

# Ward Identities for Non-Commutative Geometry\*

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*Dedicated to Res Jost and Arthur Wightman*

**Abstract.** We interpret the cocycle condition in quantum field theory as a set of integrated Ward identities for non-commutative geometry.

## I. Basic Notions

The Wightman functions of a super-symmetric quantum field theory given by a super-trace functional have a geometric or cohomological interpretation. This was shown in joint work of the authors with Lesniewski [JLO1]. This property is summarized by the construction of a cocycle  $\tau$  for the  $\partial$ -complex of entire cyclic cohomology, namely a solution to the equation

$$\partial\tau = 0. \quad (\text{I.1})$$

Here  $\tau$  is a time average of certain Euclidean Wightman functionals, and  $\partial$  is a standard coboundary operator of non-commutative geometry, see [C, JLO1]. Furthermore, this natural cohomological interpretation can be generalized to the case when the Wightman functionals  $\tau$  are constructed from a finite-temperature functional satisfying a super-version of the KMS condition of statistical mechanics [K, JLO2, JLWis]. In this case the super-trace (associated with finite volume theories) need not exist – only the super-KMS functional has to be defined.

In this note we investigate how Eq. (I.1) has a set of identities as its consequence, which can be interpreted as a symmetry of the Wightman functions. We call these identities “Ward identities for non-commutative geometry.” Let us consider an example. Let  $\gamma$  denote the  $Z_2$  grading in our theory, equal to  $(-I)^{N_f}$  in models, and let  $A \rightarrow A^\gamma = \gamma A \gamma$  denote the action of the grading on field operators. Let  $A(t) = e^{-tH} A e^{tH}$  denote the propagation of  $A$  to Euclidean time  $t$ , and let  $\langle A \rangle$  denote the expectation

$$\langle A \rangle = \text{Str}(A e^{-H}), \quad (\text{I.2})$$

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