

## On the Polynomial Hamiltonian Structure Associated with the $L(1, g + 2; g)$ Type

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**1. Introduction.** In the paper [3], we considered the differential equation

$$(1.1) \quad \frac{d}{dx} Y = \frac{1}{x} \mathcal{A}(x, t) Y, \quad \mathcal{A}(x, t) := \sum_{k=0}^{g+1} \mathcal{A}_k x^k,$$

which satisfies the conditions:

- (i)  $\mathcal{A}_k$  are  $2 \times 2$  matrices,
- (ii) the eigenvalues of  $\mathcal{A}_0$  are distinct up to additive integers,
- (iii) the eigenvalues of  $\mathcal{A}_{g+1}$  are distinct.

This equation can be reduced to the equation

$$(1.2) \quad \frac{d^2}{dx^2} y + p_1(x, t) \frac{d}{dx} y + p_2(x, t) y = 0,$$

which satisfies the three conditions:

- (iv) The Riemann scheme of (1.1) is

$$\left\{ \begin{array}{ccccccc} x = 0 & x = \lambda_1 & \cdots & x = \lambda_g & & & \\ \begin{array}{ccccccc} 0 & 0 & \cdots & 0 & & & \\ \kappa_0 & 2 & & 2 & & & \end{array} & & & & & & \\ & & & & x = \infty & & \\ & 0 & 0 & 0 & \cdots & 0 & -\kappa_\infty \\ & \frac{1}{g+1} & \frac{t_g}{g} & \frac{t_{g-1}}{g-1} & \cdots & t_1 & \kappa_\infty - \kappa_0 + 1 \end{array} \right\},$$

- (v)  $\kappa_0$  and  $\kappa_\infty$  are not integer,
- (vi)  $x = \lambda_k$  ( $k = 1, \dots, g$ ) are non-logarithmic singular points.

The equation (1.2) is called  $L(1, g + 2; g)$  type.

Let  $\mu_k$  ( $k = 1, \dots, g$ ) be the residue of  $p_2(x, t)$  at  $x = \lambda_k$  and let

$$h_j := \frac{\partial^{g-j}}{\partial x^{g-j}} \left( x p_2 - \sum_{k=1}^g \frac{\lambda_k \mu_k}{x - \lambda_k} \right) \Big|_{x=0}.$$

By the assumption (vi), we remark that these  $h_j$  are uniquely determined as rational functions in  $\lambda_k, \mu_k, t_k$  ( $k = 1, \dots, g$ ).

Using these notations  $\lambda_k, \mu_k$  and  $h_k$ , we state that the holonomic deformation of the linear equation (1.2) is governed by the Hamiltonian system:

$$\frac{\partial \lambda_i}{\partial t_j} = \frac{\partial \tilde{K}_j}{\partial \mu_i}, \quad \frac{\partial \mu_i}{\partial t_j} = \frac{\partial \tilde{K}_j}{\partial \lambda_i} \quad (i, j = 1, \dots, g),$$

where the Hamiltonian  $\tilde{K}_j$  are

$$\begin{bmatrix} \tilde{K}_1 \\ \tilde{K}_2 \\ \vdots \\ \tilde{K}_g \end{bmatrix} = \begin{bmatrix} 1 & & & 0 \\ & 2 & & \\ & & \ddots & \\ & & & g \\ 0 & & & & 0 \end{bmatrix}^{-1} \begin{bmatrix} 1 & & & 0 \\ t_g & 1 & & \\ \vdots & t_g & 1 & \\ t_3 & & \ddots & \ddots \\ t_2 & t_3 & \cdots & t_g & 1 \end{bmatrix}^{-1} \begin{bmatrix} h_1 \\ h_2 \\ \vdots \\ h_g \end{bmatrix}.$$

It is known that, if  $g = 1$  the holonomic deformation of (1.2) is governed by the fourth Painlevé equation. So the Hamiltonian system  $(\lambda, \mu, \tilde{K}, t)$  is an extension of the fourth Painlevé equation.

Now, we assume that  $g < 8$ . The purpose of this paper is to transform the Hamiltonian system  $(\lambda, \mu, \tilde{K}, t)$  into the Hamiltonian system  $(q, p, H, \xi)$  with the conditions:

- (C<sub>1</sub>)  $H_j$  ( $j = 1, \dots, g$ ) are polynomial in  $q_k, p_k$  ( $k = 1, \dots, g$ ),
- (C<sub>2</sub>)  $\frac{\partial H_k}{\partial \xi_j} = \frac{\partial H_j}{\partial \xi_k}$ .

We will state that if  $\kappa_\infty = 0$ , a special solution of this system is written by the multivalued Hermite function.

By the condition (C<sub>2</sub>), we can introduce the function  $\tau_{IV}^{(g)}$  as follows:

$$\frac{\partial}{\partial \xi_k} \log \tau_{IV}^{(g)} = H_k.$$

In [7], Okamoto defined the  $\tau$  function associated with the fourth Painlevé transcendental function. This function is equivalent to  $\tau_{IV}^{(1)}$ .

### 2. Polynomial Hamiltonian structure.

**Theorem 2.1.** Put  $\sigma_k, \rho_k$  ( $k = 1, \dots, g$ ) as follows:

$\sigma_k$  = the  $k$ -th elementary symmetric function of  $\lambda_1, \lambda_2, \dots, \lambda_g$ ,

$$\rho_k = (-1)^{k-1} \sum_{l=1}^g \lambda_l^{g-k} \mu_l \prod_{\substack{j=1 \\ j \neq l}}^g (\lambda_l - \lambda_j)^{-1}.$$