# Representations and Inequalities for Ising Model Ursell Functions 

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#### Abstract

We describe and investigate representations for the Ursell function $u_{n}$ of a family of $n$ random variables $\left\{\sigma_{i}\right\}$. The representations involve independent but identically distributed copies of the family. We apply one of these representations in the case that the random variables are spins of a finite ferromagnetic Ising model with quadratic Hamiltonian to show that $(-1)^{\frac{n}{2}+1} u_{n}\left(\sigma_{1}, \ldots, \sigma_{n}\right) \geqq 0$ for $n=2,4$, and 6 by proving the stronger statement $\left.(-1)^{\frac{n}{2}+1} \frac{\partial^{m}}{\partial J_{i_{1} j_{1}} \cdots \partial J_{i_{m} j_{m}}} Z^{\frac{n}{2}} u_{n}\right|_{J=0} \geqq 0$ for $n=2,4$, and 6 , the $J_{i}$, being coupling constants in the Hamiltonian and $Z$ the partition function. For general $n$ we combine this result with various reductions to show that sufficiently simple derivatives of $(-1)^{\frac{n}{2}+1} Z^{\frac{n}{2}} u_{n}$, evaluated at zero coupling, are nonnegative. In particular, we conclude that $(-1)^{\frac{n}{2}+1} u_{n} \geqq 0$ if all couplings are nonzero and the inverse temperature $\beta$ is sufficiently small or sufficiently large, though this result is not uniform in the order $n$ or the system size. In an appendix we give a simple proof of recent inequalities which bound $n$-spin expectations by sums of products of simpler expectations.


## 1. Introduction

The Ursell function $u_{n}\left(\sigma_{1}, \ldots, \sigma_{n}\right)$ of a family $\left\{\sigma_{i}\right\}$ of $n$ arbitrary random variables may be defined by means of a generating function as

$$
\begin{equation*}
u_{n}\left(\sigma_{1}, \ldots, \sigma_{n}\right)=\left.\frac{\partial^{n}}{\partial \lambda_{1} \cdots \partial \lambda_{n}} \ln \mathscr{E}\left(\exp \left[\sum_{i=1}^{n} \lambda_{i} \sigma_{i}\right]\right)\right|_{\lambda=0} \tag{1.1}
\end{equation*}
$$

Here $\mathscr{E}$ is the expectation integral; we assume all the necessary expectations are finite. The Ursell function may be defined recursively by

$$
\begin{equation*}
\mathscr{E}\left(\sigma_{1} \sigma_{2} \cdots \sigma_{n}\right)=\sum_{\mathscr{P}} \prod_{P \in \mathscr{P}} u_{|P|}\left(\sigma_{p_{a}}, \sigma_{p_{b}}, \ldots\right) \tag{1.2}
\end{equation*}
$$

Here $\mathscr{P}$ is a partition of $\{1, \ldots, n\}$, a set $P \in \mathscr{P}$ has elements $p_{a}, p_{b}$, etc., and $|P|$ denotes the cardinality of $P$. Finally, $u_{n}\left(\sigma_{1}, \ldots, \sigma_{n}\right)$ may be defined explicitly by

$$
\begin{equation*}
u_{n}\left(\sigma_{1}, \ldots, \sigma_{n}\right)=\sum_{\mathscr{P}}(-1)^{|\mathscr{P}|-1}(|\mathscr{P}|-1)!\prod_{P \in \mathscr{P}} \mathscr{E}\left(\prod_{p \in P} \sigma_{p}\right) \tag{1.3}
\end{equation*}
$$

where again $\mathscr{P}$ is a partition of $\{1, \ldots, n\}$.

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