

# A Boson Representation for SU ( $N$ ) Lattice Gauge Theories

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**Abstract.** SU( $N$ ) lattice gauge theories are reformulated in terms of fields varying over non-compact spaces  $\mathbb{C}^N$ , transforming as  $N$  dimensional representations of SU( $N$ ) and integrated with Gaussian measure. This reformulation is equivalent to a boson operator representation. Strong coupling expansions based on this formalism do not involve SU( $N$ ) vector coupling coefficients.

## 1. Introduction

In pure Euclidean Yang–Mills field theories on a lattice field, variables range over the group manifold itself. This manifold is compact and a non-trivial Riemannian space. The gauge groups we will consider are SU( $N$ ),  $N = 2, 3$  but our results can be immediately generalized to any  $N$ . In this article we reformulate such theories in an equivalent fashion in terms of fields taken from the flat non-compact space  $\mathbb{C}^N$ . They transform as  $N$  dimensional representations of SU( $N$ ). We will therefore call them “bosonic spinorial variables” for the gauge field. The integration is over a Gaussian measure instead of a Haar measure. A straightforward change of notation leads then to a boson operator formulation of Yang–Mills lattice field theories.

Our approach is based on Bargmann’s realization of group representations of SU( $N$ ) [1], which makes use of Hilbert spaces of entire analytic functions over  $\mathbb{C}^N$  or powers of  $\mathbb{C}^N$ . This formalism is equivalent to the so-called boson operator calculus [2]. For technical reasons and for the sake of mathematical clarity we prefer to use spaces of analytic functions in this article.

The lattice  $\Lambda$  is assumed to be hypercubic, to have dimension  $D$  and the boundary conditions are presumed to be periodic. Let  $\ell$  denote the links and  $p$  the plaquettes of  $\Lambda$ . We define the partition function by the standard ansatz

$$Z = \int \prod_{\ell \in \Lambda} (du_\ell) e^{S(u_i)},$$

$$u_\ell \in \text{SU}(N) \tag{1}$$

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