

# The Boltzmann Equation with a Soft Potential

## II. Nonlinear, Spatially-Periodic

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**Abstract.** The results of Part I are extended to include linear spatially periodic problems—solutions of the initial value are shown to exist and decay like  $e^{-\lambda t^\beta}$ . Then the full non-linear Boltzmann equation with a soft potential is solved for initial data close to equilibrium. The non-linearity is treated as a perturbation of the linear problem, and the equation is solved by iteration.

### 1. Introduction

The linear Boltzmann equation with a soft intermolecular potential was solved globally in time in Part I [1], if the initial density is a spatially homogeneous perturbation of a global Maxwellian. Moreover it was proven that this perturbation decays in  $\mathcal{L}^2$  or sup norm like  $e^{-\lambda t^\beta}$ , with  $\lambda > 0, 1 > \beta > 0$ , if it is initially bounded by a Maxwellian. We will refer to formulas or results from Part I by preceding their numbers with an “I” as in (I1.7).

In this paper we find the same result even if the initial perturbation is *spatially dependent* in the cube with periodic boundary conditions. In addition we can solve the spatially periodic *nonlinear* problem globally in time if the initial perturbation is small enough, and we find that the solution decays to the Maxwellian equilibrium.

The linear, spatially-dependent Boltzmann equation is

$$\frac{\partial}{\partial t} f + \xi \cdot \frac{\partial}{\partial \mathbf{x}} f + Lf = 0, \tag{1.1}$$

$$f(t=0) = f_0 \in \mathcal{N}, \tag{1.2}$$

where  $f_0$  and  $f = f(t, \mathbf{x}, \xi)$  are periodic in  $\mathbf{x} \in T^3 = [0, 2\pi]^3, t \geq 0, \xi \in \mathbb{R}^3$ , and  $\mathcal{N} = \left\{ g(\mathbf{x}, \xi) : \int_{T^3} \int_{\mathbb{R}^3} \psi(\xi) g(\mathbf{x}, \xi) d\xi d\mathbf{x} = 0 \text{ for } \psi(\xi) = 1, \xi_i, \text{ or } \xi_i^2 \right\}$ . The requirement that  $f_0 \in \mathcal{N}$  just means that we have chosen the right Maxwellian equilibrium to

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