

KMS Conditions and Local Thermodynamical Stability of Quantum Lattice Systems. II

Geoffrey L. Sewell

Department of Physics, Queen Mary College, London E1 4NS, England

Abstract. We prove that local thermodynamical stability (LTS), as defined in [1], implies the KMS conditions in quantum lattice systems, without any assumption of translational invariance. This result, together with those of [1], establishes the equivalence between the LTS and the KMS conditions for such systems.

Section 1

In an article by Araki and the author [1], the concept of local thermodynamical stability (LTS) was defined for quantum lattice systems; and it was shown that, if the forces were suitably tempered, the LTS conditions were implied by, and in the case of translationally invariant states equivalent to, those of Kubo-Martin-Schwinger (KMS). In the present article, we prove that the LTS conditions imply those of KMS, without any assumption of translational invariance, and thus establish the following theorem.

Theorem 1. *The LTS and KMS conditions are mutually equivalent for quantum lattice systems, subject to the same assumptions on the interactions as in [1].*

Comment. It has already been observed (cf. Note following Definition 2.2 in [2]) that, for classical lattice and hard-core continuous systems, the LTS conditions are equivalent to those of Dobrushin-Lanford-Ruelle (DLR). Thus, in view of the above theorem, we now conclude that $\text{LTS} \equiv \text{KMS}$ for quantum lattice systems, and $\equiv \text{DLR}$ for classical lattice and hard-core continuous ones.

Our notation will be based on that of [1]. We take Γ to be the lattice on which the system is situated; here it suffices to consider Γ as a denumerably infinite point set. The family $\{A\}$ of finite point subsets of Γ will be denoted by L . The algebra of observables, $\mathcal{A}(A)$, for the region $A (\in L)$ will be assumed to be a finite-dimensional, type-I factor, that is isotonic with respect to A and commutes with $\mathcal{A}(A')$ if $A \cap A' = \emptyset$; and the C^* -algebra of observables, \mathcal{A} , for the system will be taken to be the norm completion of $\mathcal{A}_L \equiv \bigcup_{A \in L} \mathcal{A}(A)$. The state space, Ω , of the system will be