

# On the Uniqueness of the Equilibrium State for Ising Spin Systems

JOEL L. LEBOWITZ\*

Belfer Graduate School of Science, Yeshiva University, New York, N.Y.

ANDERS MARTIN-LÖF

Rockefeller University, New York, N.Y.

Received November 10, 1971

**Abstract.** We show that for an Ising spin system of arbitrary spin with a ferromagnetic pair interaction and a “periodic” external magnetic field there is a unique equilibrium state if and only if the magnetization is continuous with respect to a uniform change in the external field. Hence, if the critical temperature  $T_c$  is defined as the temperature where the spontaneous magnetization (which is a non-increasing function of the temperature) becomes positive, then the equilibrium state is unique for  $T > T_c$  and is non-unique for  $T < T_c$  (when the external field is zero). This implies that the correlation functions have a cluster property for  $T > T_c$ .

We also show that for an anti-ferromagnet consisting of two sublattices there is a unique equilibrium state if and only if the staggered magnetization is continuous with respect to a change in the staggered field.

## I. Introduction

We consider an Ising spin system with a ferromagnetic pair interaction in a finite box  $A$  on a  $d$ -dimensional lattice  $\mathbb{Z}^d$ , i.e. at each point  $p$  of the lattice there is a spin  $\sigma_p = \pm 1$ , and the conditional probability of a spin configuration in the box given a configuration outside it is proportional to

$$\exp(-E_A) = \exp\left(\frac{1}{2} \sum_{p \neq q \in A} J_{p-q} \sigma_p \sigma_q + \sum_{p \in A} \sigma_p \left(H_p + h + \sum_{q \notin A} J_{p-q} \sigma_q\right)\right). \quad (1)$$

$J_p = J_{-p} \geq 0$  is the pair interaction,  $\sum_{p \in \mathbb{Z}^d} J_p < \infty$ , and  $H_p + h$  an external magnetic field. The reciprocal temperature  $\beta$  has been included in the Hamiltonian. The external field consists of a uniform part  $h$  and a periodic part  $H_p$ , i.e.  $H_p = H_{p+g}$  when  $g$  is contained in some subgroup  $G$  of  $\mathbb{Z}^d$ . A boundary condition for the box  $A$  is specified by giving a probability distribution  $b_A(d\sigma)$  for the configurations outside  $A$ .

The (equilibrium) state of the system in  $A$  is the probability distribution for configurations in  $A$  defined by (1) together with  $b_A$  or equivalently

\* Supported in part by U.S.A.F.O.S.R. under contract F 44620-71-C-0013, P001.