

149. A Necessary Unitary Field Theory as a Non-Holonomic Parabolic Lie Geometry Realized in the Three-Dimensional Cartesian Space. II

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The purpose of the present paper consists in the following five points: to deduce (i), (ii), (ii'), (iii), (iii') mentioned below from Part I (these Proc., 29 (1953).

Since the three-dimensional non-holonomic Laguerre (parabolic Lie) geometry is in law a four-dimensional teleparallelism geometry keeping the Riemann (Weyl) metric, it is remarkable that the following conjecture of Prof. Einstein of 1928, which seems to be now scarcely considered, must acquire a renaissance: "Es ist denkbar, dass diese Theorie die ursprüngliche Fassung der allgemeinen Relativitätstheorie verdrängen wird".

(i) A Unitary Field Theory of a Single Particle

6. *A Necessary Unitary Field Theory of a Single Particle Charged with Rest-mass m_0 and Constant Electricity $-e$.* In Art. 4, we have solved a two particles problem stated in Art. 2 and the resulting generalizations of the Maxwell's equations were (4.24), (4.25), (4.26) and (4.27). Thereby the continuity condition (4.23) was assumed. Now in the case of a single particle P , we have

$$(6.1) \quad \bar{m}_0 = \bar{e} = \bar{\epsilon}^i = \bar{\epsilon}^4 = \bar{\sigma}^i = \bar{\sigma}^4 = \bar{\mathcal{X}}^i = \bar{X}^i = \bar{a}^i = \bar{a}^4 = 0.$$

Hence (4.24), (4.25), (4.26) and (4.27) become the *necessary-unitary-field-theoretical generalization of the Maxwell's equations*:

$$(6.2) \quad \frac{\partial}{\partial \omega^i} (\mathcal{X}^i + eX^i) = \epsilon^i + \sigma^i, \frac{\partial}{\partial \omega^4} (a^i + ea^i) + \frac{\partial}{\partial \omega^j} (\mathcal{X}^k + eX^k) - \frac{\partial}{\partial \omega^k} (\mathcal{X}^j + eX^j) = 0,$$

$$(6.2') \quad \frac{\partial}{\partial \omega^j} (a^k + ea^k) - \frac{\partial}{\partial \omega^k} (a^j + ea^j) - \frac{\partial}{\partial \omega^4} (\mathcal{X}^i + eX^i) = \epsilon^i + \sigma^i, \frac{\partial}{\partial \omega^i} (a^i + ea^i) = 0.$$

7. *General-Relativistic and Necessary-Unitary-Field-Theoretical Generalization of the Dirac Equation in the Case of a Single Particle.* In the case of a single particle P , the *general-relativistic and necessary-unitary-field-theoretical generalizations*

$$(5.3) \quad \left[\gamma_i \left(\frac{\hbar}{2\pi i} E \frac{\partial}{\partial \omega^i} + e\phi^i + \frac{\hbar}{2\pi i} \bar{E} \frac{\partial}{\partial \omega^i} + \bar{e}\bar{\phi}^i \right) + \gamma_4 \left(\frac{\hbar}{2\pi i} E \frac{\partial}{\partial \omega^4} + e\phi^4 + \bar{m}_0 \bar{E}^2 \right) \right. \\ \left. + \gamma_5 \left(\frac{\hbar}{2\pi i} \bar{E} \frac{\partial}{\partial \omega^5} + \bar{e}\bar{\phi}^5 + m_0 E^2 \right) \right] \psi = 0,$$