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## Periodic Solution of Classical Hamiltonian Systems

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## Introduction

Let  $x = (x_1, x_2, \dots, x_n)$ ,  $y = (y_1, y_2, \dots, y_n) \in \mathbb{R}^n$  and  $H = H(x, y): \mathbb{R}^{2n} \to \mathbb{R}$ be a smooth  $(\mathbb{C}^{\infty})$  function.

We consider a Hamiltonian system of n degrees of freedom

(1)

$$\dot{x}_i = H_{y_i}$$
,  $\dot{y}_i = -H_{x_i}$ ;  $i = 1, 2, \dots, n$ .

P. Rabinowitz [3] proved that "if an energy surface

$$H^{-1}(e) = \{(x, y) \in \mathbb{R}^{2n}; H(x, y) = e\}$$

is star-shaped, then there exists at least one periodic solution of (1) on it".

"star-shaped" implies "diffeomorphic to the sphere  $S^{2n-1}$ ", but it is not known whether the condition "star-shaped" can be replaced by "diffeo. to  $S^{2n-1}$ " or not. This is a generalized Hamiltonian version of the Seifert Conjecture (Has any sufficiently smooth flow on  $S^3$  periodic orbit ?) in the theory of dynamical systems.

Classically, the system (1) is derived from the Lagrangian system and, in the time-independent case, the Hamiltonian H is the sum of the kinetic energy and the potential. So we define

DEFINITION. A Hamiltonian H=H(x, y) is called *classical* if it has the form

(2) 
$$H = \sum_{i,j=1}^{n} a^{ij}(x) y_i y_j + U(x) ,$$

where  $a^{ij}$ ,  $U: \mathbb{R}^n \to \mathbb{R}$  are smooth functions and for any  $x \in \mathbb{R}^n$ , the matrix  $(a^{ij}(x))$  is symmetric and positive definite.

The system (1) with classical Hamiltonian H is called a *classical* Hamiltonian system.

In this paper we have

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