

The Newtonian Limit of the Spherically Symmetric Vlasov–Einstein System

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Abstract. We prove that spherically symmetric solutions of the Vlasov–Einstein system with a fixed initial value converge to the corresponding solution of the Vlasov–Poisson system if the speed of light c is taken as a parameter and tends to infinity. The convergence is uniform on compact time intervals with convergence rate $1/c^2$. Thus the classical Vlasov–Poisson system appears as the Newtonian limit of the general relativistic Vlasov–Einstein system in a spherically symmetric setting.

1. Introduction

Consider an ensemble of particles (stars in a galaxy, galaxies in a galaxy cluster, etc.) which are all of the same mass and interact only by the gravitational field which they generate collectively. The ensemble is described by a time dependent density function f on phase space, and this function satisfies a continuity equation, the Vlasov or Liouville equation, which is coupled to the field equations with source terms generated by f . In the classical, Newtonian setting we obtain the Vlasov–Poisson system (VP)

$$\begin{aligned} \partial_t f + v \cdot \partial_x f - K(t, x) \cdot \partial_v f &= 0, \\ K(t, x) &:= \frac{x}{r} \frac{1}{r^2} \int_{|y| \leq r} \int_{\mathbb{R}^3} f(t, y, v) dv dy, \end{aligned}$$

where $t \geq 0$ denotes time, $x \in \mathbb{R}^3$ position, and $v \in \mathbb{R}^3$ momentum, $r := |x|$, and we have assumed that the system is spherically symmetric, i.e. $f(t, Ax, Av) = f(t, x, v)$, $K(t, Ax) = AK(t, x)$ for any orthogonal matrix A and $t \geq 0, x, v \in \mathbb{R}^3$.

If we wish to describe the above situation in the setting of general relativity, we obtain the Vlasov–Einstein system (VE γ) in the following form:

$$\begin{aligned} \partial_t f + e^{\mu-\lambda} \frac{v}{\sqrt{1+\gamma v^2}} \cdot \partial_x f - \left(\lambda \frac{\dot{x} \cdot v}{r} + e^{\mu-\lambda} \frac{1}{\gamma} \mu' \sqrt{1+\gamma v^2} \right) \frac{x}{r} \cdot \partial_v f &= 0, \\ e^{-2\lambda}(2r\lambda' - 1) + 1 &= 8\pi\gamma r^2 \rho, \\ e^{-2\lambda}(2r\mu' + 1) - 1 &= 8\pi\gamma^2 r^2 p, \end{aligned}$$