# THE PATH-CONNECTIVITY OF MRA WAVELETS IN $L^{2}\left(\mathbb{R}^{d}\right)$ 

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#### Abstract

We show that for any $d \times d$ expansive matrix $A$ with integer entries and $|\operatorname{det}(A)|=2$, the set of all $A$-dilation MRA wavelets is path-connected under the $L^{2}\left(\mathbb{R}^{d}\right)$ norm topology. We do this through the application of $A$-dilation wavelet multipliers, namely measurable functions $f$ with the property that the inverse Fourier transform of $(f \widehat{\psi})$ is an $A$-dilation wavelet for any $A$ dilation wavelet $\psi$ (where $\widehat{\psi}$ is the Fourier transform of $\psi$ ). In this process, we have completely characterized all $A$-dilation wavelet multipliers for any integral expansive matrix $A$ with $|\operatorname{det}(A)|=2$.


## 1. Introduction

Let $A$ be a $d \times d$ real expansive matrix, i.e., a matrix with real entries whose eigenvalues are all of modules greater than one. An $A$-dilation wavelet is a function $\psi \in L^{2}\left(\mathbb{R}^{d}\right)$ such that the set

$$
\left\{|\operatorname{det} A|^{\frac{n}{2}} \psi\left(A^{n} \mathbf{t}-\ell\right): n \in \mathbb{Z}, \ell \in \mathbb{Z}^{d}\right\}
$$

forms an orthonormal basis for $L^{2}\left(\mathbb{R}^{d}\right)$. For any function $f(\mathbf{t}) \in L^{1}\left(\mathbb{R}^{d}\right) \cap$ $L^{2}\left(\mathbb{R}^{d}\right)$, its Fourier transform is defined by

$$
\begin{equation*}
\mathcal{F}(f(\mathbf{t}))=\widehat{f}(\mathbf{s})=\frac{1}{(2 \pi)^{\frac{d}{2}}} \int_{\mathbb{R}^{d}} f(\mathbf{t}) e^{-i \mathbf{t} \circ \mathbf{s}} d \mathbf{t} \tag{1.1}
\end{equation*}
$$

where $\mathbf{t} \circ \mathbf{s}$ is the standard inner product of the vectors $\mathbf{s}, \mathbf{t} \in \mathbb{R}^{d}$. The inverse Fourier transform will be denoted by $\mathcal{F}^{-1}$.

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One of the many problems in wavelet theory concerns the construction of different wavelets. Naturally, one may attempt to construct new wavelets from an existing one. This approach leads to the concept of wavelet multipliers [3]. A measurable function $f$ is called an $A$-dilation wavelet multiplier if the inverse Fourier transform of $(f \widehat{\psi})$ is an $A$-dilation wavelet for any $A$-dilation wavelet $\psi$. Our study in this paper concerns the case where the dilation matrix $A$ is an expansive matrix with integer entries such that $|\operatorname{det}(A)|=2$. For such matrices, the wavelet multipliers have been studied extensively and are completely characterized for dimension $1[11,14]$ and dimension $2[9,10]$. These results then lead to the characterization of the phases of $A$-dilation MRA wavelets and the establishment of the path-connectedness of the set of all $A$-dilation MRA wavelets under the $L^{2}\left(\mathbb{R}^{1}\right)$ or $L^{2}\left(\mathbb{R}^{2}\right)$ norm topology for dimension $d=1$ and dimension $d=2$. Here in this paper, we will generalize the above mentioned results to all $d \times d$ expansive matrices with integer entries and determinant $\pm 2$. It is important for us to point out that the approach used for the case $d=2$ depends on certain special properties a $2 \times 2$ integral expansive matrix $A$ (with $|\operatorname{det}(A)|=2$ ) possesses [9, 10]. This is no longer the case for $d \geq 3$ so a different approach has to be used.

The rest of the paper is organized as follows. In the next section, we introduce the notations and terms needed for this paper, with some preliminary results needed in the later sections. In Section 3, we state and prove our results on wavelet multipliers on $L^{2}\left(\mathbb{R}^{d}\right)$. In Section 4, we prove the pathconnectivity of the set of all $A$-dilation MRA wavelets. In the last section, we compare our approach in this paper with the ones used in lower dimensions ( $d=1$ and $d=2$ ). We show that in the higher dimensions $d \geq 3$ there exist integral expansive matrices that would prohibit the direct generalizations of the approaches used to solve the path-connectivity of $A$-dilation MRA wavelets in $d=1$ and $d=2$ ([9]-[11], [14]).

## 2. Notations, definitions and preliminary results

Let $M_{d}^{(2)}(\mathbb{Z})$ be the set of all $d \times d$ expansive integral matrices (i.e., matrices with integer entries) whose determinants are $\pm 2$. Throughout this paper, we will limit our discussion to matrices $A \in M_{d}^{(2)}(\mathbb{Z})$. We will use $T, D_{A}$ as the translation and dilation unitary operators acting on $L^{2}\left(\mathbb{R}^{d}\right)$ defined by $\left(T^{\ell} f\right)(\mathbf{t})=f(\mathbf{t}-\ell),\left(D_{A} f\right)(\mathbf{t})=|\operatorname{det}(A)|^{\frac{1}{2}} f(A \mathbf{t}), \forall f \in L^{2}\left(\mathbb{R}^{d}\right), \mathbf{t} \in \mathbb{R}^{d}$ and $\ell \in \mathbb{Z}^{d}$. A measurable function $\psi \in L^{2}\left(\mathbb{R}^{d}\right)$ is called an $A$-dilation wavelet if $\left\{D_{A}^{n} T^{\ell} \psi: n \in \mathbb{Z}, \ell \in \mathbb{Z}^{d}\right\}$ is an orthonormal basis for $L^{2}\left(\mathbb{R}^{d}\right)$. Whenever we state that two functions $f, g \in L^{2}\left(\mathbb{R}^{d}\right)$ are equal, it is understood that $f(\mathbf{s})=g(\mathbf{s})$ for almost all $\mathbf{s} \in \mathbb{R}^{d}$. Furthermore, we say that $E=F$ for two measurable sets $F$ and $E$ in $\mathbb{R}^{d}$ if $(F \backslash F) \cup(E \backslash F)$ is a measure zero set.

Definition 2.1. A sequence $\left\{V_{j}: j \in \mathbb{Z}\right\}$ of closed subspaces of $L^{2}\left(\mathbb{R}^{d}\right)$ is called an $A$-dilation multi-resolution analysis (or $A$-dilation MRA for short) if the following hold:
(i) $V_{j} \subset V_{j+1}, \forall j \in \mathbb{Z}$;
(ii) $\bigcap_{j \in \mathbb{Z}} V_{j}=\{0\}, \overline{\bigcup_{j \in \mathbb{Z}} V_{j}}=L^{2}\left(\mathbb{R}^{d}\right)$;
(iii) $f(\mathbf{t}) \in V_{j}$ if and only if $f\left(A^{-j} \mathbf{t}\right) \in V_{0}$ for $j \in \mathbb{Z}$; and
(iv) There exists $\phi(\mathbf{t})$ in $V_{0}$ such that $\left\{\phi(\mathbf{t}-\ell): \ell \in \mathbb{Z}^{d}\right\}$ is an orthonormal basis for $V_{0}$.

The function $\phi(\mathbf{t})$ defined in (iv) above is called an $A$-dilation scaling function for the MRA. In our case, it is known that a single $A$-dilation wavelet can be derived from the above $A$-dilation MRA [12] (due to the fact that $|\operatorname{det}(A)|-1=1)$. An $A$-dilation wavelet $\psi \in V_{1} \cap V_{0}^{\perp}$ so obtained is called an MRA wavelet. For any $f \in V_{1}, f\left(A^{-1} \mathbf{t}\right) \in V_{0}$, hence we have

$$
\begin{equation*}
f(\mathbf{t})=|\operatorname{det}(A)| \sum_{\ell \in \mathbb{Z}^{d}} c_{\ell} \phi(A \mathbf{t}-\ell) \tag{2.1}
\end{equation*}
$$

If we define $m_{f}(\mathbf{s})=\sum_{\ell \in \mathbb{Z}^{d}} c_{\ell} e^{-i \ell o s}$, then by taking Fourier transform on both sides of (2.1) we obtain $\widehat{f}\left(A^{\tau} \mathbf{s}\right)=m_{f}(\mathbf{s}) \widehat{\phi}(\mathbf{s})$, where $A^{\tau}$ is the transpose of $A$. In particular, we have

$$
\begin{equation*}
\widehat{\phi}\left(A^{\tau} \mathbf{s}\right)=m(\mathbf{s}) \widehat{\phi}(\mathbf{s}) \tag{2.2}
\end{equation*}
$$

for some function $m(s)$ of the form similar to (2.1). $m(s)$ is called the low pass $A$-dilation filter of the MRA.

A measurable function $f(\mathbf{t}) \in L^{2}\left(\mathbb{R}^{d}\right)$ is called a $2 \pi \mathbb{Z}^{d}$ periodic if $f(\mathbf{t}+$ $2 \pi \ell)=f(\mathbf{t})$ on $\mathbb{R}^{d}$ for any $\ell \in \mathbb{Z}^{d}$.

The following lemmas are well-known results and can be easily obtained by standard textbook arguments $[2,4,7]$.

Lemma 2.1. $\psi$ is an A-dilation wavelet iff the following conditions hold:
(i) $\|\psi\|_{2}=1$;
(ii) $\sum_{j \in \mathbb{Z}}\left|\widehat{\psi}\left(\left(A^{\tau}\right)^{j} \mathbf{S}\right)\right|^{2}=1 /(2 \pi)^{d}$ and
(iii) $\sum_{j=0}^{\infty} \widehat{\psi}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \widehat{\psi}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)=0 \quad \forall \ell \in \mathbb{Z}^{d} \backslash A^{\tau} \mathbb{Z}^{d}$.

Lemma 2.2. An $A$-dilation wavelet $\psi$ is an $A$-dilation MRA wavelet iff

$$
\begin{equation*}
D_{\psi}(\mathbf{s})=\sum_{n=1}^{\infty} \sum_{\ell \in \mathbb{Z}^{d}}\left|\widehat{\psi}\left(\left(A^{\tau}\right)^{n}(\mathbf{s}+2 \pi \ell)\right)\right|^{2}=\frac{1}{(2 \pi)^{d}} \tag{2.3}
\end{equation*}
$$

Lemma 2.3. $\phi$ is an $A$-dilation scaling function for an MRA iff the following conditions hold:
(i) $\sum_{\ell \in \mathbb{Z}^{d}}|\widehat{\phi}(\mathbf{s}+2 \pi \ell)|^{2}=1 /(2 \pi)^{d}$;
(ii) $\lim _{j \rightarrow \infty}\left|\widehat{\phi}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)\right|=1 /(2 \pi)^{\frac{d}{2}}$ and
(iii) there exists a $2 \pi \mathbb{Z}^{d}$ periodic function $m(\mathbf{s}) \in L^{2}\left([-\pi, \pi)^{d}\right)$ such that $\widehat{\phi}\left(A^{\tau} \mathbf{s}\right)=m(\mathbf{s}) \widehat{\phi}(\mathbf{s})$.

Lemma 2.4 below describes some special properties a matrix in the set $M_{d}^{(2)}(\mathbb{Z})$ possesses. Its proof is elementary and our reader can also refer to [5].

LEmmA 2.4. Let $A \in M_{d}^{(2)}(\mathbb{Z})$. Then the group $\mathbb{Z}^{d} / A^{\tau} \mathbb{Z}^{d}$ is isomorphic to $\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d} / \mathbb{Z}^{d}$ and the order of $\mathbb{Z}^{d} / A^{\tau} \mathbb{Z}^{d}$ is equal to 2 . In particular, for any $\mathbf{h}_{1} \in \mathbb{Z}^{d} \backslash A^{\tau} \mathbb{Z}^{d}, \mathbb{Z}^{d}=A^{\tau} \mathbb{Z}^{d} \cup\left(A^{\tau} \mathbb{Z}^{d}+\mathbf{h}_{1}\right)$ and $\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d}=\mathbb{Z}^{d} \cup\left(\mathbb{Z}^{d}+\mathbf{h}_{2}\right)$ where $\mathbf{h}_{2}=\left(A^{\tau}\right)^{-1} \mathbf{h}_{1}$.

REMARK 2.1. Since $A \in M_{d}^{(2)}(\mathbb{Z})$, any noninteger entry in $\left(A^{\tau}\right)^{-1}$ is a rational number with denominator 2 (namely a number of the form $\frac{1}{2}(2 r+1)$ with $r \in \mathbb{Z}$. It follows that $\mathbf{h}_{2}=\left(A^{\tau}\right)^{-1} \mathbf{h}_{1} \notin \mathbb{Z}^{d}$ has at least one noninteger entry and all noninteger entries are rational numbers with denominator 2. We will use $\rho\left(\mathbf{h}_{2}\right)$ to denote the index of the first such noninteger entry in $\mathbf{h}_{2}$.

Remark 2.2. Notice that for $\mathbf{h}_{1}, \mathbf{h}_{1}^{\prime} \in \mathbb{Z}^{d} \backslash A^{\tau}\left(\mathbb{Z}^{d}\right)$ and $\mathbf{h}_{2}=\left(A^{\tau}\right)^{-1} \mathbf{h}_{1}$, $\mathbf{h}_{2}^{\prime}=\left(A^{\tau}\right)^{-1} \mathbf{h}_{1}^{\prime}$, we have $\rho\left(\mathbf{h}_{2}\right)=\rho\left(\mathbf{h}_{2}^{\prime}\right)$ since $\mathbf{h}_{2}-\mathbf{h}_{2}^{\prime} \in \mathbb{Z}^{d}$. Thus the index $\rho\left(\mathbf{h}_{2}\right)$ only depends on $A$. Hence, it is appropriate to denote such an index by $\rho(A)$. Let $\mathbf{u} \in \mathbb{R}^{d}$ be the vector with all of its entries being zero except that at its $\rho(A)$ th coordinate, where it has 1 as its entry. Then $e^{ \pm i 2 \pi \mathbf{h}_{2} \circ \mathbf{u}}=-1$. We leave it to our reader to verify that there is a unique element $\mathbf{h}_{1} \in \mathbb{Z}^{d} \backslash A^{\tau} \mathbb{Z}^{d}$ such that $\left(A^{\tau}\right)^{-1} \mathbf{h}_{1}=\mathbf{h}_{2}$ is a nonzero vector whose entries are either $1 / 2$ or 0 . In this case $\mathbf{h}_{2} \circ \mathbf{u}=1 / 2$. From now on, $\mathbf{h}_{1}$ and $\mathbf{h}_{2}$ will be understood to be these two uniquely determined vectors to avoid any confusion.

Lemma 2.5. Suppose that $\psi$ is an A-dilation MRA wavelet with scaling function $\phi$ and low pass filter $m(\mathbf{s})$, then

$$
\begin{gather*}
|\widehat{\phi}(\mathbf{s})|^{2}=\sum_{j=1}^{\infty}\left|\widehat{\psi}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2}  \tag{2.4}\\
|m(\mathbf{s})|^{2}+\left|m\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right|^{2}=1 \tag{2.5}
\end{gather*}
$$

A measurable set $F \in \mathbb{R}^{d}$ is called a $2 \pi \mathbb{Z}^{d}$ translation tiling domain (or just a tiling domain for short) if $\left\{F+2 \pi \ell: \ell \in \mathbb{Z}^{d}\right\}$ is a partition of $\mathbb{R}^{d}$. It is shown in [6] that there exist special translation tiling domains $F$ with the property that $\mathbf{0} \in F, F \subset A^{\tau} F$ and $\bigcup_{n \geq 0}\left(A^{\tau}\right)^{n} F=\mathbb{R}^{d}$. The set $E=A^{\tau} F \backslash F$ is called a generalized Shannon-type wavelet set (from which an $A$-dilation MRA wavelet can be obtained). A function $f$ with the property $|f|=1$ is called a unimodular function.

Proposition 2.1. Let $\phi \in L^{2}\left(\mathbb{R}^{d}\right)$ be a scaling function for an $A$-dilation $M R A\left\{V_{j}\right\}$ and let $m$ be its associated low pass filter. Let $\psi \in W_{0}=V_{1} \cap V_{0}^{\perp}$,
then $\left\{\psi(\mathbf{t}-\ell): \ell \in \mathbb{Z}^{d}\right\}$ is an orthonormal basis for $W_{0}$ iff there exists a $2 \pi \mathbb{Z}^{d}$ periodic, measurable and unimodular function $v: \mathbb{R}^{d} \rightarrow \mathbb{C}$ such that

$$
\begin{equation*}
\widehat{\psi}\left(A^{\tau} \mathbf{s}\right)=e^{i(\mathbf{s o u})} v\left(A^{\tau} \mathbf{s}\right) \overline{m\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)} \widehat{\phi}(\mathbf{s}) \tag{2.6}
\end{equation*}
$$

where $\mathbf{u}$ is the vector defined in Remark 2.2.
Let us give an outline of the proof for Proposition 2.1. From the discussion following (2.1), we have $\widehat{\psi}\left(A^{\tau} \mathbf{s}\right)=m_{\psi}(\mathbf{s}) \widehat{\phi}(\mathbf{s})$ for some $2 \pi \mathbb{Z}^{d}$ periodic function $m_{\psi}$. Again, standard arguments show that $\left\{\psi(\mathbf{t}-\ell): \ell \in \mathbb{Z}^{d}\right\}$ is an orthonormal basis for $W_{0}$ iff the equations $|m(\mathbf{s})|^{2}+\left|m\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right|^{2}=1$, $\left|m_{\psi}(\mathbf{s})\right|^{2}+\left|m_{\psi}\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right|^{2}=1$ and $m(\mathbf{s}) \overline{m_{\psi}(\mathbf{s})}+m\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right) \overline{m_{\psi}\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)}=$ 0 hold. The reader can verify that the solution for $m_{\psi}(\mathbf{s})$ (in terms of $m(\mathbf{s})$ ) is of the form $e^{i(\mathbf{s o u} \mathbf{u})} v\left(A^{\tau} \mathbf{s}\right) \overline{m\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)}$.

Proposition 2.2. Let $\psi$ be an A-dilation MRA wavelet, then $e^{i\left(\mathbf{s o} A^{-1} \mathbf{u}\right)}|\widehat{\psi}(\mathbf{s})|$ is the Fourier transform of an A-dilation MRA wavelet.

Proof. Let $\phi$ be the corresponding scaling function with low pass filter $m$, then $\mathcal{F}^{-1}(|\widehat{\phi}|)$ is also an $A$-dilation scaling function whose associated low pass filter is $|m|$ by Lemma 2.3. Thus, the function $\psi_{1}$ defined by $\widehat{\psi_{1}}\left(A^{\tau} \mathbf{s}\right)=$ $e^{i(\mathbf{s o u})}\left|\overline{m\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)} \widehat{\phi}(\mathbf{s})\right|=e^{i(\mathbf{s o u})}\left|\psi\left(A^{\tau} \mathbf{s}\right)\right|$ is an $A$-dilation MRA wavelet by Proposition 2.1.

A measurable subset $E$ of $\mathbb{R}^{d}$ is simple if $E(\Omega)=\bigcup_{\ell \in \mathbb{Z}^{d}}(E \cap(\Omega+2 \pi \ell)-$ $2 \pi \ell) \subset \Omega$ is a disjoint union (where $\Omega=[-\pi, \pi)^{d}$ ). Two simple sets $E, F \subset \mathbb{R}^{d}$ are said to be $2 \pi \mathbb{Z}^{d}$ translation congruent (or just translation congruent for short) to each other if $E(\Omega)=F(\Omega)$.

Proposition 2.3. Let $F$ be a $2 \pi \mathbb{Z}^{d}$ translation tiling domain of $\mathbb{R}^{d}$, then there exist disjoint subsets $F_{0}$ and $F_{1}$ of $F$ such that (1) $F=F_{0} \cup F_{1}$; (2) each of $F_{0}, F_{1}$ is a $2 \pi\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d}$ translation tiling domain of $\mathbb{R}^{d} ;(3) F_{1}-2 \pi \mathbf{h}_{2}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to $F_{0}$.

Proposition 2.3 is key to our proof for the path-connectivity of the $A$ dilation MRA wavelets so we provide a detailed proof here.

Proof of Proposition 2.3. Since $F$ is a $2 \pi \mathbb{Z}^{d}$ translation tiling domain of $\mathbb{R}^{d}$, $\left(A^{\tau}\right)^{-1} F$ is a $2 \pi\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d}$ translation tiling domain of $\mathbb{R}^{d}$. For each $\ell \in \mathbb{Z}^{d}$, define $F_{\ell}=F \cap\left(\left(\left(A^{\tau}\right)^{-1}\right) F+2 \pi \ell\right) . F_{\ell}$ is measurable since it is the intersection of two measurable sets. Now define $F_{1}=\bigcup_{\ell \in \mathbb{Z}^{d}} F_{\ell} . F_{1}$ is measurable since it is the union of countably many measurable sets. It follows that $F_{0}=F \backslash F_{1}$ is also measurable and conditions (1) is met. We now proceed to prove (2) and (3).

Claim 1. For any two points, $\mathbf{x}, \mathbf{y} \in F_{1}, \mathbf{x}-\mathbf{y} \notin 2 \pi\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d}$. For if not, then $\exists \mathbf{x}, \mathbf{y} \in F_{1}$ such that $\mathbf{x} \neq \mathbf{y}$ but $\mathbf{x}-\mathbf{y}=2 \pi\left(A^{\tau}\right)^{-1} \ell$ for some $\ell \in \mathbb{Z}^{d}$. By the definition of $F_{1}, \mathbf{x}=\left(A^{\tau}\right)^{-1} \mathbf{x}_{0}+2 \pi \ell_{1}$ for some $\mathbf{x}_{0} \in F$ and $\ell_{1} \in \mathbb{Z}^{d}$,
$\mathbf{y}=\left(A^{\tau}\right)^{-1} \mathbf{y}_{0}+2 \pi \ell_{2}$ for some $\mathbf{y}_{0} \in F$ and $\ell_{2} \in \mathbb{Z}^{d}$. It follows that $\mathbf{x}-\mathbf{y}=$ $\left(A^{\tau}\right)^{-1}\left(\mathbf{x}_{0}-\mathbf{y}_{0}\right)+2 \pi\left(\ell_{1}-\ell_{2}\right)=2 \pi\left(A^{\tau}\right)^{-1} \ell$ hence $\mathbf{x}_{0}-\mathbf{y}_{0} \in 2 \pi \mathbb{Z}^{d}$. Thus, $\mathbf{x}_{0}=\mathbf{y}_{0}$ since $F$ is a tiling domain. This in turn implies that $\mathbf{x}-\mathbf{y} \in 2 \pi \mathbb{Z}^{d}$ so $\mathbf{x}=\mathbf{y}$ (again because $F$ is a tiling domain). This contradicts to our assumption that $\mathbf{x} \neq \mathbf{y}$.

Claim 2. $F_{1}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to $\left(A^{\tau}\right)^{-1} F . \forall \mathbf{x} \in F_{1}$, then $\mathbf{x}=\left(A^{\tau}\right)^{-1} \mathbf{y}+2 \pi \ell$ for some $\mathbf{y} \in F$ and $\ell \in 2 \pi \mathbb{Z}^{d}$. Thus, $\mathbf{x}-\left(A^{\tau}\right)^{-1} \mathbf{y} \in 2 \pi \mathbb{Z}^{d}$ and it follows that $F_{1}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to a subset of $\left(A^{\tau}\right)^{-1} F$. Similarly, for any $\mathbf{y} \in\left(A^{\tau}\right)^{-1} F, \mathbf{y}=\left(A^{\tau}\right)^{-1} \mathbf{x}$ for some $\mathbf{x} \in F$. Since $F$ is a tiling of $\mathbb{R}^{d}, \mathbf{y}=\mathbf{z}+2 \pi \ell$ for some $\mathbf{z} \in F$ and $\ell \in \mathbb{Z}^{d}$. It follows that $\mathbf{z}=\mathbf{y}-2 \pi \ell=\left(A^{\tau}\right)^{-1} \mathbf{x}-2 \pi \ell \in F_{-\ell_{1}} \subset F_{1}$. Thus, $\left(A^{\tau}\right)^{-1} F$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to a subset of $F_{1}$. Therefore, $F_{1}$ is translation congruent to $\left(A^{\tau}\right)^{-1} F$.

Since $\left(A^{\tau}\right)^{-1} F$ is a tiling of $\mathbb{R}^{d}$ under the $2 \pi\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d}$ translation, it follows that $F_{1}$ is a tiling of $\mathbb{R}^{d}$ under the $2 \pi\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d}$ translation. This proves half of (2).

Claim 3. $F_{1}-2 \pi \mathbf{h}_{2}$ is translation congruent to $F_{0} . \forall \mathbf{x} \in F_{1}$, then $\mathbf{x}-$ $2 \pi \mathbf{h}_{2}=\mathbf{y}+2 \pi \ell_{1}$ for some $\mathbf{y} \in F, \ell_{1} \in \mathbb{Z}^{d}$ since $F$ is a tiling domain. If $\mathbf{y} \in F_{1}$, then $\mathbf{y}=\left(A^{\tau}\right)^{-1} \mathbf{y}_{1}+2 \pi \ell_{2}$, and $\mathbf{x}=\left(A^{\tau}\right)^{-1} \mathbf{x}_{1}+2 \pi \ell_{3}$, where $\mathbf{y}_{1}, \mathbf{x}_{1} \in F$. So $\left(A^{\tau}\right)^{-1} \mathbf{x}_{1}+2 \pi \ell_{3}-2 \pi \mathbf{h}_{2}=\left(A^{\tau}\right)^{-1} \mathbf{y}_{1}+2 \pi \ell_{2}+2 \pi \ell_{1}$. Multiplying both sides of this by $A^{\tau}$ shows that $\mathbf{x}_{1}-\mathbf{y}_{1} \in 2 \pi \mathbb{Z}^{d}$ hence $\mathbf{x}_{1}=\mathbf{y}_{1}$. But then we have $\mathbf{h}_{2}=\ell_{3}-\ell_{1}-\ell_{2} \in \mathbb{Z}^{d}$, which contradicts to $\mathbf{h}_{2} \notin \mathbb{Z}^{d}$. So we must have $\mathbf{y} \in F_{0}$ and this shows that $F_{1}-2 \pi \mathbf{h}_{2}$ is translation congruent to a subset of $F_{0}$. For any $\left.\mathbf{x}_{0} \in F_{0}, \mathbf{x}_{0}=\mathbf{x}_{1}+2 \pi\left(A^{\tau}\right)^{-1}\right) \ell_{1}$ for some $\mathbf{x}_{1} \in F_{1}$ and $\ell_{1} \in \mathbb{Z}^{d}$. It is necessary $\left.\left(A^{\tau}\right)^{-1}\right) \ell_{1} \notin \mathbb{Z}^{d}$, otherwise $\mathbf{x}_{0}=\mathbf{x}_{1}$ (again because $F$ is a tiling domain), which contradicts to $\mathbf{x}_{0} \notin F_{1}$. So $\left.2 \pi\left(A^{\tau}\right)^{-1}\right) \ell_{1}=2 \pi \mathbf{h}_{2}+2 \pi \ell_{2}$ for some $\ell_{2} \in \mathbb{Z}^{d}$. It follows that $\mathbf{x}_{0}+2 \pi \mathbf{h}_{2}$ is congruent equivalent to $\mathbf{x}_{1}$ modulo $2 \pi \mathbb{Z}^{d}$. Thus, $F_{0}$ is translation congruent to a subset of $F_{1}-2 \pi \mathbf{h}_{2}(3)$ is proven. The other half of (2) now follows easily: for any $\mathbf{z} \in \mathbb{R}^{d}, \mathbf{z}=\mathbf{x}+2 \pi\left(A^{\tau}\right)^{-1} \ell$ for some $\mathbf{x} \in F_{0}$. But $\mathbf{x}=\mathbf{x}_{1}-2 \pi \mathbf{h}_{2}+2 \pi \ell_{1}$ for some $\ell_{1} \in \mathbb{Z}^{d}$, so $\mathbf{z}=\mathbf{x}_{1}-2 \pi \mathbf{h}_{2}+$ $2 \pi \ell_{1}+2 \pi\left(A^{\tau}\right)^{-1} \ell=\mathbf{x}_{1}+2 \pi\left(A^{\tau}\right)^{-1}\left(\ell+\ell_{1}-A^{\tau} \mathbf{h}_{2}\right)=\mathbf{x}_{1}+2 \pi\left(A^{\tau}\right)^{-1} \ell_{2}$ where $\ell_{2}=\ell+\ell_{1}-A^{\tau} \mathbf{h}_{2} \in \mathbb{Z}^{d}$ and since $A^{\tau} \mathbf{h}_{2} \in \mathbb{Z}^{d}$.

A direct consequence of Proposition 2.3 is the following corollary. Its proof is elementary and is left to the reader.

Corollary 2.1. Let $A \in M_{d}^{(2)}(\mathbb{Z}), F$ a tiling domain and $\mathbf{h}_{1}, \mathbf{h}_{2}$ be as defined in Remark 2.2. Then for any $m \in L^{1}\left(\mathbb{R}^{d}\right)$ that is $2 \pi \mathbb{Z}^{d}$ periodic, we have:
(i) $\int_{F} m\left(\left(A^{\tau}\right)^{-1} \mathbf{s}\right) d \mathbf{s}=2 \int_{F_{1}} m(\mathbf{s}) d \mathbf{s}=2 \int_{F_{1}(\Omega)} m(\mathbf{s}) d \mathbf{s}$ and
(ii) $\int_{F_{j}} m(\mathbf{s}) d \mathbf{s}=\int_{F_{j}(\Omega)} m(\mathbf{s}) d \mathbf{s}=\int_{F_{1-j}(\Omega)} m\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right) d \mathbf{s}=\int_{F_{1-j}} m(\mathbf{s}+$ $\left.2 \pi \mathbf{h}_{2}\right) d \mathbf{s}$ where $j=0,1$.

## 3. $A$-dilation wavelet multipliers

A measurable function $f$ is called an $A$-dilation wavelet multiplier (or wavelet multiplier for short) if the inverse Fourier transform of $f \hat{\psi}$ is an $A$ dilation wavelet whenever $\psi$ is an $A$-dilation wavelet. In this section, we characterize the $A$-dilation wavelet multipliers. By an argument similar to the one used in [9], it can be shown that any wavelet multiplier $f$ has to be unimodular. Thus, in the following we will limit our discussion to such functions. Instead of trying to characterize the scaling function multiplier or the low pass filter multiplier (which is the approach used in [9]), we will use a different approach.

THEOREM 3.1. A unimodular function $f \in L^{\infty}\left(\mathbb{R}^{d}\right)$ is an $A$-dilation wavelet multiplier iff the function $k(\mathbf{s})=f\left(A^{\tau} \mathbf{s}\right) / f(\mathbf{s})$ is $2 \pi \mathbb{Z}^{d}$ periodic.

Proof. " $\Longleftarrow " ~ A s s u m e ~ t h a t ~ f \in L^{\infty}\left(\mathbb{R}^{d}\right)$ is a unimodular function and that $k(\mathbf{s})=f\left(A^{\tau} \mathbf{s}\right) / f(\mathbf{s})$ is $2 \pi \mathbb{Z}^{d}$ periodic. To show that $f$ is a wavelet multiplier, we need to show that for any $A$-dilation wavelet $\psi, \eta=\mathcal{F}^{-1}(f \widehat{\psi})$ is also a wavelet. It suffices to verify that $\widehat{\eta}$ satisfies conditions (ii) and (iii) in Lemma 2.1. It is easy to see that (ii) holds for $\widehat{\eta}$ since $|\widehat{\eta}|=|\widehat{\psi}|$ and (ii) holds for $\widehat{\psi}$. Applying the relation $f\left(A^{\tau} \mathbf{s}\right)=k(\mathbf{s}) f(\mathbf{s})$ repeatedly, for any $j \geq 1$ and $\ell \in \mathbb{Z}^{d}$, we obtain

$$
\begin{equation*}
f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)=k\left(\left(A^{\tau}\right)^{j-1} \mathbf{s}\right) \cdots k\left(A^{\tau} \mathbf{s}\right) k(\mathbf{s}) f(\mathbf{s}) \tag{3.1}
\end{equation*}
$$

and

$$
\begin{aligned}
f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right) & =k\left(\left(A^{\tau}\right)^{j-1}(\mathbf{s}+2 \pi \ell)\right) k\left(\left(A^{\tau}\right)^{j-2}(\mathbf{s}+2 \pi \ell)\right) \\
& \cdots k\left(A^{\tau}(\mathbf{s}+2 \pi \ell)\right) k(\mathbf{s}+2 \pi \ell) f(\mathbf{s}+2 \pi \ell) \\
& =k\left(\left(A^{\tau}\right)^{j-1} \mathbf{s}\right) \cdots k\left(A^{\tau} \mathbf{s}\right) k(\mathbf{s}) f(\mathbf{s}+2 \pi \ell)
\end{aligned}
$$

Since $k(\mathbf{s})$ is unimodular, this leads to

$$
\begin{aligned}
& f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \cdot \overline{f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)} \\
& \quad=k\left(\left(A^{\tau}\right)^{j-1} \mathbf{s}\right) \cdots k\left(A^{\tau} \mathbf{s}\right) k(\mathbf{s}) f(\mathbf{s}) \cdot \overline{k\left(\left(A^{\tau}\right)^{j-1} \mathbf{s}\right) \cdots k\left(A^{\tau} \mathbf{s}\right) k(\mathbf{s}) f(\mathbf{s}+2 \pi \ell)} \\
& \quad=f(\mathbf{s}) \overline{f(\mathbf{s}+2 \pi \ell)}
\end{aligned}
$$

for any $j \geq 0$ and $\ell \in \mathbb{Z}^{d}$. Thus,

$$
\begin{aligned}
& \sum_{j=0}^{\infty} \widehat{\eta}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \overline{\widehat{\eta}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)} \\
& \quad=f(\mathbf{s}) \overline{f(\mathbf{s}+2 \pi \ell)} \sum_{j=0}^{\infty} \widehat{\psi}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \overline{\widehat{\psi}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)}=0
\end{aligned}
$$

for any $\ell \in \mathbb{Z}^{d} \backslash A^{\tau} \mathbb{Z}^{d}$. So condition (iii) of Lemma 2.1 holds for $\widehat{\eta}$ as well.
$" \Longrightarrow "$ We need to show that $k(\mathbf{s})=f\left(A^{\tau} \mathbf{s}\right) / f(\mathbf{s})$ is $2 \pi \mathbb{Z}^{d}$ periodic. Let $\psi$ be any $A$-dilation MRA wavelet such that $\operatorname{supp}(\widehat{\psi})=\mathbb{R}^{d}$ (the existence of such $\psi$ is proven in [5, Example 5.14]). By Proposition 2.2, the function $\psi_{1}(t)$ defined by

$$
\begin{equation*}
\widehat{\psi_{1}}=e^{i\left(A^{\tau}\right)^{-1} \mathbf{s} \circ \mathbf{u}}|\widehat{\psi}(\mathbf{s})|=e^{i \mathbf{s} \circ A^{-1} \mathbf{u}}\left|\widehat{\psi_{1}}(\mathbf{s})\right| \tag{3.2}
\end{equation*}
$$

is an $A$-dilation wavelet. Since $\mathcal{F}^{-1}\left(f \widehat{\psi_{1}}\right)$ is also an $A$-dilation wavelet, $\widehat{\psi_{1}}$ and $f \widehat{\psi_{1}}$ both satisfy condition (iii) of Lemma 2.1, that is,

$$
\begin{align*}
& \sum_{j=0}^{\infty} \widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \cdot \widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)=0 \quad \text { and }  \tag{3.3}\\
& \sum_{j=0}^{\infty} f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)  \tag{3.4}\\
& \quad \times \overline{f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)} \widehat{\psi_{1}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)}=0
\end{align*}
$$

for any $\ell \in \mathbb{Z}^{d} \backslash A^{\tau} \mathbb{Z}^{d}$. Since $\ell \in \mathbb{Z}^{d} \backslash A^{\tau} \mathbb{Z}^{d}$, there exists $\ell_{1} \in \mathbb{Z}^{d}$ such that $\ell=A^{\tau} \ell_{1}+\mathbf{h}_{1}=A^{\tau}\left(\ell_{1}+\mathbf{h}_{2}\right)$ by Lemma 2.4. Thus,

$$
\begin{aligned}
\widehat{\psi_{1}}(\mathbf{s}) \widehat{\hat{\psi}_{1}}(\mathbf{s}+2 \pi \ell) & =e^{i \mathbf{s} \circ A^{-1} \mathbf{u}}\left|\widehat{\psi_{1}}(\mathbf{s})\right| \cdot e^{-i(\mathbf{s}+2 \pi \ell) \circ A^{-1} \mathbf{u}}\left|\widehat{\psi_{1}}(\mathbf{s}+2 \pi \ell)\right| \\
& =e^{-i 2 \pi\left(\ell_{1}+\mathbf{h}_{2}\right) \circ \mathbf{u}}\left|\widehat{\psi_{1}}(\mathbf{s})\right| \cdot\left|\widehat{\psi_{1}}(\mathbf{s}+2 \pi \ell)\right| \\
& =e^{-i 2 \pi \mathbf{h}_{2} \circ \mathbf{u}}\left|\widehat{\psi_{1}}(\mathbf{s})\right| \cdot\left|\widehat{\psi_{1}}(\mathbf{s}+2 \pi \ell)\right| \\
& =-\left|\widehat{\psi_{1}}(\mathbf{s})\right| \cdot\left|\widehat{\psi_{1}}(\mathbf{s}+2 \pi \ell)\right|
\end{aligned}
$$

since $\ell_{1} \circ \mathbf{u}$ is an integer and $e^{-i 2 \pi \mathbf{h}_{2} \circ \mathbf{u}}=-1$ (see Remark 2.2). On the other hand, for any $j>0$,

$$
\begin{aligned}
& \widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)
\end{aligned}
$$

$$
\begin{aligned}
& =\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right| \cdot\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)\right| .
\end{aligned}
$$

Thus, (3.3) and (3.4) can be rewritten as

$$
\begin{align*}
& \left|\widehat{\psi_{1}}(\mathbf{s})\right| \cdot\left|\widehat{\psi_{1}}(\mathbf{s}+2 \pi \ell)\right|  \tag{3.5}\\
& =\sum_{j=1}^{\infty}\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right| \cdot\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)\right| \quad \text { and } \\
& f(\mathbf{s}) \overline{f(\mathbf{s}+2 \pi \ell)} \cdot\left|\widehat{\psi_{1}}(\mathbf{s})\right| \cdot\left|\widehat{\psi_{1}}(\mathbf{s}+2 \pi \ell)\right|  \tag{3.6}\\
& =\sum_{j=1}^{\infty} f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right) \overline{f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)}\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right| \\
& \quad \times\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)\right| .
\end{align*}
$$

Since $f$ is unimodular, $\bar{f}=1 / f$. Hence, (3.6) can be rewritten as

$$
\begin{align*}
& \frac{f(\mathbf{s})}{f(\mathbf{s}+2 \pi \ell)}\left|\widehat{\psi_{1}}(\mathbf{s})\right| \cdot\left|\widehat{\psi_{1}}(\mathbf{s}+2 \pi \ell)\right|  \tag{3.7}\\
& \quad=\sum_{j=1}^{\infty} \frac{f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)}{f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)}\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right| \cdot\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)\right|
\end{align*}
$$

Combining this with (3.5) then leads to

$$
\begin{align*}
& \sum_{j=1}^{\infty}\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right| \cdot\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)\right|  \tag{3.8}\\
& =\sum_{j=1}^{\infty} \frac{f(\mathbf{s}+2 \pi \ell)}{f(\mathbf{s})} \frac{f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)}{f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)}\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right| \\
& \quad \times\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)\right| .
\end{align*}
$$

Finally, since $\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right| \cdot\left|\widehat{\psi_{1}}\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)\right|>0$ by the choice of $\psi_{1}$ and $\left|\frac{f(\mathbf{s}+2 \pi \ell)}{f(\mathbf{s})} \frac{f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)}{f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)}\right|=1$, it follows that $\frac{f(\mathbf{s}+2 \pi \ell)}{f(\mathbf{s})} \frac{f\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)}{f\left(\left(A^{\tau}\right)^{j}(\mathbf{s}+2 \pi \ell)\right)}=1$. In particular, $k(\mathbf{s})=\frac{f\left(A^{\tau} \mathbf{s}\right)}{f(\mathbf{s})}=\frac{f\left(A^{\tau}(\mathbf{s}+2 \pi \ell)\right)}{f(\mathbf{s}+2 \pi \ell)}=k(\mathbf{s}+2 \pi \ell) \forall \ell \in \mathbb{Z}^{d} \backslash A^{\tau} \mathbb{Z}^{d}$. If $\ell \in A^{\tau} \mathbb{Z}^{d}$, then $\ell-\mathbf{h}_{1} \notin A^{\tau} \mathbb{Z}^{d}$, and $k(\mathbf{s}+2 \pi \ell)=k\left(\mathbf{s}+2 \pi \mathbf{h}_{1}+2 \pi\left(\ell-\mathbf{h}_{1}\right)\right)=k\left(\mathbf{s}+2 \pi \mathbf{h}_{1}\right)=$ $k(\mathbf{s})$ as well. Therefore, $k(\mathbf{s})=f\left(A^{\tau} \mathbf{s}\right) / f(\mathbf{s})$ is $2 \pi \mathbb{Z}^{d}$ periodic.

Combining Lemma 2.2 and Theorem 3.1 then leads to the following corollary.

Corollary 3.1. A unimodular function $f \in L^{\infty}\left(\mathbb{R}^{d}\right)$ is an A-dilation MRA wavelet multiplier iff the function $k(\mathbf{s})=f\left(A^{\tau} \mathbf{s}\right) / f(\mathbf{s})$ is $2 \pi \mathbb{Z}^{d}$ periodic.

Next, we show that all $A$-dilation wavelet multipliers can be constructed in the way described in the following theorem. A measurable set $E \subset \mathbb{R}^{d}$ is called an $A$-dilation wavelet set if $\mathcal{F}^{-1}\left((2 \pi)^{-\frac{d}{2}} \chi_{E}\right)$ is an $A$-dilation wavelet. It is known that $E$ is an $A$-dilation wavelet set iff both the sets $\left\{A^{n} E: n \in \mathbb{Z}\right\}$ and $\left\{E+2 \pi \ell: \ell \in \mathbb{Z}^{d}\right\}$ are partitions of $\mathbb{R}^{d}$ [1].

Theorem 3.2. Let $E$ be an A-dilation wavelet set, $k(\mathbf{s})$ be a measurable unimodular $2 \pi \mathbb{Z}^{d}$ periodic function and $g(\mathbf{s})$ be a measurable unimodular function defined on $E$. Define

$$
f(\mathbf{s})= \begin{cases}g(\mathbf{s}), & \mathbf{s} \in E \\ k\left(\left(A^{\tau}\right)^{-1} \mathbf{s}\right) \cdots k\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right) \cdot g\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right), & \mathbf{s} \in\left(A^{\tau}\right)^{n} E, n \geq 1 \\ k(\mathbf{s}) k\left(A^{\tau} \mathbf{s}\right) \cdots k\left(\left(A^{\tau}\right)^{n-1} \mathbf{s}\right) \cdot g\left(\left(A^{\tau}\right)^{n} \mathbf{s}\right), & \mathbf{s} \in\left(A^{\tau}\right)^{-n} E, n \geq 1 \\ 1, & \mathbf{s}=0\end{cases}
$$

Then $f$ is an $A$-dilation wavelet multiplier. Moreover, any $A$-dilation wavelet multiplier can be constructed this way.

Proof. Since $k(\mathbf{s})$ is $2 \pi \mathbb{Z}^{d}$ periodic, it suffices to show that $f\left(A^{\tau} \mathbf{s}\right)=k(\mathbf{s})$. $f(\mathbf{s})$ in order to show that $f$ is an $A$-dilation wavelet multiplier by Theorem 3.1.

Case 1. $\mathbf{s} \in E$. Then $A^{\tau} \mathbf{s} \in A^{\tau} E$ and

$$
f\left(A^{\tau} \mathbf{s}\right)=k\left(\left(A^{\tau}\right)^{-1} A^{\tau} \mathbf{s}\right) g\left(\left(A^{\tau}\right)^{-1} A^{\tau} \mathbf{s}\right)=k(\mathbf{s}) g(\mathbf{s})=k(\mathbf{s}) f(\mathbf{s})
$$

Case 2. $\mathbf{s} \in\left(A^{\tau}\right)^{n} E$ where $n \geq 1$. Then $A^{\tau} \mathbf{s} \in\left(A^{\tau}\right)^{n+1} E$ and

$$
f\left(A^{\tau} \mathbf{s}\right)=k(\mathbf{s}) k\left(\left(A^{\tau}\right)^{-1} \mathbf{s}\right) \cdots k\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right) g\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)=k(\mathbf{s}) f(\mathbf{s})
$$

Case 3. $\mathbf{s} \in\left(A^{\tau}\right)^{-1} E$. Then $A^{\tau} \mathbf{s} \in E$ and $f(\mathbf{s})=\overline{k(\mathbf{s})} g\left(A^{\tau} \mathbf{s}\right)$, so $f\left(A^{\tau} \mathbf{s}\right)=$ $g\left(A^{\tau} \mathbf{s}\right)=k(\mathbf{s}) f(\mathbf{s})$.

Case 4. $\mathbf{s} \in\left(A^{\tau}\right)^{-n} E$ where $n>1$. Then $A^{\tau} \mathbf{s} \in\left(A^{\tau}\right)^{-(n-1)} E$ and

$$
f\left(A^{\tau} \mathbf{s}\right)=k(\mathbf{s}) \overline{k(\mathbf{s}) k\left(A^{\tau} \mathbf{s}\right) \cdots k\left(\left(A^{\tau}\right)^{n-1} \mathbf{s}\right)} g\left(\left(A^{\tau}\right)^{n} \mathbf{s}\right)=k(\mathbf{s}) f(\mathbf{s}) .
$$

Since $\left\{\left(A^{\tau}\right)^{n} E: n \in \mathbb{Z}\right\}$ is a partition of $\mathbb{R}^{d}$ modulo a null set, the above four cases have exhausted all possibilities for $\mathbf{s} \in \mathbb{R}^{d}$ in the sense.

Now suppose that $f(\mathbf{s})$ is an $A$-dilation wavelet multiplier. Let $g(\mathbf{s})=$ $f(\mathbf{s})$ for $\mathbf{s} \in E$, and $k(\mathbf{s})=f\left(A^{\tau} \mathbf{s}\right) / f(\mathbf{s})$. Then $k(\mathbf{s})$ is $2 \pi \mathbb{Z}^{d}$ periodic and is unimodular. We leave it to our reader to verify that $f(\mathbf{s})$ has the form given in the theorem.

## 4. Path-connectivity of the set of $A$-dilation MRA wavelets

In this section, we prove the main result of this paper, namely that the set of all $A$-dilation MRA wavelets is path-connected under the $L^{2}\left(\mathbb{R}^{d}\right)$ norm topology. For more discussions and related results on this topic, interested reader may refer to $[11,13,14]$.

Theorem 4.1. For any two $A$-dilation $M R A$ wavelets $\psi_{0}$ and $\psi_{1}$, there exists a continuous map $\gamma:[0,1] \longrightarrow L^{2}\left(\mathbb{R}^{d}\right)$ such that $\gamma(0)=\psi_{0}, \gamma(1)=\psi_{1}$ and $\gamma(t)$ is an $A$-dilation MRA wavelet for $\forall t \in[0,1]$.

We will prove the theorem by directly constructing a continuous path connecting the two MRA wavelets. Since the proof is of constructive nature and is fairly long, we will break it into several lemmas. For a given $A$-dilation MRA wavelet $\psi_{0}$, we will associate with it three special subsets of $A$-dilation MRA wavelets denoted by $\mathcal{M}_{\psi_{0}}, \mathcal{W}_{\psi_{0}}$ and $\mathcal{S}_{\psi_{0}}: \mathcal{M}_{\psi_{0}}$ contains all $A$-dilation MRA wavelets $\psi$ such that $\widehat{\psi}=v \widehat{\psi_{0}}$ for some $A$-dilation wavelet multiplier $v, \mathcal{W}_{\psi_{0}}$ contains all $A$-dilation MRA wavelets $\psi$ such that $|\widehat{\psi}|=\left|\widehat{\psi_{0}}\right|$ and $\mathcal{S}_{\psi_{0}}$ contains all $A$-dilation MRA wavelets $\psi$ such that $|\widehat{\phi}|=\left|\widehat{\phi_{0}}\right|$, where $\phi$ and $\phi_{0}$ are the corresponding scaling functions of $\psi$ and $\psi_{0}$.

Lemma 4.1. $\mathcal{S}_{\psi_{0}}=\mathcal{M}_{\psi_{0}}=\mathcal{W}_{\psi_{0}}$ for any $A$-dilation MRA wavelet $\psi_{0}$.

Proof. Let $\psi \in \mathcal{W}_{\psi_{0}}$, then $|\widehat{\psi}|=\left|\widehat{\psi_{0}}\right|$. By (2.4), we have

$$
|\widehat{\phi}(\mathbf{s})|^{2}=\sum_{j=1}^{\infty}\left|\widehat{\psi}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2}=\sum_{j=1}^{\infty}\left|\widehat{\psi_{0}}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2}=\left|\widehat{\phi_{0}}(\mathbf{s})\right|^{2}
$$

so $\psi \in \mathcal{S}_{\psi_{0}}$. This shows that $\mathcal{W}_{\psi_{0}} \subseteq \mathcal{S}_{\psi_{0}}$. $\mathcal{M}_{\psi_{0}} \subseteq \mathcal{W}_{\psi_{0}}$ by definition. $\mathcal{S}_{\psi_{0}} \subseteq$ $\mathcal{M}_{\psi_{0}}$ follows from an argument similar to the one used in the proof of [9, Theorem 1.2] and Proposition 2.1. Thus, $\mathcal{M}_{\psi_{0}} \subseteq \mathcal{W}_{\psi_{0}} \subseteq \mathcal{S}_{\psi_{0}} \subseteq \mathcal{M}_{\psi_{0}}$. Therefore, $\mathcal{S}_{\psi_{0}}=\mathcal{M}_{\psi_{0}}=\mathcal{W}_{\psi_{0}}$.

Lemma 4.2. Let $\psi_{0}$ be an $A$-dilation MRA wavelet. Then $\mathcal{M}_{\psi_{0}}$ is pathconnected.

Proof. This is proved in [9, Theorem 1.3] for two special cases of $A$. However, that proof can be easily modified for the general case and is left to the reader.

Let $F$ be a tiling domain with the property that $\mathbf{0} \in F, F \subset A^{\tau} F$ and $\bigcup_{n \geq 0}\left(A^{\tau}\right)^{n} F=\mathbb{R}^{d}$ (the existence of such sets is shown in [6] as we mentioned in Section 2). In this case, the function $\psi_{0}$ defined by

$$
\begin{equation*}
\widehat{\psi_{0}}(\mathbf{s})=(2 \pi)^{-\frac{d}{2}} e^{i \mathbf{s} \circ A^{-1} \mathbf{u}^{\prime}} \chi_{E}(\mathbf{s}) \tag{4.1}
\end{equation*}
$$

is an $A$-dilation MRA wavelet (where $E=A^{\tau} F \backslash F$ ) and is called a generalized Shannon-type wavelet. To show that any two $A$-dilation MRA wavelets are connected by a continuous path, it suffices to show any $A$-dilation MRA wavelet $\psi$ can be path-connected to $\psi_{0}$. We will do this by showing that there exists a $\psi_{1} \in \mathcal{S}_{\psi}$ such that $\psi_{1}$ is path-connected to $\psi_{0}$. By Lemmas 4.1 and 4.2 above, $\psi_{1}$ is path-connected to $\psi$ hence $\psi_{0}$ is path-connected to $\psi$ as well.

Note that the corresponding scaling function and low pass filter of $\psi_{0}$ are given by is $\hat{\phi}_{0}(\mathbf{s})=(2 \pi)^{-\frac{d}{2}} \chi_{F}$ and $m_{0}(\mathbf{s})=\chi_{\left(A^{\tau}\right)^{-1} F}$. Keep in mind that $F$ is a $2 \pi \mathbb{Z}^{d}$ tiling domain so it is $2 \pi \mathbb{Z}^{d}$ translation congruent to $\Omega=[-\pi, \pi]^{d}$.

By Lemma 2.3 and Proposition 2.2, there exists $\psi_{1} \in \mathcal{S}_{\psi}$ such that its corresponding scaling function $\phi_{1}$ and low pass filter $m_{1}$ satisfy $\widehat{\phi_{1}} \geq 0, m_{1} \geq 0$ and

$$
\begin{equation*}
\widehat{\psi_{1}}(\mathbf{s})=e^{i \mathbf{s} \circ A^{-1} \mathbf{u}^{\prime}} m_{1}\left(\left(A^{\tau}\right)^{-1} \mathbf{s}+2 \pi \mathbf{h}_{2}\right) \widehat{\phi_{1}}\left(\left(A^{\tau}\right)^{-1} \mathbf{s}\right) \tag{4.2}
\end{equation*}
$$

We will now show that this particular choice of $\psi_{1}$ suffices for our purpose. In the following, we will first build a path that connects the low-pass filters, then use this filter path to construct the path for the scaling functions and ultimately the path that connects wavelet functions $\psi_{1}$ and $\psi_{0}$.

Keep in mind that $m_{0}$ and $m_{1}$ are $2 \pi \mathbb{Z}^{d}$ periodic functions. Let $F_{1}=$ $\left(A^{\tau}\right)^{-1} F \subset F$ and $F_{0}=F \backslash F_{1}$ (thus, the support of $m_{0}$ on $F$ is simply $F_{1}$ ). Recall from Proposition 2.3 (and its proof) that $F_{0}$ and $F_{1}$ are both $2 \pi\left(A^{\tau}\right)^{-1} \mathbb{Z}^{d}$ tiling domains of $\mathbb{R}^{d}$ and $F_{1} \pm 2 \pi \mathbf{h}_{2}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to $F_{0}$. So the measure of $F_{0}$ and $F_{1}$ are both $(2 \pi)^{d} / 2$ and the measure of
$\left(A^{\tau}\right)^{-k} F_{1}$ is $(2 \pi)^{d} / 2^{k+1}$ for any $k \geq 0$. For any $t \in(0,1)$, there exists an integer $k_{0}(t) \geq 0$ such that $2^{-\left(k_{0}(t)+1\right)}<1-t \leq 2^{-k_{0}(t)}$. If $1-t=2^{-k_{0}(t)}$, define $H_{t}=\left(A^{\tau}\right)^{-k_{0}(t)} F_{1}$. If $2^{-\left(k_{0}(t)+1\right)}<1-t<2^{-k_{0}(t)}$, then there exists a positive real number $r_{t}>0$ such that the set $G_{t}=B\left(r_{t}\right) \cap\left(A^{\tau}\right)^{-k_{0}(t)}\left(F_{1} \backslash\left(A^{\tau}\right)^{-1} F_{1}\right)$ has measure $\left(1-t-2^{-\left(k_{0}(t)+1\right)}\right)(2 \pi)^{d} / 2=\left(1-t-2^{-\left(k_{0}(t)+1\right)}\right) \mu\left(F_{1}\right)$, where $B\left(r_{t}\right)=\left\{\mathbf{x} \in \mathbb{R}^{d}:|\mathbf{x}|^{2} \leq r_{t}^{2}\right\}$ and $\mu$ denotes the Lebesgue measure in $\mathbb{R}^{d}$. And we will define $H_{t}=G_{t} \cup\left(A^{\tau}\right)^{-\left(k_{0}(t)+1\right)} F_{1}$. Under this definition, observe that the measure of $H_{t}$ is precisely $(1-t) \mu\left(F_{1}\right)$. Furthermore, for any $\mathbf{s} \in F$, it is obvious that $\left(A^{\tau}\right)^{-k} \mathbf{s} \in H_{t}$ for any $k \geq k_{0}(t)+2$. For $t=0$ and $t=1, H_{t}$ is defined as $F_{1}$ and the empty set respectively. Since $H_{t}$ is a subset of $F_{1}$ and $F_{1} \pm 2 \pi \mathbf{h}_{2}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to $F_{0}, H_{t} \pm 2 \pi \mathbf{h}_{2}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to a subset (uniquely determined by $H_{t}$ modular a zero measure set) $J_{t} \subset F_{0}$. Equivalently, $J_{t} \pm 2 \pi \mathbf{h}_{2}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to $H_{t}$. The low-pass filter path $m_{t}(\mathbf{s})$ is then defined for any $\mathbf{s} \in F$ first as below and then extended to $\mathbb{R}^{d}$ as a $2 \pi \mathbb{Z}^{d}$ periodic function (this is possible since $F$ is a tiling domain).

$$
m_{t}(\mathbf{s})= \begin{cases}\sqrt{1-t+t m_{1}^{2}(\mathbf{s})}, & \mathbf{s} \in F_{1} \backslash H_{t}  \tag{4.3}\\ 1, & \mathbf{s} \in H_{t} \\ \sqrt{t} m_{1}(\mathbf{s}), & \mathbf{s} \in F_{0} \backslash J_{t} \\ 0, & \mathbf{s} \in J_{t}\end{cases}
$$

For $t=0$ and $t=1, m_{t}(\mathbf{s})$ is just $m_{0}(\mathbf{s})$ and $m_{1}(\mathbf{s})$. Furthermore, $m_{t}(\mathbf{s})$ satisfies the equation $\left|m_{t}(\mathbf{s})\right|^{2}+\left|m_{t}\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right|^{2}=1$. These are left to our reader to verify.

Lemma 4.3. For each $t \in[0,1]$, the function $\phi_{t}$ defined by

$$
\begin{equation*}
\widehat{\phi}_{t}(\mathbf{s})=(2 \pi)^{-\frac{d}{2}} \prod_{j=1}^{\infty} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right) \tag{4.4}
\end{equation*}
$$

is an MRA A-dilation scaling function hence $\psi_{t}$ defined by

$$
\begin{equation*}
\widehat{\psi_{t}}(\mathbf{s})=e^{i \mathbf{s} \circ A^{-1} \mathbf{u}} m_{t}\left(\left(A^{\tau}\right)^{-1} \mathbf{s}+2 \pi \mathbf{h}_{2}\right) \widehat{\phi}_{t}\left(\left(A^{\tau}\right)^{-1} \mathbf{s}\right) \tag{4.5}
\end{equation*}
$$

is an MRA A-dilation wavelet.
Proof. The statement holds trivially for $t=0$ and 1 , so we only need to consider the case $0<t<1$. $\widehat{\phi}_{t}$ is well defined since $0 \leq m_{t}(\mathbf{s}) \leq 1$, so is $\widehat{\psi_{t}}$. From the definition of $\widehat{\phi}_{t}$, we have $\widehat{\phi}_{t}\left(A^{\tau} \mathbf{s}\right)=m_{t}(\mathbf{s}) \widehat{\phi}_{t}(\mathbf{s})$, So $\widehat{\phi}_{t}$ satisfies condition (iii) of Lemma 2.3. On the other hand, $\widehat{\phi}_{t}(\mathbf{s})=(2 \pi)^{-\frac{d}{2}}, \forall \mathbf{s} \in H_{t}$, also by the definition of $\widehat{\phi}_{t}$. Since $0<t<1,\left(A^{\tau}\right)^{-\left(k_{0}(t)+1\right)} F_{1} \subset H_{t}$ so $H_{t}$ contains a neighborhood of $\mathbf{0}$ since $F$ (hence $F_{1}$ ) contains a neighborhood of $\mathbf{0}$. From this condition (ii) of Lemma 2.3 follows. We now prove that $\phi_{t}$ satisfies condition (i) of Lemma 2.3 as well, which then implies that $\phi_{t}$ is a scaling function and $\psi_{t}$ is an MRA $A$-dilation wavelet.

As we observed earlier, $\forall \mathbf{s} \in F,\left(A^{\tau}\right)^{-k} \mathbf{s} \in H_{t}, \forall k>k_{0}(t)+1=k_{0}^{\prime}$. So by the definition of $m_{t}(\mathbf{s})$, we have $m_{t}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right)=1$ for all such $k$. On the other hand, if $1 \leq k \leq k_{0}^{\prime}$, we have $\left(A^{\tau}\right)^{-k} \mathbf{s} \in F_{1}$ hence $m_{t}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right) \geq \sqrt{1-t}$. It follows that for any $\mathbf{s} \in F$ we have

$$
\begin{aligned}
\widehat{\phi}_{t}(\mathbf{s}) & =(2 \pi)^{-\frac{d}{2}} \prod_{k=1}^{k_{0}^{\prime}} m_{t}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right) \prod_{k=k_{0}^{\prime}+1}^{\infty} m_{t}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right) \\
& =(2 \pi)^{-\frac{d}{2}} \prod_{k=1}^{k_{0}^{\prime}} m_{t}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right) \geq(2 \pi)^{-\frac{d}{2}}(1-t)^{k_{0}^{\prime} / 2}
\end{aligned}
$$

This implies that $\chi_{F}(\mathbf{s}) \leq(2 \pi)^{\frac{d}{2}} \widehat{\phi}_{t}(\mathbf{s}) /(1-t)^{k_{0}^{\prime} / 2}$. For each $k \geq 1$, define

$$
p_{t, k}(\mathbf{s})=(2 \pi)^{-\frac{d}{2}} \chi_{F}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right) \cdot \prod_{j=1}^{k} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)
$$

Then

$$
p_{t, k}(\mathbf{s}) \leq \frac{\widehat{\phi}_{t}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right)}{(1-t)^{k_{0}^{\prime} / 2}} \prod_{j=1}^{k} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)=\frac{\widehat{\phi_{t}}(\mathbf{s})}{(1-t)^{k_{0}^{\prime} / 2}} .
$$

For $k \geq 2, \quad \int_{\mathbb{R}^{d}}\left|p_{t, k}(\mathbf{s})\right|^{2} e^{-i \mathbf{n o s}} d \mathbf{s}$ can be rewritten as (substituting $\mathbf{s}$ for $\left.\left(A^{\tau}\right)^{-k} \mathbf{s}\right)$

$$
\begin{aligned}
& \frac{1}{(2 \pi)^{d}} \int_{\mathbb{R}^{d}}\left|\chi_{F}\left(\left(A^{\tau}\right)^{-k} \mathbf{s}\right)\right|^{2} \cdot \prod_{j=1}^{k}\left|m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)\right|^{2} \cdot e^{-i \mathbf{n o s}} d \mathbf{s} \\
& =\frac{2^{k}}{(2 \pi)^{d}} \int_{F} \prod_{j=0}^{k-1}\left|m_{t}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k} \mathbf{s}\right)} d \mathbf{s} \\
& =\frac{2^{k}}{(2 \pi)^{d}}\left(\int_{F_{0}}\left|m_{t}(\mathbf{s})\right|^{2} \prod_{j=1}^{k-1}\left|m_{t}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k} \mathbf{s}\right)} d \mathbf{s}\right. \\
& \left.\quad+\int_{F_{1}}\left|m_{t}(\mathbf{s})\right|^{2} \prod_{j=1}^{k-1}\left|m_{t}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k} \mathbf{s}\right)} d \mathbf{s}\right)
\end{aligned}
$$

Since $F_{0} \pm 2 \pi \mathbf{h}_{2}$ is $2 \pi \mathbb{Z}^{d}$ translation congruent to $F_{1}$ and $e^{-i n \mathbf{s}}, m_{t}(\mathbf{s})$ are $2 \pi \mathbb{Z}^{d}$ periodic, the above equality becomes (keep in mind that $F_{1}=\left(A^{\tau}\right)^{-1} F$ )

$$
\begin{aligned}
& \frac{2^{k}}{(2 \pi)^{d}}\left(\int_{F_{1}}\left|m_{t}\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right|^{2} \prod_{j=1}^{k-1}\left|m_{t}\left(\left(A^{\tau}\right)^{j}\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k}\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right)} d \mathbf{s}\right. \\
& \left.\quad+\int_{F_{1}}\left|m_{t}(\mathbf{s})\right|^{2} \prod_{j=1}^{k-1}\left|m_{t}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k} \mathbf{s}\right)} d \mathbf{s}\right)
\end{aligned}
$$

$$
\begin{aligned}
& =\frac{2^{k}}{(2 \pi)^{d}} \int_{F_{1}}\left(\left|m_{t}(\mathbf{s})\right|^{2}+\left|m_{t}\left(\mathbf{s}+2 \pi \mathbf{h}_{2}\right)\right|^{2}\right) \prod_{j=1}^{k-1}\left|m_{t}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k} \mathbf{s}\right)} d \mathbf{s} \\
& =\frac{2^{k}}{(2 \pi)^{d}} \int_{F_{1}} \prod_{j=1}^{k-1}\left|m_{t}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k} \mathbf{s}\right)} d \mathbf{s} \\
& =\frac{2^{k-1}}{(2 \pi)^{d}} \int_{F} \prod_{j=0}^{k-2}\left|m_{t}\left(\left(A^{\tau}\right)^{j} \mathbf{s}\right)\right|^{2} e^{-i \mathbf{n} \circ\left(\left(A^{\tau}\right)^{k-1} \mathbf{s}\right)} d \mathbf{s}=\int_{\mathbb{R}^{d}}\left|p_{t, k-1}(\mathbf{s})\right|^{2} e^{-i \mathbf{n} \circ \mathbf{s}} d \mathbf{s}
\end{aligned}
$$

Repeating the above procedure then leads to

$$
\begin{aligned}
& \int_{\mathbb{R}^{d}}\left|p_{t, k}(\mathbf{s})\right|^{2} e^{-i \mathbf{n o s}} d \mathbf{s} \\
& \quad=\int_{\mathbb{R}^{d}}\left|p_{t, 1}(\mathbf{s})\right|^{2} e^{-i \mathbf{n o s}} d \mathbf{s} \\
& \quad=\frac{2}{(2 \pi)^{d}} \int_{F}\left|m_{t}(\mathbf{s})\right|^{2} e^{-i \mathbf{n} \circ\left(A^{\tau} \mathbf{s}\right)} d \mathbf{s}=\frac{2}{(2 \pi)^{d}} \int_{F_{1}} e^{-i \mathbf{n} \circ\left(A^{\tau} \mathbf{s}\right)} d \mathbf{s} \\
& \quad=\frac{1}{(2 \pi)^{d}} \int_{F} e^{-i \mathbf{n o s}} d \mathbf{s}=\frac{1}{(2 \pi)^{d}} \int_{\Omega} e^{-i \mathbf{n} \mathbf{0} \mathbf{s}} d \mathbf{s}=\delta_{\mathbf{n}, \mathbf{0}}
\end{aligned}
$$

So $\left\|p_{t, k}\right\|^{2}=1$. Clearly, $\lim _{k \rightarrow \infty} p_{t, k}(\mathbf{s})=\widehat{\phi_{t}}(\mathbf{s})$ for all $\mathbf{s} \in \mathbb{R}^{d}$. Thus, $\phi_{t} \in$ $L^{2}\left(\mathbb{R}^{d}\right)$ by Fatou's lemma. Since $p_{t, k}(\mathbf{s})$ is dominated by $\widehat{\phi_{t}}(\mathbf{s}) /(1-t)^{k_{0}^{\prime} / 2}$,

$$
\lim _{k \rightarrow \infty} \int_{\mathbb{R}^{d}}\left|p_{t, k}(\mathbf{s})\right|^{2} e^{-i \mathbf{n o s}} d \mathbf{s}=\int_{\mathbb{R}^{d}}\left|\widehat{\phi}_{t}(\mathbf{s})\right|^{2} e^{-i \mathbf{n o s}} d \mathbf{s}=\delta_{\mathbf{n}, \mathbf{0}}
$$

by Lebesgue's dominated convergence theorem. This is equivalent to

$$
\sum_{\ell \in \mathbb{Z}^{d}}\left|\widehat{\phi}_{t}(\mathbf{s}+2 \pi \ell)\right|^{2}=1 /(2 \pi)^{d}
$$

By Lemma 2.3, $\phi_{t}$ is an MRA scaling function.
Lemma 4.4. For each fixed $t_{0} \in[0,1]$, the mapping $t \mapsto m_{t}(\mathbf{s})$ is continuous at $t_{0}$ for almost all $\mathbf{s} \in \mathbb{R}^{d}$. That is, there exists a measure zero set $N_{t_{0}} \subset \mathbb{R}^{d}$ such that for each $\mathbf{s} \in \mathbb{R}^{d} \backslash N_{t_{0}}, \forall \varepsilon>0, \exists \delta>0$ such that $\left|m_{t}(\mathbf{s})-m_{t_{0}}(\mathbf{s})\right|<\varepsilon$ for all $t \in[0,1] \cap\left(t_{0}-\delta, t_{0}+\delta\right)$.

Proof. We shall only prove the case for $0<t_{0}<1$ and leave the cases $t_{0}=0$ and $t_{0}=1$ (which are simpler than the case of $0<t_{0}<1$ ) to the reader. Recall from the definition of the set $H_{t}$ that $H_{t_{2}} \subset H_{t_{1}}$ (hence $F_{1} \backslash H_{t_{1}} \subset F_{1} \backslash H_{t_{2}}$ ) whenever $t_{2} \geq t_{1}$ and that $\mu\left(H_{t}\right)=(1-t) \mu\left(F_{1}\right)$.

For any $t_{0} \in(0,1)$, define

$$
M_{t_{0}}=\left(\bigcap_{t>t_{0}}\left(\left(F_{1} \backslash H_{t}\right) \cap H_{t_{0}}\right)\right) \cup\left(\bigcap_{t<t_{0}}\left(H_{t}\right) \cap\left(F_{1} \backslash H_{t_{0}}\right)\right)
$$

It is easy to check that for any $t>t_{0}$, we have $\mu\left(\left(F_{1} \backslash H_{t}\right) \cap H_{t_{0}}\right)=\mu\left(H_{t_{0}} \backslash\right.$ $\left.H_{t}\right)=\left(t-t_{0}\right) \mu\left(F_{1}\right)$. Similarly, for any $t<t_{0}, \mu\left(H_{t} \cap\left(F_{1} \backslash H_{t_{0}}\right)\right)=\mu\left(H_{t} \backslash\right.$ $\left.H_{t_{0}}\right)=\left(t_{0}-t\right) \mu\left(F_{1}\right)$. Hence, $\mu\left(M_{t_{0}}\right) \leq 2\left|t-t_{0}\right| \mu\left(F_{1}\right)$ for any $t \neq t_{0}$. Since $t$ is arbitrary and $M_{t_{0}}$ is a fixed measurable set, we must have $\mu\left(M_{t_{0}}\right)=0$. Now define $N_{t_{0}}=\bigcup_{\ell \in \mathbb{Z}^{d}}\left(M_{t_{0}}+\pi \ell\right)$, which is also a measure zero set.

For any $\mathbf{s} \in \mathbb{R}^{d} \backslash N_{t_{0}}$, there exists a (unique) $\mathbf{s}_{0} \in F$ such that $\mathbf{s}-\mathbf{s}_{0} \in$ $2 \pi \mathbb{Z}^{d}, \mathbf{s}_{0} \in F_{1} \backslash M_{t_{0}}$ or $\mathbf{s}_{0} \in F_{0} \backslash M_{t_{0}}^{\prime}$ where $M_{t_{0}}^{\prime}$ is the subset in $F_{0}$ that is $2 \pi \mathbb{Z}^{d}$ translation congruent to $M_{t_{0}}+2 \pi \mathbf{h}_{2}$. Since $m_{t}$ is $2 \pi \mathbb{Z}^{d}$ translation periodic, we have $m_{t}(\mathbf{s})=m_{t}\left(\mathbf{s}_{0}\right)$. Consider first the case that $\mathbf{s}_{0} \in F_{1} \backslash M_{t_{0}}$. By the definition of $M_{t_{0}}$, there exist $t_{1}<t_{0}<t_{2}$ such that $\mathbf{s}_{0} \notin H_{t_{1}} \backslash H_{t_{2}}$. Choose $\delta>0$ small enough so that $\delta<\min \left\{t_{0}-t_{1}, t_{2}-t_{0}\right\}$ and $\mid \sqrt{1-t+b t}-$ $\sqrt{1-t_{0}+b t_{0}} \mid<\varepsilon$ for any $t \in(0,1) \cap\left(t_{0}-\delta, t_{0}+\delta\right)$ and any $0 \leq b \leq 1$. Then for any such $t$ we have $\mathbf{s}_{0} \in H_{t_{2}} \subset H_{t_{0}}$ or $\mathbf{s}_{0} \in F_{1} \backslash H_{t_{1}} \subset F_{1} \backslash H_{t_{0}}$. In the first case, $\left|m_{t}\left(\mathbf{s}_{0}\right)-m_{t_{0}}\left(\mathbf{s}_{0}\right)\right|=0$ and in the second case $\left|m_{t}\left(\mathbf{s}_{0}\right)-m_{t_{0}}\left(\mathbf{s}_{0}\right)\right|=$ $\left|\sqrt{1-t+t m_{1}^{2}\left(\mathbf{s}_{0}\right)}-\sqrt{1-t_{0}+t_{0} m_{1}^{2}\left(\mathbf{s}_{0}\right)}\right|<\varepsilon$ by the choice of $\delta$. The case of $\mathbf{s}_{0} \in F_{0} \backslash M_{t_{0}}^{\prime}$ can be similarly proven.

The result of Lemma 4.4 can be easily extended to the following corollary. We leave the proof to the reader.

Corollary 4.1. For each fixed $t_{0} \in[0,1]$ and each fixed $k \geq 1$, the mapping $t \mapsto \prod_{j=1}^{k} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)$ is continuous at $t_{0}$ for almost all $\mathbf{s} \in \mathbb{R}^{d}$. That is, there exists a measure zero set $N_{t_{0}}^{k} \subset \mathbb{R}^{d}$ such that for each $\mathbf{s} \in \mathbb{R}^{d} \backslash N_{t_{0}}^{k}$, $\forall \varepsilon>0, \exists \delta>0$ such that $\left|\prod_{j=1}^{k} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)-\prod_{j=1}^{k} m_{t_{0}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)\right|<\varepsilon$ for all $t \in[0,1] \cap\left(t_{0}-\delta, t_{0}+\delta\right)$.

Lemma 4.5. For any $t_{0} \in[0,1], \lim _{t \rightarrow t_{0}} \widehat{\phi_{t}}(\mathbf{s})=\widehat{\phi_{t_{0}}}(\mathbf{s})$ for $\mathbf{s} \in \mathbb{R}^{d}$. More precisely, for each fixed $t_{0} \in[0,1]$, there exists a measure zero set $N_{t_{0}}^{\prime} \subset \mathbb{R}^{d}$ such that for any $\mathbf{s} \in \mathbb{R}^{d} \backslash N_{t_{0}}^{\prime}$, we have $\lim _{t \rightarrow t_{0}} \widehat{\phi_{t}}(\mathbf{s})=\widehat{\phi_{t_{0}}}(\mathbf{s})$.

Proof. Since $\widehat{\phi_{1}} \geq 0, \lim _{j \rightarrow \infty} \widehat{\phi_{1}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)=1 /(2 \pi)^{\frac{d}{2}}$. For any given $\varepsilon>0$ and $\mathbf{s} \in \mathbb{R}^{d}$ (modular a zero measure set), there exists a positive integer $n_{0}$ such that $\widehat{\phi_{1}}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)>1 /(2 \pi)^{\frac{d}{2}}-\varepsilon / 2$ and $\left(A^{\tau}\right)^{-n} \mathbf{S} \subset\left(A^{\tau}\right)^{-1} F$ for any $n \geq n_{0}$. It follows that $m_{t}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)$ is either 1 or $\sqrt{(1-t)+t m_{1}^{2}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)}$ for any $t \in[0,1]$. In either case, $m_{t}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right) \geq m_{1}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)$, thus $\widehat{\phi}_{t}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right) \geq$ $\widehat{\phi_{1}}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)$ for any $t \in[0,1]$. Since $\widehat{\phi_{t}}\left(\mathbf{s}^{\prime}\right) \leq 1 /(2 \pi)^{\frac{d}{2}}$ for any $\mathbf{s}^{\prime} \in \mathbb{R}^{d}$ by its definition, it follows that for any $t_{1}, t_{2} \in[0,1]$, we have

$$
\begin{equation*}
\left|\widehat{\phi_{t_{1}}}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)-\widehat{\phi_{t_{2}}}\left(\left(A^{\tau}\right)^{-n} \mathbf{s}\right)\right|<\varepsilon / 2 . \tag{4.6}
\end{equation*}
$$

On the other hand, by Corollary 4.1, for each $t_{0} \in[0,1]$, there exists a zero measure set $N_{t_{0}}^{n_{0}} \subset \mathbb{R}^{d}$ such that for each $\mathbf{s} \in \mathbb{R}^{d} \backslash N_{t_{0}}^{n_{0}}, \exists \delta>0$ such that $\left|\prod_{j=1}^{n_{0}} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)-\prod_{j=1}^{n_{0}} m_{t_{0}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)\right|<\varepsilon / 2$ for all $t \in[0,1] \cap\left(t_{0}-\delta\right.$,
$\left.t_{0}+\delta\right)$. Hence, $\left|\widehat{\phi_{t}}(\mathbf{s})-\widehat{\phi_{t_{0}}}(\mathbf{s})\right|$ is bounded by

$$
\begin{aligned}
& \left|\prod_{j=1}^{n_{0}} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right) \widehat{\phi_{t}}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right)-\prod_{j=1}^{n_{0}} m_{t_{0}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right) \widehat{\phi_{t_{0}}}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right)\right| \\
& \quad=\mid \prod_{j=1}^{n_{0}} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right) \cdot \widehat{\phi}_{t}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right)-\prod_{j=1}^{n_{0}} m_{t_{0}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right) \widehat{\phi_{t}}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right) \\
& \quad+\prod_{j=1}^{n_{0}} m_{t_{0}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right) \widehat{\phi_{t}}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right)-\prod_{j=1}^{n_{0}} m_{t_{0}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right) \widehat{\phi_{t_{0}}}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right) \mid \\
& \quad \leq(2 \pi)^{-\frac{d}{2}}\left|\prod_{j=1}^{n_{0}} m_{t}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)-\prod_{j=1}^{n_{0}} m_{t_{0}}\left(\left(A^{\tau}\right)^{-j} \mathbf{s}\right)\right| \\
& \quad+\left|\widehat{\phi}_{t}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right)-\widehat{\phi_{t_{0}}}\left(\left(A^{\tau}\right)^{-n_{0}} \mathbf{s}\right)\right|<\frac{\varepsilon}{2}+\frac{\varepsilon}{2}=\varepsilon .
\end{aligned}
$$

So $\lim _{t \rightarrow t_{0}} \widehat{\phi_{t}}(\mathbf{s})=\widehat{\phi_{t_{0}}}(\mathbf{s})$.
By the continuity of $m_{t}(\mathbf{s})$ and $\widehat{\phi}_{t}$, we now have $\lim _{t \rightarrow t_{0}} \widehat{\psi_{t}}(\mathbf{s})=\widehat{\psi_{t_{0}}}(\mathbf{s})$ for almost every $\mathbf{s} \in \mathbb{R}^{d}$.

Lemma 4.6. For $t_{0}, t \in[0,1], \lim _{t \rightarrow t_{0}}\left\|\widehat{\psi_{t}}-\widehat{\psi_{t_{0}}}\right\|^{2}=0$.
Proof. Since $\left\|\widehat{\psi_{t}}\right\|^{2}=\left\|\widehat{\psi_{t_{0}}}\right\|^{2}=1,\left\|\widehat{\psi_{t}}-\widehat{\psi_{t_{0}}}\right\|^{2}=\left\langle\widehat{\psi_{t}}-\widehat{\psi_{t_{0}}}, \widehat{\psi_{t}}-\widehat{\psi_{t_{0}}}\right\rangle=2-$ $\left\langle\widehat{\psi_{t}}, \widehat{\psi_{t_{0}}}\right\rangle-\left\langle\widehat{\psi_{t_{0}}}, \widehat{\psi_{t}}\right\rangle$. Thus, it suffices to show that $\lim _{t \rightarrow t_{0}}\left\langle\widehat{\psi_{t}}, \widehat{\psi_{t_{0}}}\right\rangle=1$.

Since $\widehat{\psi_{t_{0}}} \in L^{2}\left(\mathbb{R}^{d}\right)$, for any $\varepsilon>0$, there exists a sufficiently large number $r>0$ such that $\left(\int_{|\mathbf{s}|>r}\left|\widehat{\psi_{t_{0}}}(\mathbf{s})\right|^{2} d \mathbf{s}\right)^{\frac{1}{2}}<\varepsilon / 4$. By Hölder's Inequality, we then have $\int_{|\mathbf{s}|>r}\left|\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right| \cdot\left|\widehat{\psi_{t_{0}}}(\mathbf{s})\right| d \mathbf{s} \leq\left\|\widehat{\psi_{t}}(\widehat{\mathbf{s}})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right\|\left(\int_{|\mathbf{s}|>r}\left|\widehat{\psi_{t_{0}}}(\mathbf{s})\right|^{2} d \mathbf{s}\right)^{\frac{1}{2}}<$ $\varepsilon / 2$ since $\left\|\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right\| \leq\left\|\widehat{\psi_{t}}(\mathbf{s})\right\|+\left\|\widehat{\psi_{t_{0}}}(\mathbf{s})\right\|=2$. On the other hand, $\left|\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}(\mathbf{s})}\right| \leq 1 /(\pi)^{d / 2}$ since $\left|\widehat{\psi_{t}}(\mathbf{s})\right| \leq 1 /(2 \pi)^{\frac{d}{2}}$ and $\left|\widehat{\psi_{t_{0}}(\mathbf{s})}\right| \leq 1 /(2 \pi)^{\frac{d}{2}}$ by (4.4), (4.5) and the fact that $\left|m_{t}\right| \leq 1$ for any $t$. Thus, by the dominated convergence theorem, we have $\lim _{t \rightarrow t_{0}} \int_{|\mathbf{s}| \leq r}\left|\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right| d \mathbf{s}=0$. Therefore, there exists a number $\delta>0$ such that $\int_{|\mathbf{s}| \leq r}\left|\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right| d \mathbf{s}<\varepsilon / 2$ whenever $\left|t-t_{0}\right|<\delta$. Combining the above leads to

$$
\begin{aligned}
\left|\left\langle\widehat{\psi_{t}}, \widehat{\psi_{t_{0}}}\right\rangle-1\right|= & \left|\left\langle\widehat{\psi_{t}}, \widehat{\psi_{t_{0}}}\right\rangle-\left\langle\widehat{\psi_{t_{0}}}, \widehat{\psi_{t_{0}}}\right\rangle\right|=\left|\int_{\mathbb{R}^{d}}\left(\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right) \cdot \widehat{\widehat{\psi_{0}}}(\mathbf{s}) d \mathbf{s}\right| \\
\leq & \int_{|\mathbf{s}| \leq r}\left|\left(\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right) \widehat{\psi_{t_{0}}}(\mathbf{s})\right| d \mathbf{s} \\
& +\int_{|\mathbf{s}|>r}\left|\left(\widehat{\psi_{t}}(\mathbf{s})-\widehat{\psi_{t_{0}}}(\mathbf{s})\right) \widehat{\widehat{\psi_{t_{0}}}(\mathbf{s})}\right| d \mathbf{s}<\varepsilon .
\end{aligned}
$$

So $\lim _{t \rightarrow t_{0}}\left\|\widehat{\psi_{t}}-\widehat{\psi_{t_{0}}}\right\|^{2}=0$.

Since the inverse Fourier transform preserves norm, it follows that the mapping $t \mapsto \psi_{t}$ is continuous in the $L^{2}\left(\mathbb{R}^{d}\right)$ norm topology. This completes the proof of the connectedness theorem.

## 5. Further discussions

The path-connectivity for the set of all 2-dilation MRA wavelets was first established in $[11,14]$ for the one dimensional case. In [10], the authors solved the path-connectivity problem for all matrices in $M_{2}^{(2)}(\mathbb{Z})$. The approach used in [10] depends on a special property that a matrix in $M_{2}^{(2)}(\mathbb{Z})$ possesses. The purpose of this section is to show that a matrix in $M_{3}^{(2)}(\mathbb{Z})$ may no longer possess this property. Consequently, the argument used in this paper to establish the path-connectivity of all $A$-dilation MRA wavelets for any $A \in M_{d}^{(2)}(\mathbb{Z})$ is not a simple generalization of the earlier approaches.

Two $d \times d$ integral matrices $B$ and $C$ are said to be integrally similar if there exists an integral $d \times d$ matrix $P$ such that $|\operatorname{det}(P)|=1$ and $P^{-1} B P=C$. The integral similarity then defines an equivalent relation among matrices of $M_{d}^{2}(\mathbb{Z})$. For $d=2$, there are exactly six integrally similar classes in $M_{2}^{(2)}(\mathbb{Z})$ [8]. A representative from each of these classes is listed below.

$$
\left(\begin{array}{ll}
0 & 2 \\
1 & 0
\end{array}\right),\left(\begin{array}{cc}
0 & 2 \\
-1 & 0
\end{array}\right),\left(\begin{array}{cc}
1 & 1 \\
-1 & 1
\end{array}\right),\left(\begin{array}{cc}
-1 & -1 \\
1 & -1
\end{array}\right),\left(\begin{array}{cc}
0 & 2 \\
-1 & 1
\end{array}\right),\left(\begin{array}{ll}
0 & -2 \\
1 & -1
\end{array}\right) .
$$

An important point about the above representatives is that each of them has the property that $\left(A^{\tau}\right)^{-1} \Omega \subset \Omega$ where $\Omega=[-\pi, \pi)^{2}$. For an expansive matrix with this property, one could then employ the nice geometric structures of $\Omega$ and $\left(A^{\tau}\right)^{-1} \Omega$ in the construction of $m_{t}$ (and $\left.\widehat{\phi_{t}}(\mathbf{s}), \widehat{\psi_{t}}(\mathbf{s})\right)$. In particular, a general Shannon-type wavelet can be visualized in these cases. For a matrix in $M_{2}^{(2)}(\mathbb{Z})$ that does not have this property, one can use its representative in the above list (in its equivalent class) and use the following useful theorem (proved in [10]).

Theorem 5.1 ([10]). For any $2 \times 2$ integral matrix $P$ with $|\operatorname{det} P|=1$, let $\Phi_{P}: L^{2}\left(\mathbb{R}^{2}\right) \longrightarrow L^{2}\left(\mathbb{R}^{2}\right)$ be the operator defined by $\Phi_{P}(g(\mathbf{t}))=g(P \mathbf{t})$. If $B$ and $C$ are two $2 \times 2$ integral, expansive matrices such that $P^{-1} B P=C$, then the following statements hold:
(i) $\psi$ is a $B$-dilation wavelet iff $\Phi_{P}(\psi)$ is a $C$-dilation wavelet;
(ii) A function $f \in L^{2}\left(\mathbb{R}^{2}\right)$ is a $B$-dilation wavelet multiplier iff the function $\Phi_{\left(P^{\tau}\right)^{-1}}(f)$ is a $C$-dilation wavelet multiplier.

The linear operator $\Phi_{P}: L^{2}\left(\mathbb{R}^{2}\right) \longrightarrow L^{2}\left(\mathbb{R}^{2}\right)$ defined above is obviously continuous and unitary (since $|\operatorname{det} P|=1$ ). In the case that $P$ is also integral and $P^{-1} B P=C$, then Theorem 5.1 asserts that $\Phi_{P}: \mathcal{W}_{B} \longrightarrow \mathcal{W}_{C}$ is a continuous and bijective mapping, where $\mathcal{W}_{B}$ is the set of all $B$-dilation wavelets and $\mathcal{W}_{C}$ is the set of all $C$-dilation wavelets. $\Phi_{P}$ is also a bijection between the set of all $B$-dilation MRA wavelets and the set of all $C$-dilation MRA wavelets. These observations imply that at least for the two dimensional case, the path-connectivity problem of all $B$-dilation MRA wavelets for any $B \in M_{2}^{(2)}(\mathbb{Z})$ is equivalent to the path-connectivity problem of all $A$-dilation MRA wavelets where $A$ is one of the six matrices listed above. This is precisely what these authors did in [10]: they established the path-connectivity of all $A$-dilation MRA wavelets where $A$ is one of the six matrices listed above, which then implied that all $B$-dilation MRA wavelets are path-connected for any $B \in M_{2}^{(2)}(\mathbb{Z})$. A different way to interpret the above theorem (or the classifications of $M_{2}^{(2)}(\mathbb{Z})$ matrices into the 6 integrally similar equivalent classes) is that for any matrix $A$ in $M_{2}^{(2)}(\mathbb{Z})$, a general Shannon-type $A$-dilation MRA wavelet can be constructed using a set $F$ where $F$ has the form $P \Omega$ for some $P \in M_{2}^{(1)}(\mathbb{Z})$. Of course, such explicit expression for an MRA wavelet would be desirable. In some sense, the approach in [10] made the maximum use of this property of $M_{2}^{(2)}(\mathbb{Z})$.

Let us now consider the possibility of extending that approach in the case of $d=3$. In other words, one would like to establish the following:
(1) Identify all integrally similar equivalent classes of matrices in $M_{3}^{(2)}(\mathbb{Z})$;
(2) Show that in each such class there is a representative matrix $A$ with the property that $\left(A^{\tau}\right)^{-1} \Omega \subset \Omega$ where $\Omega=[-\pi, \pi)^{3}$;
(3) Explore the structure of $\left(A^{\tau}\right)^{-1} \Omega$ and the possibility of using this structure in the definition of $m_{t}(\mathbf{s})$ (and $\left.\phi_{t}, \psi_{t}\right)$ that would lead to the establishment of path-connectivity of all $A$-dilation MRA wavelets for such matrix $A$.

Unfortunately, the first task seems to be a very difficult problem. The authors failed to find an answer to this question in the literature. The third task, even when successfully carried out, can only solve the problem for some matrices in $M_{3}^{(2)}(\mathbb{Z})$. The biggest problem turned out to be task 2: there are integrally similar equivalent classes in $M_{3}^{(2)}(\mathbb{Z})$ which do not have any representative $A$ with the property $\left(A^{\tau}\right)^{-1} \Omega \subset \Omega$. The authors do not intend to elaborate the details here. Instead, we will just list some of our findings about $M_{3}^{(2)}(\mathbb{Z})$.

Firstly, there are 14 similar equivalent classes (but we are not sure if they are all the integrally similar classes) in $M_{3}^{(2)}(\mathbb{Z})$. A representative from each
class is listed below.

$$
\begin{gathered}
\left(\begin{array}{ccc}
0 & -1 & -1 \\
0 & -1 & 1 \\
1 & 0 & 0
\end{array}\right),\left(\begin{array}{ccc}
0 & 0 & -1 \\
2 & -1 & 0 \\
0 & 1 & 0
\end{array}\right),\left(\begin{array}{ccc}
0 & -1 & 1 \\
0 & -1 & -1 \\
1 & 0 & 0
\end{array}\right),\left(\begin{array}{lll}
0 & 0 & 1 \\
2 & 0 & 0 \\
0 & 1 & 0
\end{array}\right), \\
\left(\begin{array}{lll}
0 & 1 & 0 \\
0 & 0 & 1 \\
2 & 1 & 0
\end{array}\right),\left(\begin{array}{ccc}
0 & 1 & 0 \\
0 & 0 & -1 \\
2 & -1 & 0
\end{array}\right),\left(\begin{array}{ccc}
1 & -1 & 0 \\
0 & 0 & -1 \\
1 & 1 & 0
\end{array}\right),\left(\begin{array}{ccc}
1 & 0 & 1 \\
1 & 0 & -1 \\
0 & -1 & 0
\end{array}\right) \\
\left(\begin{array}{lll}
0 & 0 & 1 \\
2 & 1 & 0 \\
0 & 1 & 0
\end{array}\right),\left(\begin{array}{ccc}
0 & 2 & 0 \\
0 & 0 & -1 \\
1 & 0 & 0
\end{array}\right),\left(\begin{array}{ccc}
-1 & -1 & -1 \\
0 & -1 & 1 \\
1 & 0 & 0
\end{array}\right),\left(\begin{array}{ccc}
-1 & -1 & -1 \\
-1 & 0 & 0 \\
0 & -1 & 1
\end{array}\right), \\
\left(\begin{array}{ccc}
-1 & -1 & 0 \\
0 & 0 & -1 \\
1 & -1 & 1
\end{array}\right),\left(\begin{array}{ccc}
0 & -1 & -1 \\
0 & 1 & -1 \\
1 & 0 & 1
\end{array}\right) .
\end{gathered}
$$

Secondly, of these 14 classes, 4 of them do not have any representative $A$ with the property $\left(A^{\tau}\right)^{-1} \Omega \subset \Omega$. The last four matrices listed above are the representatives of these four classes. It would be interesting to know if one can find a bounded set $F$ (as used in the definition of a general Shannon type wavelet) that is also geometrically nice and simple.

Thirdly, the 10 other representatives in the above list all have the property $\left(A^{\tau}\right)^{-1} \Omega \subset \Omega$. It is interesting to note that within each of these 10 classes, there are exactly 24 matrices with this property. For any other matrix in one of these classes, it is not clear whether it is integrally similar to one of these 24 representatives. We suspect that is the case.

Finally, for $d \geq 4$, the situation will be even more complicated and it is plausible that similar equivalent classes like the last 4 in the above list exist. Since our findings above are obtained through exhaustive search, the method cannot be easily generalized to handle the more general cases in higher dimension. This would be a problem for future study.

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