

106. Smooth Global Solutions for the One-Dimensional Equations in Magnetohydrodynamics

By Shuichi KAWASHIMA*) and Mari OKADA**)

(Communicated by Kôzaku YOSIDA, M. J. A., Nov. 12, 1982)

§ 1. Introduction. The motion of electrically conducting fluids on one space coordinate is described by the equations in Lagrangian representation ([1]):

$$(1) \quad \begin{cases} (1/\rho)_t - u_x = 0, \\ u_t + (p + |B|^2/2\mu_0)_x = (\nu\rho u_x)_x, & v_t - (\bar{B}^1 B/\mu_0)_x = (\mu\rho v_x)_x, \\ \theta_t + (\theta p_\theta/e_\theta)u_x = (1/e_\theta)\{(\kappa\rho\theta_x)_x + \nu\rho u_x^2 + \mu\rho|v_x|^2 + (\rho/\sigma\mu_0^2)|B_x|^2\}, \\ (B/\rho)_t - (\bar{B}^1 v)_x = \{(\rho/\sigma\mu_0)B_x\}_x. \end{cases}$$

Here $\rho > 0$, $\mathbf{u} = (u^1, u^2, u^3) \in \mathbf{R}^3$, $\theta > 0$ and $\mathbf{B} = (\bar{B}^1, B^2, B^3) \in \mathbf{R}^3$ represent the mass density, the velocity, the absolute temperature and the magnetic induction, where we write $u = u^1$, $v = (u^2, u^3)$, $B = (B^2, B^3)$, and \bar{B}^1 is a constant.

We assume that the pressure p and the internal energy e are smoothly related to ρ and θ by the equations of state

$$(2) \quad p_\rho > 0, \quad e_\theta > 0, \quad de = \theta dS - p d(1/\rho),$$

where $S = S(\rho, \theta)$ is the entropy; the coefficients of viscosity μ, ν , the coefficient of heat conductivity κ and the coefficient of electrical resistivity $1/\sigma$ (σ : the coefficient of electrical conductivity) are all smooth functions of ρ and θ , and are positive or identically zero; μ_0 is the magnetic permeability, now a positive constant.

In this paper, we seek smooth solutions of (1) in a small neighborhood of a constant state $(\rho, \mathbf{u}, \theta, B) = (\bar{\rho}, 0, \bar{\theta}, \bar{B})$ where $\bar{\rho} > 0$, $\bar{\theta} > 0$ and $\bar{B} \in \mathbf{R}^2$ are arbitrary fixed constants. To obtain the a priori estimates for the solutions, we use the following energy form ([4]):

$$\mathcal{E} = e - \bar{e} + \bar{p}(1/\rho - 1/\bar{\rho}) - \bar{\theta}(S - \bar{S}) + |\mathbf{u}|^2/2 + |B - \bar{B}|^2/2\mu_0\rho,$$

where $\bar{e} = e(\bar{\rho}, \bar{\theta})$ and so on. Note that if $|\rho - \bar{\rho}, \theta - \bar{\theta}|$ is small, \mathcal{E} is equivalent to the quadratic form $|\rho - \bar{\rho}, \mathbf{u}, \theta - \bar{\theta}, B - \bar{B}|^2$. This is based on the strict convexity of the internal energy e as a function of $1/\rho$ and S .

From (1) and (2), we have the energy conservation law:

$$(3) \quad \begin{aligned} & (e + |\mathbf{u}|^2/2 + |B|^2/2\mu_0\rho)_t + \{(p + |B|^2/2\mu_0)u - (\bar{B}^1 B/\mu_0) \cdot v\}_x \\ & = \{\nu\rho u u_x + \mu\rho v \cdot v_x + \kappa\rho\theta_x + (\rho/\sigma\mu_0^2)B \cdot B_x\}_x, \end{aligned}$$

and the equation of entropy:

*) Department of Mathematics, Nara Women's University.

**) Department of Mathematics, Kyoto University.