

Twistless KAM Tori

Giovanni Gallavotti¹

CNRS-CPT Luminy, case 907, F-13288 Marseille, France

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Abstract: A selfcontained proof of the KAM theorem in the Thirring model is discussed.

I shall particularize the Eliasson method, [E], for KAM tori to a special model, of great interest, whose relevance for the KAM problem was pointed out by Thirring, [T] (see [G] for a short discussion of the model). The idea of exposing Eliasson's method through simple particular cases appears in [V], where results of the type of the ones discussed here, and more general ones, are announced.

The connection between the methods of [E] and the tree expansions in the renormalization group approaches to quantum field theory and many body theory can be found also in [G]. The connection between the tree expansions and the breakdown of invariant tori is discussed in [PV].

The Thirring model is a system of rotators interacting via a potential. It is described by the hamiltonian (see [G] for a motivation of the name):

$$\frac{1}{2} J^{-1} \vec{A} \cdot \vec{A} + \varepsilon f(\vec{\alpha}), \quad (1)$$

where J is the (diagonal) matrix of the inertia moments, $\vec{A} = (A_1, \dots, A_l) \in \mathbb{R}^l$ are their angular momenta and $\vec{\alpha} = (\alpha_1, \dots, \alpha_l) \in T^l$ are the angles describing their positions: the matrix J will be supposed nonsingular; but we only suppose that $\min_{j=1, \dots, l} J_j = J_0 > 0$, and *no assumption* is made on the size of the *twist rate* $T = \min J_j^{-1}$: the results will be uniform in T (hence the name “twistless”: this is not a contradiction with the necessity of a twist rate in the general problems, see problems 1, 16, 17 in Sect. 5.11 of [G2], and [G]). We suppose f to be an even

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¹ E-mail: gallavotti%40221.hepnet@lbl.gov; permanent address: Dipartimento di Fisica, Università di Roma, “La Sapienza,” P. Moro 2, I-00185 Roma, Italy