

## CHERN-SIMONS PERTURBATION THEORY. II

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### Abstract

In a previous paper [2], we used superspace techniques to prove that perturbation theory (around a classical solution with no zero modes) for Chern-Simons quantum field theory on a general 3-manifold  $M$  is finite. We conjectured (and proved for the case of 2-loops) that, after adding counterterms of the expected form, the terms in the perturbation theory defined topological invariants. In this paper we prove this conjecture. Our proof uses a geometric compactification of the region on which the Feynman integrand of Feynman diagrams is smooth as well as an extension of the basic propagator of the theory.

### 1. Introduction

In a previous paper [2], we considered the perturbative expansion for three-dimensional Chern-Simons quantum field theory about a solution  $A_0$  to the equations of motion. We defined what we meant by the perturbative expansion and showed perturbation theory was finite. We showed that the first term in the perturbative expansion beyond the semiclassical limit defines a geometric invariant precisely in the manner one would expect based on Witten's exact solution [10]. We conjectured and gave strong evidence that the higher terms in the expansion were geometric invariants of the same type. In this paper we prove this conjecture.

More specifically, we take  $A_0$  to be a flat connection on a principal bundle  $P$  with a compact structure group  $G$  and a closed, oriented, three-dimensional base  $M$ . We also assume that  $A_0$  has no zero modes, i.e., that the cohomology of the exterior derivative operator  $D: \Omega^*(M, \mathfrak{g}) \rightarrow \Omega^{*+1}(M, \mathfrak{g})$ , coupled to the adjoint bundle  $\mathfrak{g}$  of  $P$  and  $A_0$ , vanishes. By rewriting the Lorentz gauge fixed theory as a superspace theory in [2], we were able to obtain Feynman rules that could be translated succinctly into the language of differential forms. To define the gauge fixing it was necessary to choose a Riemannian metric  $g$  on  $M$ . For  $l \geq 2$ , the

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