Commun. Math. Phys. 71, 95-130 (1980)

Ultraviolet Stability in Euclidean Scalar Field Theories

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Abstract. We develop a technique for reducing the problem of the ultraviolet divergences and their removal to a free field problem. This work is an example of a problem to which a rather general method can be applied. It can be thought as an attempt towards a rigorous version (in 2 or 3 space-time dimensions) of the analysis of the structure of the functional integrals developed in [9], the underlying mechanism being essentially the same as in [11, 3].

1. Introduction

The free euclidean field in \mathbb{R}^d is the gaussian field with covariance operator:

$$C = (1 - D)^{-1}, (1.1)$$

where D is the Laplace operator in \mathbb{R}^d .

We call $(\varphi_{\xi})_{\xi \in \mathbb{R}^d}$ the random field with the above covariance and we shall represent φ as a sum of independent, identically distributed up to scale factors, random fields $(\varphi_{\xi}^{(N)})_{\xi \in \mathbb{R}^d}$, $N = 0, 1, \ldots$ The field $\varphi^{(N)}$ has, by definition, the following covariance operator:

$$C^{(N)} = (\gamma^{2N} - D)^{-1} - (\gamma^{2N+2} - D)^{-1}$$
(1.2)

and γ will be appropriately chosen close to 1. From (1.2) it follows that if $d \leq 3$ the kernel $C_{\xi\eta}^{(N)}$ of $C^{(N)}$ in \mathbb{R}^d is finite when $\xi = \eta$:

$$C_{\xi\xi}^{(N)} \equiv C_{00}^{(N)} = \gamma^{(d-2)N} \frac{\gamma^2 - 1}{(2\pi)^d} \int \frac{d^d k}{(1+k^2)(\gamma^2 + k^2)} \equiv c_\gamma \gamma^{(d-2)N}.$$
(1.3)

Hence it will be convenient to introduce the normalized field:

$$z_{\xi}^{(N)} = \frac{\varphi_{\xi}^{(N)}}{\sqrt{2C_{\xi\xi}^{(N)}}} = \frac{\varphi_{\xi}^{(N)}}{\sqrt{2c_{\gamma}\gamma^{(d-2)N}}}$$
(1.4)

^{*} Supported by IHES, through financial support from the "Stiftung Volkswagenwerk"