

Tracy-Widom asymptotics for a random polymer model with gamma-distributed weights

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Abstract

We establish Tracy-Widom asymptotics for the partition function of a random polymer model with gamma-distributed weights recently introduced by Seppäläinen. We show that the partition function of this random polymer can be represented within the framework of the geometric RSK correspondence and consequently its law can be expressed in terms of Whittaker functions. This leads to a representation of the law of the partition function which is amenable to asymptotic analysis. In this model, the partition function plays a role analogous to the smallest eigenvalue in the Laguerre unitary ensemble of random matrix theory.

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1 Introduction

Denote by $\Phi_{m,n}$ the set of ‘paths’ of the form $\phi = \{(1, j_1), (2, j_2), \dots, (m, j_m)\}$, where $1 \leq j_1 \leq \dots \leq j_m \leq n$, as shown in Figure 1. Let g_{ij} be independent gamma-distributed random variables with common parameter γ , i.e.

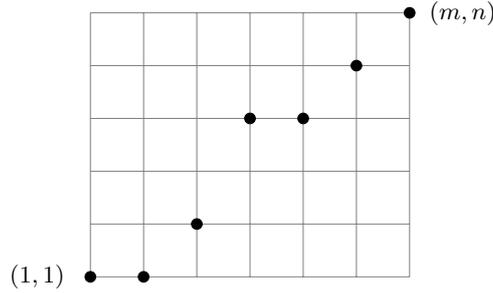
$$\mathbb{P}\{g_{ij} \in dx\} = \frac{1}{\Gamma(\gamma)} x^{\gamma-1} e^{-x} dx$$

and set

$$Z_{m,n} = \sum_{\phi \in \Phi_{m,n}} \prod_{(i,j) \in \phi} g_{ij}.$$

This is the partition function of a random polymer recently introduced by Seppäläinen [16] where it was observed that this model exhibits the so-called Burke property. The analogous property for other polymer models, specifically the semi-discrete Brownian polymer introduced in [14] and the log-gamma polymer introduced in [17], has

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Figure 1: A path in $\Phi_{m,n}$.

been used to study asymptotics of the free energy [10, 14, 17, 18]. More recently, the semi-discrete and log-gamma polymer models have been shown to have an underlying integrable structure, via a remarkable connection between a combinatorial structure known as the geometric RSK correspondence and $GL(n, \mathbb{R})$ -Whittaker functions [6, 12, 13]. This integrable structure has allowed very precise (Tracy-Widom) asymptotics to be obtained [3–5]. For these models, the partition functions play a role analogous to the largest eigenvalue in the Gaussian and Laguerre unitary ensembles of random matrix theory.

In the present paper, we show that the partition function of the above random polymer can also be represented within the framework of the geometric RSK correspondence and consequently its law can be expressed in terms of Whittaker functions. For this model, the partition function plays a role analogous to the *smallest* eigenvalue in the Laguerre unitary ensemble. This leads to a representation of the law of the partition function from which we establish Tracy-Widom asymptotics for this model. A precise statement is given as follows.

Theorem 1.1. *Suppose $m/n \rightarrow \alpha > 0$ as $n \rightarrow \infty$. Set $c = 1 + \alpha$,*

$$\mu = \inf_{z>0} [c\psi'(z + \gamma) - \psi'(z)], \quad H(z) = \ln \Gamma(z) - c \ln \Gamma(z + \gamma) + \mu z,$$

where ψ is the digamma function. The infimum in the definition of μ is achieved at some $z^* > 0$ and $\bar{g} := -H'''(z^*) > 0$. For γ sufficiently small,

$$\lim_{n \rightarrow \infty} \mathbb{P} \left\{ \frac{\ln Z_{m,n} - n\mu}{n^{1/3}} \leq r \right\} = F_{GUE} \left((\bar{g}/2)^{-1/3} r \right)$$

where F_{GUE} is the Tracy-Widom distribution function.

The connection to random matrices can be further illustrated by considering the zero-temperature limit, which corresponds to letting $\gamma \rightarrow 0$. Then the collection of random variables $-\gamma \log g_{ij}$ converge weakly to a collection of independent standard exponentially distributed variables w_{ij} and so, by the principle of the largest term, the sequence $-\gamma \log Z_{m,n}$ converges weakly to the first passage percolation variable

$$f_{m,n} = \min_{\phi \in \Phi_{m,n}} \sum_{(i,j) \in \phi} w_{ij}.$$

This first passage percolation problem was previously considered in [11] where it is argued, using a representation of $f_{m,n}$ as a departure process from a series of ‘Exp/Exp/1’ queues in tandem together with the Burke property for such queues, that, almost surely,

$$\lim_{n \rightarrow \infty} f_{\alpha n, n}/n = (\sqrt{1 + \alpha} - 1)^2. \quad (1.1)$$

Moreover, it can be inferred from further results presented in [8] on a discrete version of this model with geometric weights (or alternatively from Section 2 below) that $f_{m,n}$ has the same law as the smallest eigenvalue in the Laguerre ensemble with density proportional to

$$\prod_{1 \leq i < j \leq n} (\lambda_i - \lambda_j)^2 \prod_{i=1}^n \lambda_i^{m-1} e^{-\lambda_i} d\lambda_i.$$

Given this identity in law, the asymptotic relation (1.1) can also be seen as a consequence of the Marchenko-Pastur law. As a further consistency check, one can easily verify (see Lemma 5.2 below) that

$$-\gamma\mu \rightarrow (\sqrt{1+\alpha} - 1)^2$$

as $\gamma \rightarrow 0$, where μ is defined in the statement of Theorem 1.1.

The outline of the paper is as follows. In the next section we relate the above polymer model to the geometric RSK correspondence and deduce, using results from [6, 13], an integral formula for the Laplace transform of the partition function. In Section 3 we show that this Laplace transform can be written as a Fredholm determinant, which allows us, in Section 4 to take the limit as $n \rightarrow \infty$. Section 5 contains proofs of some lemmas that we require on the way.

After the present article appeared on the arXiv, a paper by Corwin, Seppäläinen and Shen appeared there [7]. They obtain similar results, however, their approach is very different to ours. Namely they show that the polymer model can be obtained as a limit of the discrete-time q -TASEP and then use previously known formulas for that particle system.

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2 Geometric RSK, polymers and Whittaker functions

The geometric RSK correspondence is a bijective mapping

$$T : (\mathbb{R}_{>0})^{h \times n} \rightarrow (\mathbb{R}_{>0})^{h \times n}.$$

It was introduced by Kirillov [9] as a geometric lifting of the RSK correspondence, and is defined as follows. Let $W = (w_{ij}) \in (\mathbb{R}_{>0})^{h \times n}$ and write $T(W) = (t_{ij}) \in (\mathbb{R}_{>0})^{h \times n}$. For $1 \leq k \leq n$ and $1 \leq r \leq h \wedge k$,

$$t_{h-r+1, k-r+1} \cdots t_{h-1, k-1} t_{hk} = \sum_{(\pi_1, \dots, \pi_r) \in \Pi_{h,k}^{(r)}} \prod_{(i,j) \in \pi_1 \cup \dots \cup \pi_r} w_{ij}, \quad (2.1)$$

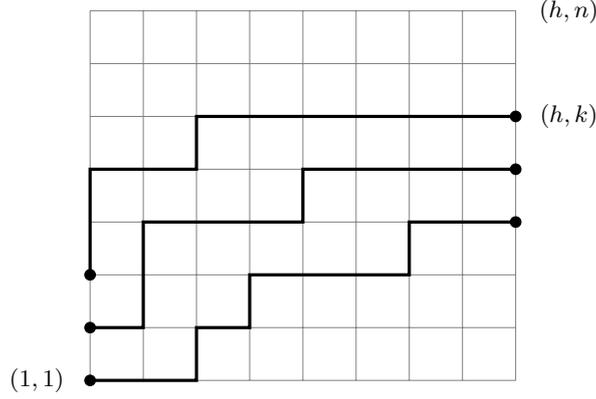
where $\Pi_{h,k}^{(r)}$ denotes the set of r -tuples of non-intersecting up/right lattice paths π_1, \dots, π_r starting at positions $(1, 1), (1, 2), \dots, (1, r)$ and ending at positions $(h, k-r+1), \dots, (h, k-1), (h, k)$, as shown in Figure 2. The remaining entries of $T(W)$ are determined by the relation $T(W^t) = T(W)^t$.

Note in particular that

$$t_{hn} = \sum_{\pi \in \Pi_{h,n}} \prod_{(i,j) \in \pi} w_{ij}, \quad (2.2)$$

where $\Pi_{h,n}$ is the set of up/right lattice paths in \mathbb{Z}^2 from $(1, 1)$ to (h, n) . This gives an interpretation of t_{hn} as a polymer partition function, providing the basis for the analysis of the log-gamma polymer developed in [6, 13].

The relation to the random polymer defined in the introduction is as follows.


 Figure 2: A 3-tuple of non-intersecting paths in $\Pi_{h,k}^{(3)}$.

Proposition 2.1. Suppose $h \geq n$ and set $m = h - n + 1$. For $1 \leq i \leq m$ and $1 \leq j \leq n$, set $g_{ij} = 1/w_{i+j-1, n-j+1}$. Then

$$\frac{1}{t_{m1}} = \sum_{\phi \in \Phi_{m,n}} \prod_{(i,j) \in \phi} g_{ij}, \quad (2.3)$$

where $\Phi_{m,n}$ the set of $\phi = \{(1, j_1), (2, j_2), \dots, (m, j_m)\}$ with $1 \leq j_1 \leq \dots \leq j_m \leq n$.

Proof. From the definition (2.1), taking $k = r = n$,

$$t_{m1} \cdots t_{h-1, n-1} t_{hn} = \sum_{(\pi_1, \dots, \pi_n) \in \Pi_{h,n}^{(n)}} \prod_{(i,j) \in \pi_1 \cup \dots \cup \pi_n} w_{ij} = \prod_{i,j} w_{ij},$$

and, taking $k = n$ and $r = n - 1$,

$$t_{m+1,2} \cdots t_{h-1, n-1} t_{hn} = \sum_{(\pi_1, \dots, \pi_{n-1}) \in \Pi_{h,n}^{(n-1)}} \prod_{(i,j) \in \pi_1 \cup \dots \cup \pi_{n-1}} w_{ij}.$$

Thus,

$$\frac{1}{t_{m1}} = \sum_{(\pi_1, \dots, \pi_{n-1}) \in \Pi_{h,n}^{(n-1)}} \prod_{(i,j) \notin \pi_1 \cup \dots \cup \pi_{n-1}} \frac{1}{w_{ij}} = \sum_{\phi \in \Phi_{m,n}} \prod_{(i,j) \in \phi} g_{ij},$$

as required. The last identity is illustrated in Figures 1 and 3. The paths in $\Phi_{m,n}$ are obtained by taking compliments of $(n - 1)$ -tuples in $\Pi_{h,n}^{(n-1)}$ (as shown in Figure 3), reflecting this picture through the horizontal, and shifting appropriately to remove the gaps - this procedure is made precise in the above definition of the g_{ij} . \square

Remark 2.2. The identity (2.3) is analogous to Theorem 5.1, equation (5.4), of the paper [8], where the corresponding identity for the usual RSK correspondence is given.

Let $a \in \mathbb{R}^n$ and $b \in \mathbb{R}^h$ be such that $a_j + b_i > 0$ for all i, j . In [6] (here we are using the notation of [13]) it was shown that, if the matrix W is chosen at random according to the probability measure

$$\mathbb{P}(dW) = \prod_{i,j} \Gamma(a_j + b_i)^{-1} e^{-1/w_{ij}} w_{ij}^{-a_j - b_i - 1} dw_{ij} \quad (2.4)$$

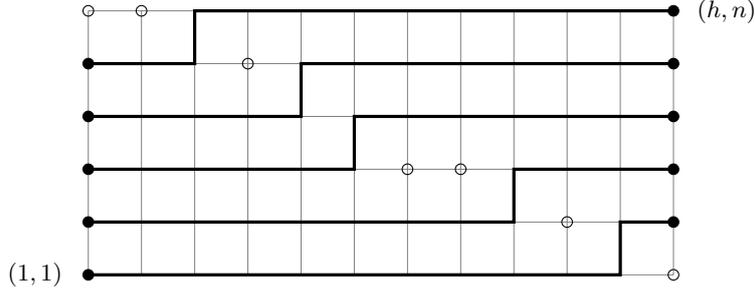


Figure 3: An $(n-1)$ -tuple of non-intersecting paths in $\Pi_{h,n}^{(n-1)}$, and its compliment. The corresponding path in $\Phi_{h-n+1,n}$ is shown in Figure 1.

then the law of the vector (t_{h1}, \dots, t_{m1}) under \mathbb{P} is given by

$$\mu_n(dx) = \prod_{i,j} \Gamma(a_j + b_i)^{-1} \Psi_a^n(x) \Psi_{b;1}^n(x) \prod_{j=1}^n \frac{dx_j}{x_j},$$

where Ψ_a^n and $\Psi_{b;1}^n$ are (generalised) Whittaker functions, as defined in [13]. Observe that with (2.4) the w_{ij} are independent random variables with inverse gamma distribution, so that the g_{ij} follow the gamma distribution. Without loss of generality we can assume that $a_j > 0$ and $b_i > 0$ for each i, j and deduce the following.

Proposition 2.3. For $s \in \mathbb{C}$ with $\Re s > 0$,

$$\mathbb{E} e^{-s/t_{m1}} = \int_{(\mathbb{R}_{>0})^n} e^{-s/x_n} \mu_n(dx) = \int_{(i\mathbb{R})^n} \prod_{i,j=1}^n \Gamma(a_i - \lambda_j) \prod_{j=1}^n \frac{s^{\lambda_j} \prod_{i=1}^h \Gamma(b_i + \lambda_j)}{s^{a_j} \prod_{i=1}^h \Gamma(b_i + a_j)} s_n(\lambda) d\lambda \quad (2.5)$$

where s_n is the density of the Sklyanin measure

$$s_n(\lambda) = \frac{1}{(2\pi i)^n n!} \prod_{i,j=1}^n \frac{1}{\Gamma(\lambda_i - \lambda_j)}. \quad (2.6)$$

Proof. By [13, Corollary 3.8] the functions $\Psi_{a;s}^n(x) \equiv e^{-s/x_n} \Psi_a^n(x)$ and $\Psi_{b;1}^n$ are both in $L_2((\mathbb{R}_{>0})^n, \prod_{j=1}^n dx_j/x_j)$ and, by [13, Corollary 3.5], for $\lambda \in (i\mathbb{R})^n$, we have

$$\int_{(\mathbb{R}_{>0})^n} \Psi_{b;1}^n(x) \Psi_{\lambda}^n(x) \prod_{j=1}^n \frac{dx_j}{x_j} = \prod_{i=1}^h \prod_{j=1}^n \Gamma(b_i + \lambda_j)$$

and

$$\int_{(\mathbb{R}_{>0})^n} \Psi_{a;s}^n(x) \Psi_{-\lambda}^n(x) \prod_{j=1}^n \frac{dx_j}{x_j} = s^{\sum_{j=1}^n (\lambda_j - a_j)} \prod_{i,j=1}^n \Gamma(a_i - \lambda_j).$$

The claim now follows from the Plancherel theorem for $GL(n)$ -Whittaker functions due to Wallach, noting that $\overline{\Psi_{\lambda}^n(x)} = \Psi_{-\lambda}^n(x)$ (see for example [13, Section 2]). \square

The Laplace transform of the partition function $Z_{m,n}$ of the random polymer (defined in the introduction) is obtained by setting $a_i = \epsilon$ and $b_j = \gamma - \epsilon$, where $0 < \epsilon < \gamma$, for in this case $Z_{m,n}$ is given by $1/t_{m1}$. We remark that the partition function for the log-gamma polymer can be defined on the same probability space as t_{hn} , given by the formula (2.2). The joint law of the two partition functions is thus given in terms of the joint first and last marginal of the probability measure μ_n .

3 Fredholm determinant representation

The first step in the proof of Theorem 1.1 is to write the right-hand side of (2.5) as a Fredholm determinant. A similar algebraic identity is proved in [5] (Theorem 2). However, this result is proved for what corresponds to the case $h = n$ in our notation, and moreover would require the poles of F_s to lie to the right of ℓ_{δ_2} . Our argument is an extension of the argument in [5] and follows the same main steps. We will sketch it below.

For $s \in \mathbb{R}$ we define a function F_s by

$$F_s(w) = s^w \prod_{j=1}^h \Gamma(b_j + w) \quad (3.1)$$

where s is a parameter to be chosen later. For $\delta > 0$ define $\ell_\delta = \delta + i\mathbb{R}$ and let C_δ be the circle centred at the origin of radius δ . Observe that the Proposition 3.1 holds for a general range of parameters a_j, b_j and not just the special choice we made at the end of section 2.

Proposition 3.1. *Let $\delta_1, \delta_2 > 0$ such that $\delta_1 < \delta_2 \wedge (1 - \delta_2)$. Suppose also that $|a_j| < \delta_1$ and $b_j > \delta_2$ for all j . Then*

$$\int e^{-s/x_n} \mu_n(dx) = \det(I + K_{n,r}^{LT})_{L^2(C_{\delta_1})} \quad (3.2)$$

where

$$K_{n,r}^{LT}(v, \tilde{v}) = \frac{1}{2\pi i} \int_{\ell_{\delta_2}} \frac{dw}{w - \tilde{v}} \frac{\pi}{\sin(\pi(v-w))} \frac{F_s(w)}{F_s(v)} \prod_{j=1}^n \frac{\Gamma(v - a_j)}{\Gamma(w - a_j)}. \quad (3.3)$$

The rest of this section is devoted to the proof of this proposition. We will begin with the right-hand side of (3.2) and show that it equals to the right-hand side of (2.5).

Step 1: Of course the right-hand side of (3.2) should be interpreted as a Fredholm series, namely

$$\det(I + K_{n,r}^{LT})_{L^2(C_{\delta_1})} = 1 + \sum_{k=1}^{\infty} \frac{1}{k!} \int \cdots \int_{C_{\delta_1}^k} dv_1 \dots dv_k \det [K_{n,r}^{LT}(v_\ell, v_r)]_{\ell,r=1}^k. \quad (3.4)$$

We need to check that this series converges. Using estimates similar to those performed in [5], in particular the estimate

$$\lim_{|y| \rightarrow \infty} \frac{|\Gamma(x + iy)|}{\sqrt{2\pi}} e^{\frac{\pi|y|}{2}} |y|^{1/2-x} = 1. \quad (3.5)$$

from Abramowitz–Stegun [1, (6.1.45)] leads to the following bound for the integrand in the definition of $K_{n,r}^{LT}(v, \tilde{v})$:

$$C_3^n |\Im(w)|^{(h-n)(\delta_2 - \frac{1}{2})} e^{-\frac{\pi}{4}(h-n)|\Im(w)|} \quad (3.6)$$

for some $C_3 > 0$. This is easily seen to be integrable over any vertical line. It follows that $|K_{n,r}^{LT}(v, \tilde{v})|$ can be bounded by C_4^n for some $C_4 > 0$, uniformly over v, \tilde{v} . Together with Hadamard's bound (see for example [15], section 2) it follows that

$$\left| \det(K_{n,r}^{LT}(v_j, v_\ell))_{j,\ell=1}^k \right| \leq k^{k/2} C_4^{nk} \quad (3.7)$$

It now follows immediately that the right hand side of (3.4) is absolutely convergent.

Step 2: Cyclic property of determinants. We now re-write the kernel defining the Fredholm determinant above by using the identity $\det(I + AB) = \det(I + BA)$ for suitable kernels A, B . Just as in [5] the operator defined by $K_{n,r}^{LT}$ can be written as a composition AB where the kernels defining the operators A, B are given by

$$K_A: C_{\delta_1} \times \ell_{\delta_2} \longrightarrow \mathbb{R}, \quad K_A(v, w) = \frac{\pi}{\sin(\pi(v-w))} \frac{F_s(w)}{F_s(v)} \prod_{j=1}^n \frac{\Gamma(v-a_j)}{\Gamma(w-a_j)}$$

$$K_B: \ell_{\delta_2} \times C_{\delta_1} \longrightarrow \mathbb{R}, \quad K_B(w, v) = \frac{1}{w-v}.$$

By the same bounds as above these define operators A from $L^2(\ell_{\delta_2})$ to $L^2(C_{\delta_1})$ and B from $L^2(C_{\delta_1})$ to $L^2(\ell_{\delta_2})$. Note that the integrals

$$\int_{C_{\delta_1}} dv K_B(w_1, v) K_A(v, w_2), \quad \int_{\ell_{\delta_2}} dw K_A(v_1, w) K_B(w, v_2)$$

are finite for all $v_1, v_2 \in C_{\delta_1}$ and $w_1, w_2 \in \ell_{\delta_2}$ (we checked one of them above, the other is similar). Thus we can write the right hand side of (3.2) as $\det\left(I + \tilde{K}_{n,r}^{LT}\right)_{L^2(\ell_{\delta_2})}$ where

$$\tilde{K}_{n,r}^{LT}(w, \tilde{w}) = \int_{C_{\delta_1}} \frac{dv}{2\pi i} \frac{1}{w-v} \frac{\pi}{\sin(\pi(v-\tilde{w}))} \frac{G(\tilde{w})}{G(v)} \quad (3.8)$$

and we have defined

$$G_s(v) = F_s(v) \prod_{j=1}^n \frac{1}{\Gamma(v-a_j)} = F_s(v) \prod_{j=1}^n \frac{v-a_j}{\Gamma(v-a_j+1)}. \quad (3.9)$$

Step 3: The integral in (3.8) can be evaluated using residue calculus: the only singularities of the integrand inside the closed contour C_{δ_1} are simple poles of the form $\frac{1}{v-a_j}$. Since ℓ_{δ_2} is a positive distance away from C_{δ_1} there are no other poles, and the fact that $\delta_1 < \delta_2 \wedge (1-\delta_2)$ implies that the fraction involving the sine does not have any singularities¹ inside C_{δ_1} . Just as in [5] it is sufficient to treat the case where the a_j are all distinct. By computing the residues at the n simple poles we see that

$$\tilde{K}_{n,r}^{LT}(w_1, w_2) = \frac{1}{2\pi i} \sum_{j=1}^n f_j(w_1) g_j(w_2)$$

with $f_j(w) = \frac{1}{w-a_j}$ and $g_j(w) = C_j G(w) \frac{\pi}{\sin(\pi(a_j-w))}$. Here the constant $C_j \in \mathbb{R}$ is given by

$$C_j = \frac{1}{F_s(a_j)} \prod_{\ell \neq j} \Gamma(a_j - a_\ell).$$

Step 4: Applying once more the cyclic property of determinants, analogously to [5], we obtain

$$\det\left(I + K_{n,r}^{LT}\right)_{L^2(C_{\delta_1})} = \det\left[I_n + \int_{\ell_{\delta_2}} \frac{dw}{2\pi i} f_j(w) g_\ell(w)\right]_{j,\ell=1}^n \quad (3.10)$$

where I_n is the $n \times n$ identity matrix.

¹These poles lie inside of the contour thanks to our assumption that $|a_j| < \delta_1$ for all j .

Step 5: We now shift the integration contour on the right-hand side of (3.10) from ℓ_{δ_2} to $-\ell_{\delta_1}$. On the way we will encounter some poles whose residues we will need to evaluate. There is sufficient decay at infinity to justify moving the contours thanks to (3.5).

The poles are different to those in [5], but the outcome is analogous: the only singularity we cross is $w = a_\ell$ from $f_j(w)$, for each j . When $j \neq \ell$ this turns out to be a removable singularity, whereas for $j = \ell$ it is a simple pole with residue given by

$$\text{Res}_{w=a_j} f_j(w)g_\ell(w) = -F_s(a_j) C_j \frac{1}{\Gamma(1)} \prod_{r \neq j} \Gamma(a_j - a_r) = -1$$

and we obtain

$$\det(I + K_{n,r}^{LT})_{L^2(C_{\delta_1})} = \det \left[\int_{-\ell_{\delta_1}} \frac{dw}{2\pi i} f_j(w)g_\ell(w) \right]_{j,\ell=1}^n \quad (3.11)$$

$$= \frac{1}{n! (2\pi i)^n} \int_{(-\ell_{\delta_1})^n} d\vec{w} \det [f_j(w_\ell)]_{j,\ell=1}^n \det [g_j(w_\ell)]_{j,\ell=1}^n \quad (3.12)$$

where the last equality follows from the Andréiev identity [2].

Step 6: It remains to show that the integrand in (3.12) is identical to that in (2.5). But this follows exactly in the same way as in [5], see the paragraphs surrounding equation (3.11) there. This completes our proof.

4 Asymptotics

In the previous section (Proposition 3.1) we saw that

$$\int_{(\mathbb{R}_{>0})^n} e^{-s/x_n} \mu_n(dx) = \det(I + K_{n,r}^{LT})_{L^2(C_{\delta_1})}$$

where

$$K_{n,r}^{LT}(v_1, v_2) = \int_{\ell_{\delta_2}} \frac{dw}{2\pi i} \frac{\pi}{\sin(\pi(v_1 - w))} \frac{F_s(w)}{F_s(v_1)} \frac{1}{w - v_2} \prod_{j=1}^n \frac{\Gamma(v_1 - a_j)}{\Gamma(w - a_j)} \quad (4.1)$$

and the function F_s was defined in (3.1). From now on we choose $a_j = 0$ and $b_j = \gamma$ for all j , where $\gamma > 0$. Then $1/x_n$ has the same law under μ_n as the partition function $Z_{m,n}$ of the random polymer defined in the introduction, taking $m = h - n + 1$. We will set $h = \lceil cn \rceil$ for some fixed $c > 1$. The correct choice of the parameter s will turn out to be $s = e^{-n\mu - rn^{-1/3}}$ with μ defined in (4.6) below. Then

$$e^{-sZ_{m,n}} = f_{n,r} \left(\frac{\ln Z_{m,n} - n\mu}{n^{1/3}} \right) \quad (4.2)$$

where $f_{n,r}(x) = \exp \left\{ -e^{n^{1/3}(x-r)} \right\}$. In this section we show that the expectation of the left-hand side above converges, as $n \rightarrow \infty$, to a rescaled version of the Tracy–Widom GUE distribution function. Observe that with our choice of parameter s this expectation equals $\det(I + K_{n,r})$ where

$$K_{n,r}(v_1, v_2) = \frac{1}{2\pi i} \int_{\ell_{\delta_2}} \frac{dw}{w - v_2} \frac{\pi}{\sin(\pi(v_1 - w))} \exp \left\{ n(H_{n,c,\gamma}(v_1) - H_{n,c,\gamma}(w)) - rn^{1/3}(w - v_1) \right\} \quad (4.3)$$

and, recalling that $h = \lceil cn \rceil$

$$H_{n,c,\gamma}(z) = \ln \Gamma(z) - \tilde{c}_n \ln \Gamma(\gamma + z) + \mu z. \quad (4.4)$$

and $\tilde{c}_n = \frac{\lceil cn \rceil}{n}$.

Theorem 4.1. *For γ sufficiently small we have*

$$\lim_{n \rightarrow \infty} \det(I + K_{n,r})_{L^2(C_{\delta_1})} = F_{GUE} \left((\bar{g}/2)^3 r \right)$$

where \bar{g} was defined in Theorem 1.1.

The proof of Theorem 1.1 is completed by noting that $f_{n,r}(x) = f_{n,0}(x - r)$ for all r and that $(f_n := f_{n,0} : n \in \mathbb{N})$ and $p := F_{GUE}$ satisfy the conditions of Lemma 4.2, whose proof is elementary and can be found in [3, Lemma 4.1.39].

Lemma 4.2. *For each $n \in \mathbb{N}$ let $f_n : \mathbb{R} \rightarrow [0, 1]$ be f_n strictly decreasing and converge to 0 at ∞ and 1 at $-\infty$. Suppose further that for each $\delta > 0$, $(f_n : n \in \mathbb{N})$ converges uniformly to $\mathbf{1}_{(-\infty, 0]}$. Let $(X_n : n \in \mathbb{N})$ be real-valued random variables such that for each $r \in \mathbb{R}$,*

$$\lim_{n \rightarrow \infty} \mathbb{E}(f_n(X_n - r)) = p(r)$$

where p is a continuous probability distribution function. Then $(X_n : n \in \mathbb{N})$ converges in distribution to a random variable with distribution function p .

It therefore remains to prove Theorem 4.1. Recall that we need to compute the $n \rightarrow \infty$ limit of $\det(I + K_{n,r})_{L^2(C_{\delta_1})}$ with $K_{n,r}$ as defined in (4.3) above.

The first step is to identify suitable steepest descent contours to which we will deform the contours C_{δ_1} and ℓ_{δ_2} . We also introduce the function $H_{c,\gamma}(z) = \ln \Gamma(z) - c \ln \Gamma(z + \gamma) + \mu z$. Observe that for $z \in \mathbb{C}$,

$$H_{c,\gamma}(z) - H_{n,c,\gamma}(z) = (\tilde{c}_n - c) \ln(\Gamma(z + \gamma)). \quad (4.5)$$

and that $\tilde{c}_n - c = O(n^{-1})$. For later use we record the first few derivatives of $H_{c,\gamma}$:

$$\begin{aligned} H'_{c,\gamma}(z) &= \psi(z) - c\psi(\gamma + z) + \mu \\ H''_{c,\gamma}(z) &= \psi_1(z) - c\psi_1(\gamma + z) \\ H'''_{c,\gamma}(z) &= \psi_2(z) - c\psi_2(\gamma + z) \end{aligned}$$

where $\psi_k(x) = \frac{d^{k+1}}{dx^{k+1}} \ln(\Gamma(x))$ is the k^{th} polygamma function (in particular $\psi = \psi_0$ is the digamma function as above). Let $\lambda_c > 0$ be small, with the precise value to be chosen later. The proof of the following calculus lemma can be found in Section 5.

Lemma 4.3. *For each $c > 0$ and $\gamma > 0$ small enough there exists unique $z_{c,\gamma}^*$ such that $H''_{c,\gamma}(z_{c,\gamma}^*) = 0$. Moreover $H'''_{c,\gamma}(z_{c,\gamma}^*) < 0$ and we can write $z_{c,\gamma}^* = \gamma \tilde{z}_{c,\gamma}^* + O(\gamma)$ with $\lim_{\gamma \rightarrow 0} \tilde{z}_{c,\gamma}^* = \frac{1}{\sqrt{c-1}}$.*

Our asymptotic analysis will consist of shifting our contours to curves that pass through or near $z_{c,\gamma}^*$ and showing that in the $n \rightarrow \infty$ limit only the parts of the contour near $z_{c,\gamma}^*$ survive. We will see that the right choice for $\mu = \mu_c$ is such that $H'_{c,\gamma}(z_{c,\gamma}^*) = 0$, i.e.

$$\mu_c = c\psi(\gamma + z_{c,\gamma}^*) - \psi(z_{c,\gamma}^*) \quad (4.6)$$

$$= \inf_{z > 0} \{c\psi(z + \gamma) - \psi(z)\}. \quad (4.7)$$

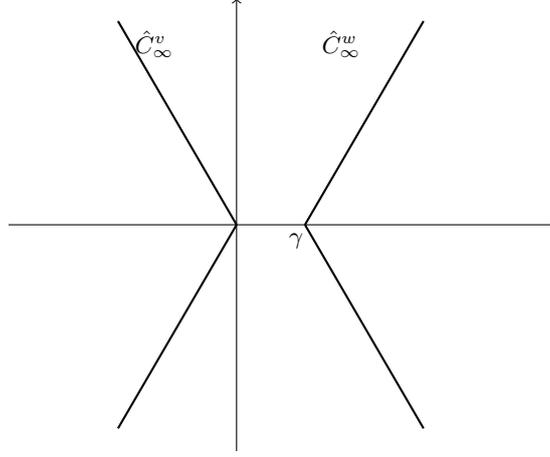


Figure 4: The contours \hat{C}_∞^v and \hat{C}_∞^w . The angle between \hat{C}_∞^v and \hat{C}_∞^w and the negative and positive x-axes respectively is given by $\frac{\pi}{3}$.

with infimum rather than supremum because $\bar{g}_c := -H_{c,\gamma}'''(z_{c,\gamma}^*) > 0$. Taylor's theorem implies therefore that, for v, w near $z_{c,\gamma}^*$,

$$H_{c,\gamma}(v_1) - H_{c,\gamma}(w) = \frac{\bar{g}_c(w - z_{c,\gamma}^*)^3}{6} - \frac{\bar{g}_c(v_1 - z_{c,\gamma}^*)^3}{6} + O\left((w - z_{c,\gamma}^*)^4\right) + O\left((v_1 - z_{c,\gamma}^*)^4\right). \quad (4.8)$$

The fact that the lowest power is a cube suggests a scaling of order $n^{1/3}$ around the critical point and we set $\tilde{v}_j = n^{1/3}(v_j - z_{c,\gamma}^*)$ and $\tilde{w} = n^{1/3}(w - z_{c,\gamma}^*)$. We will see below that only a small part of the integral around the critical point contributes to the limit which leads to

Proposition 4.4. *We have*

$$\lim_{n \rightarrow \infty} \det(I + K_{n,r})_{L^2(C^v)} = \det(1 + K_r^{LT})_{L^2(\hat{C}_\infty^v)} \quad (4.9)$$

where

$$K_r^{LT}(\tilde{v}_1, \tilde{v}_2) = \frac{1}{2\pi i} \int_{\hat{C}_\infty^w} \frac{d\tilde{w}}{\tilde{w} - \tilde{v}_2} \frac{1}{\tilde{v}_1 - \tilde{w}} \exp\left\{ \frac{\bar{g}_c(\tilde{w}^3 - \tilde{v}_2^3)}{6} + r(\tilde{v}_1 - \tilde{w}) \right\} \quad (4.10)$$

and further $\hat{C}_\infty^v = e^{2i\pi/3}\mathbb{R}_{\geq 0} \cup e^{4i\pi/3}\mathbb{R}_{\geq 0}$ and $\hat{C}_\infty^w = \gamma + (e^{i\pi/3}\mathbb{R}_{\geq 0} \cup e^{-i\pi/3}\mathbb{R}_{\geq 0})$, see Figure 4.

Setting now $v = \left(\frac{\bar{g}_c}{2}\right)^{1/3} \tilde{v}$ and similarly $w = \left(\frac{\bar{g}_c}{2}\right)^{1/3} \tilde{w}$ we obtain $\det(I + \widetilde{K}_r^{LT})_{L^2(C^v)}$ where

$$\widetilde{K}_r^{LT}(v_1, v_2) = \frac{1}{2\pi i} \int_{\hat{C}_\infty^w} \frac{dw}{w - v_2} \frac{1}{v_1 - w} \exp\left\{ \frac{-\frac{v^3}{3} + \left(\frac{\bar{g}_c}{2}\right)^{-1/3} rv}{-\frac{w^3}{3} + \left(\frac{\bar{g}_c}{2}\right)^{-1/3} rw} \right\}$$

But this is exactly one of the definitions of the Tracy-Widom GUE distribution, see for example Lemma 8.6 in [4].

We begin by deforming the contours C_{δ_1} and ℓ_{δ_2} to suitable steepest descent contours. In fact, for γ small enough we will be able to do this without passing through any pole.

The integrand on the right hand side of (4.1) has the following poles in the integration variable w :

- $w = v_2$
- $w = -M - \gamma$ for $M \in \mathbb{Z}_{\geq 0}$ (these are the poles of F)
- $w = v_1 + 2p\pi$ for all $p \in \mathbb{Z}$

On the other hand the poles of the kernel in v_1, v_2 are given by

- $v_1 = w + 2p\pi$ for all $p \in \mathbb{Z}$
- $v_2 = w$
- $v_2 = 0$
- $v_1 = a_j - M$ for all $M \in \mathbb{Z}_{\geq 0}$

We would like to move the contours C_{δ_1} and ℓ_{δ_2} to the following contours, which are illustrated in Figure 5.

Denote by $C^{w,\pm}$ the line segments of length $\ell - n^{-1/3}$ starting at $z_{c,\gamma}^* + \gamma n^{-1/3}$ making angles $\frac{\pi}{3}$ and $-\frac{\pi}{3}$ respectively with the positive x -axis and let $C^w = (z_{c,\gamma}^* + \gamma \ell e^{-i\pi/3} + i\mathbb{R}_{\geq 0}) \cup C^{w,+} \cup C^{w,-} \cup (z_{c,\gamma}^* + \gamma \ell e^{i\pi/3} + i\mathbb{R}_{\geq 0})$, oriented to have increasing imaginary part.

The closed contour C^v is defined differently according to whether c is larger than $\frac{5}{2}$ or not. For $c > \frac{5}{2}$ let C^v be the union of the line segments of length $\frac{6\gamma}{5(\sqrt{c-1})}$ making angles $\pm \frac{2\pi}{3}$ with the positive x -axis and the circular segment, centred at $z_{c,\gamma}^*$, that connects the end-points of these two segments. For $c \leq \frac{5}{2}$ we define C^v to be the union of the following four line segments: those starting at $z_{c,\gamma}^*$ of length $\frac{2\gamma}{c-1}$ making angles $\pm \frac{2\pi}{3}$ with the positive x -axis and those connecting the end-points of the former with the point $-\frac{2\gamma}{c-1}$. In both cases we give C^v the positive orientation.

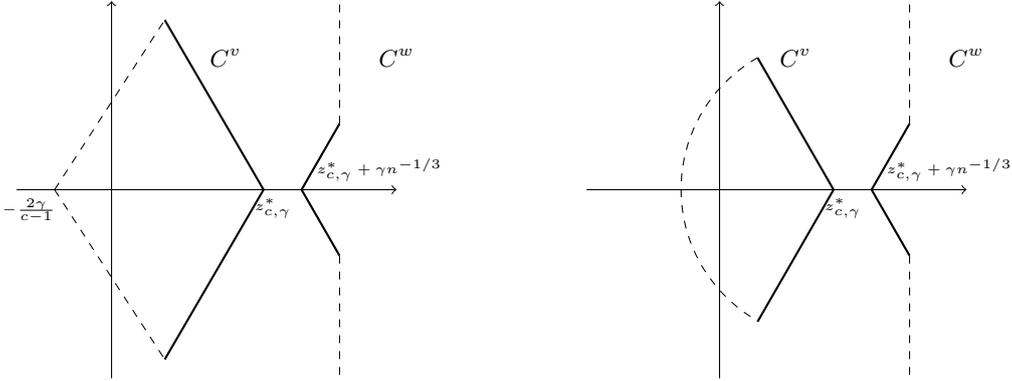


Figure 5: Contours C^v and C^w for large $c \leq \frac{5}{2}$ (on the left) and $c > \frac{5}{2}$ (on the right). The parts C_{irrel}^v and C_{irrel}^w of the contours are drawn as dashed lines.

It is easy to see that we do not cross any poles of the integrand, further the estimate (3.5) gives sufficient decay at infinity to justify moving the infinite w -contour. It follows that $\mathbb{E}e^{-s/Z_n} = \det(I + K_{n,r}^{LT})_{C^v}$ where

$$K_{n,r}^{LT}(v, \tilde{v}) = \frac{1}{2\pi i} \int_{C^w} \frac{dw}{w - \tilde{v}} \frac{\pi}{\sin(\pi(v - \tilde{v}))} \frac{F(w)}{F(v)} \prod_{j=1}^n \frac{\Gamma(v - a_j)}{\Gamma(w - a_j)}. \quad (4.11)$$

The proof in the rigorous steepest descent analysis now goes along similar lines as, for example, [4, 5]. Fix $\epsilon > 0$. We will show that the difference between our formula for the Laplace transform of $\mathbb{E}e^{-s/Z_N}$ and the right hand side of (4.10) can be bounded by ϵ for large enough n .

Lemma 4.5. *There exists $M^* > 0$ such that for $M > M^*$,*

$$\left| \det (I + K_{r,M}^{trunc})_{L^2(\hat{C}_M^v)} - \det (I + K_r^{LT})_{L^2(\hat{C}_\infty^v)} \right| < \frac{\epsilon}{3}$$

where $\hat{C}_M^v = \{z \in \hat{C}_\infty^v : |z| \leq M\}$,

$$K_{r,M}^{trunc}(\tilde{v}_1, \tilde{v}_2) = \frac{1}{2\pi i} \int_{\hat{C}_M^w} \frac{d\tilde{w}}{\tilde{w} - \tilde{v}_2} \frac{1}{\tilde{v}_1 - \tilde{w}} \exp \left\{ \frac{\bar{g}(\tilde{w}^3 - \tilde{v}_2^3)}{6} + r(\tilde{v}_1 - \tilde{w}) \right\}$$

and similarly $\hat{C}_M^w = \{z \in \hat{C}_\infty^w : |z| \leq M\}$.

From now on we assume that $M > M^*$. Denote by C_{rel}^v the part of C^v consisting of the two line segments starting at $z_{c,\gamma}^*$. Similarly let C_{rel}^w be the corresponding part of C^w . Further define $C_{\text{irrel}}^v = C^v \setminus C_{\text{rel}}^v$ and $C_{\text{irrel}}^w = C^w \setminus C_{\text{rel}}^w$ (see Figure 5).

Lemma 4.6. *There exist $\gamma^* > 0$ and $\ell > 0$ such that for $\gamma < \gamma^*$ and n sufficiently large the following hold*

(i) *There exists $C_1 > 0$ such that for $v \in C_{\text{irrel}}^v$,*

$$\Re(H_{n,c,\gamma}(v) - H_{n,c,\gamma}(z_{c,\gamma}^*)) \leq -C_1. \quad (4.12)$$

(ii) *There is $C_2 > 0$ such that for all $v \in C_{\text{rel}}^v$ with $|v| \geq \ell$,*

$$\Re[H_{n,c,\gamma}(v) - H_{n,c,\gamma}(z_{c,\gamma}^*)] \leq -C_2 \quad (4.13)$$

(iii) *There is $C_3 > 0$ such that for all $v \in C_{\text{rel}}^v$ with $|v| \leq \ell$,*

$$\Re[H_{n,c,\gamma}(v) - H_{n,c,\gamma}(z_{c,\gamma}^*)] \leq -C_3 \Re[(v - z_{c,\gamma}^*)^3] \quad (4.14)$$

(iv) *There is $C_4 > 0$ such that for all $w \in C_{\text{rel}}^w$*

$$\Re[H_{n,c,\gamma}(z_{c,\gamma}^*) - H_{n,c,\gamma}(w)] \leq -C_4 \Re[(z_{c,\gamma}^* - w)^3]. \quad (4.15)$$

(v) *There exists $C_5 > 0$ such that for all $\gamma < \gamma^*$ and $w \in C_{\text{irrel}}^w$*

$$\Re[H_{n,c,\gamma}(z_{c,\gamma}^*) - H_{n,c,\gamma}(w)] \leq -C_5. \quad (4.16)$$

Further there exists $L = L_{c,\gamma} > 0$ such that if additionally $|w| > L$ then

$$\Re[H_{n,c,\gamma}(z_{c,\gamma}^*) - H_{n,c,\gamma}(w)] \leq \frac{(1-c)\pi}{4} |\Im(w)| \quad (4.17)$$

The proof of Lemmas 4.5 and 4.6 can be found in Section 5. From now on we assume that $\gamma < \gamma^*$.

By observing that the estimates of Lemma 4.6 are uniform in n and applying Hadamard's bound in exactly the same way as in Step 1 of the proof of Proposition 3.1 we deduce that the series defining $\det(I + K_{n,r}^{LT})_{C^v}$ is uniformly convergent in n . Thus we may interchange the $n \rightarrow \infty$ limit with the series in k .

Thanks to (v) the contribution of the w integral along C_{irrel}^w becomes negligible as n tends to infinity. That is, uniformly in $v_1, v_2 \in C^v$, as $n \rightarrow \infty$,

$$\int_{C_{\text{irrel}}^w} \frac{dw}{w - v_1} \frac{\pi}{\sin(\pi(v_2 - w))} e^{n(H_{n,c,\gamma}(z_{c,\gamma}^*) - H_{n,c,\gamma}(w))} \rightarrow 0 \quad (4.18)$$

Similarly it follows from (4.12) and uniform convergence that only the ‘relevant’ part of the v -contour survives in the limit. That is, there exists $N \in \mathbb{N}$ such that for all $n \geq N$,

$$\left| \det(I + K_{n,r}^{LT})_{L^2(C^v)} - \det(I + K_{n,r}^{LT})_{L^2(C_{\text{rel}}^v)} \right| < \frac{\epsilon}{3}. \quad (4.19)$$

The estimates from (4.14) and (4.15) now allow us to further discard the parts of C_{rel}^v and C_{rel}^w which are further than $Mn^{-1/3}$ away from $z_{c,\gamma}^*$ and $z_{c,\gamma}^* + n^{-1/3}$ respectively. Now we make the change of variables $v_j = n^{1/3}\tilde{v}_j + z_{c,\gamma}^*$ and $w_j = n^{1/3}\tilde{w} + z_{c,\gamma}^*$, and write $K_{n,r}^{LT}(v_1, v_2) = \tilde{K}_{n,r}^{LT}(\tilde{v}_1, \tilde{v}_2)$ for $\tilde{v}_1, \tilde{v}_2 \in \hat{C}_M^v$. Then

$$\lim_{n \rightarrow \infty} \det(I + K_{n,r}^{LT})_{L^2(C_{\text{rel}}^v)} = \lim_{n \rightarrow \infty} \det(I + \tilde{K}_{n,r}^{LT})_{L^2(\hat{C}_M^v)}.$$

We will show that $\tilde{K}_{n,r}^{LT}$ converges pointwise to $K_{r,M}^{\text{trunc}}$. Once this has been established we can conclude by the DCT and uniform convergence that $\det(I + K_{n,r}^{LT})_{L^2(C^v)}$ converges to $\det(I + K_{r,M}^{\text{trunc}})_{L^2(\hat{C}_M^v)}$. But by Lemma 4.5 this differs only by $\frac{\epsilon}{3}$ from $\det(1 + K_r^{LT})_{L^2(\hat{C}_\infty^v)}$. So we have shown that for N sufficiently large, γ sufficiently small and $M > M^*$,

$$\left| \det(I + K_{n,r}^{LT})_{L^2(C^v)} - \det(1 + K_r^{LT})_{L^2(\hat{C}_\infty^v)} \right| < \epsilon$$

subject to establishing pointwise convergence of $\tilde{K}_{n,r}^{LT}$ to $K_{r,M}^{\text{trunc}}$. For this observe that

$$\frac{dw}{w - v_2} = \frac{d\tilde{w}}{\tilde{w} - \tilde{v}_2}, \quad n^{-1/3} \frac{\pi}{\sin(\pi(v_1 - w))} = \frac{1}{\tilde{v}_1 - \tilde{w}} + O(n^{-1/3}), \quad rn^{1/3}(w - v_1) = r(\tilde{w} - \tilde{v}_1)$$

and, thanks to (4.8) and the fact that $\tilde{c}_n = c + O(\frac{1}{n})$,

$$n(H_{n,c,\gamma}(v_1) - H_{n,c,\gamma}(w)) = \frac{\bar{g}_c}{6} \left[(\tilde{w} - z_{c,\gamma}^*)^3 - (\tilde{v}_1 - z_{c,\gamma}^*)^3 \right] + O(n^{-1/3}).$$

This concludes the proof of Proposition 4.4

5 Proof of Lemmas

This section is devoted to proving the auxiliary results from Section 4 above.

5.1 Proof of Lemma 4.5

The last lemma to prove replaces the finite contour \hat{C}_M^v by \hat{C}_∞^v . By the Dominated Convergence Theorem and continuity of the determinant we have, for sufficiently large M ,

$$\left| \det(I + K_{r,M}^{\text{trunc}})_{L^2(\hat{C}_M^v)} - \det(K_r^{LT})_{L^2(\hat{C}_M^v)} \right| < \frac{\epsilon}{6}. \quad (5.1)$$

The following useful result can be found as Lemma 8.4 in [4].

Lemma 5.1. *Let Λ be an infinite complex curve and K an integral operator on Λ . Suppose that there exists $C_1, C_2, C_3 > 0$ such that $|K(v_1, v_2)| \leq C_1$ for all $v_1, v_2 \in \Lambda$ and that*

$$|K(\Lambda(s_1), \Lambda(s_2))| \leq C_2 e^{-C_3 |s_1|} \quad (5.2)$$

for all $s \in \mathbb{R}$ (here, $\Lambda(s)$ denotes the parametrisation of Λ by arc length). Then the Fredholm series defining $\det(I + K)_{L^2(\Lambda)}$ is well defined, and for any $\epsilon > 0$ there exists $M_\epsilon > 0$ such that for all $M > M_\epsilon$,

$$\left| \det(I + K)_{L^2(\Lambda)} - \det(I + K)_{L^2(\Lambda_M)} \right| \leq \epsilon \quad (5.3)$$

where $\Gamma_M = \{\Gamma(s) : |s| \leq M\}$.

The proof of Lemma 4.5 is therefore complete if we can find $C_1, C_2 > 0$ such that $|K_r^{LT}(v_1, v_2)| \leq C_1 e^{-C_2 v_1}$ for all $v_1, v_2 \in \hat{C}_\infty^v$. But this follows immediately from (4.10).

5.2 Proof of Lemma 4.3

Convexity considerations show that if there exists a zero of $H''_{c,\gamma}$ then it is unique. Let us write $z = \gamma \tilde{z}$ then

$$H''_{c,\gamma}(z) = \gamma^{-2} \left(\frac{1}{\tilde{z}^2} - \frac{c}{(1 + \tilde{z})^2} \right) + \frac{\pi^2}{6} (1 - c) + O(\gamma)$$

with the error being uniform in \tilde{z} over compact intervals. Hence, for γ small enough we have $H''_{c,\gamma}\left(\frac{\gamma}{\sqrt{c}-1}\right) < 0$ and $H''_{c,\gamma}\left(\frac{\gamma}{\sqrt{c}-1} - \lambda_c\right) > 0$, from which the result follows.

5.3 Proof of Lemma 4.6

The following small γ estimates will be useful. Throughout we set $z = \gamma \tilde{z}$.

Lemma 5.2. *There exist $(\tilde{\mu}_{c,\gamma}^* : \gamma > 0)$ such that*

$$\mu_{c,\gamma}^* = \frac{\tilde{\mu}_{c,\gamma}^*}{\gamma} + O(\gamma) \quad (5.4)$$

and $\tilde{\mu}_{c,\gamma}^* \rightarrow -(\sqrt{c}-1)^2$ as $\gamma \rightarrow 0$.

Proof. We have

$$\begin{aligned} \mu_{c,\gamma}^* &= c\psi(\gamma(\tilde{z}_{c,\gamma}^* + 1)) - \psi(\gamma\tilde{z}_{c,\gamma}^*) \\ &= \frac{1}{\gamma} \left[\frac{1}{\tilde{z}_{c,\gamma}^*} - \frac{c}{1 + \tilde{z}_{c,\gamma}^*} \right] + O(\gamma). \end{aligned}$$

The claim now follows from Lemma 4.3. \square

We also record the following small γ expansions.

$$\bar{g}_c = -H''''_{c,\gamma}(z_{c,\gamma}^*) = \gamma^{-3} \left(\frac{2c}{(1 + \tilde{z}_{c,\gamma}^*)^3} - \frac{2}{(\tilde{z}_{c,\gamma}^*)^3} \right) \quad (5.5)$$

$$\begin{aligned} H_{c,\gamma}(z) - H_{c,\gamma}(v) &= c \log(1 + \tilde{v}) - c \log(1 + \tilde{z}) - \log(\tilde{v}) + \log(\tilde{z}) \\ &\quad + \tilde{\mu}_{c,\gamma}^* (\tilde{v} - \tilde{z}) + O(\gamma) \end{aligned} \quad (5.6)$$

Proof of Lemma 4.6. (i) Because v varies over a compact set it follows from (4.5) that there exists some $C > 0$ such that

$$|H_{n,c,\gamma}(v) - H_{n,c,\gamma}(z_{c,\gamma}^*) - [H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*)]| < \frac{C}{n}$$

holds for all $v \in C^v$. Therefore we may as well prove the claim with $H_{n,c,\gamma}$ replaced by $H_{c,\gamma}$, which is what we will do.

Since the contours are different we will consider the cases $c > \frac{5}{2}$ and $c \leq \frac{5}{2}$ separately.

Case I: $c > \frac{5}{2}$. Fix $\epsilon > 0$ to be chosen later and write $v = \gamma\tilde{v}$. Recall that $z_{c,\gamma}^* = \gamma\tilde{z}_{c,\gamma}^*$ and $\mu_{c,\gamma}^* = \frac{\tilde{\mu}_{c,\gamma}^*}{\gamma}$. By (5.6),

$$H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*) = c \ln \left(\frac{\tilde{v} + 1}{\tilde{z}_{c,\gamma}^* + 1} \right) - \ln \left(\frac{\tilde{v}}{\tilde{z}_{c,\gamma}^*} \right) + \tilde{\mu}_{c,\gamma}^* (\tilde{v} - \tilde{z}_{c,\gamma}^*) + O(\gamma)$$

where the error term is uniform in \tilde{v} (because the latter varies over a compact contour). Now $\tilde{v} = \tilde{z}_{c,\gamma}^* + r_c e^{i\theta}$ where $r_c = \frac{6}{5} \frac{1}{\sqrt{c-1}}$ and $\theta \in [\frac{2\pi}{3}, \frac{4\pi}{3}]$, so we obtain, for γ small enough,

$$\Re(H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*)) \leq c \ln \left| 1 + \frac{r_c}{\tilde{z}_{c,\gamma}^* + 1} e^{i\theta} \right| - \ln \left| 1 + \frac{r_c}{\tilde{z}_{c,\gamma}^*} e^{i\theta} \right| + \Re(\tilde{\mu}_{c,\gamma}^* r_c e^{i\theta}) + \frac{\epsilon}{2}.$$

By Lemmas 4.3 and 5.2 we can now ensure, by choosing γ small enough, that

$$\begin{aligned} \Re(H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*)) &\leq c \ln \left| 1 + \frac{6}{5\sqrt{c}} e^{i\theta} \right| - \ln \left| 1 + \frac{6}{5} e^{i\theta} \right| - \frac{6}{5} (\sqrt{c} - 1) \cos(\theta) + \epsilon \\ &= \frac{c}{2} \ln \left[\left(1 + \frac{6}{5} \sqrt{c} \alpha \right)^2 + \frac{36}{25c} (1 - \alpha^2) \right] \\ &\quad - \frac{1}{2} \ln \left[\left(1 + \frac{6}{5} \alpha \right)^2 + \frac{36}{25} (1 - \alpha) \right] - \frac{6}{5} (\sqrt{c} - 1) \alpha + \epsilon \end{aligned}$$

where we have written $\alpha = \cos(\theta) \in [-1, -\frac{1}{2}]$. Denote by $f(c, \alpha)$ the last expression above, with $\epsilon = 0$. For each $c > \frac{5}{2}$ the function $\alpha \mapsto f(c, \alpha)$ has a unique critical point on the interval $[-1, -\frac{1}{2}]$ which turns out to be a minimum. Thus we are reduced to consider the end-points. Now $f(\cdot, -\frac{1}{2})$ is strictly decreasing and clearly $C_{11} = f(\frac{5}{2}, -\frac{1}{2}) < 0$. On the other hand $f(\cdot, -1)$ is strictly increasing and tends to $C_{12} = \ln(5) - \frac{48}{25} < 0$ as $c \rightarrow \infty$. Taking now $C_1 = \epsilon = -\frac{1}{2} \min\{C_{11}, C_{12}\}$ completes the proof for the case $c > \frac{5}{2}$.

Case II: $c \leq \frac{5}{2}$. The contour in question is the union of the (complex) line segments $\left[\frac{2\gamma}{c-1} e^{2i\pi/3}, -\frac{2\gamma}{c-1} \right]$ and $\left[\frac{2\gamma}{c-1} e^{-2\pi i/3}, -\frac{2\gamma}{c-1} \right]$. By symmetry it suffices to consider the former. Thus, writing $v = \gamma\tilde{v}$,

$$\tilde{v} = t + i \frac{\sqrt{3}}{\sqrt{c} + 2} \left(t + \frac{2}{c-1} \right), \quad t \in \left[-\frac{2}{c-1}, \frac{\sqrt{c}}{c-1} \right]. \quad (5.7)$$

Fix $\epsilon > 0$. By Lemmas 4.3 and 5.2 as well as (5.6) and (5.7) we have, for γ small enough and then n large enough,

$$\begin{aligned} \Re[H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*)] &\leq (c-1) \ln(\sqrt{c}-1) + \ln(\sqrt{c}) - \frac{1}{(\sqrt{c}-1)^2} \left(t - \frac{1}{\sqrt{c}-1} \right) \\ &\quad - \ln \left[t^2 + \frac{3}{(\sqrt{c}+2)^2} \left(t + \frac{2}{c-1} \right)^2 \right] \\ &\quad + c \ln \left[(t+1)^2 + \frac{3}{(\sqrt{c}+2)^2} \left(t + \frac{2}{c-1} \right)^2 \right] + \epsilon \end{aligned}$$

Temporarily denote the right hand side above by $F(c, t, \epsilon)$. For any fixed $c \in (1, \frac{5}{2}]$ the function $F(c, \cdot, 0)$ has a unique critical point on the interval $[-\frac{2}{c-1}, \frac{\sqrt{c}}{c-1}]$, at which point the second derivative is positive. Furthermore it is easy to verify that the end points $t_0 = -\frac{2}{c-1}$ and $t_1 = \frac{\sqrt{c}}{c-1}$ satisfy $F(c, t_0, 0) < 0$ and $F(c, t_1, 0) < 0$. This completes the proof for $c \leq \frac{5}{2}$ and hence for part (i) of the lemma.

(ii) By the same argument as in part (i) we may replace $H_{n,c,\gamma}$ by $H_{c,\gamma}$.

Consider first the case where $c > \frac{5}{2}$, so that we have $v = \gamma \left(\tilde{z}_{c,\gamma}^* + \frac{re^{2i\pi/3}}{\sqrt{c-1}} \right)$ for $r \in [0, \frac{6}{5}]$. Then

$$\Re [H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*)] = c \ln \left| 1 + \frac{re^{2i\pi/3}}{(\sqrt{c}-1)(\tilde{z}_{c,\gamma}^*+1)} \right| - \ln \left| 1 + \frac{re^{2i\pi/3}}{\tilde{z}_{c,\gamma}^*(\sqrt{c}-1)} \right| - \frac{1}{2} r \tilde{\mu}_{c,\gamma}^*$$

Fix $\epsilon > 0$. Using (5.6) and Lemmas 4.3 and 5.2 as above we have, for γ suitably small,

$$\begin{aligned} \Re [H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*)] &\leq c \ln \left[\left(1 + \frac{r}{2\sqrt{c}} \right)^2 + \frac{3r^2}{4\sqrt{c}} \right] - \ln \left[\left(\frac{r}{2} \right)^2 + \frac{3r^2}{4} \right] \\ &\quad + \frac{r}{2} (\sqrt{c}-1) + \epsilon. \end{aligned}$$

Now for any fixed r the right hand side is decreasing in c , so it is enough to consider the case where $c = \frac{5}{2}$, for which it is easy to see that the quantity above is bounded above away from zero (for small enough ϵ), uniformly in $r \in [0, \frac{6}{5}]$. Taking γ small enough deals with the error term (which is uniform in the other variables involved).

For the case $c \leq \frac{5}{2}$ we set $v = \gamma \left(\tilde{z}_{c,\gamma}^* + \frac{2r}{c-1} e^{2i\pi/3} \right)$ with $r \in [0, 1]$. A similar computation as in the case $c > \frac{5}{2}$ shows that for this choice of v (at any fixed $r \in [0, 1]$) the function $c \mapsto \Re [H_{c,\gamma}(v) - H_{c,\gamma}(z_{c,\gamma}^*)]$ is strictly increasing in c and converges to zero as $c \rightarrow 1$. Thus the claim holds for any fixed $c > 1$, as required.

Parts (iii) and (iv) follow from Taylor's theorem and the fact that $H'_{n,c,\gamma}(z_{c,\gamma}^*) = H''_{n,c,\gamma}(z_{c,\gamma}^*) = 0$.

It remains to prove part (v). For the first assertion observe first that by (4.5) we have, uniformly in $w \in C_{\text{irrel}}^w$,

$$\begin{aligned} \Re [H_{n,c,\gamma}(z_{c,\gamma}^*) - H_{n,c,\gamma}(w) - (H_{c,\gamma}(z_{c,\gamma}^*) - H_{c,\gamma}(w))] \\ &= \underbrace{\Re [H_{n,c,\gamma}(z_{c,\gamma}^*) - H_{c,\gamma}(z_{c,\gamma}^*)]}_{O(n^{-1})} + \Re [H_{c,\gamma}(w) - H_{n,c,\gamma}(w)] \\ &= (\tilde{c}_n - c) \ln |\Gamma(w + \gamma)| + O(n^{-1}) \end{aligned}$$

Now $|\Gamma(w + \gamma)| < 1$ for $w \in C_{\text{irrel}}^w$ and we have chosen $\tilde{c}_n > c$, so the first summand above is negative and we can once more reduce to the case where $H_{n,c,\gamma}$ is replaced by $H_{c,\gamma}$. Next, write $w = \gamma \tilde{w}$ so that $\tilde{w} = \tilde{z}_{c,\gamma}^* + le^{i\pi/3} + iy$ for $y \geq 0$ or $\tilde{w} = \tilde{z}_{c,\gamma}^* + le^{-i\pi/3} + iy$ for $y \leq 0$. By symmetry it is enough to consider the former case. Fix $\epsilon > 0$. Applying once more Lemmas 4.3 and 5.2 and (5.6) as well as the fact that $e^{i\pi/3} = \frac{1}{2} + i\frac{\sqrt{3}}{2}$ we get, for

suitably small γ ,

$$\begin{aligned} \Re [H_{c,\gamma}(z_{c,\gamma}^*) - H_{c,\gamma}(w)] &\leq \frac{1}{2} \ln \left[\left(1 + \frac{\ell(\sqrt{c}-1)}{2} \right)^2 + (\sqrt{c}-1)^2 \left(y + \frac{\sqrt{3}\ell}{2} \right)^2 \right] \\ &\quad - \frac{c}{2} \ln \left[\left(1 + \frac{\ell(\sqrt{c}-1)}{2\sqrt{c}} \right)^2 + \frac{(\sqrt{c}-1)^2 \left(y + \frac{\sqrt{3}\ell}{2} \right)^2}{c} \right] \\ &\quad - \frac{(\sqrt{c}-1)^2 \ell}{2} + \epsilon \end{aligned}$$

Let us denote by $F(c, \ell, y, \epsilon)$ the last expression above. It is straightforward to check that the map $y \mapsto F(c, \ell, y, 0)$ is strictly decreasing on $[0, \infty)$. Furthermore the map $\ell \mapsto F(c, \ell, 0, 0)$ is strictly decreasing on $[0, \infty)$ and moreover $F(c, 0, 0, 0) = 0$. Since $\ell > 0$ it follows that there exists $\tilde{C}_5 > 0$ such that $F(c, \ell, y, 0) \leq -\tilde{C}_5$ for all $c > 1$ and $y \geq 0$. The first assertion now follows by choosing $\epsilon = C_5 = \frac{1}{2}\tilde{C}_5$.

For the second assertion we will apply the bound (3.5): for any $\eta > 0$,

$$\begin{aligned} \Re [H_{n,c,\gamma}(z_{c,\gamma}^*) - H_{n,c,\gamma}(w)] &\leq C - \ln |\Gamma(x + iy)| + \tilde{c}_n \ln |\Gamma(\gamma + x + iy)| \\ &\leq C - \ln \left(\sqrt{2\pi}(1 + \eta)e^{-\pi|y|/2} |y|^{x-\frac{1}{2}} \right) \\ &\quad + \tilde{c}_n \ln \left(\sqrt{2\pi}(1 - \eta)e^{-\pi|y|/2} |y|^{x+\gamma-\frac{1}{2}} \right) \\ &\leq C + (1 - \tilde{c}_n) \frac{\pi}{2} |y| \end{aligned}$$

from which the estimate follows by observing that $\tilde{c}_n \in [c, c + \frac{1}{n}]$. \square

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