x) = \begin{cases} f(x) & \text{if } x \in E, \\ 0 & \text{if } x \notin E. \end{cases}$$

*Proof.* By Lemma 4.1,  $\|W_{P_{\Pi_n}}\|_{\infty} = \|P_{\Pi_n}\|_{\infty} = 1$ . Hence it suffices to prove that  $W_{P_{\Pi_n}}f \to P_E f$  for all f in some dense subset of  $L^2((0,1)^N)$ ,

say for  $f \in C([0,1]^N)$ . Thus, fix  $f \in C([0,1]^N)$  and  $\varepsilon > 0$ . For  $i = (i_1, \ldots, i_N) \in \{0, 1, \ldots, n\}^N$ , put

$$Q_i = \left[\frac{i_1}{n+1}, \frac{i_1+1}{n+1}\right) \times \cdots \times \left[\frac{i_N}{n+1}, \frac{i_N+1}{n+1}\right).$$

We have

(26) 
$$\|W_{P_{\Pi_n}}f - P_E f\|^2 = \sum_{i_1, \dots, i_N=0}^n \int_{Q_i} |W_{P_{\Pi_n}}f(x) - P_E f(x)|^2 dx.$$

If  $x = (x_1, ..., x_N) \in Q_i$ , then  $[(n+1)x_1] = i_1, ..., [(n+1)x_N] = i_N$  and hence

$$(W_{P_{\Pi_n}}f)(x)$$

$$= (n+1)^N \int_{(0,1)^N} (P_{\Pi_n})_{[(n+1)x_1],\dots,[(n+1)x_N],[(n+1)y_1],\dots,[(n+1)y_N]} \times f(y) dy$$

$$= (n+1)^N \int_{(0,1)^N} (P_{\Pi_n})_{i_1,\dots,i_N,[(n+1)y_1],\dots,[(n+1)y_N]} f(y) dy.$$

By virtue of (25),

$$(W_{P_{\Pi_n}}f)(x) = (n+1)^N \int_{O_i} (P_{\Pi_n})_{i_1,\dots,i_N,i_1,\dots,i_N} f(y) dy.$$

If  $(i_1/(n+1),\ldots,i_N/(n+1)) \notin E$  then  $(i_1/n,\ldots,i_N/n) \notin E$ . Thus, in this case we have  $i=(i_1,\ldots,i_N) \notin \Pi_n$  and  $Q_i \subset [0,1]^N \setminus E$ . Consequently,  $(P_{\Pi_n})_{i,i}=0$  and  $(P_Ef)(x)=0$  for  $x\in Q_i$ , which implies that

$$\int_{Q_i} |W_{P_{\Pi_n}} f(x) - P_E f(x)|^2 dx = 0.$$

Suppose  $((i_1+1)/n, ..., (i_N+1)/n) \in E$ . Then  $i = (i_1, ..., i_N) \in \Pi_n$  and since  $((i_1+1)/(n+1), ..., (i_N+1)/(n+1)) \in E$ , it follows that  $Q_i \subset E$ . Thus,  $(P_{\Pi_n})_{i,i} = 1$  and  $(P_E f)(x) = f(x)$  for  $x \in Q_i$ . This

gives

$$\int_{Q_{i}} |(W_{P_{\Pi_{n}}}f)(x) - (P_{E}f)(x)|^{2} dx 
\leq \frac{1}{(n+1)^{N}} \sup_{x \in Q_{i}} |(W_{P_{\Pi_{n}}}f)(x) - (P_{E}f)(x)|^{2} 
= \frac{1}{(n+1)^{N}} \sup_{x \in Q_{i}} |(n+1)^{N} \int_{Q_{i}} (f(y) - f(x)) dy|^{2} 
\leq \frac{1}{(n+1)^{N}} \sup_{x,y \in Q_{i}} |f(y) - f(x)|^{2} \leq \frac{\varepsilon}{2 n^{N}}$$

if only  $n \geq n_1 = n_1(\varepsilon)$ . It follows that the sum of the terms

$$\int_{Q_i} |W_{P_{\Pi_n}} f(x) - P_E f(x)|^2 dx$$

over  $((i_1+1)/n, \ldots, (i_N+1)/n) \in E$  is at most  $n^N(\varepsilon/(2n^N)) = \varepsilon/2$  for all  $n \geq n_1$ . We are left with the case where

$$\left(\frac{i_1}{n+1},\ldots,\frac{i_N}{n+1}\right) \in E, \quad \left(\frac{i_1+1}{n},\ldots,\frac{i_N+1}{n}\right) \notin E.$$

These points  $(i_1, \ldots, i_N)$  are all in a small shell around the boundary of nE and hence their number is  $O(n^{N-1})$ . Summing up over these points we get

$$\sum_{i} \int_{Q_{i}} |(W_{P_{\pi_{n}}} f)(x) - (P_{E} f)(x)|^{2} dx$$

$$\leq \sum_{i} \left( 2 \max_{x \in [0,1]^{N}} |f(x)| \right)^{2} \frac{1}{(n+1)^{N}}$$

$$= O(n^{N-1}) \left( 2 \max_{x \in [0,1]^{N}} |f(x)| \right)^{2} \frac{1}{(n+1)^{N}},$$

which is smaller than  $\varepsilon/2$  if  $n \geq n_2 = n_2(\varepsilon)$ . In summary, (26) is smaller that  $\varepsilon$  if  $n \geq \max(n_1, n_2)$ .

We are now in a position to prove (11) and (12) of Theorem 1.1. It is well known that if K is a compact operator,  $||K_n - K||_{\infty} \to 0$ ,

and  $C_n^* \to C^*$  strongly (the asterisk denoting the adjoint), then  $||K_nC_n - KC||_{\infty} \to 0$  (see, e.g., [5, Lemma 2.8]). Put

$$K_n = (n+1)^{-\sigma|\nu|} W_{D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N}},$$

$$K = T_{\nu_1,\alpha_1} \otimes \cdots \otimes T_{\nu_N,\alpha_N}$$

$$C_n = C_n^* = W_{P_{\Pi_n}}, \quad C = C^* = P_E.$$

It is easily seen that all  $T_{\nu_j,\alpha_j}$  are Hilbert-Schmidt operators. (Here we are using our assumption that  $\nu_j \geq 1$  for all j.) This implies that  $T_{\nu_1,\alpha_1} \otimes \cdots \otimes T_{\nu_N,\alpha_N}$  is also Hilbert-Schmidt and thus compact. From (24) we know that  $||K_n - K||_{\infty} \to 0$ , and Lemma 4.3 states that  $C_n^* \to C^*$  strongly. Consequently,

$$\begin{aligned} \|(n+1)^{-\sigma|\nu|} W_{D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N}} W_{P_{\Pi_n}} \\ &- (T_{\nu_1,\alpha_1} \otimes \cdots \otimes T_{\nu_N,\alpha_N}) P_E \|_{\infty} \to 0. \end{aligned}$$

This yields that

$$(n+1)^{-\sigma|\nu|} \|W_{D_{\alpha_1,n}^{\nu_1} \otimes \cdots \otimes D_{\alpha_N,n}^{\nu_N}} W_{P_{\Pi_n}}\|_{\infty}$$

$$\longrightarrow \|(T_{\nu_1,\alpha_1} \otimes \cdots \otimes T_{\nu_N,\alpha_N}) P_E\|_{\infty}.$$

From Lemma 4.1 we infer that

$$\|W_{D^{\nu_1}_{\alpha_1,n}\otimes\cdots\otimes D^{\nu_N}_{\alpha_N,n}}W_{P_{\Pi_n}}\|_{\infty}=\|(D^{\nu_1}_{\alpha_1,n}\otimes\cdots\otimes D^{\nu_N}_{\alpha_N,n})P_{\Pi_n}\|_{\infty},$$

and since  $n/(n+1) \to 1$ , Lemma 4.2 gives the desired result.

We are left with the case where some of the numbers  $\nu_1, \ldots, \nu_N$  are zero. We assume without loss of generality that  $\nu_1, \ldots, \nu_k \geq 1$  and  $\nu_{k+1} = \ldots = \nu_N = 0$ . Note that then

$$E_{\nu} = \{(x_1, \dots, x_k) \in [0, 1]^k : (x_1, \dots, x_k, 0, \dots, 0) \in E\}.$$

For the sake of definiteness, we consider the Laguerre case. Let

$$C_n = \lambda(\partial_1^{\nu_1} \cdots \partial_k^{\nu_k} \mid \mathcal{P}_n(E_\nu))$$

be the norm of  $\partial_1^{\nu_1} \dots \partial_k^{\nu_k}$  on  $\mathcal{P}_n(E_{\nu})$ . We have already proved that

(27) 
$$C_n \sim n^{|\nu|} \|L_{\nu_1,\alpha_1} \otimes \cdots \otimes L_{\nu_k,\alpha_k} |L^2(E_{\nu})\|_{\infty}.$$

There exists a polynomial  $g \in \mathcal{P}_n(E_{\nu})$  such that  $\|\partial_1^{\nu_1} \dots \partial_k^{\nu_k} g\| = C_n\|g\|$ . Define  $f \in \mathcal{P}_n(E)$  by  $f(t_1, \dots, t_N) = g(t_1, \dots, t_k)$ . Then  $\partial_1^{\nu_1} \dots \partial_N^{\nu_N} f = \partial_1^{\nu_1} \dots \partial_k^{\nu_k} g$  and hence

$$\int_{(0,\infty)^k} |(\partial_1^{\nu_1} \cdots \partial_N^{\nu_N} f)(t_1, \dots, t_N)|^2 t_1^{\alpha_1} \cdots t_k^{\alpha_k} e^{-t_1} \cdots e^{-t_k} dt_1 \cdots dt_k$$

$$= C_n^2 \int_{(0,\infty)^k} |f(t_1, \dots, t_N)|^2 t_1^{\alpha_1} \cdots t_k^{\alpha_k} e^{-t_1} \cdots e^{-t_k} dt_1 \cdots dt_k$$

for each point  $(t_{k+1}, \ldots, t_N) \in (0, \infty)^{N-k}$ . Multiplying this equality by

(28) 
$$t_{k+1}^{\alpha_{k+1}} \cdots t_N^{\alpha_N} e^{-t_{k+1}} \cdots e^{-t_N}$$

and integrating the result over  $(0,\infty)^{N-k}$  we get  $\|\partial_1^{\nu_1}\cdots\partial_N^{\nu_N}f\|^2=C_n^2\|f\|^2$ . Thus,  $\lambda(\partial_1^{\nu_1}\cdots\partial_N^{\nu_N}\mid\mathcal{P}_n(E))\geq C_n$ .

On the other hand, every polynomial  $f \in \mathcal{P}_n(E)$  may be written as

(29) 
$$f(t_1, \dots, t_N) = \sum p_{i_1, \dots, i_k}(t_{k+1}, \dots, t_N) t_1^{i_1} \cdots t_k^{i_k}$$

where

$$p_{i_1,\dots,i_k}(t_{k+1},\dots,t_N) = \sum p_{i_1,\dots,i_k,\ell_{k+1},\dots,\ell_N} t_{k+1}^{\ell_{k+1}} \cdots t_N^{\ell_N}$$

and  $(i_1/n, \ldots, i_k/n, \ell_{k+1}/n, \ldots, \ell_N/n) \in E$ . This implies that  $(i_1/n, \ldots, i_k/n) \in E_{\nu}$  and that hence the polynomial (29) belongs to  $\mathcal{P}_n(E_{\nu})$  for each fixed point  $(t_{k+1}, \ldots, t_N) \in (0, \infty)^{N-k}$ . We obtain that, for fixed  $(t_{k+1}, \ldots, t_N) \in (0, \infty)^{N-k}$ ,

$$\int_{(0,\infty)^k} |(\partial_1^{\nu_1} \cdots \partial_N^{\nu_N} f)(t_1, \dots, t_N)|^2 t_1^{\alpha_1} \cdots t_k^{\alpha_k} e^{-t_1} \cdots e^{-t_k} dt_1 \cdots dt_k 
\leq C_n^2 \int_{(0,\infty)^k} |f(t_1, \dots, t_N)|^2 t_1^{\alpha_1} \cdots t_k^{\alpha_k} e^{-t_1} \cdots e^{-t_k} dt_1 \cdots dt_k,$$

which after multiplication by (28) and integration over  $(0,\infty)^{N-k}$  becomes the inequality  $\|\partial_1^{\nu_1}\cdots\partial_N^{\nu_N}f\|^2 \leq C_n^2\|f\|^2$ . This proves that  $\lambda(\partial_1^{\nu_1}\cdots\partial_N^{\nu_N}\mid \mathcal{P}_n(E))\leq C_n$ .

In summary, we have  $\lambda(\partial_1^{\nu_1}\cdots\partial_N^{\nu_N}\mid\mathcal{P}_n(E))=C_n$ , which in conjunction with (27) completes the proof of Theorem 1.1.

Here is the proof of Theorem 1.2. The operator

$$U: L^2(E^2) \to L^2(E), \quad (Uf)(t_1, \dots, t_k) = 2^{k/2} t_1^{1/2} \cdots t_k^{1/2} f(t_1^2, \dots, t_k^2)$$

is an isometry and the inverse operator acts by the rule

$$(U^{-1}g)(x_1,\ldots,x_k)=2^{-k/2}x_1^{-1/4}\cdots x_k^{-1/4}g(x_1^{1/2},\ldots,x_k^{1/2}).$$

The kernel of the integral operator  $G_{\nu,\alpha} \mid L^2(E)$  is

$$\kappa(x_1, \dots, x_k, y_1, \dots, y_k) = \prod_{i=1}^k \frac{x_j^{1/2 + \alpha_j} y_j^{1/2 - \alpha_j}}{2^{\nu_j - 1} \Gamma(\nu_j)} (y_j^2 - x_j^2)^{\nu_j - 1} \chi(y_j - x_j)$$

where  $\chi(\xi) = 1$  for  $\xi > 0$  and  $\chi(\xi) = 0$  for  $\xi < 0$ . Thus, for  $x = (x_1, \dots, x_k) \in E$ ,

$$(U^{-1}(G_{\nu,\alpha} \mid L^{2}(E))Uf)(x)$$

$$= 2^{-k/2} \prod_{j=1}^{k} x_{j}^{-1/4} \int_{E} \kappa(x_{1}^{1/2}, \dots, x_{k}^{1/2}, t_{1}, \dots, t_{k})(Uf)(t_{1}, \dots, t_{k}) dt$$

and the integral equals

$$\begin{split} \int_{E^2} \kappa(x_1^{1/2}, \dots, x_k^{1/2}, y_1^{1/2}, \dots, y_k^{1/2}) (Uf) (y_1^{1/2}, \dots, y_k^{1/2}) \, \frac{1}{2^k} \\ & \times \prod_{j=1}^k y_j^{-1/2} \, dy \\ &= \int_{E^2} \kappa(x_1^{1/2}, \dots, x_k^{1/2}, y_1^{1/2}, \dots, y_k^{1/2}) \, 2^{k/2} \\ & \times \prod_{j=1}^k y_j^{1/4} f(y_1, \dots, y_k) \, \frac{1}{2^k} \prod_{j=1}^k y_j^{-1/2} \, dy. \end{split}$$

Consequently,  $U^{-1}G_{\nu,\alpha}U$  is the integral operator on  $L^2(E^2)$  with the kernel

$$\begin{split} \frac{1}{2^k} \bigg( \prod_{j=1}^k x_j^{-1/4} y_j^{-1/4} \bigg) \kappa \big( x_1^{1/2}, \dots, x_k^{1/2}, y_1^{1/2}, \dots, y_k^{1/2} \big) \\ &= \frac{1}{2^k} \prod_{j=1}^k \frac{1}{2^{\nu_j - 1} \Gamma(\nu_j)} \, x_j^{-1/4} y_j^{-1/4} x_j^{1/4 + \alpha_j/2} y_j^{1/4 - \alpha_j/2} \\ &\qquad \qquad \times (y_j - x_j)^{\nu_j - 1} \chi \big( y_j^{1/2} - x_j^{1/2} \big) \\ &= \frac{1}{2^{|\nu|}} \prod_{j=1}^k \frac{1}{\Gamma(\nu_j)} \, x_j^{\alpha_j/2} y_j^{-\alpha_j/2} \big( y_j - x_j \big)^{\nu_j - 1} \chi \big( y_j - x_j \big). \end{split}$$

As this is just  $1/2^{|\nu|}$  times the kernel of  $L_{\nu,\alpha}$ , the proof of Theorem 1.2 is complete.

## **5. Bounds.** In this section we prove Theorem 1.3.

Thus, let  $\nu=(\nu_1,\ldots,\nu_k)$  with natural numbers  $\nu_j\geq 1$  and  $\alpha=(\alpha_1,\ldots,\alpha_k)$  with real numbers  $\alpha_j>-1$ . The operator  $L_{\nu,\alpha}$  has the kernel

$$\prod_{j=1}^k rac{x_j^{lpha_j/2} y_j^{-lpha_j/2}}{\Gamma(
u_j)} (y_j - x_j)^{
u_j - 1} \chi(y_j - x_j).$$

The adjoint of the operator  $L_{\nu,\alpha} \mid L^2(E)$  is  $L_{\nu,\alpha}^* \mid L^2(E)$  where  $L_{\nu,\alpha}^*$  is the integral operator with the kernel

$$\prod_{j=1}^k rac{x_j^{-lpha_j/2}y_j^{lpha_j/2}}{\Gamma(
u_j)} (x_j-y_j)^{
u_j-1} \chi(x_j-y_j).$$

Recall that  $\Omega_k^{\delta} = \{x \in [0,1]^k : x_1^{1/\delta} + \dots + x_k^{1/\delta} \leq 1\}$ . In what follows we will make frequent use of Euler's formula

$$\int_0^1 t^{p-1} (1-t)^{q-1} dt = \frac{\Gamma(p)\Gamma(q)}{\Gamma(p+q)}$$

and of Dirichlet's formula

(30) 
$$\int_{\Omega_k^{\delta}} x_1^{p_1-1} \cdots x_k^{p_k-1} dx = \frac{\delta^k \Gamma(\delta p_1) \cdots \Gamma(\delta p_k)}{\Gamma(\delta p_1 + \cdots + \delta p_k + 1)}.$$

The Hilbert-Schmidt norm  $||K|| L^2(E)||_2$  of an integral operator is the  $L^2$  norm of its kernel over  $E \times E$ . We therefore have

(31) 
$$\left( \prod_{j=1}^{k} \Gamma(\nu_{j})^{2} \right) \|L_{\nu,\alpha} | L^{2}(\Omega_{k}^{\delta})\|_{2}^{2}$$

$$= \left( \prod_{j=1}^{k} \Gamma(\nu_{j})^{2} \right) \|L_{\nu,\alpha}^{*} | L^{2}(\Omega_{k}^{\delta})\|_{2}^{2}$$

$$= \int_{\Omega_{k}^{\delta}} \int_{\Omega_{k}^{\delta}} \prod_{j=1}^{k} x_{j}^{-\alpha_{j}} y_{j}^{\alpha_{j}} (x_{j} - y_{j})^{2\nu_{j} - 2} \chi(x_{j} - y_{j}) \, dy \, dx.$$

If  $(x_1,\ldots,x_k)\in\Omega_k^\delta$  and  $y_j< x_j$  for all j, then  $(y_1,\ldots,y_j)$  is automatically in  $\Omega_k^\delta$ . Thus, (31) is

$$\begin{split} \int_{\Omega_k^{\delta}} \bigg( \prod_{j=1}^k x_j^{-\alpha_j} \int_0^{x_j} y_j^{\alpha_j} (x_j - y_j)^{2\nu_j - 2} dy_j \bigg) \, dx \\ &= \int_{\Omega_k^{\delta}} \bigg( \prod_{j=1}^k x_j^{-\alpha_j} \int_0^1 x_j^{\alpha_j} t^{\alpha_j} x_j^{2\nu_j - 2} (1 - t)^{2\nu_j - 2} x_j \, dt \bigg) \, dx \\ &= \bigg( \prod_{j=1}^k \frac{\Gamma(\alpha_j + 1) \Gamma(2\nu_j - 1)}{\Gamma(\alpha_j + 2\nu_j)} \bigg) \int_{\Omega_k^{\delta}} \prod_{j=1}^k x_j^{2\nu_j - 1} \, dx \end{split}$$

and formula (30) now implies the equality asserted in Theorem 1.3.

To prove the inequality stated in Theorem 1.3, note first that the kernel of the integral operator  $L_{\nu,\alpha}^* L_{\nu,\alpha} \mid L^2(\Omega_k^{\delta})$  is

$$\int_{\Omega_k^{\delta}} g(t, x) g(t, y) dt$$

where

$$g(t,z) = \prod_{j=1}^{k} \frac{t_j^{\alpha_j/2} z_j^{-\alpha_j/2}}{\Gamma(\nu_j)} (z_j - t_j)^{\nu_j - 1} \chi(z_j - t_j).$$

Consequently,

$$(32) \quad \|L_{\nu,\alpha}^* L_{\nu,\alpha} \mid L^2(\Omega_k^{\delta})\|_2^2 = \int_{\Omega_k^{\delta}} \int_{\Omega_k^{\delta}} \left( \int_{\Omega_k^{\delta}} g(t,x) g(t,y) \, dt \right)^2 dx \, dy.$$

The inner integral in (32) is

$$\begin{split} I_1 &:= \bigg(\prod_{j=1}^k \frac{x_j^{-\alpha_j/2} y_j^{-\alpha_j/2}}{\Gamma(\nu_j)^2}\bigg) \\ &\times \int_{\Omega_k^\delta} \bigg(\prod_{j=1}^k t_j^{\alpha_j} (y_j - t_j)^{\nu_j - 1} (x_j - t_j)^{\nu_j - 1} \chi(y_j - t_j) \chi(x_j - t_j)\bigg) \, dt. \end{split}$$

Put  $u_j = \min(x_j, y_j)$ . If x and y are in  $\Omega_k^{\delta}$ , then  $(u_1, \ldots, u_k)$  is also in  $\Omega_k^{\delta}$ . Therefore the integral in (33) is

$$I_2 := \prod_{j=1}^k \int_0^{u_j} t_j^{lpha_j} (y_j - t_j)^{
u_j - 1} (x_j - t_j)^{
u_j - 1} dt_j.$$

If  $y_i \leq x_j$  then  $u_j = y_j$  and

$$(34) \int_{0}^{u_{j}} t_{j}^{\alpha_{j}} (y_{j} - t_{j})^{\nu_{j}-1} (x_{j} - t_{j})^{\nu_{j}-1} dt_{j}$$

$$= \int_{0}^{y_{j}} t_{j}^{\alpha_{j}} (y_{j} - t_{j})^{\nu_{j}-1} (x_{j} - t_{j})^{\nu_{j}-1} dt_{j}$$

$$= \int_{0}^{1} y_{j}^{\alpha_{j}} s^{\alpha_{j}} y_{j}^{\nu_{j}-1} (1 - s)^{\nu_{j}-1} (x_{j} - y_{j}s)^{\nu_{j}-1} y_{j} ds.$$

Since  $x_j - y_j s \ge x_j - x_j s$  for  $0 \le y_j \le x_j$  and  $0 \le s \le 1$ , integral (34) is at least

(35) 
$$\int_{0}^{1} y_{j}^{\alpha_{j}+\nu_{j}} x_{j}^{\nu_{j}-1} s^{\alpha_{j}} (1-s)^{2\nu_{j}-2} ds$$

$$= \frac{\Gamma(\alpha_{j}+1)\Gamma(2\nu_{j}-1)}{\Gamma(\alpha_{j}+2\nu_{j})} y_{j}^{\alpha_{j}+\nu_{j}} x_{j}^{\nu_{j}-1}$$

$$= \frac{\Gamma(\alpha_{j}+1)\Gamma(2\nu_{j}-1)}{\Gamma(\alpha_{j}+2\nu_{j})} u_{j}^{\alpha_{j}+\nu_{j}} w_{j}^{\nu_{j}-1},$$

where  $w_j = \max(x_j, y_j)$ . The integral  $I_2$  is symmetric in x and y and hence (35) will also result in the case  $y_j > x_j$ . Thus,

$$I_2 \ge \Gamma_{\nu,\alpha} \prod_{j=1}^k u_j^{\alpha_j + \nu_j} w_j^{\nu_j - 1},$$

$$\Gamma_{\nu,\alpha} := \prod_{j=1}^k \frac{\Gamma(\alpha_j + 1) \Gamma(2\nu_j - 1)}{\Gamma(\alpha_j + 2\nu_j)}.$$

Letting  $C_{\nu} = \prod_{j=1}^{k} (1/\Gamma(\nu_j))$ , we obtain that

$$\begin{split} I_1^2 &\geq C_{\nu}^4 \Gamma_{\nu,\alpha}^2 \prod_{j=1}^k x_j^{-\alpha_j} y_j^{-\alpha_j} u_j^{2\alpha_j + 2\nu_j} w_j^{2\nu_j - 2} \\ &= C_{\nu}^4 \Gamma_{\nu,\alpha}^2 \prod_{j=1}^k u_j^{\alpha_j + 2\nu_j} w_j^{2\nu_j - \alpha_j - 2}. \end{split}$$

Consequently, a lower bound for (32) is

(36) 
$$C_{\nu}^{4}\Gamma_{\nu,\alpha}^{2} \int_{\Omega_{k}^{\delta}} \int_{\Omega_{k}^{\delta}} \left( \prod_{j=1}^{k} u_{j}^{\alpha_{j}+2\nu_{j}} w_{j}^{2\nu_{j}-\alpha_{j}-2} \right) dx dy.$$

For each j, we have, up to sets of measure zero, the two possibilities  $y_j < x_j$  or  $y_j > x_j$ . This gives  $2^k$  possibilities for all j. Accordingly, we may partition  $\Omega_k^{\delta} \times \Omega_k^{\delta}$  into  $2^k$  domains  $B_i$   $(i = 1, \ldots, 2^k)$ . For a fixed j, we have either  $u_j = x_j$  and  $w_j = y_j$  or  $u_j = y_j$  and  $w_j = x_j$  throughout each domain  $B_i$ . Thus,

$$dx_1 \cdots dx_k dy_1 \cdots dy_k = du_1 \cdots du_k dw_1 \cdots dw_k$$

in each  $B_i$ . This implies that each  $B_i$  makes the same contribution to (36) and that hence (36) equals

(37) 
$$2^k C_{\nu}^4 \Gamma_{\nu,\alpha}^2 \int_{B_1} \left( \prod_{j=1}^k u_j^{\alpha_j + 2\nu_j} w_j^{2\nu_j - \alpha_j - 2} \right) du \, dw$$

where  $B_1 = \{(u, w) \in \Omega_k^{\delta} \times \Omega_k^{\delta} : u_j < w_j \text{ for all } j\}$ . If  $w \in \Omega_k^{\delta}$  and  $u_j < w_j$  for all j, then u is automatically in  $\Omega_k^{\delta}$ . This shows that (37) is equal to

$$\begin{split} 2^k C_{\nu}^4 \Gamma_{\nu,\alpha}^2 \int_{\Omega_k^{\delta}} \left( \prod_{j=1}^k \int_0^{w_j} u_j^{\alpha_j + 2\nu_j} w_j^{2\nu_j - \alpha_j - 2} du_j \right) dw \\ &= 2^k C_{\nu}^4 \Gamma_{\nu,\alpha}^2 \int_{\Omega_k^{\delta}} \left( \prod_{j=1}^k \frac{w_j^{4\nu_j - 1}}{\alpha_j + 2\nu_j + 1} \right) dw. \end{split}$$

From (30) we infer that

$$\int_{\Omega_k^{\delta}} \prod_{i=1}^k w_j^{4\nu_j - 1} dw = \frac{\delta^k}{\Gamma(4\delta|\nu| + 1)} \prod_{i=1}^k \Gamma(4\delta\nu_i).$$

Putting things together we arrive at the inequality in Theorem 1.3.

Here is, for the sake of completeness, a proof of (14). The squares of the bounds  $b_1 := b_1(\nu, \alpha, \Omega_k^{\delta})$  and  $b_2 := b_2(\nu, \alpha, \Omega_k^{\delta})$  are

$$b_2^2 = \frac{\delta^k}{\Gamma(2\delta|\nu|+1)} \prod_{j=1}^k \frac{\Gamma(\alpha_j+1)\Gamma(2\nu_j-1)\Gamma(2\delta\nu_j)}{\Gamma(\alpha_j+2\nu_j)\Gamma(\nu_j)^2},$$

$$b_1^2 = \frac{2^k\Gamma(2\delta|\nu|+1)}{\Gamma(4\delta|\nu|+1)} \prod_{j=1}^k \frac{\Gamma(\alpha_j+1)\Gamma(2\nu_j-1)\Gamma(4\delta\nu_j)}{\Gamma(\alpha_j+2\nu_j)(\alpha_j+2\nu_j+1)\Gamma(\nu_j)^2\Gamma(2\delta\nu_j)},$$

and the quotient of these two bounds is

(38) 
$$\frac{b_2^2}{b_1^2} = \left(\frac{\delta}{2}\right)^k \frac{\Gamma(4\delta|\nu|+1)}{\Gamma(2\delta|\nu|+1)^2} \prod_{j=1}^k (\alpha_j + 2\nu_j + 1) \frac{\Gamma(2\delta\nu_j)^2}{\Gamma(4\delta\nu_j)}.$$

Suppose  $k, \delta, \alpha_1, \ldots, \alpha_k$  remain fixed and  $\nu_j \to \infty$  for all j. Taking into account that

$$\frac{\Gamma(2\mu)}{\Gamma(\mu)^2} \sim 2^{2\mu} \, \sqrt{\frac{\mu}{4\pi}} \quad \text{as } \mu \to \infty,$$

we see that (38) is

$$\frac{\delta^{k}}{2^{k}} \frac{4\delta|\nu| \Gamma(4\delta|\nu|)}{4\delta^{2}|\nu|^{2} \Gamma(2\delta|\nu|)^{2}} \prod_{j=1}^{k} (\alpha_{j} + 2\nu_{j} + 1) \frac{\Gamma(2\delta\nu_{j})^{2}}{\Gamma(4\delta\nu_{j})}$$

$$\sim \frac{\delta^{k}}{2^{k}} \frac{1}{\delta|\nu|} 2^{4\delta|\nu|} \left(\frac{2\delta|\nu|}{4\pi}\right)^{1/2} 2^{k} \prod_{j=1}^{k} \nu_{j} \frac{1}{2^{4\delta\nu_{j}}} \left(\frac{4\pi}{2\delta\nu_{j}}\right)^{1/2}$$

$$= (2\pi\delta)^{(k-1)/2} \frac{1}{|\nu|^{1/2}} \prod_{j=1}^{k} \nu_{j}^{1/2},$$

which is the same as (14).

## **6. Asymptotics.** In this section we present a proof of Theorem 1.4.

Let  $\varrho_1, \ldots, \varrho_k$  be positive real numbers. For a real number  $\tau > 0$ , we consider the integral operator  $L_{\tau}$  with the kernel

$$g_{\tau}(x,y) = \prod_{j=1}^{k} x_{j}^{-\alpha_{j}/2} y_{j}^{\alpha_{j}/2} (x_{j} - y_{j})^{\varrho_{j}\tau - 1} \chi(x_{j} - y_{j})$$

on  $L^2(\Omega_k^{\delta})$ . Clearly,  $L_{\tau}$  is nothing but

$$\bigg(\prod_{j=1}^k \Gamma(\varrho_j\tau)\bigg)(L_{\varrho_1\tau,\alpha_1}^*\otimes\cdots\otimes L_{\varrho_k\tau,\alpha_k}^*)\mid L^2(\Omega_k^\delta).$$

The function  $\prod_{j=1}^k x_j^{\varrho_j}$  attains its maximum on  $\Omega_k^{\delta}$  at the point  $(p_1, \ldots, p_k)$  and nowhere else. This point lies on the boundary of  $\Omega_k^{\delta}$  and is given by

$$p_j = \left(\frac{\varrho_j}{|\varrho|}\right)^{\delta}, \quad |\varrho| := \varrho_1 + \dots + \varrho_k.$$

We denote by  $H_{\tau}$  the integral operator on  $L^{2}(\Omega_{k}^{\delta})$  whose kernel is

$$h_{\tau}(x,y) = \prod_{j=1}^{k} x_{j}^{-\alpha_{j}/2} y_{j}^{\alpha_{j}/2} x_{j}^{\varrho_{j}\tau-1} \left(1 - \frac{y_{j}}{p_{j}}\right)^{\varrho_{j}\tau-1} \chi(p_{j} - y_{j}).$$

As  $h_{\tau}(x,y)$  is of the form a(x)b(y), the operator norm of  $H_{\tau}$  is the product of the  $L^2$  norms of a and b. Thus, using the Dirichlet and Euler formulas,

$$(39) \quad \|H_{\tau}\|_{\infty}^{2} = \int_{\Omega_{k}^{\delta}} \prod_{j=1}^{k} x_{j}^{2\varrho_{j}\tau - \alpha_{j} - 2} dx \prod_{j=1}^{k} \int_{0}^{p_{j}} y_{j}^{\alpha_{j}} \left(1 - \frac{y_{j}}{p_{j}}\right)^{2\varrho_{j}\tau - 2} dy_{j}$$

$$= \frac{\delta^{k} \prod_{j=1}^{k} \Gamma(2\delta\varrho_{j}\tau - (\alpha_{j} + 1)\delta)}{\Gamma(2\delta|\varrho|\tau - \sum(\alpha_{j} + 1)\delta + 1)}$$

$$\times \prod_{j=1}^{k} p_{j}^{\alpha_{j} + 1} \frac{\Gamma(\alpha_{j} + 1)\Gamma(2\varrho_{j}\tau - 1)}{\Gamma(\alpha_{j} + 2\varrho_{j}\tau)},$$

the sum over j = 1, ..., k. From Theorem 1.3 we infer that

$$||L_{\tau}||_{2}^{2} = \frac{\delta^{k}}{\Gamma(2\delta|\varrho|\tau+1)} \prod_{j=1}^{k} \frac{\Gamma(\alpha_{j}+1)\Gamma(2\varrho_{j}\tau-1)\Gamma(2\delta\varrho_{j}\tau)}{\Gamma(\alpha_{j}+2\varrho_{j}\tau)}.$$

Consequently, by Stirling's formula,

$$\begin{split} \frac{\|H_{\tau}\|_{\infty}^{2}}{\|L_{\tau}\|_{2}^{2}} &= \frac{\Gamma(2\delta|\varrho|\tau+1)}{\Gamma(2\delta|\varrho|\tau-\sum(\alpha_{j}+1)\delta+1)} \\ &\times \prod_{j=1}^{k} p_{j}^{\alpha_{j}+1} \frac{\Gamma(2\delta\varrho_{j}\tau-(\alpha_{j}+1)\delta)}{\Gamma(2\delta\varrho_{j}\tau)} \\ &\sim (2\delta|\varrho|\tau)^{\sum(\alpha_{j}+1)\delta} \prod_{j=1}^{k} p_{j}^{\alpha_{j}+1} (2\delta\varrho_{j}\tau)^{-(\alpha_{j}+1)\delta} \\ &= |\varrho|^{\sum(\alpha_{j}+1)\delta} \prod_{j=1}^{k} p_{j}^{\alpha_{j}+1} \varrho_{j}^{-(\alpha_{j}+1)\delta} \\ &= |\varrho|^{\sum(\alpha_{j}+1)\delta} \prod_{j=1}^{k} \left(\frac{\varrho_{j}}{|\varrho|}\right)^{\delta(\alpha_{j}+1)} \varrho_{j}^{-(\alpha_{j}+1)\delta} = 1, \end{split}$$

that is,  $||H_{\tau}||_{\infty} = ||L_{\tau}||_{2}(1 + o(1))$ . Now suppose we had shown that

$$||L_{\tau} - H_{\tau}||_{2}^{2} = o(||H_{\tau}||_{\infty}^{2}).$$

It would follow that  $||L_{\tau} - H_{\tau}||_{\infty} = o(||H_{\tau}||_{\infty})$  and hence

$$||L_{\tau}||_{\infty} = ||H_{\tau}||_{\infty} + o(||H_{\tau}||_{\infty})$$
  
=  $||H_{\tau}||_{\infty} (1 + o(1)) = ||L_{\tau}||_{2} (1 + o(1)),$ 

as desired. Thus, we are left with (40).

We have

$$\begin{split} \|L_{\tau} - H_{\tau}\|_{2}^{2} &= \int_{\Omega_{k}^{\delta}} \int_{\Omega_{k}^{\delta}} (g_{\tau}^{2} - 2 g_{\tau} h_{\tau} + h_{\tau}^{2}) \, dy \, dx \\ &= \|L_{\tau}\|_{2}^{2} - 2 \int_{\Omega_{k}^{\delta}} \int_{\Omega_{k}^{\delta}} g_{\tau} h_{\tau} \, dy \, dx + \|H_{\tau}\|_{2}^{2}. \end{split}$$

We already showed that  $||L_{\tau}||_2^2/||H_{\tau}||_{\infty}^2 \to 1$ , and since the kernel of  $H_{\tau}$  is of the form a(x)b(y) we may conclude that  $||H_{\tau}||_2^2/||H_{\tau}||_{\infty}^2 = 1$ . Consequently, estimate (40) will follow as soon as we have proved that

(41) 
$$\int_{\Omega_h^{\delta}} \int_{\Omega_h^{\delta}} g_{\tau} h_{\tau} \, dy \, dx / \|H_{\tau}\|_{\infty}^2 \longrightarrow 1.$$

The integral in (41) equals

$$\begin{split} & \int_{\Omega_k^\delta} \prod_{j=1}^k x_j^{\varrho_j \tau - \alpha_j - 1} \\ & \times \left( \prod_{i=1}^k \int_0^{p_j} y_j^{\alpha_j} \left( 1 - \frac{y_j}{p_j} \right)^{\varrho_j \tau - 1} \chi(x_j - y_j) (x_j - y_j)^{\varrho_j \tau - 1} \, dy_j \right) dx. \end{split}$$

To tackle the inner integrals, we use that if  $\alpha > -1$  and  $0 \le \beta \le 1$ , then

$$\int_0^1 t^{\alpha} (1 - \beta t)^{\lambda} (1 - t)^{\lambda} dt = \frac{\Gamma(\alpha + 1)}{(1 + \beta)^{\alpha + 1}} \frac{1}{\lambda^{\alpha + 1}} (1 + o(1))$$

as  $\lambda \to \infty$ , the o(1) being unform in  $\beta \in [0,1]$ ; see, for example, [11, Section 2.4]. If  $x_j < p_j$ , the jth inner integral is

$$\int_{0}^{x_{j}} y_{j}^{\alpha_{j}} \left(1 - \frac{y_{j}}{p_{j}}\right)^{\varrho_{j}\tau - 1} (x_{j} - y_{j})^{\varrho_{j}\tau - 1} dy_{j}$$

$$= x_{j}^{\alpha_{j} + \varrho_{j}\tau} \int_{0}^{1} t^{\alpha_{j}} \left(1 - \frac{x_{j}t}{p_{j}}\right)^{\varrho_{j}\tau - 1} (1 - t)^{\varrho_{j}\tau - 1} dt$$

$$= x_{j}^{\alpha_{j} + \varrho_{j}\tau} \frac{\Gamma(\alpha_{j} + 1)}{(1 + x_{j}/p_{j})^{\alpha_{j} + 1}} \frac{1}{(\varrho_{j}\tau)^{\alpha_{j} + 1}} (1 + o(1)),$$

while for  $x_j > p_j$  it is

$$\begin{split} \int_{0}^{p_{j}} y_{j}^{\alpha_{j}} \left(1 - \frac{y_{j}}{p_{j}}\right)^{\varrho_{j}\tau - 1} (x_{j} - y_{j})^{\varrho_{j}\tau - 1} dy_{j} \\ &= x_{j}^{\varrho_{j}\tau - 1} p_{j}^{\alpha_{j} + 1} \int_{0}^{1} t^{\alpha_{j}} (1 - t)^{\varrho_{j}\tau - 1} \left(1 - \frac{p_{j}t}{x_{j}}\right)^{\varrho_{j}\tau - 1} dt \\ &= x_{j}^{\varrho_{j}\tau - 1} p_{j}^{\alpha_{j} + 1} \frac{\Gamma(\alpha_{j} + 1)}{(1 + p_{j}/x_{j})^{\alpha_{j} + 1}} \frac{1}{(\varrho_{j}\tau)^{\alpha_{j} + 1}} (1 + o(1)) \\ &= x_{j}^{\alpha_{j} + \varrho_{j}\tau} \frac{\Gamma(\alpha_{j} + 1)}{(1 + x_{j}/p_{j})^{\alpha_{j} + 1}} \frac{1}{(\varrho_{j}\tau)^{\alpha_{j} + 1}} (1 + o(1)). \end{split}$$

Thus, the integral in (41) is asymptotically equal to

(42) 
$$\prod_{j=1}^{k} \frac{\Gamma(\alpha_{j}+1)}{(\varrho_{j}\tau)^{\alpha_{j}+1}} \int_{\Omega_{k}^{\delta}} \prod_{j=1}^{k} \frac{x_{j}^{2\varrho_{j}\tau-1}}{(1+x_{j}/p_{j})^{\alpha_{j}+1}} dx.$$

Taking into account that  $\prod x_j^{\varrho_j}$  has its maximum at  $(p_1, \ldots, p_k)$  one can employ standard methods, for example, such as in [13, Section II.4], to show that the integral in (42) is asymptotically equal to

$$(43) \qquad \int_{\Omega_k^{\delta}} \prod_{j=1}^k \frac{x_j^{2\varrho_j\tau-1}}{(1+p_j/p_j)^{\alpha_j+1}} \, dx = \prod_{j=1}^k \frac{1}{2^{\alpha_j+1}} \int_{\Omega_k^{\delta}} \prod_{j=1}^k x_j^{2\varrho_j\tau-1} dx.$$

By Dirichlet's formula (30), the integral on the right of (43) is

$$\frac{\delta^k}{\Gamma(2\delta|\varrho|\tau+1)} \prod_{j=1}^k \Gamma(2\delta\varrho_j\tau).$$

In summary, the integral in (41) equals

$$\frac{\delta^k}{\Gamma(2\delta|\varrho|\tau+1)} \prod_{j=1}^k \frac{\Gamma(\alpha_j+1)\Gamma(2\delta\varrho_j\tau)}{2^{\alpha_j+1}(\varrho_j\tau)^{\alpha_j+1}} (1+o(1)).$$

This in conjunction with (39) gives that the left-hand side of (41) is asymptotically equal to

$$\begin{split} \frac{\Gamma(2\delta|\varrho|\tau-\sum(\alpha_{j}+1)\delta+1)}{\Gamma(2\delta|\varrho|\tau+1)} \\ &\times \prod_{j=1}^{k} \frac{\Gamma(2\delta\varrho_{j}\tau)\Gamma(\alpha_{j}+2\varrho_{j}\tau)}{(2p_{j}\varrho_{j}\tau)^{\alpha_{j}+1}\Gamma(2\delta\varrho_{j}\tau-(\alpha_{j}+1)\delta)\Gamma(2\varrho_{j}\tau-1)} \\ &\sim \frac{1}{(2\delta|\varrho|\tau)^{\sum(\alpha_{j}+1)\delta}} \prod_{j=1}^{k} \frac{(2\delta\varrho_{j}\tau)^{(\alpha_{j}+1)\delta}(2\varrho_{j}\tau)^{\alpha_{j}+1}}{(2p_{j}\varrho_{j}\tau)^{\alpha_{j}+1}} \\ &= \prod_{j=1}^{k} \frac{1}{p_{j}^{\alpha_{j}+1}} \left(\frac{\varrho_{j}}{|\varrho|}\right)^{(\alpha_{j}+1)\delta} = \prod_{j=1}^{k} \frac{1}{p_{j}^{\alpha_{j}+1}} p_{j}^{\alpha_{j}+1} = 1, \end{split}$$

which completes the proof.

#### 7. Linear combinations. Here is the proof of Theorem 1.5.

Let  $m \geq 1$  be an integer and  $\alpha > -1$  a real number. We put  $T_{m,\alpha} = L_{m,\alpha}$  in the Laguerre case and  $T_{m,\alpha} = G_{m,\alpha}$  in the Gegenbauer case. Let  $M_{(2x)^{m/2}}$  be the operator of multiplication by  $(2x)^{m/2}$  on  $L^2(0,1)$ , and put  $T_{m,\alpha} = M_{(2x)^{m/2}}$  in the Hermite case; clearly, in that case  $T_{m,\alpha}$  does actually not depend on  $\alpha$ . In either case, we let  $T_{0,\alpha}$  be the identity operator. We finally use the abbreviation

(44) 
$$C_n(p) = C(p(\partial_1, \dots, \partial_N) \mid \mathcal{P}_n(E)).$$

In Section 4 we introduced the matrix representations  $D_{\alpha,n}$  of the operator of differentiation on  $\mathcal{P}_n$  in the Laguerre and Gegenbauer cases. We use the notation  $D_{\alpha,n}$  also for the matrix representation of the operator  $D: \mathcal{P}_n \to \mathcal{P}_n$  in the Hermite basis (where in fact there is no dependence on  $\alpha$ ).

## Lemma 7.1. We have

$$(n+1)^{-\sigma|\nu|}W_{D_{\alpha_1,n}^{\nu_1}}\otimes\cdots\otimes W_{D_{\alpha_N,n}^{\nu_N}}\longrightarrow T_{\nu_1,\alpha_1}\otimes\cdots\otimes T_{\nu_N,\alpha_N}$$

strongly on  $L^2((0,1)^N)$  as  $n \to \infty$ .

*Proof.* Let  $I_n$  be the  $(n+1)\times (n+1)$  identity matrix. It is easily seen that  $W_{I_n}\to I$  strongly. In the Laguerre and Gegenbauer cases, we know from (23) that  $(n+1)^{-\sigma\,m}W_{D^m_{\alpha,n}}\to T_{m,\alpha}$  in the norm provided  $m\geq 1$ . This implies the lemma in these two cases. Let us consider the Hermite case. For  $x\in Q_j:=[j/(n+1),(j+1)/(n+1))$ , we obtain from (20) that

$$\begin{split} (W_{D_{\alpha,n}^m} f)(x) &= (n+1) \sum_{i=0}^n \int_{Q_i} (D_{\alpha,n}^m)_{j,i} f(y) \, dy \\ &= (n+1) \int_{Q_{j+m}} (D_{\alpha,n}^m)_{j,j+m} f(y) \, dy \\ &= (n+1) \int_{Q_{j+m}} 2^{m/2} \sqrt{\frac{\Gamma(j+m+1)}{\Gamma(j+1)}} \, f(y) \, dy \\ &= (n+1) \int_{Q_{j+m}} 2^{m/2} \sqrt{(j+m)\cdots(j+1)} \, f(y) \, dy, \end{split}$$

provided  $j + m \leq n$ . Thus,

$$(n+1)^{-m/2}(W_{D_{\alpha,n}^m}f)(x)$$

$$= (n+1) \cdot 2^{m/2} \int_{Q_{j+m}} \sqrt{\left(\frac{j}{n+1} + \frac{m}{n+1}\right) \cdots \left(\frac{j}{n+1} + \frac{1}{n+1}\right)} f(y) dy.$$

If f is in C[0, 1] and n is large, then the right-hand side of this equality is approximately equal to  $2^{m/2}x^{m/2}f(x)$ . It is not difficult to make this precise and to show that

$$\|(n+1)^{-m/2}W_{D_{\alpha}^{m}}f-M_{(2x)^{m/2}}f\|\longrightarrow 0$$

for every  $f \in C[0,1]$ . Since, by Lemma 4.1 and the result known for the one-dimensional case,

$$\|(n+1)^{-m/2}W_{D_{\alpha,n}^m}\|_{\infty} = (n+1)^{-m/2}\|D_{\alpha,n}^m\|_{\infty}$$
$$= (n+1)^{-m/2}\eta_n(D^m) \sim 2^{m/2},$$

it follows that  $(n+1)^{-m/2}W_{D^m_{\alpha,n}}\to M_{(2x)^{m/2}}$  strongly on  $L^2(0,1)$ . Tensoring we get the assertion in the Hermite case.

Since  $(n+1)/n \to 1$ , Lemmas 4.1, 4.3 and 7.1 give that

$$\sum_{\nu_1 + \dots + \nu_N \leq M} n^{-\sigma|\nu|} p_{\nu_1, \dots, \nu_N} W_{(D_{\alpha_1, n}^{\nu_1} \otimes \dots \otimes D_{\alpha_N, n}^{\nu_N}) P_{\Pi_n}}$$

$$\longrightarrow \sum_{\nu_1 + \dots + \nu_N \leq M} p_{\nu_1, \dots, \nu_N} (T_{\nu_1, \alpha_1} \otimes \dots \otimes T_{\nu_N, \alpha_N}) P_E$$

strongly, which in turn implies that

$$(45) \quad n^{-\sigma M} \sum_{\nu_1 + \dots + \nu_N = M} p_{\nu_1, \dots, \nu_N} W_{(D_{\alpha_1, n}^{\nu_1} \otimes \dots \otimes D_{\alpha_N, n}^{\nu_N}) P_{\Pi_n}}$$

$$\longrightarrow \sum_{\nu_1 + \dots + \nu_N = M} p_{\nu_1, \dots, \nu_N} (T_{\nu_1, \alpha_1} \otimes \dots \otimes T_{\nu_N, \alpha_N}) P_E$$

strongly. From the Banach-Steinhaus theorem we therefore obtain that

$$\lim_{n \to \infty} \inf n^{-\sigma M} \left\| \sum_{\nu_1 + \dots + \nu_N = M} p_{\nu_1, \dots, \nu_N} W_{(D_{\alpha_1, n}^{\nu_1} \otimes \dots \otimes D_{\alpha_N, n}^{\nu_N}) P_{\Pi_n}} \right\|_{\infty} \\
\ge \left\| \sum_{\nu_1 + \dots + \nu_N = M} p_{\nu_1, \dots, \nu_N} (T_{\nu_1, \alpha_1} \otimes \dots \otimes T_{\nu_N, \alpha_N}) P_E \right\|_{\infty}.$$

The right-hand side of this inequality is strictly positive and the left-hand side is just  $\liminf n^{-\sigma M} C_n(p_0)$  due to Lemma 4.1. Hence

$$\liminf n^{-\sigma M} C_n(p_0) > 0.$$

**Lemma 7.2.** If p and q are any two polynomials, then  $C_n(p+q) \leq C_n(p) + C_n(q)$ .

*Proof.* Obvious.  $\Box$ 

Lemma 7.2 in conjunction with Theorem 1.1 shows that

$$\lim\sup n^{-\sigma\,M}C_n(p_0)<\infty.$$

Thus, at this point we have proved (18). From Lemma 7.2 we also get

$$C_{n}(p) \leq C_{n}(p_{0}) + \sum_{\nu_{1} + \dots + \nu_{N} \leq M-1} |p_{\nu_{1},\dots,\nu_{N}}| C(\partial_{1}^{\nu_{1}} \dots \partial_{N}^{\nu_{n}} | \mathcal{P}_{n}(E)),$$

$$C_{n}(p_{0}) \leq C_{n}(p) + \sum_{\nu_{1} + \dots + \nu_{N} \leq M-1} |p_{\nu_{1},\dots,\nu_{N}}| C(\partial_{1}^{\nu_{1}} \dots \partial_{N}^{\nu_{n}} | \mathcal{P}_{n}(E)),$$

and Theorem 1.1 therefore yields that  $C_n(p) = C_n(p_0) + O(n^{\sigma(M-1)})$ , which together with (18) gives (19) and thus completes the proof of Theorem 1.5.

Remark 7.3. The arguments used in this section reveal the difference between the Hermite case on the one hand and the Laguerre and Gegenbauer cases on the other. In contrast to the Laguerre and Gegenbauer cases, the limiting operators in the Hermite case are no longer (compact) integral operators, but multiplication operators, and secondly, in the Hermite case the convergence is no longer uniform, but only strong. In this light it comes as a fortune that the Hermite case can be disposed of by the simple reasoning presented in Section 3.

## 8. Open problems.

**Problem 8.1.** Let N=2 and consider the operator

$$p(\partial_1, \partial_2) = p_{30}\partial_1^3 + p_{21}\partial_1^2\partial_2 + p_{12}\partial_1\partial_2^2 + p_{03}\partial_2^3 + \sum_{\nu_1 + \nu_2 < 2} p_{\nu_1 \nu_2} \partial_1^{\nu_1} \partial_2^{\nu_2}.$$

The principal part is  $p_0(\partial_1, \partial_2) = p_{30}\partial_1^3 + p_{21}\partial_1^2\partial_2 + p_{12}\partial_1\partial_2^2 + p_{03}\partial_2^3$ . Theorem 1.5 tells us that, with abbreviation (44),

$$C_n(p) \sim C_n(p_0) \simeq n^{3\sigma}$$

and as long as one of the terms  $p_{30}\partial_1^3$  and  $p_{03}\partial_2^3$  is present, we cannot say more. However, if  $p_{30} = p_{03} = 0$  and if we are in the Laguerre or Gegenbauer cases, then the convergence in (45) is uniform because all occurring  $T_{\nu_i,\alpha_i}$  are compact integral operators. Consequently, in these cases there is no need in having recourse to the Banach-Steinhaus theorem, since we can rather conclude straightforwardly that

$$C_n(p_0) \sim n^{3\sigma} \|p_{21}T_{2,\alpha_1} \otimes T_{1,\alpha_2} + p_{12}T_{1,\alpha_1} \otimes T_{2,\alpha_2} \mid L^2(E)\|_{\infty}.$$

However, in the general case we must leave as an open problem the replacement of " $\simeq$ " in (18) by " $\sim$  times a constant."

**Problem 8.2.** The most embarrassing message is that our approach fails for the Laplace operator. This failure is of course connected with Problem 8.1. Let, for example, N=2 and  $\Delta=\partial_1^2+\partial_2^2$ . From Theorem 1.5 we deduce that

$$C(\Delta \mid \mathcal{P}_n(E)) \simeq n^{2\sigma},$$

but we cannot even prove that  $C(\Delta \mid \mathcal{P}_n(E))/n^{2\sigma}$  converges to a limit. On the credit side we have a few modest estimates. From (45) we get

$$(46) \qquad \liminf_{n\to\infty}\frac{C(\Delta\mid\mathcal{P}_n(E))}{n^{2\,\sigma}}\geq \|T_{2,\alpha_1}\otimes I+I\otimes T_{2,\alpha_2}\mid L^2(E)\|_{\infty},$$

while the combination of Theorem 1.1 and Lemma 7.2 yields that (47)

$$\limsup_{n \to \infty} \frac{C(\Delta \mid \mathcal{P}_n(E))}{n^{2\sigma}} \le \|T_{2,\alpha_1} \mid L^2(E_{2,0})\|_{\infty} + \|T_{2,\alpha_2} \mid L^2(E_{0,2})\|_{\infty}.$$

Suppose  $\alpha_1 = \alpha_2 = 0$  and  $E = \Omega_2$ . In that case  $E_{2,0} = E_{0,2} = [0,1]$ .

For the Hermite weight, the right-hand sides of (46) and (47) become

$$||M_{2x_1+2x_2}||L^2(\Omega_2)||_{\infty} = 2$$
 and  $2||M_{2x}||L^2(0,1)||_{\infty} = 4$ ,

respectively, which results in the estimates

$$2 \leq \liminf_{n \to \infty} \frac{\eta(\Delta \mid \mathcal{P}_n(\Omega_2))}{n} \leq \limsup_{n \to \infty} \frac{\eta(\Delta \mid \mathcal{P}_n(\Omega_2))}{n} \leq 4.$$

Let us turn to the Laguerre weight. In that case the right-hand side of (47) is

$$2 \|L_{2,0} \| L^2(0,1)\|_{\infty} = 2 \times 0.284 \dots < 0.569$$

(see [33] or [3]). On the right of (46) we now have the operator norm of the operator  $L_{2,0} \otimes I + I \otimes L_{2,0}$  on  $L^2(\Omega_2)$ . Let f be identically 1 on  $\Omega_2$ . Then  $||f||^2 = 1/2$  and for  $(x_1, x_2) \in \Omega_2$ ,

$$((L_{2,0}^* \otimes I + I \otimes L_{2,0}^*)f)(x_1, x_2) = \int_0^{x_1} (x_1 - y_1)f(y_1, x_2) dy_1$$

$$+ \int_0^{x_2} (x_2 - y_2)f(x_1, y_2) dy_2$$

$$= \int_0^{x_1} (x_1 - y_1) dy_1$$

$$+ \int_0^{x_2} (x_2 - y_2) dy_2 = \frac{x_1^2 + x_2^2}{2}.$$

Thus,

$$\begin{split} \|L_{2,0} \otimes I + I \otimes L_{2,0} \mid L^{2}(\Omega_{2})\|_{\infty}^{2} \\ &= \|L_{2,0}^{*} \otimes I + I \otimes L_{2,0}^{*} \mid L^{2}(\Omega_{2})\|_{\infty}^{2} \\ &\geq 2 \int_{\Omega_{2}} \left(\frac{x_{1}^{2} + x_{2}^{2}}{2}\right)^{2} dx_{1} dx_{2} = \frac{7}{180} > 0.197^{2}. \end{split}$$

It follows that

$$0.197 \leq \liminf_{n \to \infty} \frac{\lambda(\Delta \mid \mathcal{P}_n(\Omega_2))}{n^2} \leq \limsup_{n \to \infty} \frac{\lambda(\Delta \mid \mathcal{P}_n(\Omega_2))}{n^2} < 0.569.$$

In the Legendre case (= Gegenbauer case for  $\alpha_1 = \alpha_2 = 0$ ) we obtain analogously that

$$0.091 \leq \liminf_{n \to \infty} \frac{\gamma(\Delta \mid \mathcal{P}_n(\Omega_2))}{n^4} \leq \limsup_{n \to \infty} \frac{\gamma(\Delta \mid \mathcal{P}_n(\Omega_2))}{n^4} < 0.143.$$

We remark that for the N-dimensional Laplace operator  $\Delta = \partial_1^2 + \cdots + \partial_N^2$  the estimates obtained in this way become worse and worse: in the Hermite case the right-hand sides of (46) and (47) are 2 and 2N, respectively, whereas in the Laguerre and Gegenbauer cases the upper bounds go to infinity and the lower bounds approach zero as  $N \to \infty$ . Thus, in general we cannot answer even the question whether

$$\liminf_{n \to \infty} \frac{C(\Delta \mid \mathcal{P}_n(\Omega_N))}{n^{2\sigma}}, \quad \limsup_{n \to \infty} \frac{C(\Delta \mid \mathcal{P}_n(\Omega_N))}{n^{2\sigma}}$$

converge to zero, increase to infinity, or remain bounded and bounded away from zero as  $N \to \infty$ .

**Problem 8.3.** Consider the wave operator  $\Box = \partial_1^2 - \partial_2^2$  on  $\mathcal{P}_n(\Omega_2)$ . The linear operator

$$S: \mathcal{P}_n(\Omega_2) o \mathcal{P}_n(\Omega_2), \quad (Sf)(t_1, t_2) = f\left(rac{t_1 + t_2}{\sqrt{2}}, rac{t_1 - t_2}{\sqrt{2}}
ight)$$

is an isometry when taking the Hermite norm. Moreover, the factorization

$$\partial_1^2 - \partial_2^2 = (\partial_1 + \partial_2)(\partial_1 - \partial_2)$$

may be written as the identity  $2 \partial_1 \partial_2 S = S(\partial_1^2 - \partial_2^2)$ . It follows that

$$\begin{split} \eta(\Box \mid \mathcal{P}_n(\Omega_2)) &= \max_{\|f\|=1} \|(\partial_1^2 - \partial_2^2) f\| = \max_{\|Sf\|=1} \|S(\partial_1^2 - \partial_2^2) f\| \\ &= \max_{\|Sf\|=1} \|2 \, \partial_1 \partial_2 \, Sf\| = \max_{\|g\|=1} \|2 \, \partial_1 \partial_2 \, g\| \\ &= 2 \, \eta(\partial_1 \partial_2 \mid \mathcal{P}_n(\Omega_2)) \end{split}$$

and from Theorem 1.1 we obtain that

$$2\eta(\partial_1\partial_2\mid \mathcal{P}_n(\Omega_2))\sim 2n\max_{(x_1,x_2)\in\Omega_2}\sqrt{4x_1x_2}=2n.$$

The following problems remain open. What can be said about  $C(\Box \mid \mathcal{P}_n(\Omega_2))$  if  $C = \lambda$  or  $C = \gamma$ ? What happens if  $\Box$  is the more general operator  $\partial_1^2 - c^2 \partial_2^2$ ? Can one tackle the Laplace operator by using the factorization  $\Delta = (\partial_1 + i\partial_2)(\partial_1 - i\partial_2)$ ?

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