

## Parity Operator and Quantization of $\delta$ -Functions

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**Abstract.** In the Weyl quantization scheme, the  $\delta$ -function at the origin of phase space corresponds to the parity operator. The quantization of a function  $f(v)$  on phase space is the operator  $\int f(v/2)W(v)dvM$ , where  $M$  is the parity and  $W(v)$  the Weyl operator.

### Introduction

We are concerned here with the elementary problem of writing down an operator  $Q(f)$  which quantizes a function  $f$  on (flat) phase space. The existing solutions [1] (see also [2]) all involve, to the best of our knowledge, the performing of Fourier transforms. By contrast, our equation (10 bis) picks up local contribution from the classical function and also exhibits a rather unexpected role played by the parity operator.

### 1. Displaced Parity Operators

Let  $E$  be the phase space for  $v < \infty$  degrees of freedom, i.e. a  $2v$ -dimensional vector space over  $\mathbb{R}$ , with a symplectic form  $\sigma(v, a) \cdot (a, v \in E)$ . Let  $v \rightarrow W(v)$  ( $v \in E$ ) be a Weyl system over  $E$ , i.e. a strongly continuous family of unitary operators acting irreducibly on a separable Hilbert space  $\mathcal{H}$  and satisfying

$$W(a)W(v) = e^{i\sigma(a, v)}W(a+v). \quad (1)$$

We have introduced the abbreviation

$$e^{i\sigma(a, v)} = e^{2i\pi\sigma(a, v)}. \quad (2)$$

The family  $W'(v) = W(-v)$  also satisfies (1). By the uniqueness theorem of von Neumann, there exists in  $\mathcal{H}$  a unitary operator  $M$ , determined up to a phase, and such that  $W(v)M = MW(-v)$  for every  $v \in E$ . Since  $M^2$  commutes with the irreducible family of operators  $W(v)$ , it is a number of modulus 1, which can be adjusted to 1 by a multiplication of a suitable number  $e^{i\theta}$  to  $M$ . Then  $M = M^*$  and  $M$  is determined up to a sign.