

27. Path Integral for the Weyl Quantized Relativistic Hamiltonian^{†)}

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1. Introduction. The aim of this note is to give a path integral representation of the solution of the Cauchy problem for

$$(1.1) \quad \partial_t \psi(t, x) = -[H - mc^2] \psi(t, x), \quad t > 0, \quad x \in \mathbf{R}^d.$$

Here c is the light velocity. H is the quantum Hamiltonian via the Weyl correspondence, i.e. the pseudo-differential operator ([1], [2], [6])

$$(1.2) \quad (Hg)(x) = (2\pi)^{-d} \iint_{\mathbf{R}^d \times \mathbf{R}^d} e^{i(x-y)p} \hbar \left(p, \frac{x+y}{2} \right) g(y) dy dp, \quad g \in \mathcal{S}(\mathbf{R}^d),$$

associated with the classical Hamiltonian

$$(1.3) \quad h(p, x) = [cp - eA(x)]^2 + m^2 c^4 + e\Phi(x), \quad p \in \mathbf{R}^d, \quad x \in \mathbf{R}^d,$$

of a relativistic spinless particle of mass $m > 0$ and charge e interacting with electromagnetic vector and scalar potentials $A(x)$ and $\Phi(x)$ (e.g. [5]). The Planck constant \hbar is taken to equal 1.

The present approach is a rigorous application of the phase space path integral or Hamiltonian path integral method with the "midpoint" prescription ([6], [7]). The path space measure used is a probability measure on the space of the right-continuous paths $X: [0, \infty) \rightarrow \mathbf{R}^d$ having the left-hand limits. Each path $X(s)$ is called a d -dimensional time homogeneous Lévy process ([3], [4]). The path integral formula obtained has a close analogy with the Feynman-Kac-Itô formula for the quantum Hamiltonian of a nonrelativistic spinless particle of the same mass and charge interacting with vector and scalar potentials (e.g. [8]).

2. Path integral representation. To formulate our result we need some notions from a time homogeneous Lévy process ([3], [4]). The path space measure which we are going to use is the probability measure $\lambda_{0,x}$ on the space $D_{0,x}([0, \infty) \rightarrow \mathbf{R}^d)$ of the right-continuous paths having the left-hand limits and satisfying $X(0) = x$ whose characteristic function is given by

$$(2.1) \quad \exp\{-t[(c^2 p^2 + m^2 c^4)^{1/2} - mc^2]\} = \int e^{ip(X(t) - X(0))} d\lambda_{0,x}(X).$$

The Lévy-Khinchin formula turns out to be

$$(2.2) \quad (c^2 p^2 + m^2 c^4)^{1/2} - mc^2 = - \int_{\mathbf{R}^d \setminus \{0\}} [e^{ip \cdot y} - 1 - ipy I_{\{|y| < 1\}}(y)] n(dy).$$

Here $n(dy)$ is the Lévy measure which is a σ -finite measure on $\mathbf{R}^d \setminus \{0\}$ satis-

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